

$B \rightarrow \tau \nu$ & $B \rightarrow D^{(*)} \tau \nu$ decays at Belle

Jacek Stypuła^{*†}

H. Niewodniczański Institute of Nuclear Physics Kraków

E-mail: Jacek.Stypula@ifj.edu.pl

We present measurements of the decays $B \rightarrow \tau \nu_\tau$ and $B \rightarrow D^{(*)} \tau \nu_\tau$ in a large data sample recorded with the Belle detector at the KEKB asymmetric energy e^+e^- collider. We obtain the branching fractions for these decays. The resulting constraints on a charged Higgs boson are also discussed.

35th International Conference of High Energy Physics

July 22-28, 2010

Paris, France

^{*}Speaker.

[†]On behalf of the Belle collaboration and supported by the Polish Ministry of Education and Science grant No N N202 287138.

1. Introduction

B meson decays with τ leptons in the final state, despite of experimental difficulties, are of great importance. In the Standard Model (SM) scenario measurement of the tauonic B decay can provide direct experimental determination of B meson decay constant, which can be compared to the lattice QCD calculations. Semitauonic B decays provide access to form-factors that cannot be measured in other semileptonic B decays. Due to large τ lepton mass both decay modes are sensitive to extended Higgs sector, $B \rightarrow \tau\nu_\tau$ through branching fraction (BF) effects while $B \rightarrow D^{(*)}\tau\nu_\tau$ mostly through other observables such as *e.g.* polarizations.

These analyses are based on data samples recorded at the $\Upsilon(4S)$ resonance with the Belle detector [1] at the KEKB collider [2].

1.1 Experimental techniques

A decay with 2 or 3 neutrinos can be observed using kinematic constraints available only at B-factories which are clean sources of exclusive $B\bar{B}$ pairs. To ensure that we have missing four-momentum consistent with multi neutrino hypothesis we take the advantage of the accompanying B meson referred to as B_{tag} . The B_{tag} can be reconstructed in several exclusive modes first and then checks whether remaining particles are consistent with the signal B (B_{sig}) decay can be done. We refer to this method as an “exclusive B_{tag} reconstruction”. The B_{tag} can be also reconstructed inclusively from all the particles that remain after selecting B_{sig} candidate. We refer to this method as an “inclusive B_{tag} reconstruction”. The analyses presented here exploit both mentioned approaches depending on the final state.

2. $B \rightarrow \tau\nu$

In the SM a leptonic B decay is a W -mediated annihilation with the decay rate simply related to B meson decay constant f_B and the quark-mixing amplitude V_{ub} :

$$\mathcal{B}(B^+ \rightarrow l^+ \nu_l) \Big|_{\text{SM}} = \frac{G_F^2 m_B}{8\pi} m_l^2 \left(1 - \frac{m_l^2}{m_B^2}\right)^2 f_B^2 |V_{ub}|^2 \tau_B, \quad (2.1)$$

where G_F is the Fermi constant, m_B and m_l are the B meson and lepton masses while τ_B is the B meson lifetime. The decay is helicity suppressed thus the τ mode is favoured in comparison with the light lepton modes. This decay is also sensitive to charged Higgs which would modify the branching. *E.g.* in the type-II two-Higgs doublet model (2HDM) we have

$$\mathcal{B} = \mathcal{B} \Big|_{\text{SM}} \times r_H, \quad r_H = \left(1 - \frac{m_B^2}{m_H^2} \tan^2 \beta\right)^2, \quad (2.2)$$

where m_H is the charged Higgs mass and $\tan \beta$ is the ratio of Higgs vacuum expectation values [3].

Belle previously reported the first evidence of $B^+ \rightarrow \tau^+ \nu_\tau$ ¹ decay with the “exclusive B_{tag} reconstruction” in hadronic modes [4]. The signal is extracted from a fit to the remaining energy in the electromagnetic calorimeter (E_{ECL}), which is the sum of the energies of neutral clusters

¹Charge conjugate modes are implied throughout this report unless otherwise stated.

