



# A study of $\gamma$ -hadron correlation in p + Pb collisions at $\sqrt{s_{\rm NN}} = 5.02$ TeV

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Under the assumption that a quark-gluon plasma droplet is created in p + A collisions and partons traversing it will lose their energy, we calculate  $\gamma$ -triggered hadron correlation in p + Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV, within a next-to-leading-order perturbative QCD parton model with the medium-modified fragmentation functions. The parton energy loss can be controlled by the scaled jet transport coefficient  $\hat{q}/T^3$  within the high-twist (HT) approach. The evolution informations of such QGP medium created in p + A collisions are provided by the SuperSONIC hydrodynamics model. With the value of  $\hat{q}/T^3$  extracted via single hadron suppressions in A + A collisions with similar highest initial temperature as in p + A collisions, the  $\gamma$ -hadron spectra with  $p_T^{\gamma} = 12 - 40$  GeV/*c* show a suppression of 5%~10% in the most central 0 - 10% p + Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV. We also provide the predictions for  $\gamma$ -hadron suppression in Pb + Pb collisions at the LHC energies.

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### 1. Introduction

It is generally believed that  $\gamma$ -jet production is a golden probe for studying the parton energy loss [1] in high-energy heavy-ion collisions. If we assume that a small QGP droplet is produced in p + A collisions and its evolution can be described by hydrodynamics, we can also predict suppression of  $\gamma$ -hadron production with medium-modified fragmentation functions. We assume that partons will lose their energy when traversing such a medium and the lost energy is controlled by the jet transport coefficient  $\hat{q}$  [2] which is defined as the transverse momentum broadening squared per unit length. It depends on medium temperature T and four fluid velocity  $u^{\mu}$  in the form  $\hat{q} = \hat{q}_0 \frac{T_3^3}{T_0^3} \frac{p^{\mu} \cdot u^{\mu}}{p_0}$ , where T<sub>0</sub> is the highest temperature at the center of the medium at the initial time  $\tau_0$  for the QGP formation. The information for T and  $u^{\mu}$  are given by event-by-event simulations of the superSONIC hydrodynamic model[3, 4]. In our last work [5], we find that the scaled dimensionless initial jet transport coefficient  $\hat{q}/T^3$  is decreasing slightly with the initial temperature. So we believe that the same value of jet transport coefficient can be approximatively applied for  $\gamma$ -hadron in p + A collisions as obtained for single hadron in A + A collisions at the similar initial temperature of the created QGP mediums.

#### 2. The theory model

In p +A collisions,  $\gamma$ -hadron cross section can be expressed as,

$$\frac{d\sigma_{pA}^{\gamma h}}{dy^{\gamma} d^2 p_{\rm T}^{\gamma} dy^h d^2 p_{\rm T}^h} = \sum_{abd} \int d^2 r dz_d t_A(\vec{r}) f_{a/A}(x_a, \mu^2, \vec{r}) f_{b/p}(x_b, \mu^2) \frac{x_a x_b}{\pi z_d^2} \\
\times \frac{d\sigma_{ab \to \gamma d}}{d\hat{t}} \tilde{D}_{h/d}(z_d, \mu^2, \Delta E_d) + O(\alpha_e \alpha_s^2),$$
(1)

where  $f_b/p(x_b, \mu^2)$  and  $f_a/A(x_a, \mu^2, \vec{r})$  are the parton distribution functions and  $t_A(\vec{r})$  is the nuclear thickness function.  $d\sigma_{ab\to\gamma d}/d\hat{t}$  are the tree-level  $2 \to 2$  partonic scattering cross sections. The NLO correction at  $O(\alpha_e \alpha_s^2)$  order included in our calculation contains  $2 \to 2$  virtual diagrams and  $2 \to 3$  tree diagrams. We only focus on direct photons which come from the hard processes of the Compton  $(qg \to q\gamma)$  or annihilation  $(q\bar{q} \to g\gamma)$  scattering. With isolation cuts the contribution from fragmentation photons is only about 10%. We can ignore them here [7].

The medium-modified fragmentation function  $\tilde{D}_{h/d}(z_d, \mu^2, \Delta E_d)$  can be expressed as [6],

$$\tilde{D}_{h/d}(z_d, \mu^2, \Delta E_d) = (1 - e^{-\langle N_g^d \rangle}) \left[ \frac{z'_d}{z_d} D_{h/d}(z'_d, \mu^2) + \langle N_g^d \rangle \frac{z_g'}{z_d} D_{h/g}(z_g', \mu^2) \right] + e^{-\langle N_g^d \rangle} D_{h/d}(z_d, \mu^2).$$
(2)

With the high-wist formalism [8], the radiative energy loss  $\Delta E_d$  can be calculated as,

$$\frac{\Delta E_d}{E} = \frac{2C_A \alpha_s}{\pi} \int d\tau \int \frac{dl_{\rm T}^2}{l_{\rm T}^4} \int dz \times \left[1 + (1-z)^2\right] \hat{q}_d \sin^2 \left[\frac{l_{\rm T}^2(\tau-\tau_0)}{4z(1-z)E}\right],\tag{3}$$

where  $C_A = 3$ ,  $l_T$  is the transverse momentum of the radiated gluon.

