

Exotic Searches in ATLAS

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The latest searches for (non-SUSY) Beyond Standard Model phenomena performed with the ATLAS detector are reported. The searches have been performed with data from proton-proton collisions at a centre-of-mass energy of 7 TeV collected in 2010 and 2011. Various experimental signatures have been studied involving the reconstruction and measurement of leptons, photons, jets, missing transverse energy, as well as the reconstruction of top quarks. For most of the signatures, the experimental reach is significantly increased with respect to previous results.

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1. Introduction

The Standard Model of elementary particles (SM) has had enormous phenomenological success during the past half century. However, it is generally believed that it is a low energy limit of a more general theory, and numerous theoretical scenarios are studied in order to address at least some of the open issues such as the unification of forces, existence of dark matter or hierarchy problem. Many of these models indicate that new phenomena are likely to appear at the TeV scale, the direct probe of which is possible with the Large Hadron Collider (LHC) at CERN. Supersymmetry (SUSY) is most commonly invoked to address the opened questions of the SM, but there are a large number of important and well-motivated theoretical models that seek to answer some or all of fundamental questions. These, non-SUSY, phenomena are often collectively referred to as "Exotic" physics. The ATLAS experiment [1] explores TeV-scale territory performing a large number of searches for BSM physics. A comprehensive review is beyond the scope of this document, but for the up to date list of published ATLAS results, please consult official ATLAS Exotic page [2]. The emphasis in this conference report is given on the recent results.

This document is organized as follows. After a brief description of data used by ATLAS, searches for two-object resonances reconstructed using leptons, jets, photons and missing energy are described. This is followed by a review of the ATLAS searches for Extra Dimensions, including searches for TeV-scale gravity (a.k.a. black holes). After that, a brief description of some other exotic searches, like contact interactions, the search for right-handed W bosons and Majorana neutrinos, leptoquarks, and exotic top partners are described. It should be emphasized that searches for new physics are typically signature based, although most analyses do include interpretations in one or more models. A full review of all possible experimental signatures is also beyond the scope of this report.

2. The LHC, ATLAS Detector and Data Samples

The searches presented in this report have been performed with the data from proton-proton collisions at a centre-of-mass energy of 7 TeV collected in 2010 and 2011 with the ATLAS detector at the LHC. Excellent performance of the LHC in 2010, leading to 45 pb⁻¹ of data recorded by ATLAS, continued also in 2011. The data taking started with luminosities similar to those at the end of 2010, with instantaneous luminosity reaching 3.65×10^{33} cm⁻²s⁻¹, and on the "best day" 122 pb⁻¹ was recorded. The total integrated luminosity delivered to ATLAS was more than 5 fb⁻¹, five times larger than previously targeted for 2011. ATLAS data taking efficiency was 93.5%, and fraction of good quality data 90-100% depending on the detector subsystem. The most striking change in beam conditions between 2010 and 2011 data-taking has been increase in number of interactions per bunch crossing at a given point in an LHC fill. With up to 34 collisions per bunch crossing, multiple *pp* interactions ("pileup") in both the triggered beam-crossing ("in-time") and ones nearby in time ("out-of-time") posed a significant experimental challenge and much effort on understanding its impact on detector performance with data and simulation was needed. The performance of the ATLAS detector, in terms of the lepton reconstruction and identification efficiency, lepton and jet energy resolutions and scales, as well as the performance of the measurement of the missing transverse energy (E_T^{miss}) have been established from the data itself. The details of the

ATLAS performance are presented at this conference [3]. The searches presented in this report are based on data samples corresponding to integrated luminosities of up to $\sim 2 \text{ fb}^{-1}$.

3. Searches for Heavy Resonances

The existence of heavy two-object resonances are predicted by a variety of BSM models. The analysis basically consists of search for a bump in the background mass spectra which is either small or could be parametrized from the data control samples. Relatively simple topologies allow for a robust analysis to be done rapidly. ATLAS Collaboration has searched for dilepton (including lepton+neutrino), diphoton, dijet and jet+photon resonances. Diphoton searches are described in the section about Extra Dimensions.

3.1 Dilepton resonances

Here, a high mass resonance decaying to a pair of isolated leptons e^+e^- , or $\mu^+\mu^-$ is searched for [4]. Such resonances include spin-1 heavy neutral gauge bosons such as sequential SM-like Z' [5], as well as spin-2 Randall-Sundrum gravitons G^* [6]. The recent ATLAS results is based on the analysis of $1.1\text{--}1.2 \text{ fb}^{-1}$ of data. Good lepton reconstruction and identification with high efficiency and resolution under control is crucial for the searches for high mass dilepton (and especially lepton+neutrino) resonances. For example identification and quality criteria, as well as calorimeter isolation requirements, are imposed to suppress background from jets faking electrons. To optimize the muon momentum resolution, each muon candidate is required to pass quality cuts in the Inner Detector (ID) and to have at least three hits in each of the inner, middle, and outer layers of the Muon Spectrometer (MS). The effects of misalignments and intrinsic position resolution are included in the simulation. The muon p_T resolution at 1 TeV ranges from 15% (central) to 44% (for $|\eta| > 2$). The background is dominated by the irreducible background due to the Z/γ^* (Drell-Yan) process, characterized by the same final state as the signal. Small contributions come from $t\bar{t}$ and diboson production. QCD background consist of semi-leptonic decays of b and c quarks in the $\mu^+\mu^-$ channel, and a mixture of photon conversions, semi-leptonic heavy quark decays, and hadrons faking electrons in the e^+e^- channel. Jets accompanying W bosons may similarly produce lepton candidates. Backgrounds from cosmic rays are negligible. Level of background, except for QCD background, are initially estimated from Monte Carlo, and after that simulated backgrounds are rescaled so that the sum of the background matches the observed number of events in data around the Z peak. The advantage is that the uncertainty on the luminosity, and any mass independent uncertainty on efficiency cancel between Z' and the Z background. The dielectron mass distribution is shown in Fig. 1 (top). No significant deviation from SM expectation is observed in either e^+e^- or $\mu^+\mu^-$ invariant mass spectra, and limits¹ are set on cross section times branching ratio (σB) for a number of Z' models, Fig. 1 (center). Observed lower mass limit on SM-like Z' is 1.83 TeV from the combination of both channels and now supersedes Tevatron limit. Observed mass limit on G^* is 1.63 TeV, and also suppresses the current Tevatron result. It is worth noting that the observed upper limits on σB in the wide range of Z' masses between 0.2 TeV to 1.5 TeV are lower than the Tevatron limits, reflecting increase in new physics sensitivity at the LHC with respect to the Tevatron.

