

Charm Spectroscopy at *BABAR*

Denis Altenburg*

University of Dortmund, Germany

E-mail: d.altenburg@physik.tu-dresden.de

In this note we present results on charmed hadron decays recently obtained from data recorded with the *BABAR* detector at the Stanford Linear Accelerator Center (SLAC) PEP-II B-Factory.

1. A Dalitz plot analysis of the three body decay $D^0 \rightarrow \bar{K}^0 K^+ K^-$ has been performed and the branching ratio of $D^0 \rightarrow \bar{K}^0 K^+ K^-$ and $D^0 \rightarrow \bar{K}^0 \pi^+ \pi^-$ has been measured very accurately to be

$$BR = \frac{\Gamma(D^0 \rightarrow \bar{K}^0 K^+ K^-)}{\Gamma(D^0 \rightarrow \bar{K}^0 \pi^+ \pi^-)} = (15.8 \pm 0.1(stat.) \pm 0.5(syst.)) \times 10^{-2}$$

Moreover a partial wave analysis has been done in order to separate the S- and P-wave contributions. As a result the $a_0(980) \rightarrow KK$ lineshape has been extracted for the first time almost background free and we further found a very clean S-P-interference within the $\phi(1020)$ mass region.

2. The mass of the Λ_c^+ baryon has been obtained from the low- Q -value modes $\Lambda_c^+ \rightarrow \Sigma^0 K_s^0 K^+$ and $\Lambda_c^+ \rightarrow \Lambda K_s^0 K^+$ in order to minimize systematic uncertainties mainly due to detector effects. We determined

$$m_{\Lambda_c^+} = (2286.46 \pm 0.14) \text{ MeV}/c^2$$

3. A study of Ω_c^0 production in continuum and $\Upsilon(4S)$ events showed up a very clear first evidence for Ω_c^0 production in B meson decays.

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1. Introduction

Although the *BABAR* project is known as a *B* meson factory there is much more than *B* physics which can be done at this facility. The data sample recorded by the *BABAR* detector since October 1999 contains a large amount of charmed-hadron decays, which can be used to perform precision measurements on the charm sector.

2. The *BABAR* Detector

The following is a brief summary of the most important components for the 3 analyses presented here, which are the tracking devices and the particle identification system. A more complete overview of the detector is given elsewhere [1]. The interaction region is surrounded by a 5-layer double-sided silicon vertex tracker (SVT) followed by a 40-layer drift chamber (DCH), which is filled with a mixture of isobutane and helium. The tracking devices are embedded in a 1.5-T solenoid field. The SVT and DCH measure the energy loss dE/dx , which provides charged particle identification in the low momentum region ($p < 0.7 \text{ GeV}/c$). Higher energetic charged particles are identified by a Cherenkov detector referred as DIRC (**D**etector of **I**nternally **R**eflected **C**herenkov light), which was designed to measure the angle of Cherenkov photons emitted when a charged particle passes one of the bars of fused silica.

3. Dalitz plot analysis of $D^0 \rightarrow \bar{K}^0 K^+ K^-$

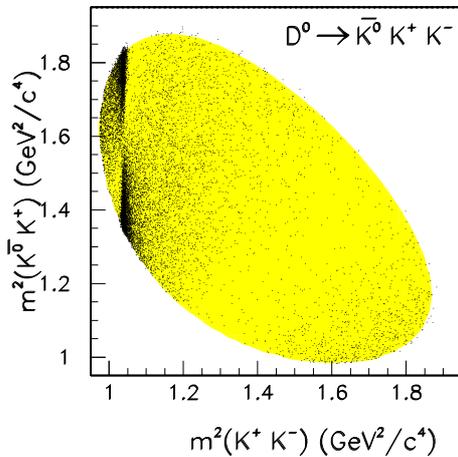


Figure 1: Dalitz plot of $D^0 \rightarrow \bar{K}^0 K^+ K^-$

$\bar{K}^0 K^+ K^-$ and $D^0 \rightarrow \bar{K}^0 \pi^+ \pi^-$ the Dalitz plane was divided into approximately 500 cells and under usage of the efficiency corrected yields we determined

$$BR = \frac{\Gamma(D^0 \rightarrow \bar{K}^0 K^+ K^-)}{\Gamma(D^0 \rightarrow \bar{K}^0 \pi^+ \pi^-)} = (15.8 \pm 0.1(\text{stat.}) \pm 0.5(\text{syst.})) \times 10^{-2}$$

