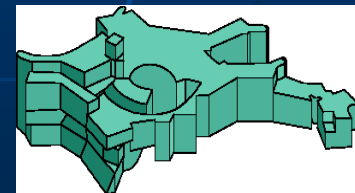


Supernova Theory

(and Cosmological Distances)

Wolfgang Hillebrandt
MPI für Astrophysik
Garching

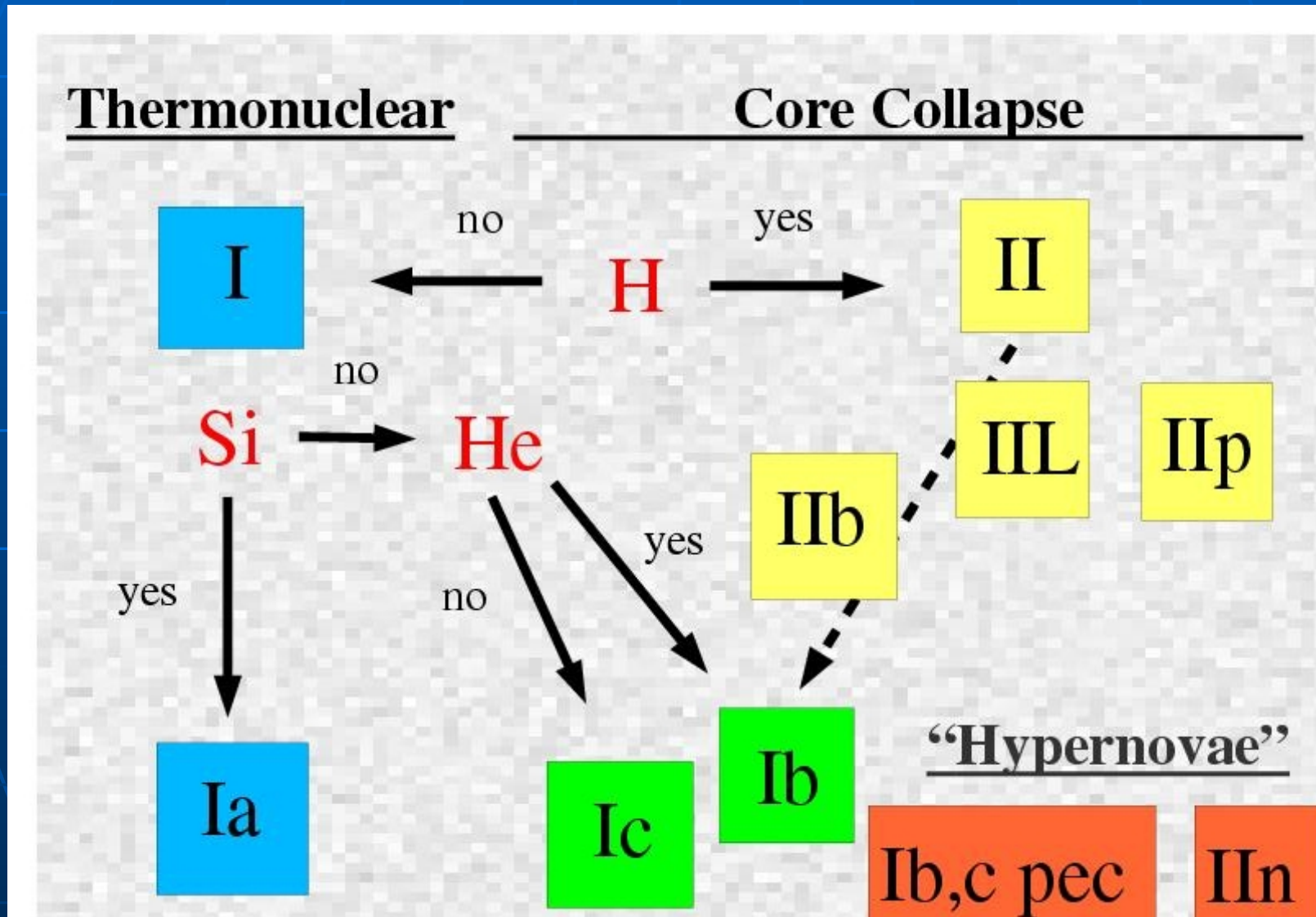
NIC School, CERN,
June 19 – 23, 2006



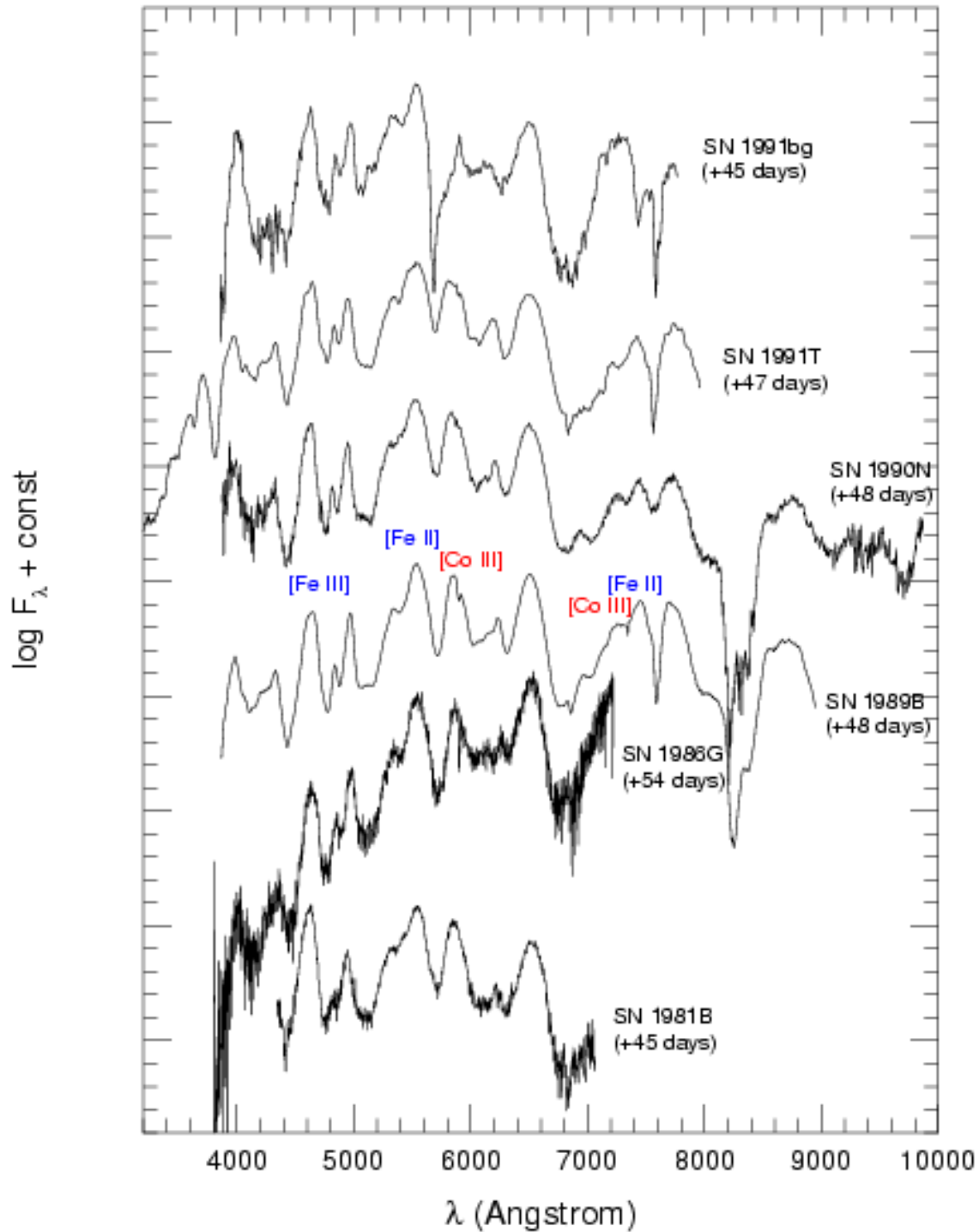
Outline of the lectures

- Supernova types and phenomenology
- Models of core-collapse supernovae (Type II; Type Ib,c; GRB's)
- Models of thermonuclear supernovae (Type Ia)
- Luminosity distances and supernova cosmology