¹ All limits in this document are at 95% confidence level (CL).

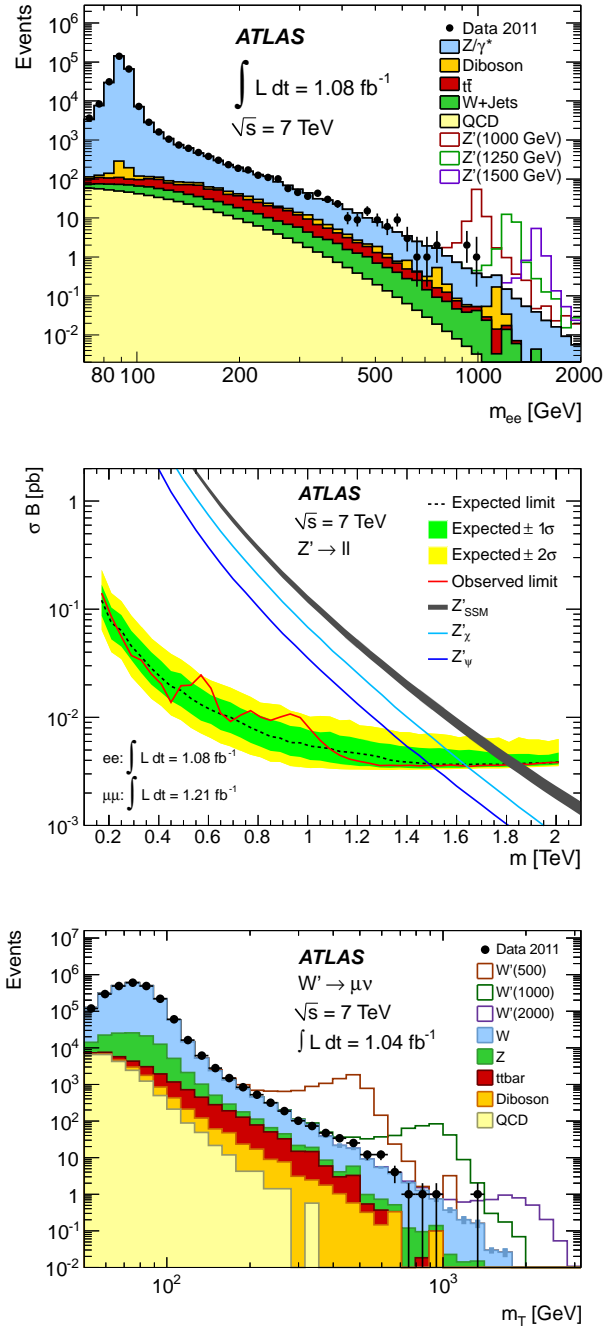


Figure 1: Top: Dielectron invariant mass distribution, compared to the expected backgrounds, with three example SSM Z' signals overlaid. Center: Expected and observed limits on σB and expected σB for SSM Z' production and the two E_6 -motivated Z' models for the combination of e^+e^- and $\mu^+\mu^-$ channel. Bottom: Spectrum of m_T for the muon channel in the W' search.

Another example of search for dilepton resonances, is the search for technivector mesons ρ_T and ω_T [7]. Technicolor (TC) models predict new technihadron states that could be copiously

produced at the LHC. The lowest mass states are the scalar technipion $\pi^{\pm,0}$ and the vector technirho $\rho^{\pm,0}$ and techniomega ω^0 . Additional motivation for this search stems from the observation by CDF of an excess in the dijet mass spectrum in Wjj events, which has led to speculation that the excess may be due to a 290 GeV ρ_T that decays into a W and a 160 GeV π_T . Since the neutral technimesons are narrow, spin 1 resonances, the search methodology is identical to that developed for the $Z' \rightarrow l^+l^-$ search. TC mesons are excluded for ρ_T masses 130 - 480 GeV and for π_T masses 50-480 GeV.

3.2 Lepton+neutrino resonances

As for heavy neutral Z' , many models predict additional heavy charged gauge bosons, commonly denoted W' . The experimental signature consist of exactly one isolated lepton (electron or muon), and high E_T^{miss} arising from neutrino escaping detection. Since the longitudinal component of the neutrino's momentum cannot be unambitiously reconstructed, transverse mass variable defined as $m_T = \sqrt{p_T E_T^{miss} (1 - \cos\phi_{lv})}$ is used as a discriminant variable. The most recent ATLAS results are based on the analysis of pp collisions corresponding to an integrated luminosity of 1.04 fb^{-1} [8]. The main background to the $W' \rightarrow lv$ signal comes from the high- m_T tail of SM W boson. This, as well as $Z \rightarrow ll$, top and diboson backgrounds are estimated using Monte Carlo simulation, while QCD and cosmic backgrounds are estimated using data driven techniques. Transverse mass in the muon channel is shown in Fig. 1 (top). No excess beyond Standard Model expectations is observed. W' with SM-like couplings is excluded at the 95% CL for masses up to 2.15 TeV using both $W' \rightarrow ev$ and $W' \rightarrow \mu\nu$ channels. The limit on σB suppress the Tevatron limits for $m_{W'} > 600 \text{ GeV}$. The obtained lower mass limit represents an increase of about 50% (approximately 700 GeV) with respect to the mass limit obtained by analyzing 36 pb^{-1} of 2010 data [9].