	$\mathcal{B}(B^+ \rightarrow \tau^+ \nu_\tau)$
Belle hadronic tag	$[1.79^{+0.56}(\text{stat})^{+0.46}(\text{syst})] \times 10^{-4}$
Belle semileptonic tag	$[1.54^{+0.38}(\text{stat})^{+0.29}(\text{syst})] \times 10^{-4}$
SM – Eq. (2.1)	$[1.2 \pm 0.25] \times 10^{-4}$
SM – CKM fitter	$[0.763^{+0.113}_{-0.061}] \times 10^{-4}$

Table 1: Summary of $B \rightarrow \tau\nu$ results at Belle with a comparison to SM predictions from Eq. (2.1) with $|V_{ub}|$ taken from [6], f_B taken from [7] and from CKM fitter results with $B \rightarrow \tau\nu$ [8].

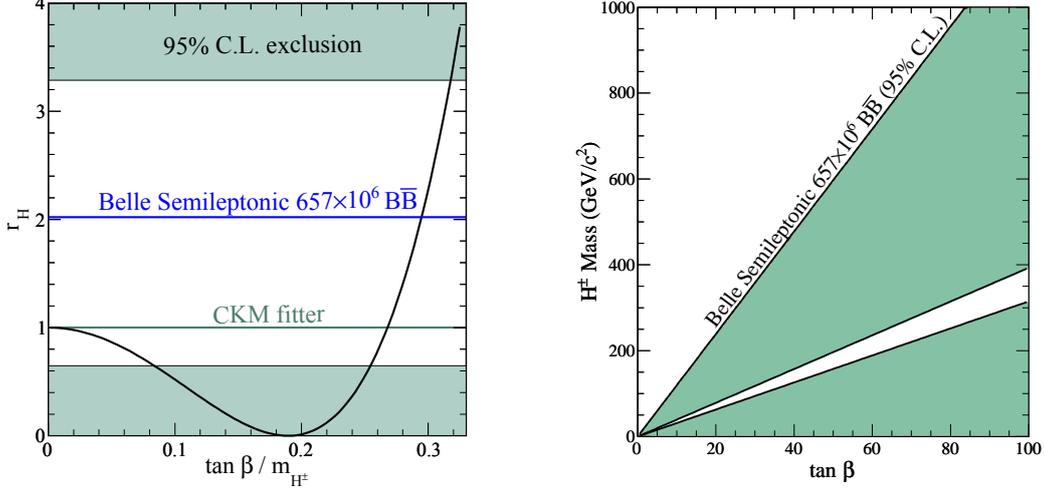


Figure 1: Constraints on the type-II two-Higgs doublet model placed by Belle semileptonic tag result.

that are not associated with either the B_{tag} or the π^0 candidate from the $\tau^+ \rightarrow \pi^+ \pi^0 \bar{\nu}_\tau$ decay. For signal events, E_{ECL} must be either zero or a small value arising from beam background hits while background events are distributed toward higher E_{ECL} due to the contribution from additional neutral clusters. To better establish this decay mode and determine the BF with greater precision we have done a measurement where B_{tag} is reconstructed in semileptonic modes [5]. This result is consistent with the previous one and with the SM predictions (Table 1). Figure 1 shows the constraints placed by our recent result on the 2HDM.

3. $B \rightarrow D^{(*)}\tau\nu$

Semitauonic B decays are complementary to and competitive with tauonic ones due to different theoretical uncertainties and more observables. In $B \rightarrow D^{(*)}\tau\nu$ we are free from f_B and $|V_{ub}|$ which have large uncertainties. We have a dependence on the formfactors and $|V_{cb}|$ instead however the latter cancels out in the ratio

$$R = \frac{\mathcal{B}(B \rightarrow D\tau\nu)}{\mathcal{B}(B \rightarrow D\nu_l)} \tag{3.1}$$

The exclusive semitauonic decay was first observed by Belle in the $B^0 \rightarrow D^{*-}\tau^+\nu_\tau$ mode [9] using “inclusive B_{tag} reconstruction”. At large missing masses most of background components behave combinatorial in the beam constrained mass $M_{\text{tag}} = \sqrt{E_{\text{beam}}^2 - p_{\text{tag}}^2}$, where E_{beam} is the