Supernova classification



(Leibundgut et al. 1993)

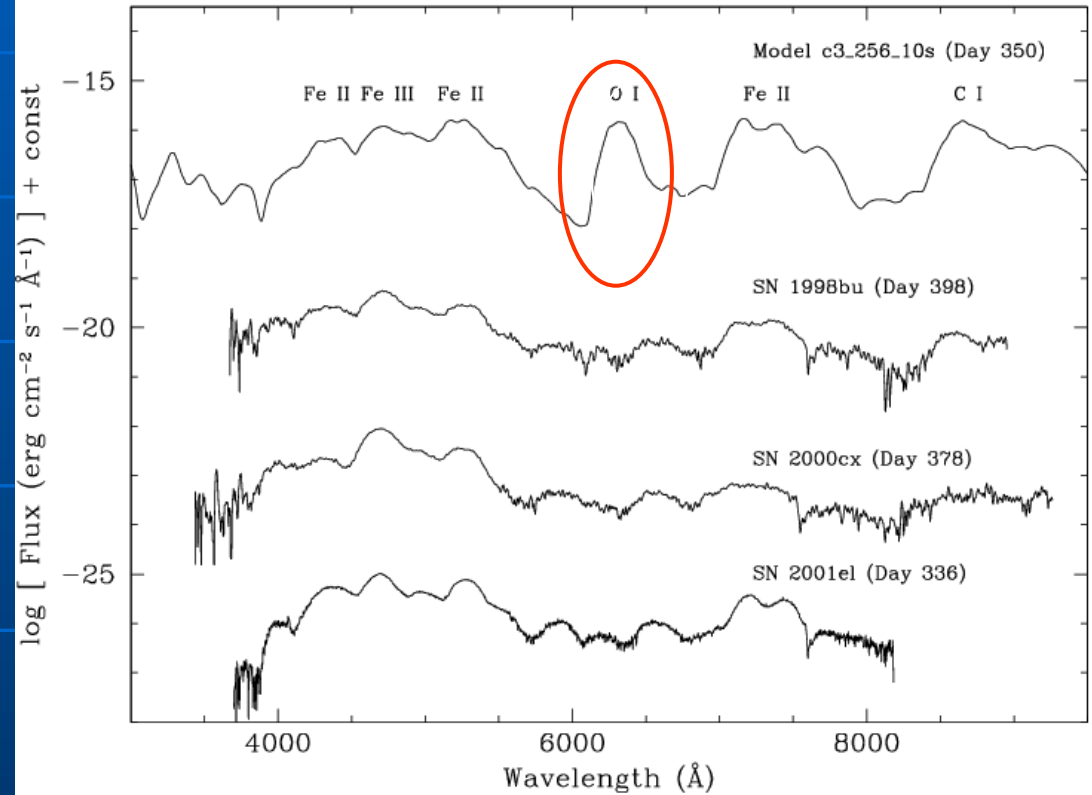
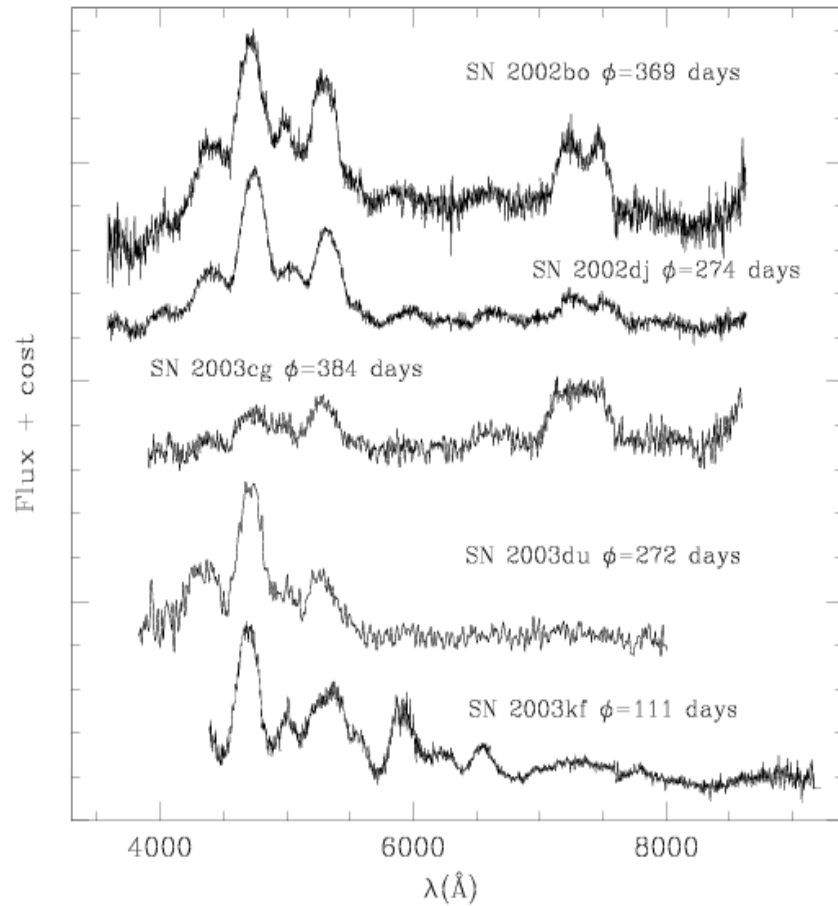


Supernova Spectroscopy

Type Ia

SN Ia: Nebular spectra

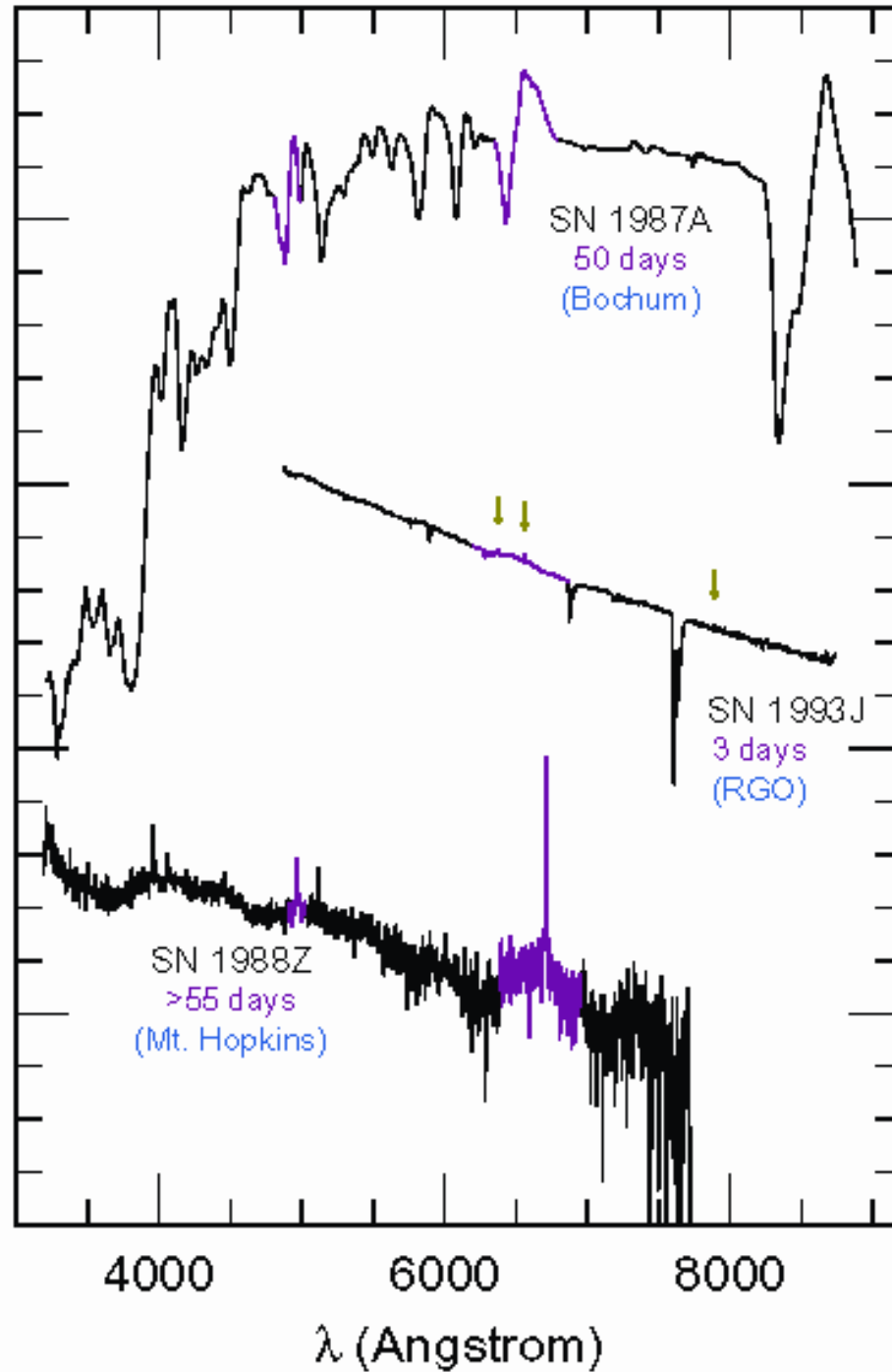
(Model: 3D Monte Carlo; Kozma et al. 2005)



