

Search for extra spatial dimensions at DØ

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1. Introduction

In this talk results from searches for Large Extra spatial Dimensions (LEDs) and Randall-Sundrum (RS) gravitons at DØ are presented, detailed information about the searches presented in this talk can be found in [1, 2, 3, 4]. The results from the search for TeV^{-1} -size extra dimensions can be found in [5]. Large Extra spatial Dimensions (LEDs) was proposed by Arkani-Hamed, Dimopoulos and Dvali (ADD) [6], they postulate that the standard model (SM) particles and gauge interactions are confined to a three-dimensional “brane” in a bulk which consists of 3 plus n additional compact spatial dimensions. This removes the hierarchy problem as gravity appears weak, ($M_{EW} \ll M_{Pl}$), because it propagates in the bulk. The reduced Planck scale M_D is the fundamental scale in the $(4+n)$ space-time. To be able to ‘dilute’ gravity in the volume of the extra dimensions, so that it appears weak in the 3D-world, these compact extra dimensions need to be ‘large’. The size of the LEDs, the compactification radius R , is given by: $R^n = M_{Pl}^2 / (8\pi M_D^{n+2})$. Gravitons propagate in the bulk and they appear as a tower of Kaluza-Klein (KK) excited modes from the point of view of the SM-brane and can be detected via virtual graviton exchange or direct graviton emission.

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Lisa Randall and Raman Sundrum has in [7] proposed a rigorous solution of the hierarchy problem, by adding a single extra dimension with a non-Euclidian, warped metric (RS-model). The size of the ED is R and the curvature of the warped metric is k , (the warp factor). The hierarchy problem is solved by localization of gravity to the Planck brane, at $\phi = 0$, while the SM fields are localised on the TeV-brane at $\phi = \pi$, where ϕ is the coordinate along the ED. Gravity originates on the Planck brane and the graviton wave function is exponentially suppressed away from the brane along the ED due to the warp-factor. The M_{Pl} -size operators on the Planck brane yield low-energy effects on the SM brane with a typical scale of $\Lambda_\pi = \bar{M}_{Pl} \exp(-k\pi R)$, where $\bar{M}_{Pl} \equiv M_{Pl}/\sqrt{8\pi}$ is the reduced Planck mass. The hierarchy problem is solved if $\Lambda_\pi \sim 1$ TeV, which can be achieved by requiring $kR \sim 10$. In the simplest RS model only gravitons propagate in the bulk, and can be detected as low mass resonances, the lowest order of these KK-modes is massless, and the searches presented here thus put a limit on M_1 , the mass of the first resonance. The parameters k and R can be expressed in M_1 and the coupling to the SM fields: k/\bar{M}_{Pl} . The theoretically preferred range of M_1 is a few hundred GeV to a few TeV, while k/\bar{M}_{Pl} is expected to be between 0.01 – 0.1.

2. Searches for LEDs via virtual graviton exchange.

Virtual graviton exchange can be detected as an enhancement at high masses of the di-lepton or di-photon mass spectrum. The change of the cross section due to the interference of gravity effects is: $d^2\sigma/dM d\cos\theta^* = f_{sm} + f_{int}\eta_G + f_{KK}\eta_G^2$, where f_{sm} is the SM-term, f_{int} is the interference and f_{KK} is the graviton terms, M is the invariant di-lepton or di-photon mass, and θ^* is the scattering angle. The effects of LEDs are parametrised by a single variable $\eta_G = F/M_s^4$. The dimensionless parameter F of order unity reflects the dependence of the cross section on the different formalisms for the effective Lagrangian: GRW [8], HLZ [9] and Hewett [10]. The only difference between the formalisms is the definitions of M_s , which is the ultra-violet cut-off required to keep the divergent sum over KK states finite. Virtual graviton effects are sensitive only to M_s and not M_D , but the two scales are expected to be of the same order. M_s could be slightly lower, in which case virtual graviton searches are more sensitive to LEDs than direct graviton searches. The analyses presented here are performed in the dielectron, diphoton and dimuon channels [1, 2]. To maximise the reconstruction efficiency for electrons and photons, the tracking information was ignored and these channels were combined in a single calorimeter based 'diEM' channel. The EM-objects were required to be isolated, to have transverse energy $E_T > 25$ GeV and to be in central or forward regions of the calorimeter, but with at least one of the 2 objects in the central region. Events with more than 2 high- p_T EM objects were rejected. The muons were identified in the muon spectrometer and were required to have a matching track in the central tracking detector, transverse momentum $p_T > 15$ GeV, be isolated and pass a cosmic veto. The main backgrounds are Drell-Yan (DY) and direct $\gamma\gamma$ production, and instrumental background from jets misidentified as EM objects ('QCD' background). The former was estimated using simulation and the QCD background was estimated from data. The resulting invariant mass spectra are shown in Fig. 1 and Fig. 2, the data is in agreement with the expected background. Limits on M_s are set using the distributions of M and $\cos\theta^*$ in a 2D likelihood fit of the parameter η_G . In the diEM channel, using 200 pb^{-1} of data, a new 95% C.L limit of M_s in the GRW convention of 1.36 TeV is set and combined with the RunI result the limit is 1.43 TeV. In the dimuon channel, using 246 pb^{-1} of data, 95% C.L limits of M_s between

0.85 and 1.27 TeV are set in the different formalisms. These are the most stringent limits on LEDs to date.

