# PoS

# Observation of the hadronic final state charge asymmetry in high $Q^2$ deep-inelastic scattering at HERA

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> A first measurement has been made of the charge asymmetry in the scattered hadronic final state from the hard interaction in high  $Q^2$  (100 <  $Q^2$  < 8,000 GeV<sup>2</sup>) deep-inelastic *ep* neutral current scattering at HERA. The difference between the event normalised distribution of the scaled momentum,  $x_p$ , for positively and negatively charged particles, measured in the current region of the Breit frame, has been studied together with its evolution as a function of Q. The results are compared to Monte Carlo models at the hadron and parton level.

XVIII International Workshop on Deep-Inelastic Scattering and Related Subjects April 19 -23, 2010 Convitto della Calza, Firenze, Italy

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### 1. Introduction

In lepton proton deep-inelastic scattering (DIS) at large Bjorken x the contribution of u valence quarks from the proton to the hard interaction dominates over that from the d valence quarks due to their larger charge and greater abundance. Hence an asymmetry in the number of positively and negatively charged particles is expected in the final state.

In recent papers, H1 has presented studies of the inclusive charged particle production in high  $Q^2$  deep-inelastic scattering (DIS) at HERA [1]. The measurement is performed in the current hemisphere of the Breit frame. In the naïve quark parton model (QPM) the momentum of the scattered parton in the Breit frame is Q/2, where  $Q^2$  is the virtuality of the exchanged boson. The main observable studied is  $x_p$ , the charged particle momentum in the current region of the Breit frame scaled to Q/2. Hadrons with small values of  $x_p$  are predominately produced by fragmentation, while hadrons at large  $x_p$  are more likely to contain a parton from the hard interaction. Therefore, a study of the  $x_p$  distribution,  $D(x_p, Q)$ , separately for positively and negatively charged particles should reveal information about the valence quarks and their fragmentation.

## 2. Data Selection and Correction

The data used in this analysis correspond to an integrated luminosity of 44  $pb^{-1}$  and were taken by H1 in the year 2000 when protons with an energy of 920 GeV collided with positrons with an energy of 27.5 GeV.

The kinematic phase space is defined as follows. The scattered positron is detected in the LAr calorimeter in the polar angular range  $10^{\circ} < \theta_e < 150^{\circ}$  and with energy greater than 11 GeV. The negative squared photon momentum is required to be in the range  $100 < Q^2 < 8000 \text{ GeV}^2$  and the inelasticity *y*, which is the fractional energy loss of the positron in the proton rest frame, to be in the range 0.05 < y < 0.6. The polar scattering angle for a massless parton, calculated from the positron kinematics in the QPM approximation, is required to be in the range  $30^{\circ} < \theta_{q,lab} < 150^{\circ}$ . This ensures that the current region of the Breit frame remains in the central region of the detector where there is high acceptance and track reconstruction efficiency. It should be noted that the kinematic phase space is defined solely from the scattered electron and can be applied in a simple way to theoretical models. Additional selections are made to reduce QED radiation effects and to suppress background events (photoproduction, beam gas and QED Compton ). The final event selection results in a data sample of about 60,000 events. Primary vertex fitted reconstructed charged tracks, which pass a variety of track quality cuts, are used to study the fragmentation process.

The data are corrected for detector acceptance, efficiency, resolution effects, and QED radiation. The total correction factor applied to the uncorrected data points is typically 1.0 - 1.2 for  $D(x_p, Q)$ . In the measurement of the charge asymmetry ratio contributions to the correction factor mostly cancel and as a result the correction factor is consistent with 1.0.

Several sources of systematic errors are considered: the scattered positron energy scale uncertainty and angular resolution, the hadronic energy scale uncertainty, the model uncertainty from the correction procedure, and the track reconstruction efficiency. All sources of error are treated as uncorrelated, apart from the positron energy scale uncertainty which is treated as fully correlated between bins. For  $D(x_p, Q)$  the largest contributions to the systematic error are the positron energy scale uncertainty, 0.5% ( $x_p \sim 0.1$ ) to 11% ( $x_p \sim 1.0$ ), and the track reconstruction uncertainty, 2.5%. For the charge asymmetry measurements the systematic errors mostly cancel and the only significant contribution is from the track reconstruction uncertainty, 2.5%.

#### 3. Phenomenology

The data are compared to predictions of the Parton Shower model (PS), as implemented in the RAPGAP Monte Carlo program, and to predictions of the Colour Dipole Model (CDM) both matched to  $O(\alpha_s)$  matrix elements. ARIADNE provides an implementation of CDM and is used in the DJANGO Monte Carlo program. Both the PS and CDM predictions use the Lund string model for hadronisation. The HERWIG Monte Carlo program uses the parton shower model to describe the fragmentation process but incorporates the cluster model of hadronisation. The data are also compared to predictions from the Soft Colour Interaction model (SCI) using the generalised area law (GAL) as implemented in LEPTO.

