

# 1 QCD Analysis of the Combined HERA Inclusive Data 2 together with HERA Jet and Charm Data

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This talk presents results of the latest combination of inclusive deep inelastic cross sections published by the H1 and ZEUS collaborations at HERA for neutral and charged current  $e^\pm p$  scattering for zero beam polarisation, including charm and jet data. The combined cross sections were used as input to QCD analyses at leading order and next-to-leading order providing a new set of parton distribution functions (PDFs), called HERAPDF2.0Jets. This set is a variation in the HERA-PDF2.0 family set of pdfs which is based on fitting DGLAP evolution equations to HERA data only. The inclusion of charm- and jet-production cross sections made a simultaneous determination of these parton distributions and the strong coupling constant possible, resulting in  $\alpha_s(M_Z^2) = 0.1183 \pm 0.0009(\text{exp}) \pm 0.0005(\text{model/parameterisation}) \pm 0.0012(\text{hadronisation})_{-0.0030}^{+0.0037}(\text{scale})$ . Comparisons of cross sections predictions based on HERAPDF2.0Jets NLO to the input data are excellent.

*The XXIII International Workshop on Deep Inelastic Scattering and Related Subjects*

*April 27 - May 1, 2015*

*Southern Methodist University*

*Dallas, Texas 75275*

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### 3 1. Introduction

4 Deep inelastic scattering (DIS) of electrons on protons at centre-of-mass energies of up to  
 5  $\sqrt{s} \simeq 320 \text{ GeV}$  at HERA has been central to the exploration of proton structure and quark–gluon  
 6 dynamics as described by perturbative quantum chromodynamics (pQCD) [1]. The H1 and ZEUS  
 7 collaborations have recently published a comprehensive QCD analysis of inclusive and various  
 8 other previously published data [2]. The data were taken at proton beam energies of 920, 820,  
 9 575 and 460 GeV and an electron (or positron) beam energy of 27.5 GeV. They correspond to an  
 10 integrated luminosity of about  $1 \text{ fb}^{-1}$  and span six orders of magnitude in negative four-momentum-  
 11 transfer squared,  $Q^2$ , and Bjorken  $x$ .

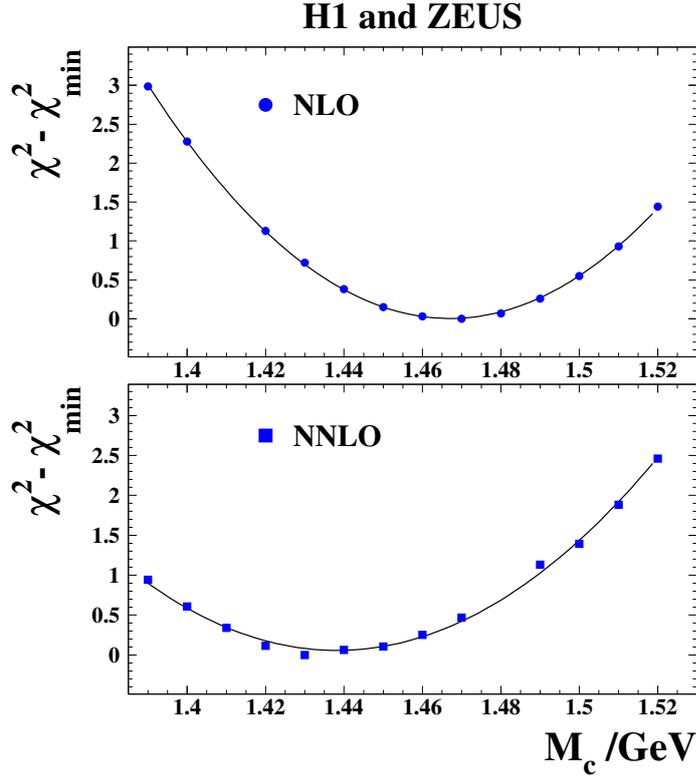
12 The name HERAPDF stands for the family of pQCD analyses based on the DGLAP formalism  
 13 using only these data. The nominal HERAPDF2.0 pdf set, based only on the inclusive cross sec-  
 14 tions of the full HERA data set, is described in Ref. [2]. The different variations of HERAPDF2.0  
 15 include different assumptions on  $Q_{min}^2$ , heavy flavour scheme, NLO and NNLO calculations and  
 16 the inclusion of additional data sets. Both H1 and ZEUS also published charm production cross  
 17 sections, some of which were combined and analysed previously [3], and jet production cross sec-  
 18 tions [4–8]. These data were included in the QCD analysis to obtain the variant HERAPDF2.0Jets,  
 19 presented in this talk.