3.3 Same sign dimuon resonances

Production of prompt same sign muon pairs is predicted by several BSM theories (SUSY, Universal Extra Dimensions, heavy Majorana neutrinos, fourth-generation quarks and doubly charged Higgs bosons). This production, however, occurs comparatively rarely in the Standard Model. Here the ATLAS search in the context of the doubly charged Higgs boson ($H^{\pm\pm} \rightarrow \mu^{\pm}\mu^{\pm}$) is presented [10]. The fiducial region is defined as two like-sign muons with $|\eta| < 2.5$ and $p_T > 20$ (10) GeV for the leading (subleading) muon that are outside hadronic jets, and have an invariant mass $m_{\mu\mu} > 15 \text{ GeV}$. This fiducial region corresponds well to the experimental requirements made in the event selection. In order to suppress background from non-prompt muon production and cosmic muons, the transverse and longitudinal impact parameters (measured with respect to the event primary vertex) of the muon candidates must be small. Muons are required to be isolated from other tracks in ID, and sign of the muon track is required to be the same in ID and MS. Background is composed of the events coming from diboson production (WZ, ZZ) and non-prompt muon production (semi-leptonic b and c decays, π/K decays, cosmic rays). Non-prompt muon background is estimated using several data-driven techniques, and good agreement between data and background estimation is obtained, as demonstrated in Fig. 2 (left). Lower mass limits of 375(295) GeV are set on left-handed (right-handed) doubly charged Higgs production assuming the branching ratio to muons is 100%. These are most stringent limit to date. Limits are calculated

in the mass - branching ratio plane as well. Further, model independent upper limits on the cross section of new physics contributing to like-sign dimuon production are placed as function of the lower dimuon mass cut assuming the same acceptance: the limit ranges from 304 fb for $m_{\mu\mu} > 15$ GeV to 5.0 fb for $m_{\mu\mu} > 300$ GeV. The limits obtained represent a significant improvement compared to those obtained from the analysis based on the 2010.

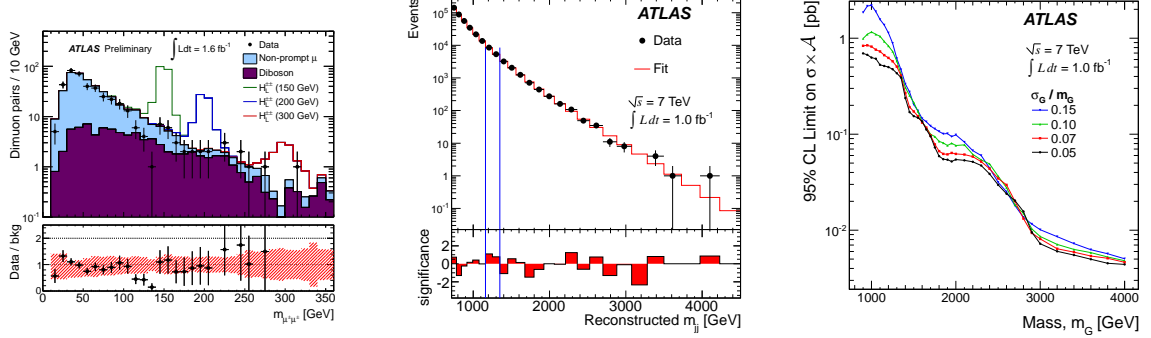


Figure 2: Left: Invariant mass distribution for $\mu^\pm\mu^\pm$ and for hypothetical left-handed $H^{\pm\pm}$. Center: The dijet mass distribution fitted with a smooth functional form describing the QCD background. Right: The upper limits on $\sigma \times A$ for a dijet Gaussian resonance m_G , with four values of σ_G/m_G .

3.4 Dijet resonances

Another search is the hunt for a bump in the dijet mass spectrum. The main observable is m_{jj} , the mass of the system of the two leading jets in p_T . Studies have shown that the dijet mass distribution are sensitive to the highest mass scales accessible in hadronic final states. Dijet mass distribution is examined in a search for resonances due to new phenomena localized near a given mass, employing a data driven background estimation that does not rely on detailed QCD calculations. ATLAS Collaboration has recently updated the search for new physics in the dijet mass distribution using 1 fb⁻¹ of 2011 data [11]. Jets are reconstructed using the anti- k_t algorithm, with the distance parameter $R=0.6$, with momentum calibration based on MC studies including full detector simulation, and validated with extensive test-beam and collision data studies. Measured dijet mass distributions are not corrected for detector resolution, which, in terms of mass smearing, is $\frac{\sigma_{m_{jj}}}{m_{jj}}=5\%$ at $m_{jj}=1$ TeV, dropping to 4.5% at 2 TeV, and asymptotically approaching 4% at m_{jj} of 5 TeV and above. For final event selection, additional kinematic criteria are applied, requiring that the two leading jets have $|\eta_j| < 2.8$, and the rapidity in the parton CM frame satisfies $|y^*| < 0.6$, where y is the jet rapidity defined as $y \equiv \frac{1}{2} \ln \frac{E+p_z}{E-p_z}$ with E being energy, and p_z z -component of the jet's momentum, and $y^* = \frac{1}{2}(y_1 - y_2)$. These criteria favour central collisions and have been shown, based on studies of expected signals and QCD background, to optimise the analysis sensitivity. The observed dijet mass distribution after all selection cuts is shown in Fig.2 (center), with highest entry > 4 TeV. m_{jj} spectrum is fitted to the smooth functional form which has been empirically shown to fit well data as well as the QCD prediction given by several MC programs. Limits are derived on cross section times acceptance ($\sigma \times A$) on several new physics models, among which excited quarks are excluded below masses of 2.99 TeV. Fig. 2 (right) shows model-independent limits on

$\sigma \times A$ for a collection of hypothetical signals that are assumed to be Gaussian-distributed in m_{jj} with mean (m_G) ranging from 0.9 to 4.0 TeV and standard deviation (σ_G) from 5% to 15% of the mean.