It is possible to turn off the hadronisation and compare data with parton level predictions using the assumption of local parton hadron duality. This has been done with the CDM predictions where a quark, with fractional charge, is taken as equivalent to a charged hadron of unit charge. The predictions are made after the main parton cascade has taken place and the gluons are ignored.

The CTEQ5L PDF is used for all model predictions. However other PDFs lead to charge asymmetries in agreement with the prediction based on CTEQ5L to within  $\pm$  0.01.

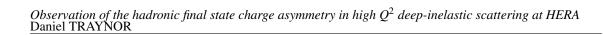
#### 4. Results

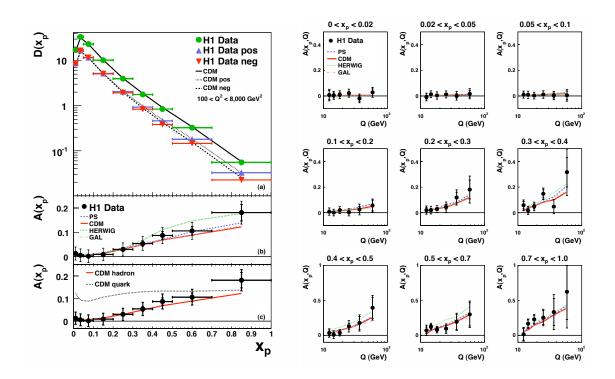
The scaled momentum distribution for all charged particles, and for positive, and negative particles separately, is shown in figure 1-Left (a). There are significantly more particles produced at low  $x_p$  than at high  $x_p$ . The scaled momentum distribution for positive and negative particles are very similar at low  $x_p$  but at high  $x_p$  there is a clear excess of positive particles.

The charge asymmetry can be as large as 20% as shown in figure 1-Left (b). The scaled momentum distribution and its asymmetry is described by the Monte Carlo models. Models using string hadronisation (PS, CDM, GAL) produce a smaller charge asymmetry at high  $x_p$  than that produced by the cluster hadronisation model (HERWIG).

In figure 1-Left (c) the charge asymmetry is compared to predictions from CDM before and after hadronisation. It is observed that at high  $x_p$  the hadron and quark levels are in good agreement and both agree with the data. However, as  $x_p$  get smaller a large difference develops with the quark level asymmetry prediction constant at ~ 12%, while the hadron level, and the data, fall to zero. This is consistent with the expectation that the hadrons at low  $x_p$  are dominantly produced by fragmentation while hadrons at high  $x_p$  retain the memory of the charge of the scattered quark from the hard interaction. It should be noted that sea quarks and gluons will produce, on average, a charge symmetric hadronic final states reducing the charge asymmetry expected by valence quarks alone.

Figures 1-Right shows the charge asymmetry as a function of Q in different  $x_p$  intervals. The charge asymmetry observed at large  $x_p$  evolves to larger values as Q increases. The largest asymmetries of up to 40% are obtained in the highest Q and highest  $x_p$  intervals. It should be noted





**Figure 1:** Left - (a) The measured normalised distributions of the scaled momentum,  $D(x_p)$ , for all charged particles and for positively (pos) and negatively (neg) charged particles separately, and (b), (c) the charge asymmetry, as a function of  $x_p$ . Right - The charge asymmetry, as a function of Q for nine different  $x_p$  regions. The data are displayed at the average value of Q.

that higher average Q corresponds to higher average x and hence the highest Q intervals are most sensitive to the valence quark distribution. The distributions are well described by the Monte Carlo models.

#### 5. Conclusions

The first measurement of the charge asymmetry of the hadronic final state at HERA is presented. The charge asymmetry is found to be dependent on the scaled momentum  $x_p$  with a larger asymmetry for larger  $x_p$ . The observed charge asymmetry at large  $x_p$  is found to increase with the scale Q corresponding at HERA to an enhancement at large Bjorken x. The results are consistent with the expectation that the asymmetry is directly related to the valence quark content of the proton. The observed charge asymmetry is reproduced by various models. The data are expected to provide useful information for the extraction of fragmentation functions and additional constraints on the valence quark distribution of the proton.

#### References

[1] F. D. Aaron *et al.* [H1 Collaboration], Phys. Lett. B **654** (2007) 148 [arXiv:0706.2456].
F. D. Aaron *et al.* [H1 Collaboration], Phys. Lett. B **681** (2009) 125 [arXiv:0907.2666].