A partial wave analysis has been done and we find a strong interference of the S- and P-wave amplitude contributions within the $\phi(1020)$ mass region. Under the assumption that the

P contribution is entirely due to the $\phi(1020)$ the $a_0(980) \rightarrow KK$ lineshape has been determined almost background free. Furthermore an unbinned maximum likelihood fit has been performed in order to determine relative phases and amplitudes of intermediate resonant and non-resonant states. Table 1 summarizes the fit results. The decay is dominated by $D^0 \rightarrow \bar{K}^0 a_0(980)^0$, $D^0 \rightarrow \bar{K}^0 \phi(1020)$ and $D^0 \rightarrow K^- a_0(980)^+$. The $f_0(980)$ and the doubly-cabbibo suppressed contribution ($D^0 \rightarrow K^+ a_0(980)^-$) are consistent with 0. The remaining contribution is non-uniform and could be interpreted by a tail of a broad resonance, e.g. $f_0(1400)$. Since the coupling of the $f_0(980)$ to the KK system is poorly known it requires more data to make statements about the properties of the $f_0(980)$ and the $a_0(980)$.

Final state	Amplitude	Phase (radians)	Fraction (%)
$\bar{K}^0 a_0(980)^0$	1.	0.	$66.4 \pm 1.6 \pm 7.0$
$\bar{K}^0 \phi(1020)$	$0.437 \pm 0.006 \pm 0.060$	$1.91 \pm 0.02 \pm 0.10$	$45.9 \pm 0.7 \pm 0.7$
$K^- a_0(980)^+$	$0.460 \pm 0.017 \pm 0.056$	$3.59 \pm 0.05 \pm 0.20$	$13.4 \pm 1.1 \pm 3.7$
$\bar{K}^0 f_0(1400)$	$0.435 \pm 0.033 \pm 0.162$	$-2.63 \pm 0.10 \pm 0.71$	$3.8 \pm 0.7 \pm 2.3$
$\bar{K}^0 f_0(980)$			$0.4 \pm 0.2 \pm 0.8$
$K^+ a_0(980)^-$			$0.8 \pm 0.3 \pm 0.8$
Sum			130.7 ± 2.2

Table 1: Results from the Dalitz plot analysis of $D^0 \rightarrow \bar{K}^0 K^+ K^-$

4. Precise measurement of the Λ_c^+ baryon mass

So far the best Λ_c^+ mass measurement was done by the CLEO collaboration using a sample of approximately 1000 events [5]. BABAR collected a very large sample of charmed-hadron decays and

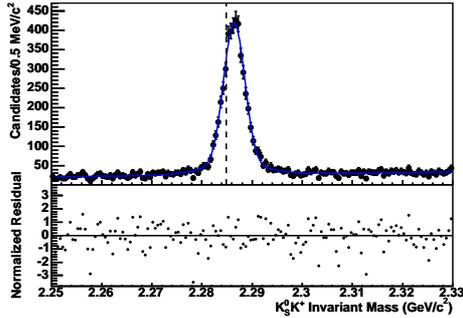


Figure 2: Invariant $\Lambda K_S^0 K^+$ mass. The lower part of the figure shows the normalized fit residuals. The dashed line indicates the present PDG value.

charged particle momenta are measured with very high precision. The challenge of such a precision mass measurement lies in the track reconstruction uncertainties, arising from energy-loss correction and the knowledge of the magnetic field. Such systematic uncertainties tend to increase with increasing Q , whereas Q denotes the 4-momentum transfer. Therefore we used low- Q -modes to reconstruct the Λ_c^+ , such as $\Lambda_c^+ \rightarrow \Lambda K_S^0 K^+$ and $\Lambda_c^+ \rightarrow \Sigma^0 K_S^0 K^+$.