### 20 2. Adding HERA Charm Data to the HERAPDF2.0 Fit

21 At HERA charm is dominantly produced in virtual photon-gluon fusion. The charm mass pa-  
 22 rameter  $M_c$  is an important scale in pQCD calculations of the process. Sensitivity to the gluon PDF  
 23 is expected since charm contributes up to 30% at high  $Q^2$ . However, in a previous QCD analysis us-  
 24 ing only HERA-I data (HERA data taken in the years up to 2000) the impact of adding charm data  
 25 to inclusive data as input to NLO QCD fits has been established to mainly result in a reduction of  
 26 the uncertainty on the charm mass parameter  $M_c$  for a given flavour scheme [3]. It was also estab-  
 27 lished that the optimal value of  $M_c$  can differ depending on the choice of a particular general-mass  
 28 variable-flavour-number scheme (GM-VFNS) chosen for the fit. The fits for all schemes considered  
 29 were found to be of similar quality in describing the HERA data. For the HERAPDF2.0 analysis,  
 30 a total of 47 data points on charm production with  $Q^2$  larger than  $Q_{min}^2 = 3.5 \text{ GeV}^2$  were added as  
 31 input to the NLO fits. The 42 sources of correlated systematic uncertainty from the H1 and ZEUS  
 32 data sets on charm production and two additional sources due to the combination procedure were  
 33 taken into account. Figure 1 shows  $\chi^2$  scans of NLO and NNLO pQCD fits to HERA inclusive  
 34 and combined charm data to determine the optimal value of  $M_c$ . The GM-VFNS RTOPT flavor  
 35 scheme was used in the fits. The optimal value  $M_c = 1.47 \text{ GeV}$  at NLO ( $M_c = 1.43 \text{ GeV}$  at NNLO)  
 36 are chosen at minimum  $\chi^2$  and are used in subsequent fits. Fits including the charm data were  
 37 proven to be consistent with fits using the inclusive data only, but only a marginal reduction in the  
 38 uncertainty on the low- $x$  gluon PDF was obtained.

### 39 3. Adding HERA Jet Data to the HERAPDF2.0fit

40 In pQCD fits to inclusive DIS data alone, the gluon PDF is determined indirectly via the  
 41 DGLAP equations using the observed scaling violations, while the value of  $\alpha_s(M_Z^2)$  is given as

