

KLOE extraction of V_{us} from kaon decays and lifetimes

KLOE collaboration *

presented by Barbara Sciascia[†]

INFN Frascati, Italy

E-mail: barbara.sciascia@lnf.infn.it

The most precise determinations of V_{us} comes from semileptonic kaon decays. The KLOE experiment at DAΦNE the Frascati ϕ factory, has measured all the experimental inputs to V_{us} for both neutral and charged kaons. Preliminary results for $BR(K^\pm \rightarrow \pi^0 e^\pm \nu)$ and $BR(K^\pm \rightarrow \pi^0 \mu^\pm \nu)$, and for τ_\pm are presented, together with measurements of the $K_L e 3$ and $K_L \mu 3$ decay BR, the $K_S e 3$ decay BR and the K_L lifetime τ_L . From our results for the 5 branching ratios and τ_L we find $V_{us}=0.2258\pm 0.0022$. We have also measured the fully inclusive $K_{\mu 2}^+(\gamma)$ absolute branching ratio for which we obtain $BR(K^+ \rightarrow \mu \nu(\gamma))=0.6366\pm 0.0017$. Combining this value with recent lattice results for f_K/f_π gives $V_{us}=0.2223\pm 0.0026$.

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*F. Ambrosino, A. Antonelli, M. Antonelli, C. Bacci, P. Beltrame, G. Bencivenni, S. Bertolucci, C. Bini, C. Bloise, V. Bocci, F. Bossi, D. Bowring, P. Branchini, R. Caloi, P. Campana, G. Capon, T. Capussela, F. Ceradini, S. Chi, G. Chiefari, P. Ciambrone, S. Conetti, E. De Lucia, A. De Santis, P. De Simone, G. De Zorzi, S. Dell’Agnello, A. Denig, A. Di Domenico, C. Di Donato, S. Di Falco, B. Di Micco, A. Doria, M. Dreucci, G. Felici, A. Ferrari, M. L. Ferrer, G. Finocchiaro, C. Forti, P. Franzini, C. Gatti, P. Gauzzi, S. Giovannella, E. Gorini, E. Graziani, M. Incagli, W. Kluge, V. Kulikov, F. Lacava, G. Lanfranchi, J. Lee-Franzini, D. Leone, M. Martini, P. Massarotti, W. Mei, L. Meola, S. Missetti, M. Moulson, S. Müller, F. Murtas, M. Napolitano, F. Nguyen, M. Palutan, E. Pasqualucci, A. Passeri, V. Patera, F. Peretto, L. Pontecorvo, M. Primavera, P. Santangelo, E. Santovetti, G. Saracino, B. Sciascia, A. Sciubba, F. Scuri, I. Sfiligoi, A. Sibidanov, T. Spadaro, M. Testa, L. Tortora, P. Valente, B. Valeriani, G. Venanzoni, S. Veneziano, A. Ventura, R. Versaci, G. Xu.

[†]Speaker.

The most precise verification of the unitarity of the CKM mixing matrix is obtained today from the V_{us} and V_{ud} values, neglecting $|V_{ub}|^2 \sim 0.00002$. With the KLOE detector we can measure all experimental inputs to V_{us} : branching ratios, lifetimes, and form factors.

1. DAΦNE and KLOE

In the DAΦNE e^+e^- collider, beams collide at a center-of-mass energy $W \sim M(\phi)$. Since 2001, KLOE has collected an integrated luminosity of $\sim 2 \text{ fb}^{-1}$. Results presented below are based on 2001-02 data for $\sim 450 \text{ pb}^{-1}$.

The KLOE detector consists of a large cylindrical drift chamber surrounded by a lead/scintillating-fiber electromagnetic calorimeter. A superconducting coil around the calorimeter provides a 0.52 T field. The drift chamber, Ref. [1], is 4 m in diameter and 3.3 m long. The momentum resolution is $\sigma(p_T)/p_T \sim 0.4\%$. Two track vertices are reconstructed with a spatial resolution of $\sim 3 \text{ mm}$. The calorimeter, Ref. [2], composed of a barrel and two endcaps, covers 98% of the solid angle. Energy and time resolution are $\sigma(E)/E = 5.7\%/\sqrt{E(\text{GeV})}$ and $\sigma(t) = 54 \text{ ps}/\sqrt{E(\text{GeV})} \oplus 50 \text{ ps}$. The KLOE trigger, Ref. [3], uses calorimeter and drift chamber information. For the present analysis only the calorimeter signals are used. Two energy deposits above threshold, $E > 50 \text{ MeV}$ for the barrel and $E > 150 \text{ MeV}$ for the endcaps, are required.