Too much
oxygen at low
velocities!

Type II Supernovae

$\log F_\lambda + \text{const.}$

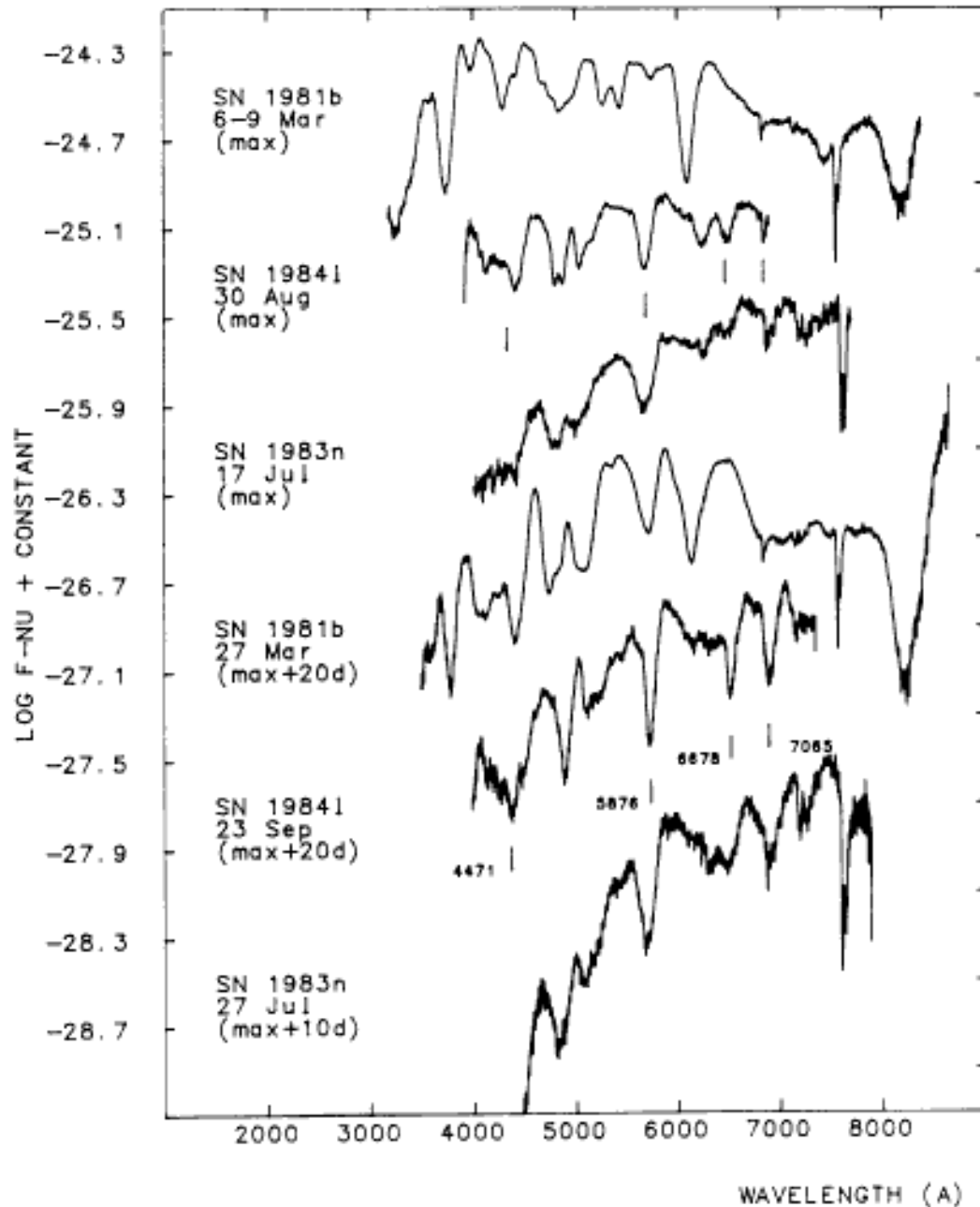


Supernova Spectroscopy

Type II

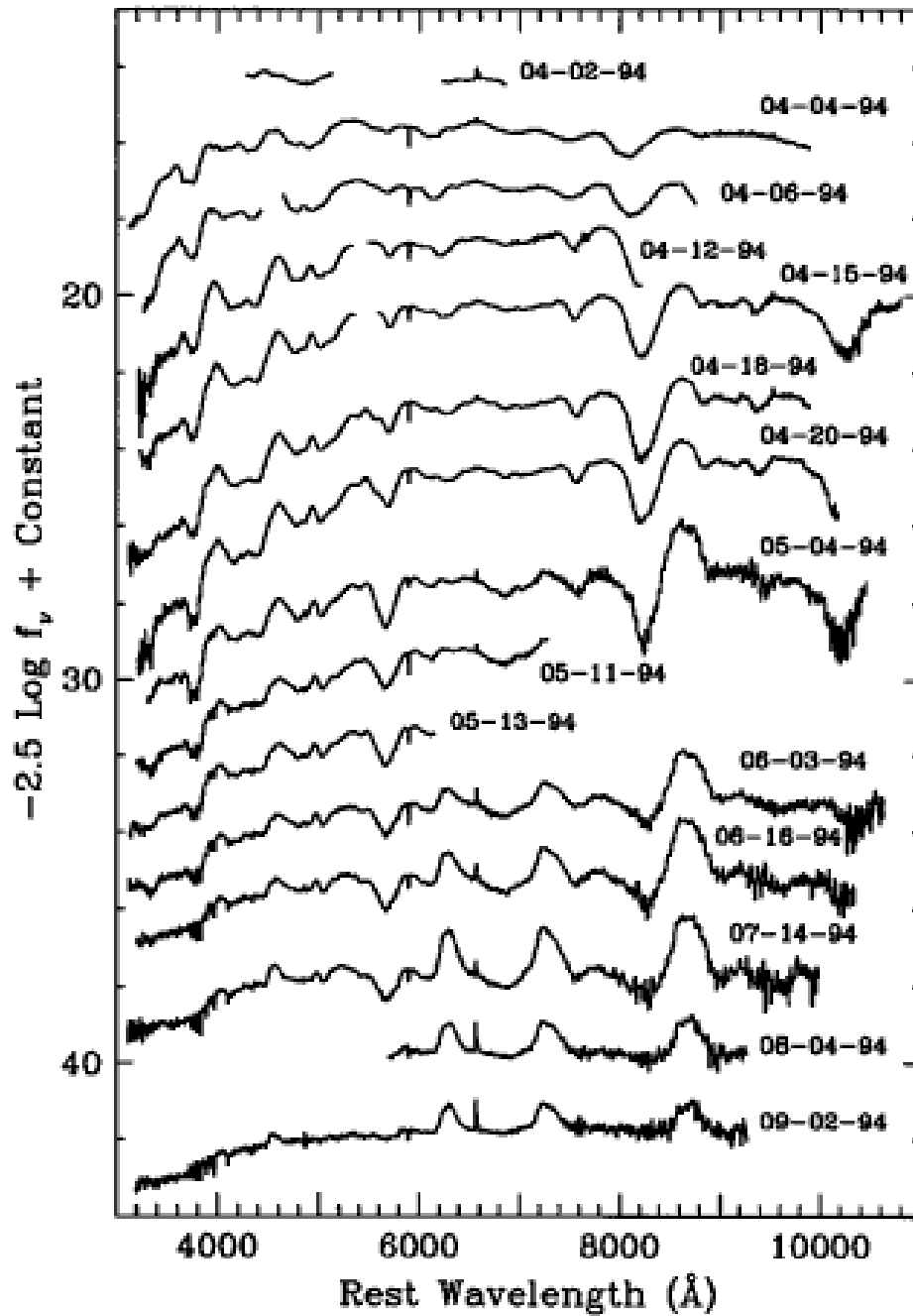
Supernova Spectroscopy

Type Ib



Supernova Spectroscopy

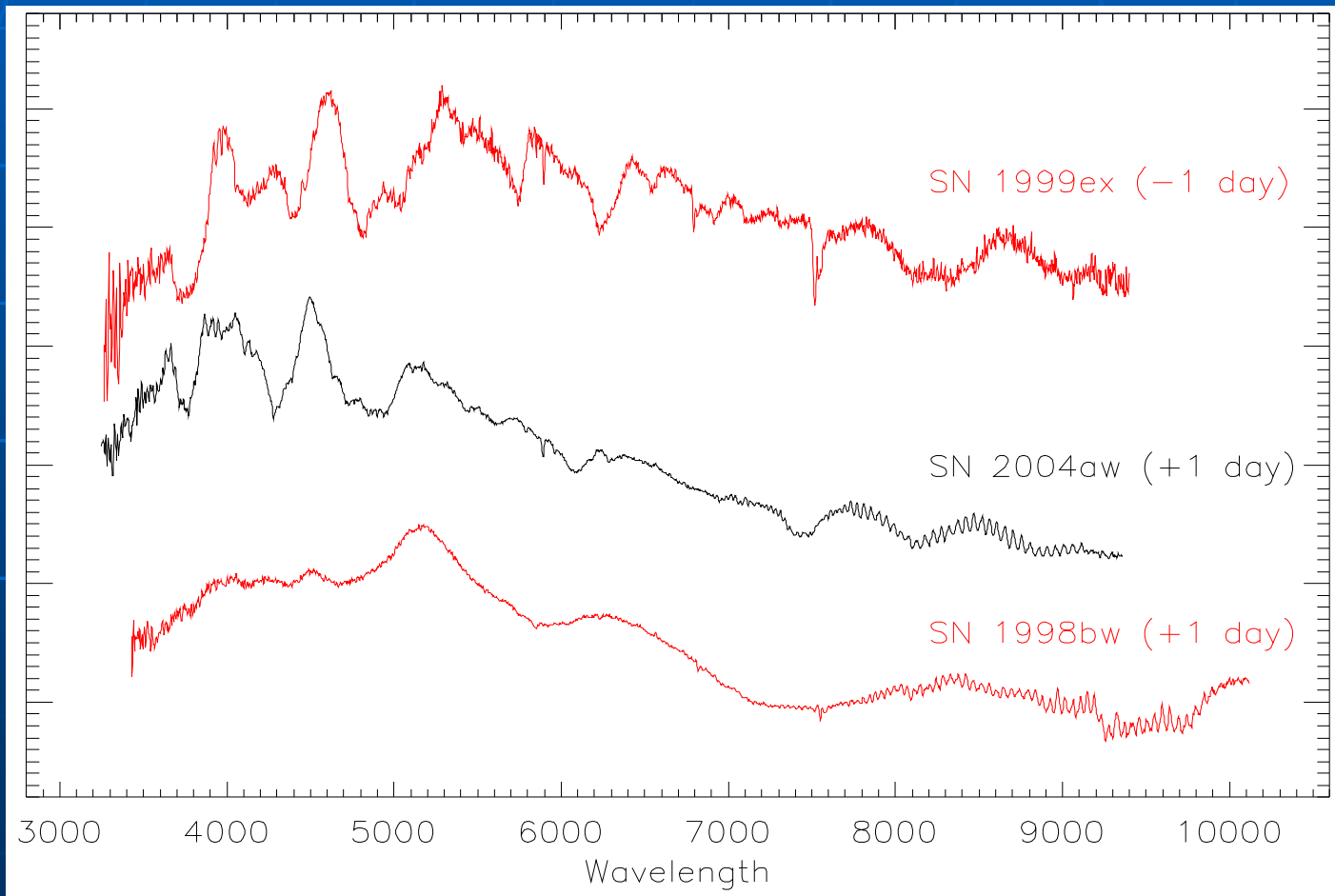
Type Ic



SN 1994I (Filippenko et al.)