3.5 γ + jet resonances

Analogue to the dijet resonance search, a search for a resonance consisting of a back-to-back jet and a isolated photon is performed with the ATLAS detector [12]. The photon-jet invariant mass ($m_{\gamma j}$) distribution is used as a main discriminant variable. Background resulting from the mixture of direct photon-jet production, from radiation of final-state quarks or from multijet processes where dijet or higher-order events produce secondary photons during fragmentation of the hard- scatter quarks and gluons, is estimated by an empirical fit in the control region. Limits are set on $\sigma B \times A$ for several exotic models, excluding excited-quark model for masses up to 2.46 TeV. Limits are also set on a generic signal with Gaussian distribution and arbitrary production cross section, assuming the signal lineshape to be a Gaussian distribution with one of three widths, $\frac{\sigma_G}{m_G}=5\%, 7\%, \text{ and } 10\%$.

4. Searches for Extra Dimensions

In order to address hierarchy problem theories with extra spatial dimensions have been developed. Also, the complete unification of particle forces may require the existence of the additional spacial dimensions, three of which are within of our senses, and others being compactified at distance of the order of 10^{-32} m. These scenarios provided exciting new ground for experimental probes, and some of the recent searches performed with ATLAS are reported in this section.

4.1 Graviton dilepton and diphoton searches

The Randall-Sundrum (RS) model of extra dimensions predicts excited Kaluza-Klein (KK) modes of the graviton, which appear as spin-2 resonances. These modes have a narrow intrinsic width when $k/\overline{M}_{Pl} < 0.1$, where k is the spacetime curvature in the extra dimension, and $\overline{M}_{Pl} = M_{Pl}/\sqrt{8\pi}$ is the reduced Planck scale. KK gravitons in this model would have a mass splitting of order 1 TeV and would appear as new resonances. The phenomenology can be described in terms of the mass of the lightest KK graviton excitation (m_G) and the dimensionless coupling to the SM fields, k/\overline{M}_{Pl} . It is theoretically preferred for k/\overline{M}_{Pl} to have a value in the range from 0.01 to 0.1. The recent ATLAS search for dilepton final states [4] is described earlier. The diphoton final state provides a sensitive channel for the search for extra dimensions due to the clean experimental signature, excellent diphoton mass resolution, and modest backgrounds, as well as a branching ratio for graviton decay to diphotons that is twice the value of that for graviton decay to any individual charged-lepton pair. Another model that could be probed with the diphoton final states, is the Arkani-Hamed, Dimopoulos and Dvali (ADD) model [13]. The ADD model postulates the existence of flat additional spatial dimensions compactified with radius R , in which only gravity propagates. The mass splitting between subsequent KK states is of order $1/R$. In this model, resolving the hierarchy problem requires typically small values of $1/R$, giving rise to an almost continuous spectrum of KK graviton states. Effects due to ADD graviton exchange would be evidenced by a non-resonant deviation from the SM background expectation.

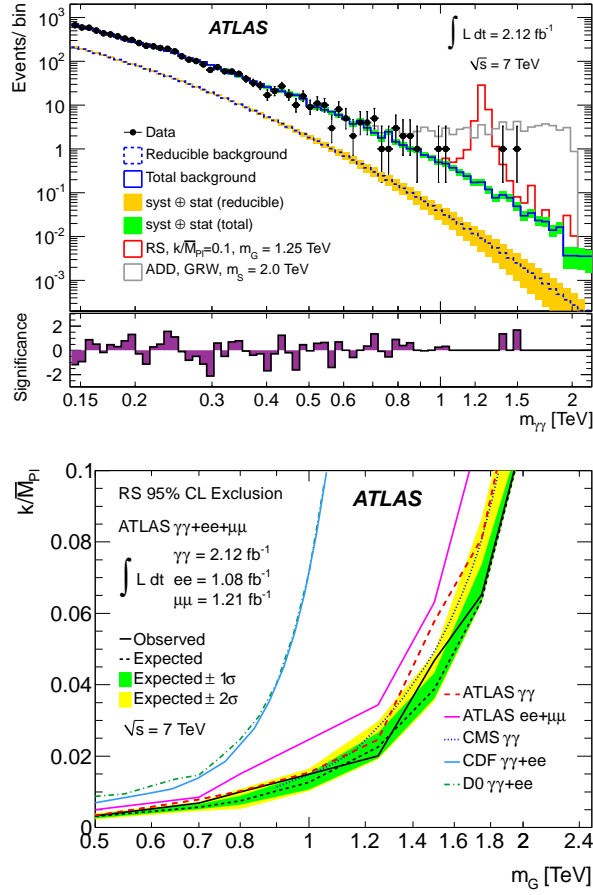


Figure 3: Top: The observed $m_{\gamma\gamma}$ distribution of diphoton events, superimposed with the predicted SM background and expected signals for ADD and RS models with indicated choices of parameters. Bottom: The RS limit interpreted in the plane of coupling versus graviton mass.

ATLAS searches for extra dimensions with $\gamma\gamma$ final states performed with 2.12 fb^{-1} of data are described in [14]. The largest background for this analysis is the irreducible background due to SM $\gamma\gamma$ production. The shape of the diphoton invariant mass spectrum from this background was estimated using MC, reweighting samples to the differential cross section predictions of state-of-the-art higher order generator. Another significant background component is the reducible background that includes events in which one or both of the reconstructed photon candidates result from a different physics object being misidentified as a photon. This background is dominated by γj and $j j$ events, with one or two jets faking photons, respectively. This background is estimated using data-driven techniques. Backgrounds with electrons faking photons, such as the Drell-Yan production of electron-positron pairs as well as $W/Z + \gamma$ and top processes, were verified to be small.