2. Kaon tagging

The ϕ meson decays mainly into kaons: 49% to K^+K^- and 34% to K_L - K_S pairs. We can thus tag K_L , K_S , K^+ , and K^- beams by detecting respectively K_S , K_L , K^- , and K^+ decays. Tagging allows measurements of absolute branching ratios. A K_S beam is tagged using events with a K_L interaction in the calorimeter. KLOE results on K_S decays are given in these proceedings [4]. K_L -mesons are tagged detecting $K_S \rightarrow \pi^+\pi^-$ decays. Charged kaons are tagged using two-body decays, $K^\pm \rightarrow \mu^\pm \nu$ and $K^\pm \rightarrow \pi^\pm \pi^0$. Perfect tagging requires that the detection efficiency of the tagging mode be independent of the decay mode of the tagged kaon. In reality, some dependency of the tagging efficiency on the decay mode of the signal kaon exists. This dependence must be carefully measured using Monte Carlo (MC) and data control samples for each BR measurement.

3. Semileptonic K^\pm decays

The measurement of the branching ratios for the K^\pm semileptonic decays is performed using four samples defined by different decay modes for the tagging kaon: $K_{\mu 2}^+$, $K_{\pi 2}^+$, $K_{\mu 2}^-$, and $K_{\pi 2}^-$. This redundancy allows the systematic effects due to the tag selection to be kept under control. Kaons are identified as tracks with momentum $70 < p < 130 \text{ MeV}$, originating from the collision point. The kaon decay vertex must be within a fiducial volume (FV) defined as a cylinder of radius $40 < r < 150 \text{ cm}$, centered at the collision point, coaxial with the beams. The decay track, extrapolated to the calorimeter, must overlap an appropriate energy deposit. $K_{\mu 2}$ and $K_{\pi 2}$ decays are selected by applying a 3σ cuts around the kaon rest-frame momentum distribution calculated in the relevant mass hypothesis. For the $K_{\pi 2}^\pm$ tag, identification of the π^0 from the vertex is also required. Finally, to reduce the dependency of the tag selection efficiency on the decay mode of the

signal kaon, the tagging decay is required to satisfy the calorimeter trigger. In the analyzed data set about 60 million tag decays were identified and divided into the four tag samples. To select a semileptonic decay on the signal side, a one-prong kaon decay vertex must be present in the FV. The daughter track has to reach the calorimeter and to overlap an energy deposit. Two-body decays are rejected by requiring that the decay momentum in the kaon frame, computed assuming the pion mass, is < 195 MeV. The lepton mass is obtained from the velocity of the lepton computed from the time of flight. The number of K_{e3} and $K_{\mu3}$ decays is then obtained by fitting the m_{lept}^2 distribution to a sum of MC distributions for the signals and various background sources, each multiplied by a free scale factor. An example of the lepton mass distribution is shown in fig. 1, left. The BR

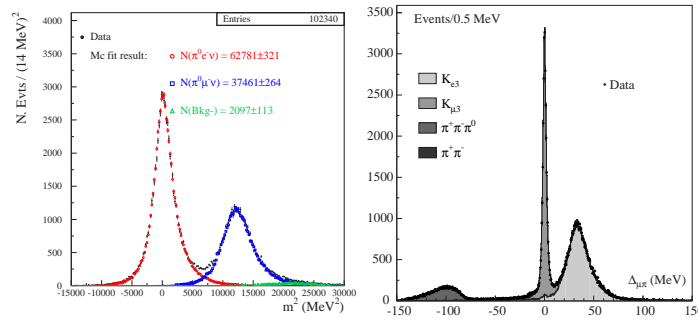


Figure 1: Left. Lepton mass distribution. Right. $\Delta_{\mu\pi}$ distribution.

is evaluated separately for each tag sample, dividing by the number of tag counts and correcting for acceptances. The latter are obtained from MC simulations. Corrections are applied to account for data-MC differences in tracking and cluster. About 190 000 K_{e3}^{\pm} and 100 000 $K_{\mu3}^{\pm}$ decays are selected. The resulting BR's are:

$$BR(K_{e3}^{\pm}) = (5.047 \pm 0.046_{Stat+Tag} \pm 0.080_{Syst})\%$$

$$BR(K_{\mu3}^{\pm}) = (3.310 \pm 0.040_{Stat+Tag} \pm 0.070_{Syst})\%.$$

These values are averages over the four different tag samples for each channel, and have been calculated with correlations carefully taken into account. The error is dominated by the error on data/MC efficiency corrections and the systematic error from the signal selection efficiency is still preliminary.

4. K^{\pm} lifetime

The K^{\pm} lifetime is an experimental input to the determination of V_{us} . The present fractional uncertainty is about 0.2%, corresponding to an uncertainty of 0.1% for V_{us} . However, see Ref. [7], there are large discrepancies between results of different experiments. The value of τ_{\pm} also affects the geometrical acceptance in BR measurements. KLOE is nearing completion of a new, high-statistics τ_{\pm} measurement using two different methods: one based on the measurement of the decay length and the other on the decay time of the kaons. Both methods have fractional errors at the few per mil level; comparison between the methods allows part of the systematic error to be assessed.