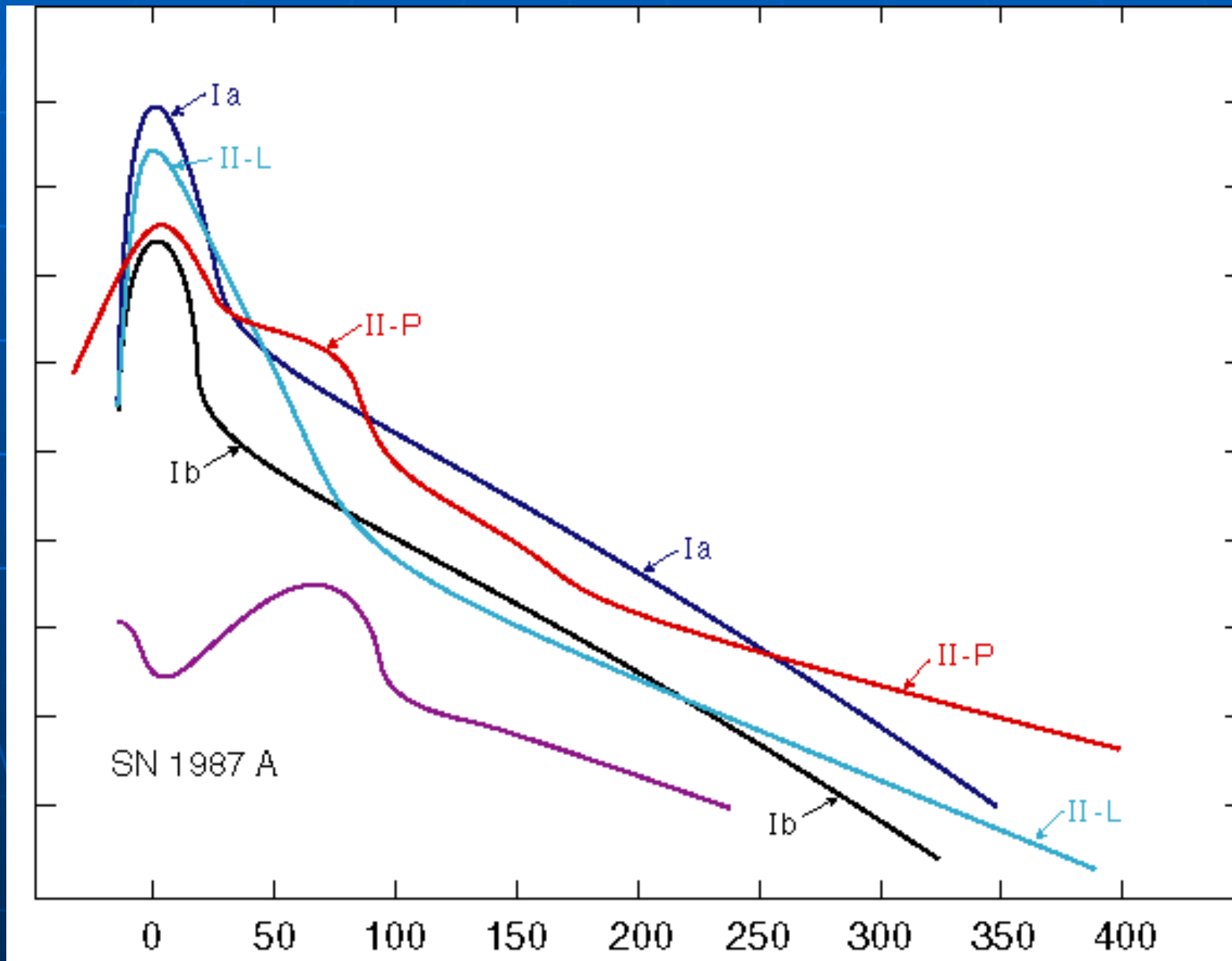
Supernova Spectroscopy

Type Ic pec



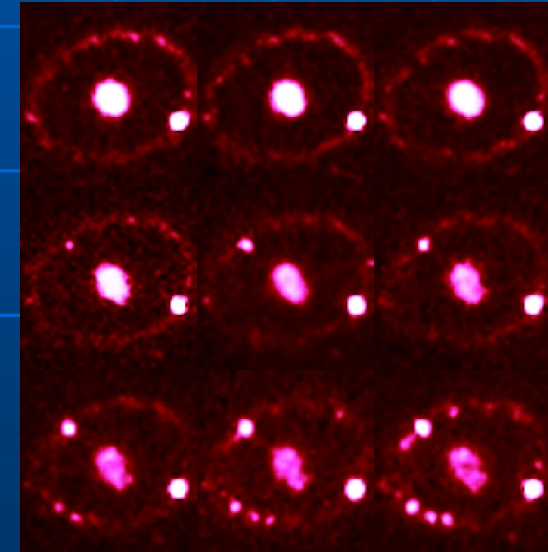
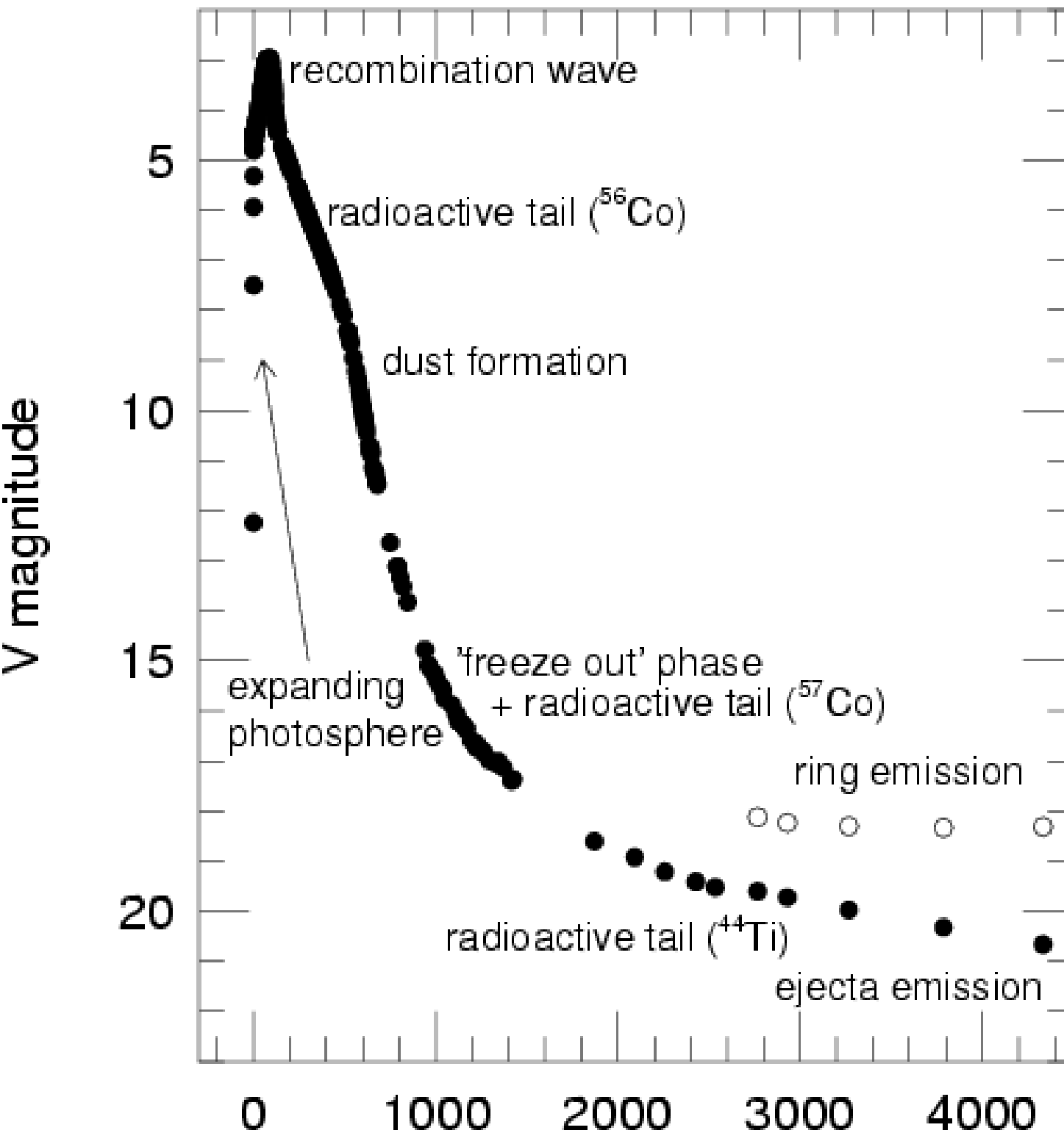