The resulting $\gamma\gamma$ spectrum is presented in Fig. 3 (top). No signal is observed, and upper limits are determined on the RS and ADD signal cross sections. The RS model results are combined with the limits obtained from the search with dilepton states, Fig. 3 (bottom). The mass of the lightest

RS graviton is excluded below the value of 1.95 TeV for $k/\overline{M}_{Pl}=0.1$. The observed upper limit on ADD model is 2.49 fb for the product of the cross section due to new physics multiplied by the acceptance and efficiency.

4.2 Universal Extra Dimensions: $\gamma\gamma + E_T^{miss}$ signature

This analysis considers the case of a single Universal Extra Dimension (UED) [15], with compactification radius $R \approx 1 \text{ TeV}^{-1}$. At the LHC, the main UED process would be the production via the strong interaction of a pair of first-level KK quarks and/or gluons. These would decay via cascades involving other KK particles until reaching the lightest KK particle (LKP), i.e. the first level KK photon γ^* decaying via $\gamma^* \rightarrow \gamma + G$, where G is a tower of eV-spaced graviton states, leading to a graviton mass between 0 and $1/R$. With two decay chains per event, the final state would contain $\gamma\gamma + E_T^{miss}$, where E_T^{miss} results from the escaping gravitons. A search for UED with this event topology has been performed using 1.07 fb^{-1} [16]. No excess of events is observed above the Standard Model prediction and upper limits are set on the production cross section for new physics, leading to the lower limit of 1.23 TeV on $1/R$. The results are also interpreted in the context of a generalised model of gauge-mediated supersymmetry breaking with a bino-like lightest neutralino [17], and in the context of a minimal model of gauge-mediated supersymmetry breaking [18].

4.3 Search for monojets

In the ADD Large Extra Dimensions (LED) model, the four-dimensional Planck scale, M_{Pl} , is related to the fundamental $4+n$ -dimensional Planck scale, M_D , by $M_{Pl}^2 \sim M_D^{2+n} R^n$, where n and R are the number and compactification radius of the extra dimensions, respectively. An appropriate choice of R for a given n allows for a value of M_D close to the electroweak scale. The compactification of the extra spatial dimensions results in a KK tower of massive graviton modes. These graviton modes are produced in association with a jet and do not interact with the detector, which results in a monojet signature in the final state. ATLAS has published an update of the previous analysis using a much larger 2011 data sample corresponding to 1.0 fb^{-1} [19]. Event topology consists of a high energy jet, and large E_T^{miss} . Three signal regions with increasing cuts on the jet $p_{T,j}$ and E_T^{miss} are defined i) 'low p_T ' $p_{T,j} > 120 \text{ GeV}$ and $E_T^{miss} > 220 \text{ GeV}$, ii) 'high p_T ' $p_{T,j} > 250 \text{ GeV}$ and $E_T^{miss} > 250 \text{ GeV}$, iii) 'very high p_T ' $p_{T,j} > 350 \text{ GeV}$ and $E_T^{miss} > 300 \text{ GeV}$. The aim is to maintain the sensitivity to a variety of models for new phenomena. Jets are required to be central $|\eta_j| < 2$, and events are vetoed if there is another jet or lepton present in the event. The expected background to the monojet signature is dominated by $Z \rightarrow \nu\bar{\nu} + \text{jets}$ and $W + \text{jets}$ production, and is estimated from data itself. The agreement between the data and the SM predictions for the total number of events in the different analyses is translated into model-independent upper limits on the $\sigma \times A$. The limits are set also in the context of the ADD model: for number of extra dimensions $n=2-6$, $M_D > 3.2 \text{ TeV} - 2.0 \text{ TeV}$.

4.4 Searches for TeV-gravity signature in final states with leptons and jets

Low-scale gravity models allow the existence of non-perturbative gravitational states such as black holes. In ADD model black holes can form if $\sqrt{s} > M_D$, and collision impact parameter

is smaller than higher dimensional Schwarzschild radius. Microscopic black holes decaying via Hawking radiation with isotropic and "democratic" decays leading to large multiplicity of particles. Searches for these states have previously been performed by investigating final states with multiple high- p_T objects with high- p_T jets only. Depending on the model 15-50% of events with black holes produced should have at least 1 charged lepton. Requiring the presence of the lepton dramatically reduces background and systematic uncertainties. ATLAS search for black holes in jets+lepton final state is presented here [20]. Scalar sum of p_T of high- p_T leptons and jets in the event is used as discriminant variable as shown in Fig. 4. Model independent limits (using true fiducial region separately defined for the electron and muon channel) are estimating total acceptance of 60-90% (40-60%) for electron (muon) channel, excluding new physics cross section 169 fb - 8.7 fb (77 fb - 4.8 fb) depending on the cut on Σp_T . As an illustration, mass limits within the ADD model are quoted: $M_D > 1.0$ TeV for $M_{Th} > 4.6$ TeV, or $M_D > 2.5$ TeV for $M_{Th} > 4.0$ TeV, where M_{Th} is the threshold for the creation of the black hole.