Decay mode	BF&significance from inclusive tag	BF&significance from exclusive tag
$B^+ \rightarrow \bar{D}^{*0}\tau^+\nu_\tau$	$[2.12^{+0.28}_{-0.27} \pm 0.29]\%$ 8.1σ	$[3.04^{+0.69+0.40}_{-0.66-0.47}]\%$ 3.9σ
$B^0 \rightarrow D^{*-}\tau^+\nu_\tau$	$[2.02^{+0.40}_{-0.37} \pm 0.37]\%$ 5.2σ	$[2.56^{+0.75+0.31}_{-0.66-0.22}]\%$ 4.7σ
$B^+ \rightarrow \bar{D}^0\tau^+\nu_\tau$	$[0.77 \pm 0.22 \pm 0.12]\%$ 3.5σ	$[1.51^{+0.41+0.24}_{-0.39-0.19}]\%$ 3.8σ
$B^0 \rightarrow D^-\tau^+\nu_\tau$	$[1.01^{+0.46+0.13}_{-0.41-0.11}]\%$ 2.6σ	—

Table 2: Summary of $B \rightarrow D^{(*)}\tau\nu$ results at Belle along with statistical and systematical uncertainties.

beam energy and p_{tag} is the momentum of B_{tag} candidate (residual particles). On the other hand the signal is visible as a well reconstructed B_{tag} and thus was extracted from a fit to M_{tag} . Recently we have done a next-step analysis which includes simultaneous extraction of signals in charged B decays to $D^*\tau\nu$ and to $D\tau\nu$ taking into account D^*D cross-feeds [10]. All the results including preliminary one using “exclusive B_{tag} reconstruction”[11] are summarized in Table 2.

These results are consistent within experimental uncertainties with SM [12] and provide constraints on 2HDM complementary to those from purely tauonic B decays [13].

4. Summary

The studies of B decays to τ at Belle brought significant advances in this field, providing the first evidence of the purely leptonic $B^+ \rightarrow \tau^+\nu_\tau$ mode, semi-tauonic $B \rightarrow D\tau^+\nu_\tau$ modes and the observation of semi-tauonic B decays in the $B \rightarrow D^*\tau^+\nu_\tau$ channels. These results are consistent with the SM but, given the uncertainties, there is still a room for a sizeable non-SM contribution. Belle II experiment on SuperKEKB Super B-factory with ≈ 50 times higher statistics should measure these modes with much higher precision.

References

- [1] A. Abashian *et al.* (Belle Collaboration), *Nucl. Instr. and Meth A* **479** 117 (2002)
- [2] S. Kurokawa, E. Kikutani, *Nucl. Instr. and Meth A* **499** 1 (2003) and references therein
- [3] W. S. Hou, *Phys. Rev. D* **48**, 2342 (1993)
- [4] K. Ikado *et al.* (Belle Collaboration), *Phys. Rev. Lett.* **97**, 251802 (2006) [hep-ex/0604018v3]
- [5] K. Hara *et al.* (Belle Collaboration), *Phys. Rev. D* **82**, 071101(R) (2010) [hep-ex/1006.4201v2]
- [6] <http://www.slac.stanford.edu/xorg/hfag/semi/ichep08/>
- [7] HPQCD Collaboration, *Phys. Rev. D* **80**, 014503 (2009) [hep-lat/0902.1815v3]
- [8] http://ckmfitter.in2p3.fr/plots_FPCP10/
- [9] A. Matyja *et al.* (Belle Collaboration), *Phys. Rev. Lett.* **99**, 191807 (2007) [hep-ex/0706.4429v2]
- [10] A. Bozek *et al.* (Belle Collaboration), *Phys. Rev. D* **82**, 072005 (2010) [hep-ex/1005.2302v1]
- [11] I. Adachi *et al.* (Belle Collaboration), hep-ex/0910.4301v1
- [12] C.-H. Chen, C.-Q. Geng, *JHEP* **0610**, 053 (2006) [hep-ph/0608166v3]
- [13] M. Tanaka, R. Watanabe, hep-ph/1005.4306v3