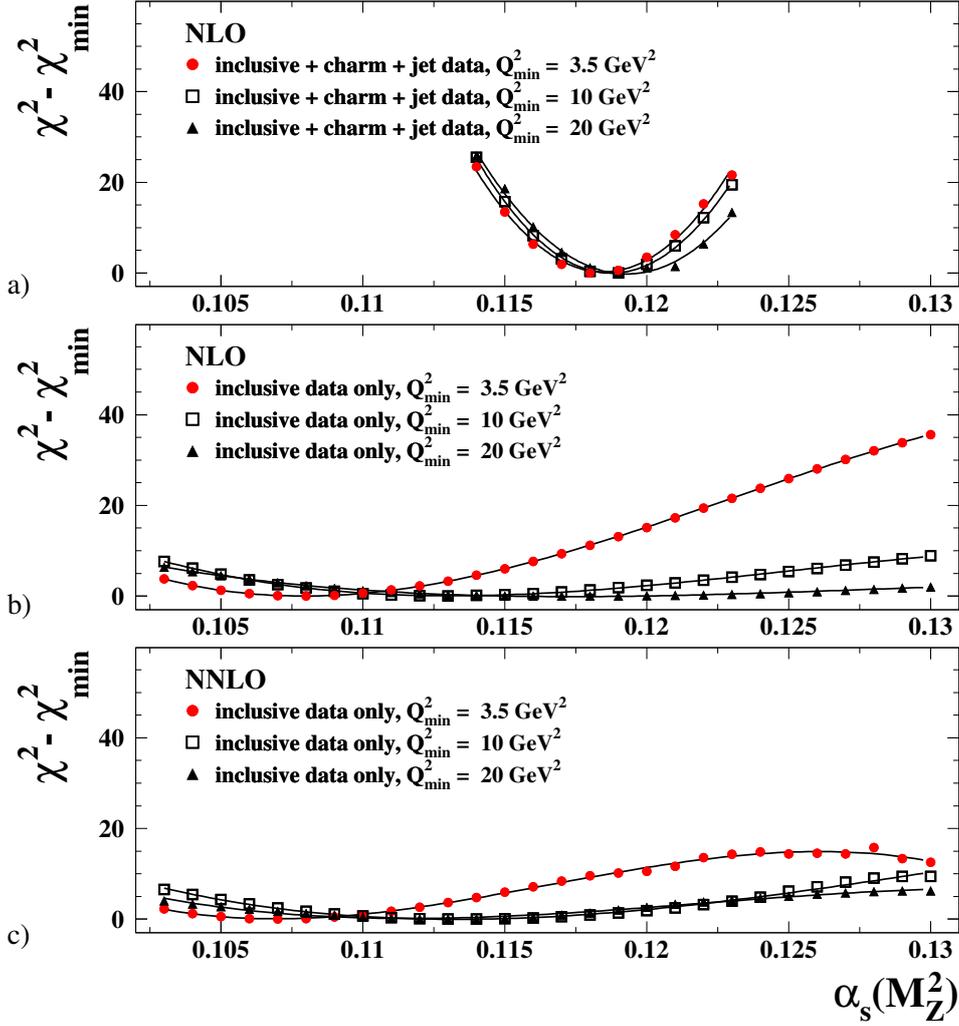


**Figure 1:** The  $\Delta\chi^2 = \chi^2 - \chi^2_{\min}$  versus the charm mass parameter  $M_c$  for NLO and NNLO fits based on the combined data on charm production in addition to the combined inclusive data.

input from external information [9]. Data on jet production cross sections provide an independent measurement of the gluon distribution. They are sensitive to the gluon distribution at lower  $Q^2$  and to the valence-quark distribution at higher  $Q^2$ . The inclusion of jet data in pQCD fits not only reduces the uncertainty on the high- $x$  gluon distribution in fits with fixed  $\alpha_s(M_Z^2)$  but also allow the accurate simultaneous determination of  $\alpha_s(M_Z^2)$  and the gluon distribution. This is demonstrated in Fig. 2 where a  $\chi^2$  fit to  $\alpha_s(M_Z^2)$  is shown. Fits including these jet data and including the combined charm data were performed with  $\alpha_s(M_Z^2) = 0.118$  fixed and with  $\alpha_s(M_Z^2)$  as a free parameter in the fit. This inclusion was first used to validate the choice of  $\alpha_s(M_Z^2) = 0.118$  for HERAPDF2.0 by investigating the dependence of the  $\chi^2$ s of the HERAPDF2.0 pQCD fits on  $\alpha_s(M_Z^2)$ . Three  $\chi^2$  scans vs. the value of  $\alpha_s(M_Z^2)$  were performed at NLO for three values of  $Q_{\min}^2$ . The result is depicted in the top panel of Fig. 2. A distinct minimum at  $\alpha_s(M_Z^2) \approx 0.118$  is observed, which is basically independent of  $Q_{\min}^2$ . This validates the choice of  $\alpha_s(M_Z^2) = 0.118$  for HERAPDF2.0 NLO. Scans at NLO and NNLO were also performed for fits to inclusive data only. The middle and bottom panels of Fig. 2 show that these scans yielded similar shallow  $\chi^2$  dependences and the minima were strongly dependent on the  $Q_{\min}^2$ . This demonstrates that the inclusive data alone cannot constrain  $\alpha_s(M_Z^2)$  reasonably.

The resulting HERAPDF variant with free  $\alpha_s(M_Z^2)$  is called HERAPDF2.0Jets. A full uncertainty analysis was performed for the HERAPDF2.0Jets variant, including model and parameterisation uncertainties and additional hadronisation uncertainties on the jet data as evaluated for the

## H1 and ZEUS



**Figure 2:**  $\Delta\chi^2 = \chi^2 - \chi_{\min}^2$  vs.  $\alpha_s(M_Z^2)$  for pQCD fits with different  $Q_{\min}^2$  using data on (a) inclusive, charm and jet production at NLO, (b) inclusive  $ep$  scattering only at NLO and (c) inclusive  $ep$  scattering only at NNLO.