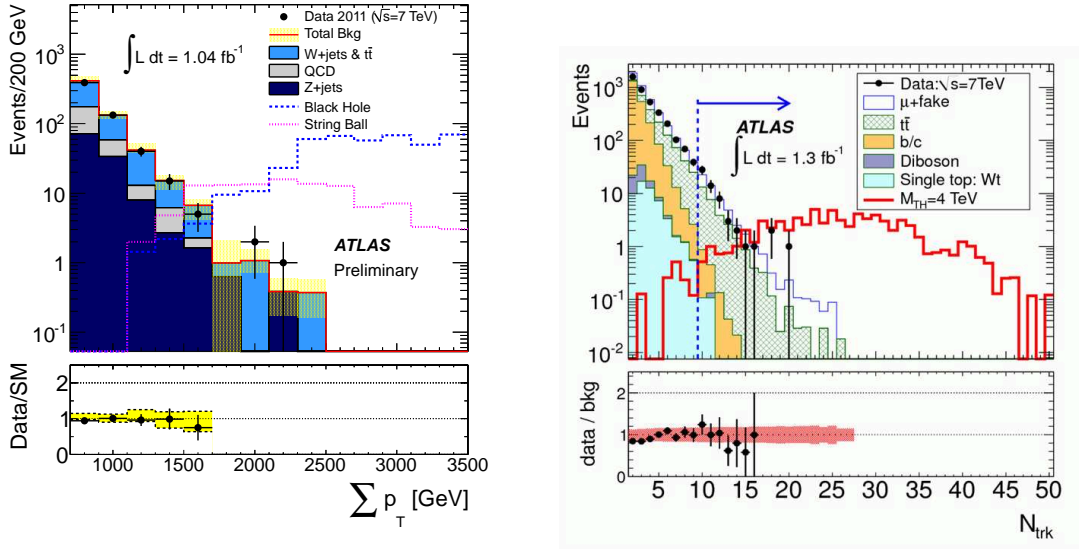


Figure 4: Left: Σp_T distributions for the signal region in the electron channel of the lepton+jets black hole search. Right: The track multiplicity distribution for same-sign dimuon events.

4.5 Search for Strong Gravity Signatures in Same-sign Dimuon Final States

In this analysis [21], events are selected containing two muons of the same charge. This channel is expected to have low Standard Model backgrounds while retaining good signal acceptance. In order to maintain optimal acceptance for a possible signal, only one of the muons is required to be isolated in this analysis, thereby typically increasing the acceptance in the signal region by 50%. Black hole events typically have a high number of tracks per event (N_{trk}), while Standard Model processes have sharply falling track multiplicity distributions. So N_{trk} is used as a discriminant variable, and the background is estimated from the control region. No excess of events over the SM prediction is observed and exclusion contours are obtained in the plane of the reduced Planck scale M_D and M_{Th} . A model independent limit on $\sigma B \times A$, where B is the branching ratio to dimuons, and A the acceptance of non Standard Model contributions, is derived to be 0.018 pb.

5. Searches for Contact Interactions, W_R and Majorana Neutrinos, and Leptoquarks

5.1 Search for Contact interaction

If quarks and leptons are composite, with at least one common constituent, the interaction of these constituents would likely be manifested through an effective four-fermion contact interaction (CI) at energies well below the compositeness scale. Such a contact interaction could also describe a new interaction with a messenger too heavy for direct observation at the LHC, in analogy with Fermi's nuclear β decay theory. A new interaction at high energy scale could be described by the effective Lagrangian, with parameter Λ interpreted as a scale below quarks and leptons constituents are bound. One investigated process is $q\bar{q}l^+l^-$, $l = e, \mu$ contact interaction [22]. The addition of the contact interaction term to the SM Lagrangian describing Drell-Yan production alters the DY production cross section. The largest deviations, either constructive or destructive, are expected at high dilepton invariant mass and are determined by the scale Λ . This analysis interprets the data in the context of the left-left isoscalar model (LLIM), which is commonly used as a benchmark for contact interaction searches [23]. Fig. 5 (top) represents DY mass spectra in the $\mu^+\mu^-$, with deviation from the spectra if as a function of Λ if constructive/destructive interference had been present. To test the consistency between the data and the standard model, a likelihood ratio test was performed. The derived p -value corresponding to the probability of observing a fluctuation corresponds to the likelihood after in the pseudoexperiments that is at least as signal-like as that seen in the data (i.e. with a maximum likelihood ratio greater or equal to that obtained in the data), is estimated to be 39% (79%) in the electron channel and 21% (5%) in the muon channel for constructive (destructive) interference. These values indicate that there is no significant evidence for contact interactions in the analyzed data and thus limits are set on the contact interaction scale Λ . The observed limits are $\Lambda^- > 10.1$ TeV ($\Lambda^+ > 9.4$ TeV) in the electron channel and $\Lambda^- > 8.0$ TeV ($\Lambda^+ > 7.0$ TeV) in the muon channel for constructive (destructive) interference.

Another type of contact interaction that has been investigated, albeit with smaller dataset, is $q\bar{q}q\bar{q}$ CI [24]. Basically, this is dijet signature, with event selection similar as in dijet resonant search described previously. The basic observable is $\chi \equiv e^{|\Delta y|}$ where Δy is the rapidity difference between the leading and the subleading jet. A derived observable is F_χ , which is the fraction of events at $\chi < e^{1.2}$. The choice of $e^{1.2}$ is based on optimization of the sensitivity to quark CI, and new physics would appear as an increase in F_χ . When F_χ is computed in bins of m_{jj} , the result is a spectrum which can indicate new physics by an increase of $F_\chi(m_{jj})$ in m_{jj} bins that contain significant amounts of signal. Fig. 5 (bottom) shows the observed and expected $F_\chi(m_{jj})$ spectra. No such signal is observed, and limit on Λ is set to be $\Lambda > 9.5$ TeV.

5.2 Search for W_R and Majorana Neutrinos

ATLAS has recently released a search for hypothetical heavy Majorana neutrinos and W_R gauge bosons [25]. The results are based on 34 pb^{-1} of data collected in 2010. The approach is guided by the Left-Right Symmetric model (LRSM) [26] in which a new gauge group with its force particles (W_R) manifesting themselves directly at LHC energies is introduced [27]. In this model, the heavy neutrinos N are produced in the decays of a W_R via $q\bar{q} \rightarrow W_R \rightarrow lN$, with N decaying subsequently to $N \rightarrow lW_R^* \rightarrow ljj$. The final state signature is two high- p_T leptons (electron or muon) and two high- p_T jets. The N and W_R invariant masses can be fully reconstructed from the