Fig. 2 shows the invariant $\Lambda K_S^0 K^+$ mass fitted to a sum of two Gaussians with a common mean for signal and a linear function for combinatorial background. The signal yield was determined to be 4627 ± 84 and the fitted peak position is $(2286.44 \pm 0.04) \text{ GeV}/c^2$.

Four control samples ($\Lambda \rightarrow p\pi^-$, $K_S^0 \rightarrow \pi^+\pi^-$, $\Lambda_c^+ \rightarrow pK^-\pi^+$ and $\Lambda_c^+ \rightarrow pK_S^0$) have been used to study mass dependencies on the amount of material assumed for the energy-loss correction and the magnetic field. As a result it was found that the Λ_c mass needs to be corrected for

underestimated energy loss. The correction is calculated by increasing the SVT material by 20%. We find $m_{\Lambda_c^+}(\Lambda K_s^0 K^+) = (2286.501 \pm 0.042(stat.) \pm 0.144(syst.)) \text{ MeV}/c^2$ and $m_{\Lambda_c^+}(\Sigma^0 K_s^0 K^+) = (2286.303 \pm 0.181(stat.) \pm 0.126(syst.)) \text{ MeV}/c^2$. We combine these results to $m_{\Lambda_c^+} = (2286.46 \pm 0.14) \text{ MeV}/c^2$ which is in agreement with the present PDG value but much more precise.

5. Production and decays of Ω_c^0 baryons

Ω_c^0 baryons are reconstructed via their decay into $\Omega^- \pi^+$. The resulting invariant mass spectra have been fitted in 11 bins of the c.m. momentum $p^*(\Omega_c^0)$. In Fig. 3 the uncorrected momentum spectrum can be seen and it shows a very clear two peak structure. The peak above 1.5 GeV/c can be identified as Ω_c^0 production in continuum events (The blue bands illustrate the MC continuum expectation). Since the peak at low momenta is not seen in data taken below the $B\bar{B}$ threshold this peak can be interpreted as the first evidence for Ω_c^0 baryons produced in B decays.

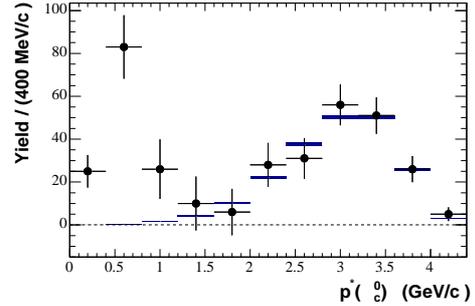


Figure 3: Ω_c^0 momentum spectrum.

6. Summary

In summary a Dalitz plot analysis of $D^0 \rightarrow \bar{K}^0 K^+ K^-$ has been performed which indicates that this decay is dominated by $D^0 \rightarrow \bar{K}^0 a_0(980)^0$, $D^0 \rightarrow \bar{K}^0 \phi(1020)$ and $D^0 \rightarrow K^- a_0(980)^+$. The Λ_c^+ mass has been measured with an improvement in precision of factor 4 compared to the PDG value using low- Q -modes. We have found a first clear evidence for Ω_c production in B meson decays.

7. Acknowledgements

We are grateful for the extraordinary contributions of our PEP-II colleagues in achieving the excellent luminosity and machine conditions that have made this work possible. The success of this project also relies critically on the expertise and dedication of the computing organizations that support BABAR. The collaborating institutions wish to thank SLAC for its support and the kind hospitality extended to them. This work is supported by the US Department of energy and national funding agencies.

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