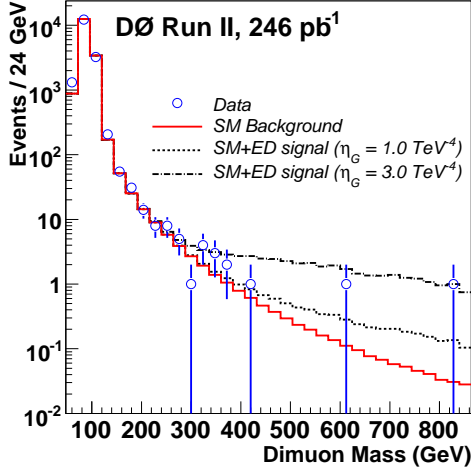


Figure 1: The dimuon mass distribution for data and SM background in the LED search, the effects of LED are shown for $\eta_G = 1 \text{ TeV}^{-4}$ and 3 TeV^{-4} in dashed lines.

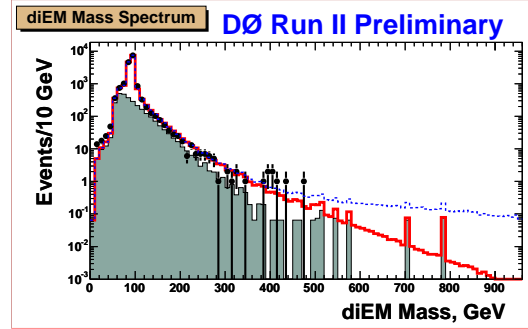


Figure 2: The diEM Mass distribution for data (errorbars), QCD background (histogram), and a fit to the SM background and QCD (full line). The dashed line show the effect of the ED signal for $\eta_G = 0.6 \text{ TeV}^{-4}$.

3. Searches for LEDs via direct graviton emission.

At the Tevatron gravitons can be produced recoiling against a quark or a gluon jet, which results in a monojet-like topology. The analysis presented here is the first preliminary result from Run II of searches for LEDs in the monojet channel [3]. The search is performed using events with single jets and acoplanar jets with missing E_T . Data was collected using a trigger specially designed for Jets + E_T topologies. The leading jet E_T and E_T were required to exceed 150 GeV, no additional jet with $E_T > 50 \text{ GeV}$ was allowed, and the angle (jet, E_T) was required to be at least 30° . The resulting missing E_T spectra is shown in Figure 3, as well as the uncertainty due to the jet energy scale (JES). The challenge in this analysis is the large instrumental background from E_T -mismeasurements and the sensitivity to uncertainty of the JES. Nevertheless, using only 85 pb^{-1} the 95% C.L. limit on M_D is 0.68 TeV for $n = 4$, which is better than the $D\emptyset$ Run I limit. This analysis is being updated using more data and improved JES.

4. Direct search for Randall-Sundrum gravitons.

In the RS-model, the excited KK-modes can decay into fermion-antifermion or diboson pairs and be detected as a resonance in the invariant mass spectrum. The analysis presented here is a direct search for RS-gravitons in the dielectron, diphoton and dimuon final states [4]. As for the LED-search, the electron and photon channel was combined into a 'diEM' channel to maximise the reconstruction efficiency. The requirements on the muons and EM-objects as well as the main

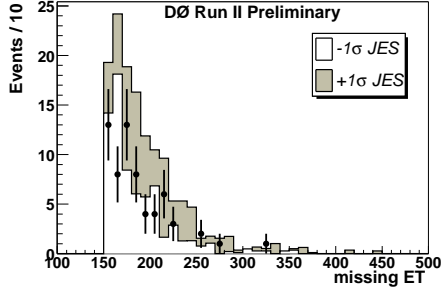


Figure 3: The E_T distribution of the jets+ E_T channel, the uncertainty due to JES is also shown.

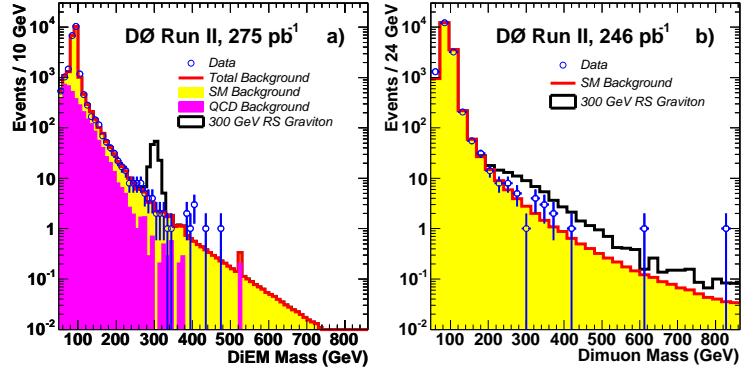


Figure 4: The resulting invariant mass spectra of the RS-graviton search in the diEM channel, (to the right) and the dimuon channel (to the left).

backgrounds are similar to the ones described in Section 2. The resulting invariant mass spectra are shown in Figure 4, the data is in agreement with the expected background, and limits on M_1 are set using the mass window method. A combination of the diEM-channel using 275 pb^{-1} of data, and the dimuon channel, using 246 pb^{-1} , excludes graviton masses up to 785 (250) GeV, for $k/M_{Pl} = 0.1$ (0.01), at 95% C.L.. These are the most stringent limits on RS-gravitons to date.

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