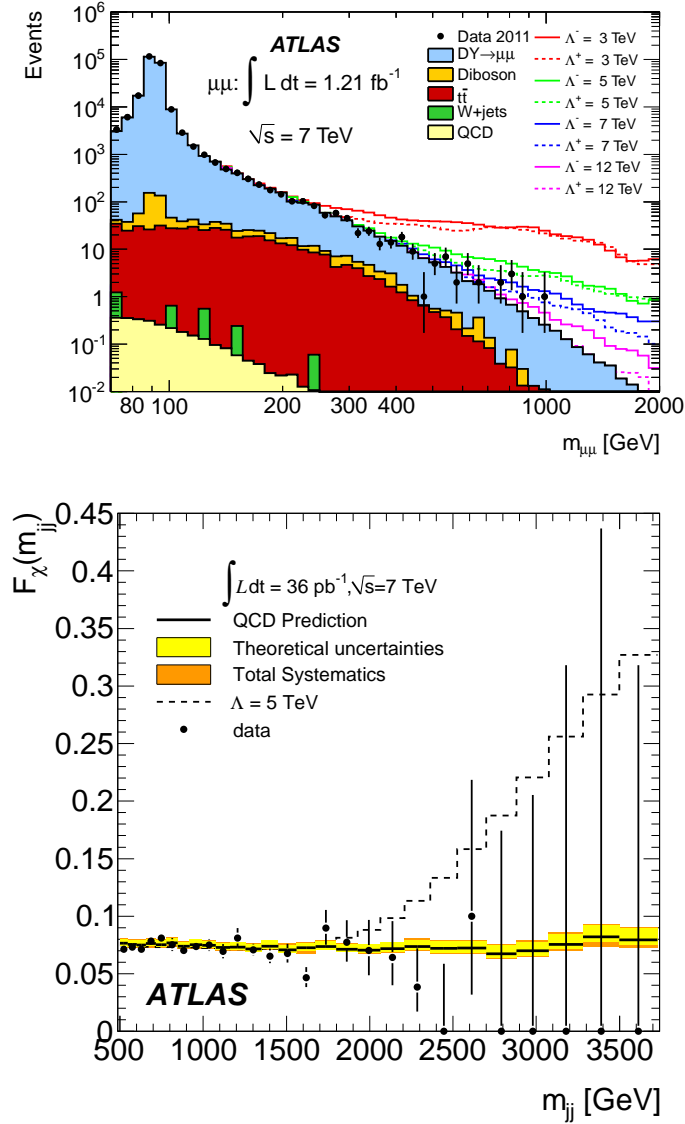


Figure 5: Top: $m_{\mu\mu}$ distribution in the contact interaction search. The open histograms correspond to the distributions expected in the presence of contact interactions with different values of Λ for constructive or destructive interference. Bottom: The $F_\chi(m_{jj})$ versus m_{jj} and with the expected signal from QCD plus a quark contact interaction as described in the text.

decay products. Similarly to SM neutrinos, heavy Majorana neutrinos can mix if their masses are different. Both the scenarios of no-mixing and 100% mixing between two generations of lepton flavour (e or μ) are investigated. Both the same-sign (SS) and opposite-sign (OS) dilepton final states are considered. If the heavy neutrinos are of Majorana type, they would contribute to both the SS and OS channels, while non-Majorana heavy neutrinos would contribute solely to the OS channel. The OS final state can be exploited for the LRSM search since the expected signal events, for the considered W_R masses, populate a kinematic region where the SM backgrounds are small.

The final selection requires $m_{ll} > 110$ GeV and $m_{lljj} > 400$ GeV in the SS search, while in OS search additionally $S_T > 400$ GeV is required, where S_T is the scalar sum of the transverse energies of the leptons and up to the two leading jets (in p_T). The numbers of observed data events and the expected SM background events in the SS and OS final states signal regions for the preselection and final selection are in good agreement. The upper limits on the σB for new physics based on the LRSM models are obtained. The limits on the heavy Majorana neutrino and W_R masses, for each mixing scenario, are determined for SS events and SS+OS events. The SS+OS limit is obtained by combining the two individual limits, and results are shown in Fig. 6.

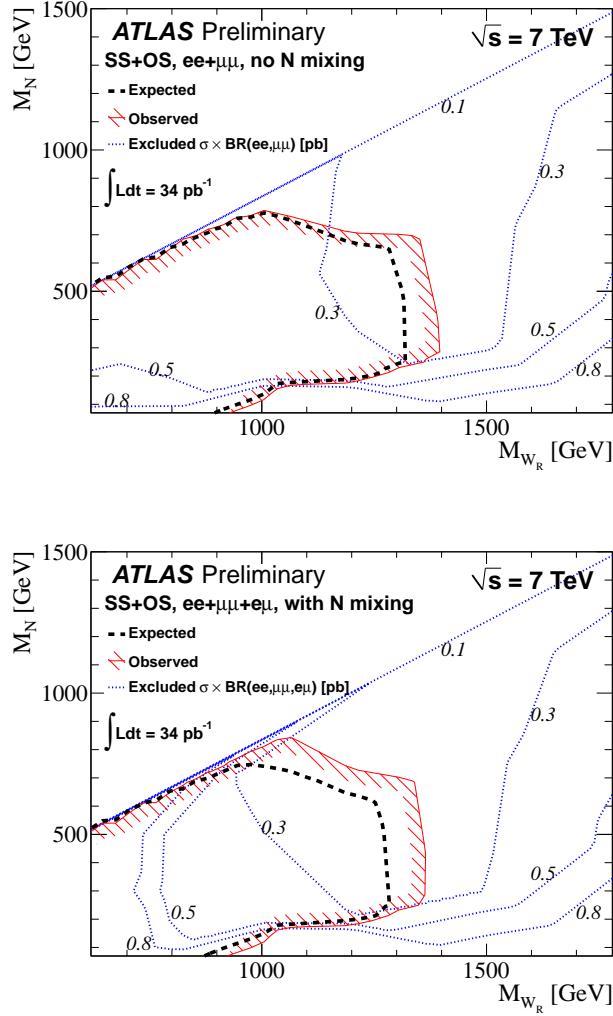


Figure 6: Observed and expected limits on the heavy Majorana neutrino and W_R masses for the combined SS+OS channels. The no-mixing scenario is shown (top), as well as 100% mixing scenario (bottom).