(Taubenberger et al. 2005)

Supernova light curves



l. (2001)

Core-collapse light curve



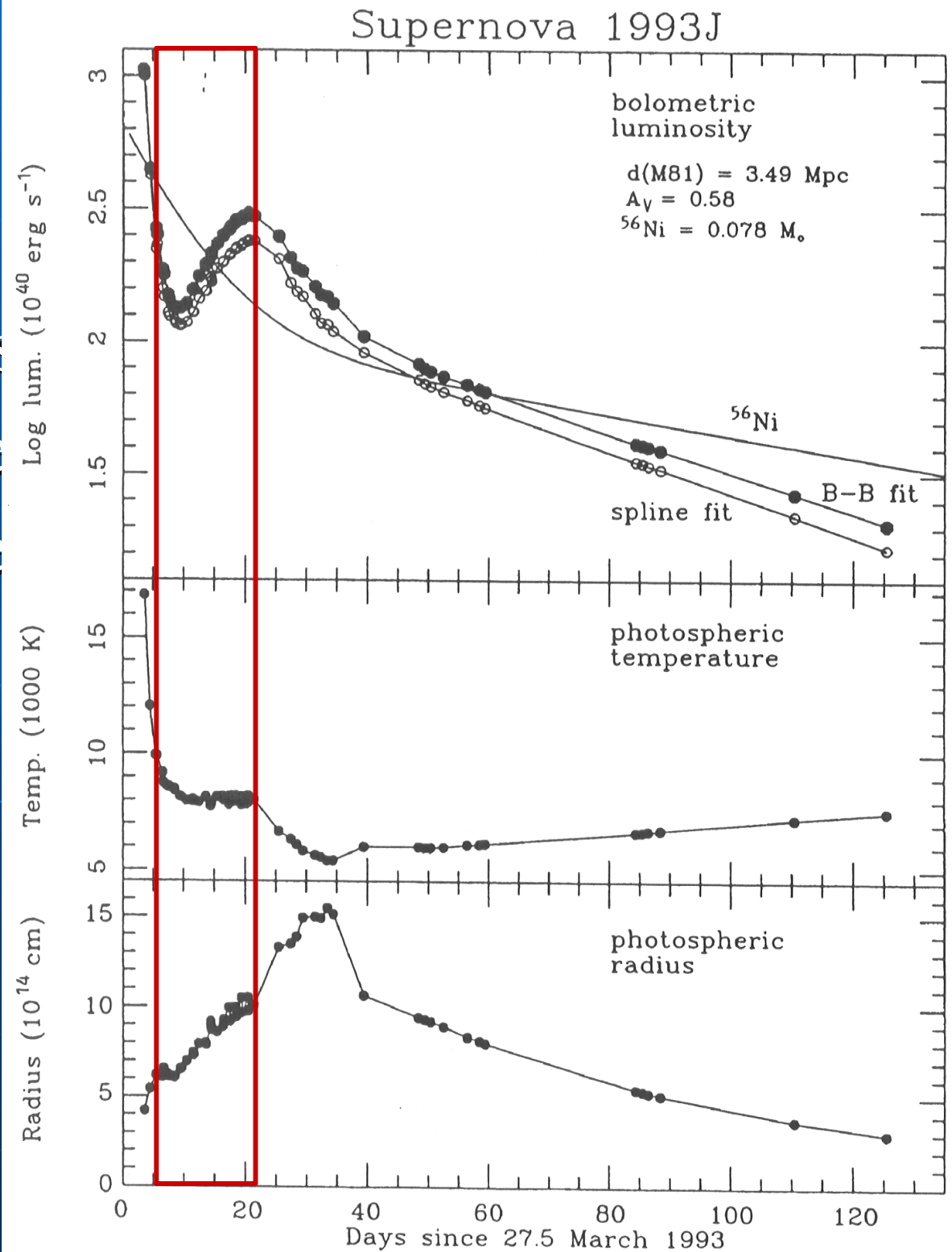
Suntzeff (2003)