5. Semileptonic K_L decays and the K_L lifetime

KLOE has also measured the dominant K_L branching ratios using the K_L beam tagged by $K_S \rightarrow \pi^+ \pi^-$ decays [6]. About 13×10^6 tagged K_L decays are used for the measurement, and $\sim 4 \times 10^6$ to evaluate efficiencies. For K_{e3} , $K_{\mu 3}$, and $\pi^+ \pi^- \pi^0$ decays we compute smallest of the two values of $\Delta = E_{\text{miss}} - p_{\text{miss}}$ assuming each charged particle to be pion and muon (muon and pion). The distribution in this variable shows three well separated peaks corresponding to the three modes mentioned, see Fig. 1, right. Fitting the distribution with MC obtained shapes for each mode, gives the number of events in each channel. To select $K_L \rightarrow 3\pi^0$ events, at least three photons are required from the K_L decay vertex. The reconstruction efficiency and purity of the selected sample are both about 99%.

The resulting BR's are: $\text{BR}(K_L \rightarrow \pi e \nu(\gamma)) = 0.4007 \pm 0.0006 \pm 0.0014_{\text{Tag+Trk}}$, $\text{BR}(K_L \rightarrow \pi \mu \nu(\gamma)) = 0.2698 \pm 0.0006 \pm 0.0014_{\text{Tag+Trk}}$, $\text{BR}(K_L \rightarrow 3\pi^0) = 0.1997 \pm 0.0005 \pm 0.0019_{\text{Tag+\gammacount}}$ and, $\text{BR}(K_L \rightarrow \pi^+ \pi^- \pi^0(\gamma)) = 0.1263 \pm 0.0005 \pm 0.0011_{\text{Tag+Trk}}$, after imposing the constraint $\sum \text{BR}(K_L) = 1$. This corresponds to also measuring the lifetime by counting the number of decays in a time interval for a beam of known intensity.

The K_L lifetime has been also measured directly [8], employing 10^7 $K_L \rightarrow 3\pi^0$ events. The result is $\tau_L = (50.92 \pm 0.17 \pm 0.25)$ ns, which together with that from the K_L BR measurements gives the KLOE average: $\tau_L = (50.84 \pm 0.23)$ ns

6. Determination of V_{us}

The BR's of the semileptonic K_L decays, together with the result $\text{BR}(K_S \rightarrow \pi e \nu(\gamma)) = (7.09 \pm 0.09) \times 10^{-4}$ and the preliminary results on the semileptonic K^\pm decays, allow five independent determinations of the observable $|V_{us} f_+(0)|$, as shown in Fig. 2, in which the new KLOE value of τ_L has been used to covert K_L BR's to partial widths. Averaging the five KLOE values gives $|V_{us} f_+(0)| = 0.2170 \pm 0.0005$, with $\chi^2/\text{DoF} = 1.7/4$. The fractional uncertainty is about 0.25%.

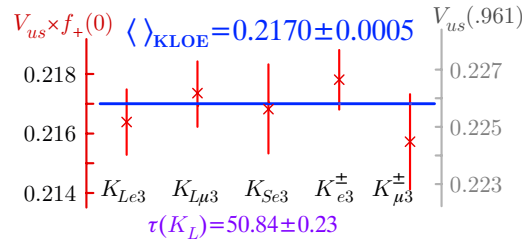


Figure 2: $|V_{us} f_+(0)|$ measurements. For K_L BR the KLOE τ_L has been used.

A precise estimate of $f_+(0)$, 0.961 ± 0.008 , was first given in 1984 [9]. Recent lattice calculations [10] give $f_+(0) = 0.960 \pm 0.009$, in agreement with Ref. [9]. Using the value from [9] and the average of our results for $|V_{us} f_+(0)|$ we find $V_{us} = 0.2258 \pm 0.0022$.

7. V_{us} from $\text{BR}(K^+ \rightarrow \mu^+ \nu(\gamma))$

KLOE has also measured the fully inclusive, absolute $K_{\mu 2}^+$ branching ratio. From about 9×10^5

$K_{\mu 2}^+$ decays obtained from a sample of about 250pb^{-1} KLOE obtains $BR(K^+ \rightarrow \mu^+ \nu(\gamma)) = 0.6366 \pm 0.0009_{Stat} \pm 0.0015_{Syst}$ [11], for an overall fractional error of 0.27%. Using recent lattice results on the decay constants of pseudoscalar mesons [5], we find $V_{us} = 0.2223 \pm 0.0026$. The fractional error of about 1% is dominated by the uncertainty in the f_K/f_π computation.

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