61 original publications [4–8].

62 The PDFs resulting from a fit with free  $\alpha_s(M_Z^2)$ , HERAPDF2.0Jets, and from a fit with fixed  
 63  $\alpha_s(M_Z^2) = 0.118$  are shown in Fig. 3. A full uncertainty analysis was performed in both cases, in-  
 64 cluding model and parameterisation uncertainties as well as additional hadronisation uncertainties  
 65 on the jet data. The PDFs are very similar, because the HERAPDF2.0Jets fit with free  $\alpha_s(M_Z^2)$   
 66 yields a value which is very close to the value used for the fit with fixed  $\alpha_s(M_Z^2)$ . The jet data deter-  
 67 mine the value of  $\alpha_s(M_Z^2)$  very well in the HERAPDF2.0Jets fit. Thus, the uncertainty on  $\alpha_s(M_Z^2)$   
 68 in this fit does not significantly increase the uncertainty on the gluon PDF with respect to the fit  
 69 with  $\alpha_s(M_Z^2)$  fixed. The difference in the  $\alpha_s(M_Z^2)$  free fit is mostly due to extra uncertainty coming  
 70 from the hadronisation corrections.

71 The PDFs from the HERAPDF2.0Jets fit with  $\alpha_s(M_Z^2) = 0.118$  fixed are also very similar to  
 72 the standard PDFs from HERAPDF2.0 NLO. This is demonstrated in Fig. 4. This is again the

73 result of the choice of  $\alpha_s(M_Z^2) = 0.118$  for HERAPDF2.0 which is also the preferred value for  
 74 HERAPDF2.0Jets. Consequently, there is only a small reduction of the uncertainty on the gluon  
 75 distribution observed for HERAPDF2.0Jets.

76 The  $\chi^2$  of the HERAPDF2.0Jets fit with free  $\alpha_s(M_Z^2)$  is the same as for the fit with fixed  
 77  $\alpha_s(M_Z^2) = 0.118$ . This is again due the fact that the value of  $\alpha_s(M_Z^2)$  obtained from the fit is very  
 78 close to the value previously fixed. The strong coupling constant obtained is

$$\alpha_s(M_Z^2) = 0.1183 \pm 0.0009(\text{exp}) \pm 0.0005(\text{model/parameterisation}) \\ \pm 0.0012(\text{hadronisation}) \begin{matrix} +0.0037 \\ -0.0030 \end{matrix}(\text{scale}) .$$

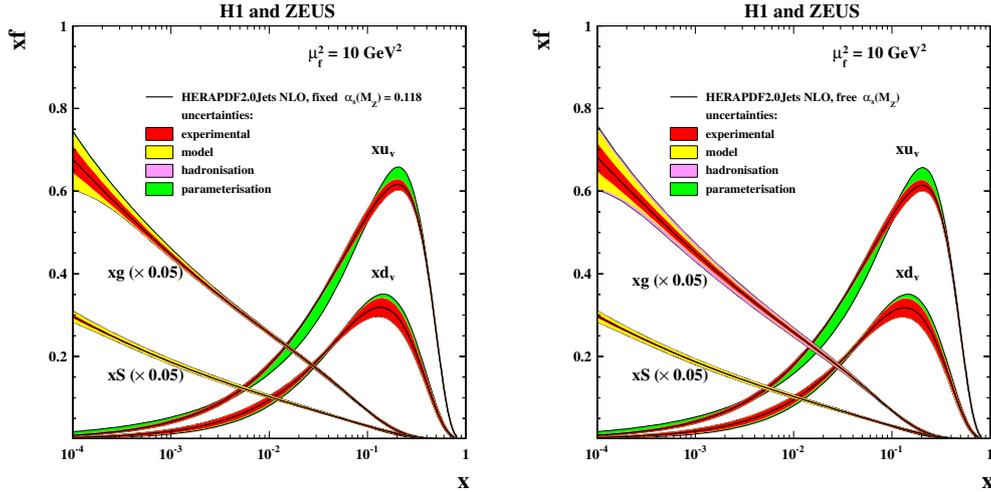
79 The uncertainty on  $\alpha_s(M_Z^2)$  due to scale uncertainties was evaluated by varying the renormalisation  
 80 and factorisation scales by a factor of two, both separately and simultaneously, and taking the  
 81 maximal positive and negative deviations. The uncertainties were assumed to be 50 % correlated  
 82 and 50 % uncorrelated between bins and data sets. This resulted in an asymmetric uncertainty of  
 83  $+0.0037$  and  $-0.0030$ . The result on  $\alpha_s(M_Z^2)$  is compatible with the world average [9] and it is  
 84 competitive with other determinations at NLO.