Energy sources for light curves

- **shock**
 - breakout
 - kinetic energy
- **cooling**
 - due to expansion of the ejecta
- **radioactivity**
 - nucleosynthesis
- **recombination**
 - of the shock-ionised material

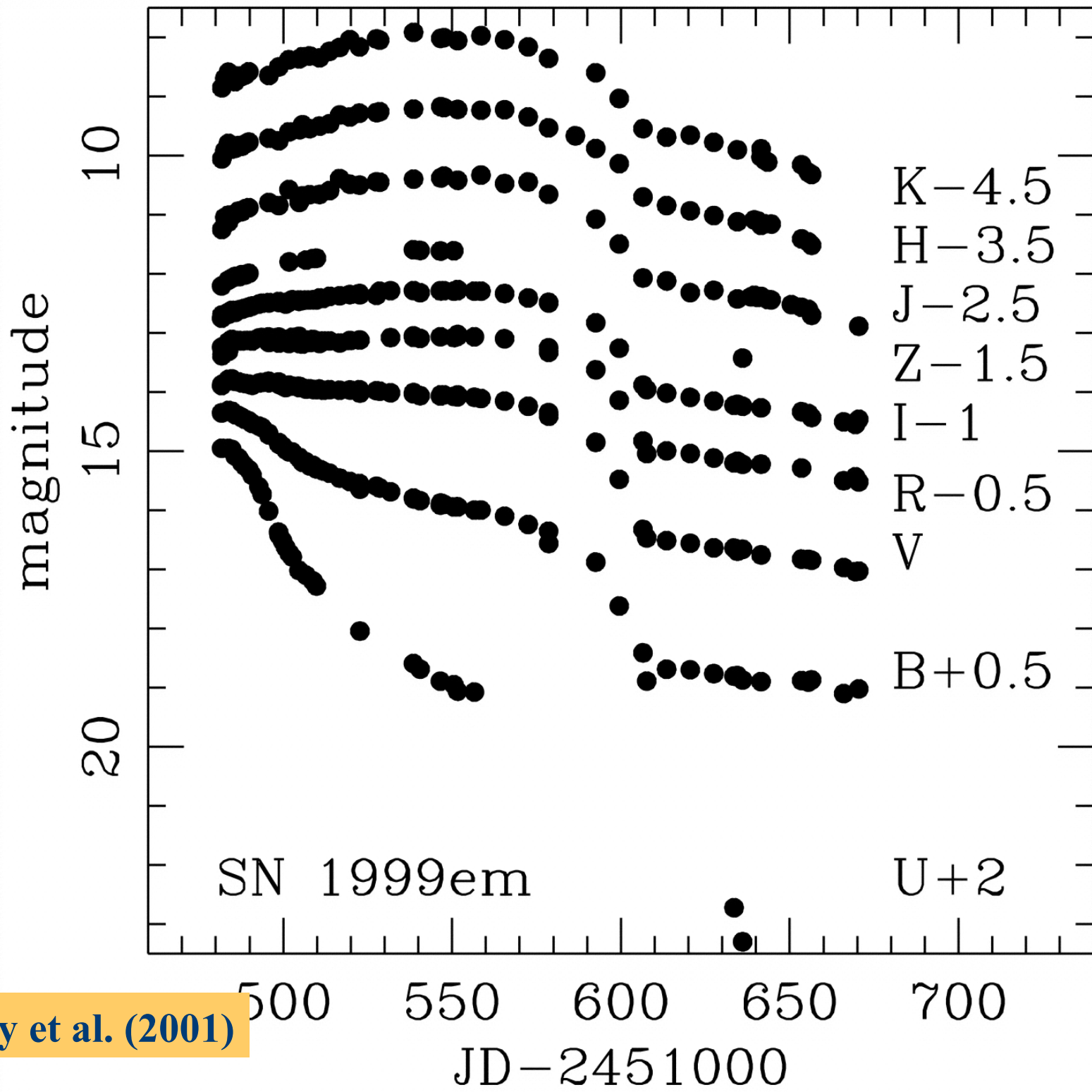
Expansion

- Brightness increases
 - increased surface area
 - slow temperature

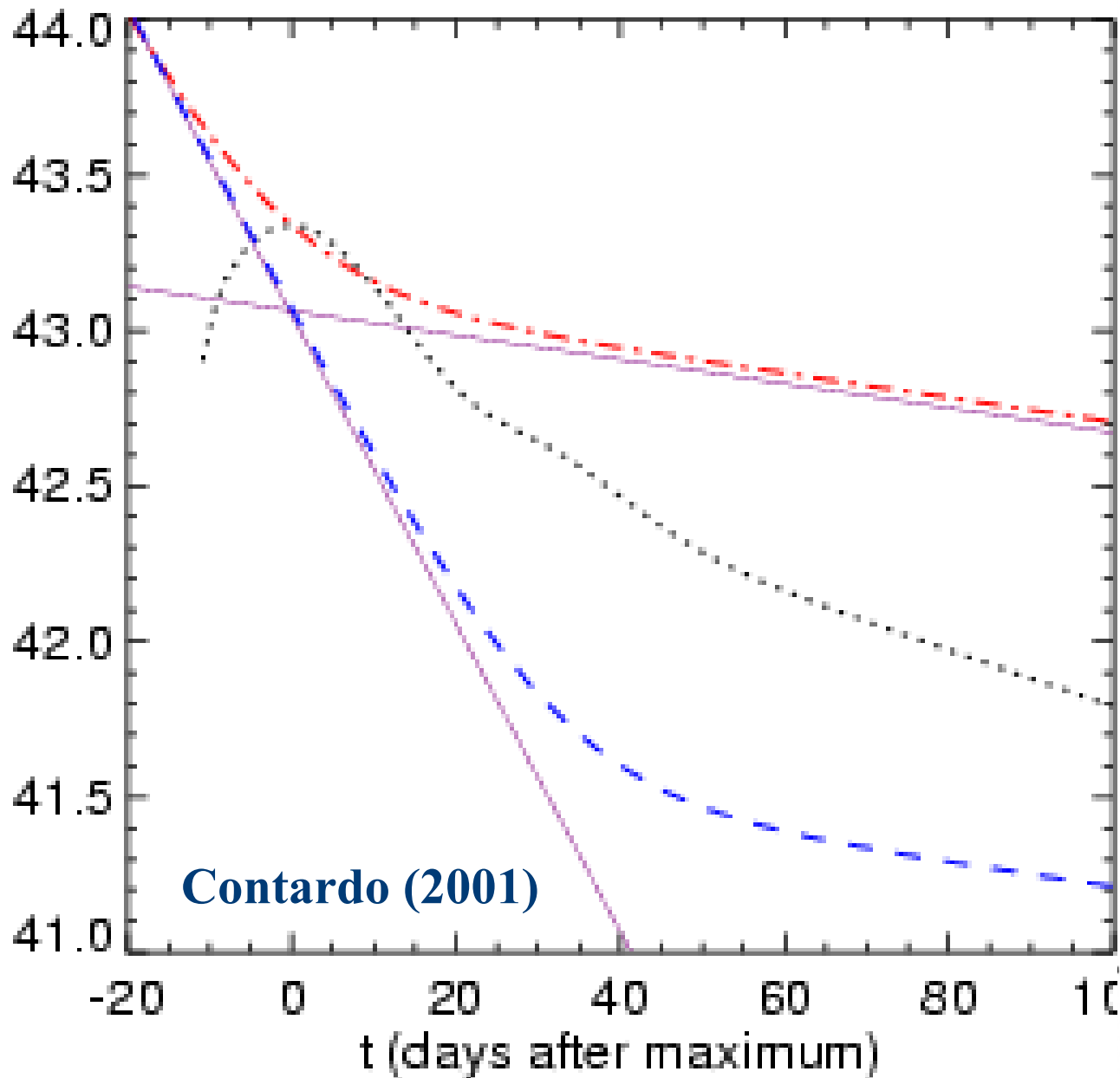
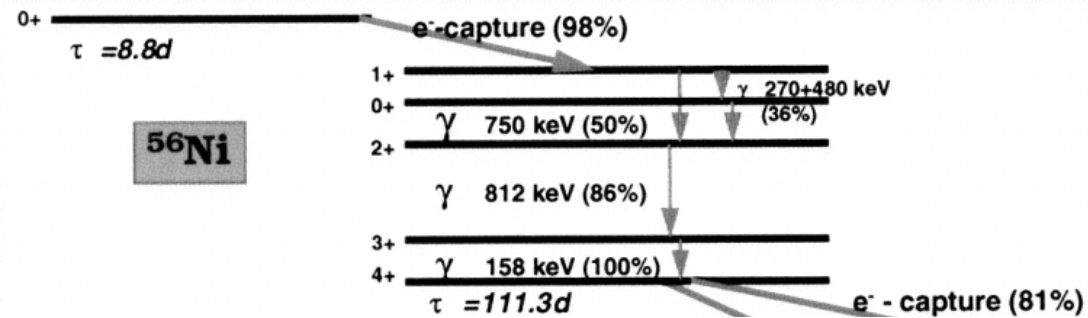


Recombination

- Balance of expansion
- leads to



Hamuy et al. (2001)



Radioactivity

- Isotopes of Ni and other elements
 - conversion of γ -rays and positrons into heat and optical photons

Supernova types: Summary

■ Thermonuclear SNe

- from low-mass stars ($<8M$)
- highly evolved stars (white dwarfs)
- explosive C and O burning
- binary systems required
- complete disruption

■ Core-collapse SNe

- high mass stars ($>8M$)
- large envelopes (still burning)