5.3 Search for the first generation of leptoquarks

A search for the first generation of scalar leptoquarks using 1.03 fb^{-1} of pp data has been

recently released by the ATLAS [28]. Leptoquarks are sought via their decay into an electron or neutrino and a quark, producing events with two oppositely charged electrons and at least two jets, or events with an electron, large E_T^{miss} and at least two jets. Control data samples are used to validate background predictions from MC. Likelihood ratio is used to separate signal from background in order to allow mass independent optimization of the search. The following discriminating variables, selected to give the best separation between signal and background, are used for the construction of the likelihood: *i*) for the $eejj$ channel scalar sum of electron and jet p_T (S_T), and the average invariant LQ mass; *ii*) for the $evjj$ channel m_T , S_T , the transverse LQ mass, and the invariant LQ mass $m_{LQ}(e, jet)$ are used. No signal is observed and limits are set on $M_{LQ} > 662$ GeV assuming 100% branching ratio to lepton+quark and $M_{LQ} > 606$ GeV assuming 50% branching ratio to lepton+quark.

6. Searches for New Phenomena in $t\bar{t}$ Events with Large Missing Energy

The top quark holds great promise as a probe for new phenomena at the TeV scale. It has the strongest coupling to the Standard Model Higgs boson, and as a consequence it is the main contributor to the quadratic divergence in the Higgs mass. Thus, assuming the "naturalness" hypothesis of effective QFT, light top partners (with masses ~ 1 TeV) should correspond to one of the most robust predictions of solutions to the hierarchy problem. There are quite a lot of searches for new physics in ATLAS involving top quarks, some of them are presented at this conference [29]. As an example, a search for pair-produced exotic top partners $T\bar{T}$, is described [30]. Each of these particles decay to top quark and a neutral weakly-interacting particle ($T\bar{T} \rightarrow t\bar{t}A_0A_0$) which in some models maybe its own antiparticle. The final state for such a process is identical to $t\bar{t}$, though with a larger amount of E_T^{miss} from undetectable A_0 . The search is performed in the semileptonic channel, resulting in a final state with an isolated lepton, four or more jets, and high E_T^{miss} . The observed yield in the signal region is compared with the SM expectation, Fig. 7. In the absence of signal an upper limit on σB is computed. In the model of exotic fourth generation up-type quarks the $T\bar{T}$ production cross-section is predicted to be approximately six times higher than for stop squarks with a similar mass, due to the multiple spin states of two T 's compared to scalar stops. For this model the cross-section limits are converted to an exclusion curve in the T vs A_0 mass parameter space. No excess is found over SM predictions, with derived limit $m_T > 420$ GeV for $m_{A_0} > 10$ GeV, Fig. 7. It should be noted that the estimated acceptance times efficiency for spin- $\frac{1}{2}$ $T\bar{T}$ models is consistent within systematic uncertainties with that for scalar models, such as pair production of stop squarks. The cross-section limits are therefore approximately valid for such models, although the predicted cross-section is typically below the current sensitivity.

7. Summary and Prospects

In this conference report, a subset of ATLAS searches for new phenomena, other than those based on Supersymmetry, are briefly described. There is unfortunately no sign of new physics yet at the LHC, though most of the limits supersede the ones from the Tevatron experiments. ATLAS however will continue to explore a wide variety of signatures, increasing the mass reach and present the results in as model-independent manner as possible. In 2012, $O(10)$ fb $^{-1}$ of data is foreseen,

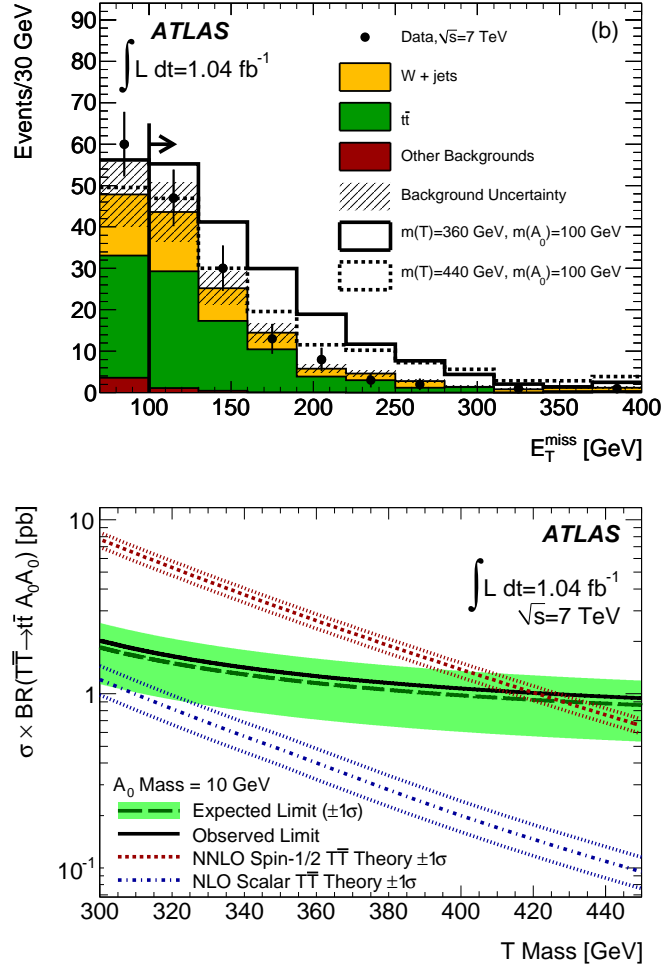


Figure 7: Top: E_T^{miss} in the search for pair-produced exotic top partners $T\bar{T}$. Bottom: σB excluded at the 95 confidence level versus T mass for $m(A_0) = 10 \text{ GeV}$.

with expected increase in collision energy to 8 TeV, which will result in significant gains in mass reach for many BSM phenomena. More significant gains might be expected by pushing towards smaller couplings, compared to the ones in common "benchmark" models.

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