Using the spectrum in p + p collisions as a baseline, the gamma-hadron nuclear modification factor  $I_{pA}^{\gamma h}(z_{\rm T}) = \frac{D_{pA}^{\gamma h}(z_{\rm T})}{D_{pA}^{\gamma h}(z_{\rm T})}$  can be expressed as a function of  $z_{\rm T} = p_{\rm T}^{h}/p_{\rm T}^{\gamma}$  [9].  $D_{pA}^{\gamma h}(z_{\rm T})$  is the  $\gamma$ -triggered fragmentation function which can be defined as the ratio of  $\gamma$ -hadron cross section to the trigger photon cross section.

#### 3. Numerical results

We first calculate the  $\gamma^{\text{dir}}$ -triggered fragmentation function in p + p collisions at  $\sqrt{s_{\text{NN}}} = 0.2$ TeV and the corresponding medium modification factor  $I_{AuAu}^{\gamma h}$  in 0 - 10% Au + Au collisions. Both of them agree well with the experimental data. The details are shown in our recent paper [10]. The predictions for  $\gamma$ -hadron suppression  $I_{PbPb}^{\gamma h}$  in Pb + Pb collisions at  $\sqrt{s_{\text{NN}}} = 2.76$  TeV and 5.02 TeV for different (0 - 5%, 20 - 30%, 40 - 50%, 60 - 70%) centralities are shown in Fig. 1. The corresponding  $\hat{q}_0 = 1.8$  GeV<sup>2</sup>/fm in Pb + Pb collisions at  $\sqrt{s_{\text{NN}}} = 2.76$  TeV and  $\hat{q}_0 = 2.0$  GeV<sup>2</sup>/fm in Pb + Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV are extracted via comparisons to single inclusive hadron suppressions in 0 - 5% Pb + Pb collisions at these two energies, respectively [10].

Shown in Fig. 2 are our predictions for  $I_{pPb}^{\gamma h}$  in p + Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV with  $\hat{q}_0 = 1.5$  GeV<sup>2</sup>/fm which is extracted from single hadron production in Au + Au collisions at  $\sqrt{s_{NN}} = 0.2$  TeV which have the similar central temperature as in p + Pb collisions. The shaded bands indicate variations of the results when one changes the initial time for parton-medium interaction between  $\tau_0 = 0.5$  and 1.0 fm/c. For  $\gamma$ -hadron spectra in p + Pb collisions, we see a suppression of about 5%~10% due to jet quenching. In both Pb + Pb and p + Pb collisions, the suppression of  $\gamma$ -triggered hadron spectra becomes weaker with a larger  $p_T$  trigger photon.



**Figure 1:**  $\gamma^{\text{dir}}$ -hadron suppression factors as a function of  $z_{\text{T}}$  in 0 - 5%, 20 - 30%, 40 - 50% and 60 - 70% Pb + Pb collisions, with  $12 < p_{\text{T}}^{\gamma} < 40 \text{ GeV}/c$ ,  $0.5 < p_{\text{T}}^{h} < 15 \text{ GeV}/c$  (upper panels) and  $40 < p_{\text{T}}^{\gamma} < 60 \text{ GeV}/c$ ,  $0.5 < p_{\text{T}}^{h} < 45 \text{ GeV}/c$  (lower panels) at  $\sqrt{s_{\text{NN}}} = 2.76 \text{ TeV}$  with  $\hat{q}_{0} = 1.8 \text{ GeV}^{2}/\text{fm}$  (left panels) and at  $\sqrt{s_{\text{NN}}} = 5.02 \text{ TeV}$  with  $\hat{q}_{0} = 2.0 \text{ GeV}^{2}/\text{fm}$  (right panels).

### 4. Summary

Under the assumption that a QGP droplet is produced in p + Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV and its evolution can be described by hydrodynamics, we predict the suppression of  $\gamma$ -triggered





**Figure 2:**  $\gamma^{\text{dir}}$ -hadron suppression factors as a function of  $z_{\text{T}}$  in 0 - 10%, 20 - 30%, 40 - 50% and 60 - 80% p + Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV with  $12 < p_{\text{T}}^{\gamma} < 40$  GeV/c,  $0.5 < p_{\text{T}}^{h} < 15$  GeV/c (upper panels) and  $40 < p_{\text{T}}^{\gamma} < 60$  GeV/c,  $0.5 < p_{\text{T}}^{h} < 45$  GeV/c (lower panels). The shaded bands indicate variations of the results when one changes the initial time for parton-medium interaction between  $\tau_0 = 0.5$  and 1.0 fm/c.

hadron spectra within NLO perturbative QCD parton model with medium modified fragmentation function due to parton energy loss. The parton energy loss is calculated with the high-twist formalism. The evolution information of the medium created in p + Pb collisions are provided by event-by-event superSONIC hydrodynamics model. Our numerical results show a suppression of about 5%~10% for  $\gamma$ -hadron spectra for 12 <  $p_T^{\gamma}$  < 40 GeV/*c* in the most 0 - 10% central p + Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV with the initial jet transport coefficient  $\hat{q}_0$  extracted from the suppression of single hadron spectra in A + A collisions. And we also predict the  $\gamma$ -hadron productions in Pb + Pb collisions at the LHC energies.

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