- burning due to compression
- single stars (binaries for SNe Ib/c?)
- neutron star

Core-collapse Supernovae

(In part “borrowed” from a lecture by Ewald Müller)



Prototype:

Crab nebula with pulsar
(constellation Orion)

Remnant of a supernova
observed in 1054

30 Doradus region in the Large Magellanic Cloud

(d ~ 160 000 light years)



© Anglo-Australian Observatory

Supernova 1987A
7:35 UT 23.2.1987

Blue Supergiant
Sanduleak 69.202

A few observational facts

(core collapse supernovae, i.e. SNe II, Ib, Ic)

very bright events:

$$L \sim 10^{10} L_{\text{sun}}$$

fast expanding ejecta:

$$v \sim 10^4 \text{ km/s}$$

energies: electromagnetic:

$$\sim 10^{49} \text{ erg}$$

kinetic:

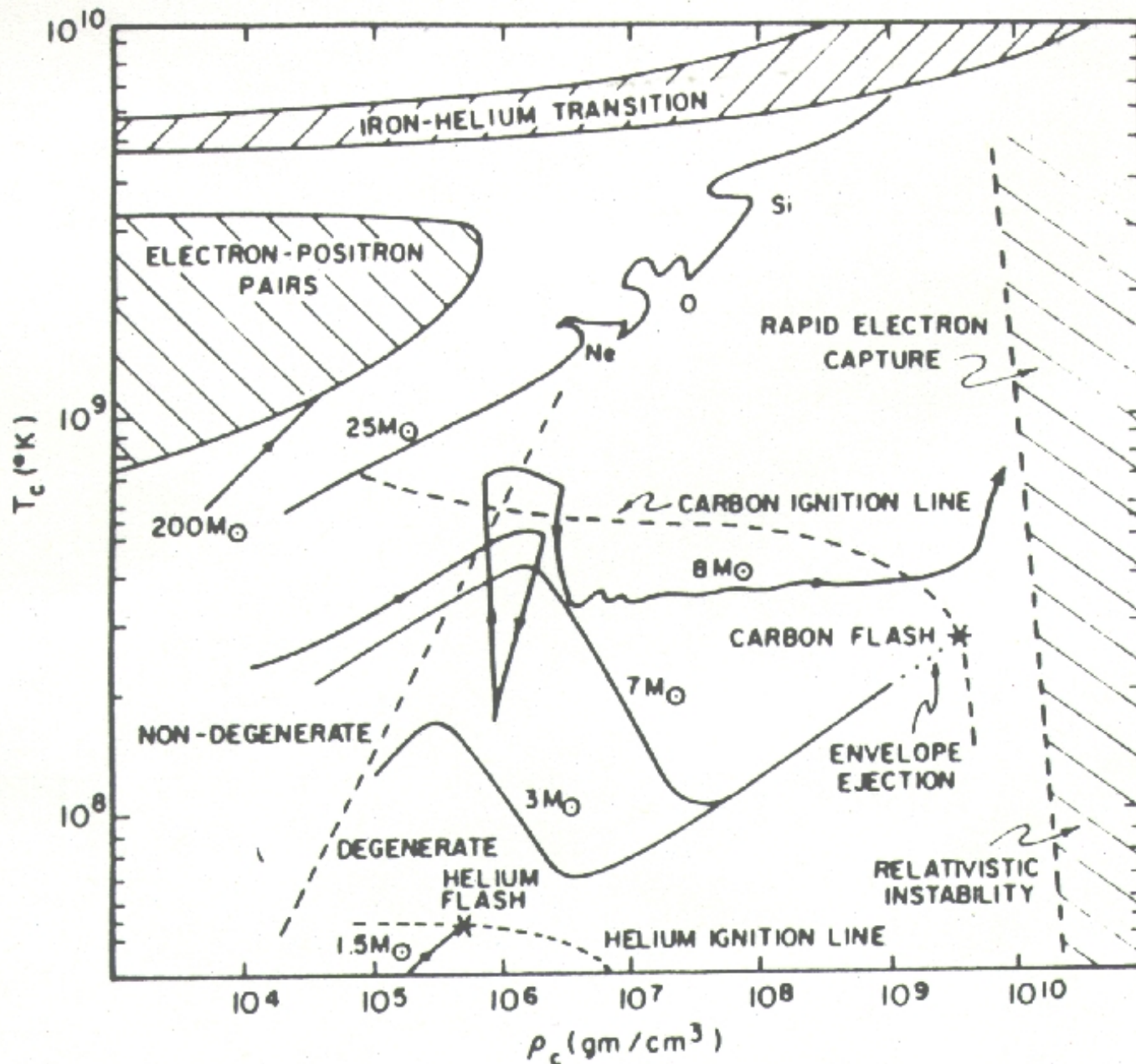
$$\sim 10^{51} \text{ erg}$$

neutrinos (SN1987A):

$$\sim 3 \cdot 10^{53} \text{ erg}$$

progenitor star destroyed (SN 1987A, SN 1993J)