### 85 3.1 Comparison of HERAPDF2.0Jets NLO to data

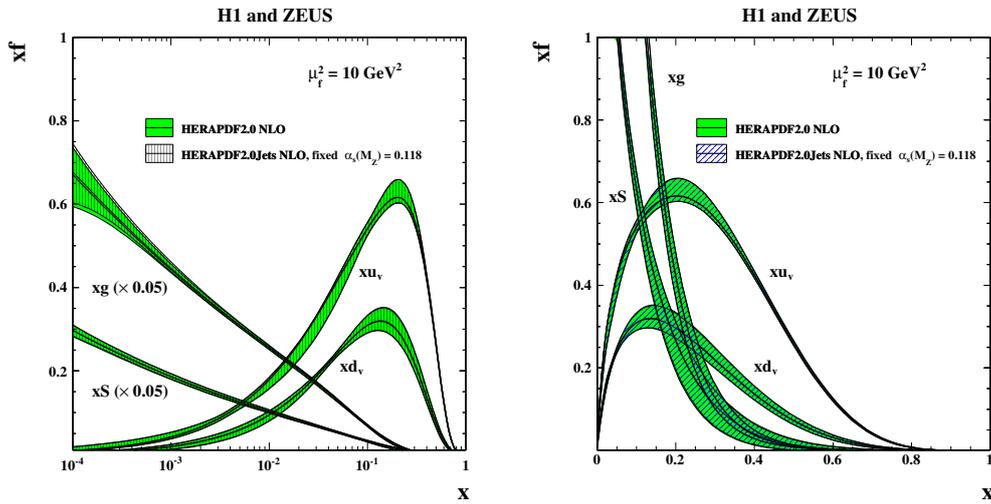
86 The predictions of HERAPDF2.0Jets NLO with free  $\alpha_s(M_Z^2)$  were compared to charm input  
 87 data [3] and the description of the data is excellent. Comparisons of the predictions of HERA-  
 88 PDF2.0Jets to the data on jet production used as input [4–8]. were also done. All analyses were  
 89 performed using the assumption of massless jets, i.e. the transverse energy,  $E_T$ , and the transverse  
 90 momentum of a jet,  $p_T$ , are equivalent. For inclusive jet analyses, each jet is entered separately with  
 91 its  $p_T$ . For dijet and trijet analyses, the average of the transverse momenta is used as  $p_T$ . These  
 92 different definitions of  $p_T$  were also used to set the the renormalisation scale to  $\mu_r^2 = (Q^2 + p_T^2)/2$   
 93 for calculating predictions. The factorisation scale was chosen as  $\mu_f^2 = Q^2$ . Scale uncertainties  
 94 were not considered for the comparisons to data. The description of all the data on jet production  
 95 by HERAPDF2.0Jets NLO is also excellent.

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**Figure 3:** The parton distribution functions  $xu_v$ ,  $xd_v$ ,  $xS = 2x(\bar{U} + \bar{D})$  and  $xg$  of HERAPDF2.0Jets NLO at  $\mu_f^2 = 10\text{GeV}^2$  with fixed  $\alpha_s(M_Z^2) = 0.118$  (top) and free  $\alpha_s(M_Z^2)$  (bottom). The experimental, model and parameterisation uncertainties are shown. The hadronisation uncertainty is also included, but it is only visible for the fit with free  $\alpha_s(M_Z^2)$ .



**Figure 4:** The parton distribution functions  $xu_v$ ,  $xd_v$ ,  $xS = 2x(\bar{U} + \bar{D})$  and  $xg$  of HERAPDF2.0Jets NLO at  $\mu_f^2 = 10\text{GeV}^2$  compared to those of HERAPDF2.0 NLO on logarithmic (top) and linear (bottom) scales. The fits were done with fixed  $\alpha_s(M_Z^2) = 0.118$ . The bands represent the total uncertainties.

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