Evolution of massive stars

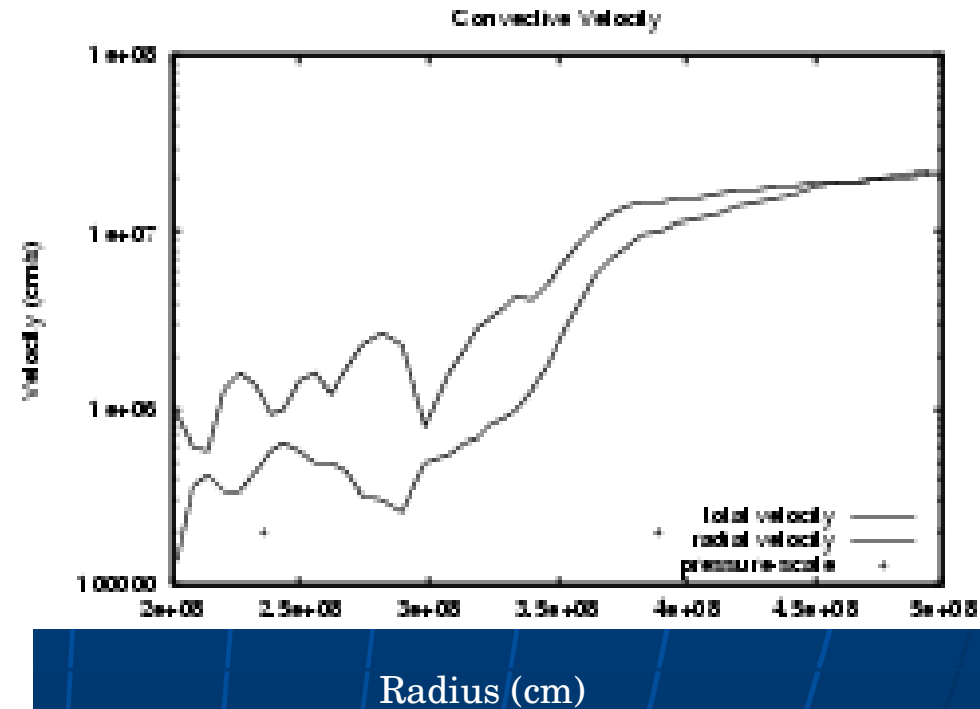
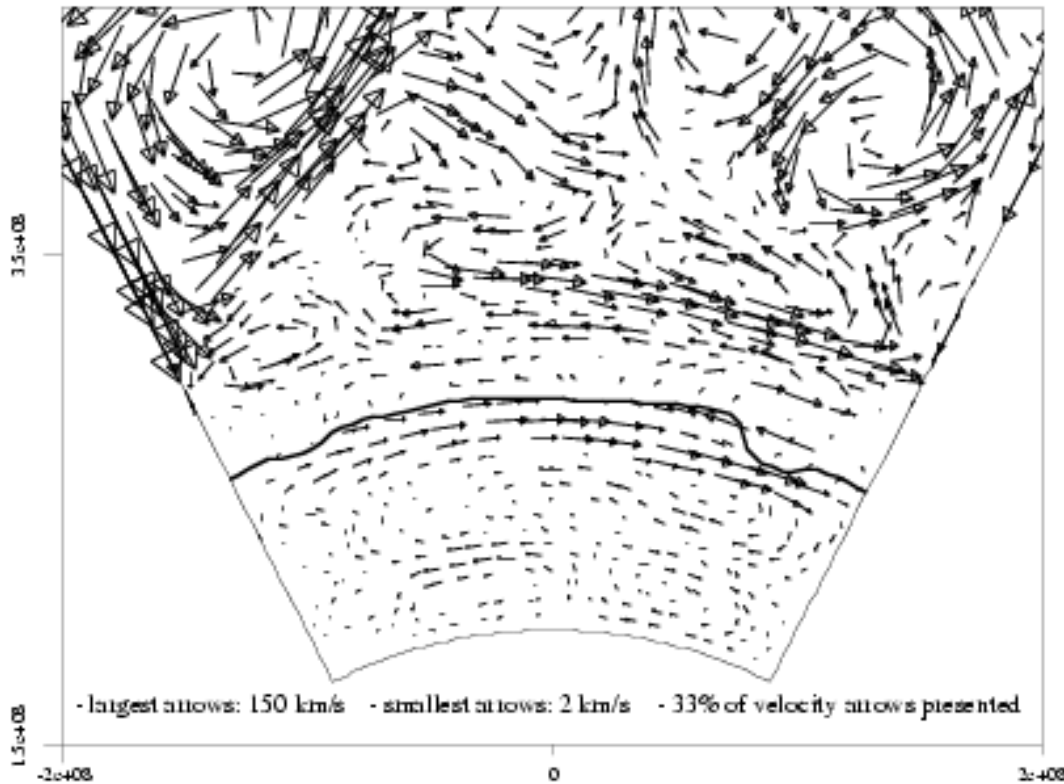


For models of the explosions:

"Fe"-core masses and their entropies have to be known!

Problems with massive star evolution:

Non-local time-dependent convection!

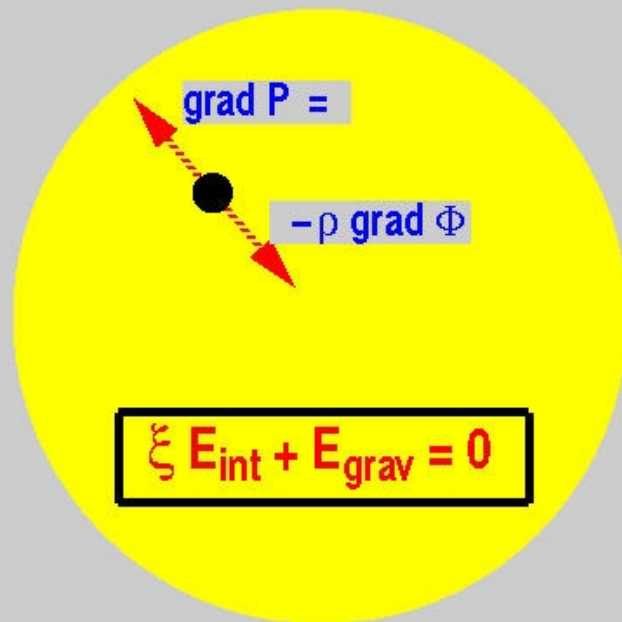


Hydrostatic O-burning
(Asida & Arnett, 2000)

(See also Brummell et al. 2002,
Rogers et al. 2003, ...)

Evolution towards gravitational collapse

stellar evolution mostly **hydrostatic**, i.e. pressure and gravitational forces are in equilibrium



virial theorem

$\xi := 3P/\rho u$ ideal gas: $P = (\gamma-1)\rho u \rightarrow \xi = 3(\gamma-1)$
 relativ. Fermi gas: $P = 1/3\rho u \rightarrow \xi = 1$

total energy:

$$W := E_{\text{int}} + E_{\text{grav}} = (1-\xi) E_{\text{int}} = (\xi-1)/\xi E_{\text{grav}}$$

if $\xi = 1 \rightarrow W = 0!$

Evolution towards gravitational collapse

gas: **finite temperature** \rightarrow star radiates

energy conservation:

$$dW / dt + L = 0$$

luminosity

$$L = (\xi-1) dE_{\text{int}} / dt = -(\xi-1)/\xi dE_{\text{grav}} / dt$$

if $L > 0 \rightarrow dE_{\text{grav}} / dt < 0 \leftrightarrow$

contraction $\rightarrow dE_{\text{int}} / dt > 0$

contraction with $\gamma = 5/3$ ($\xi = 2$):

50% of liberated energy are radiated away

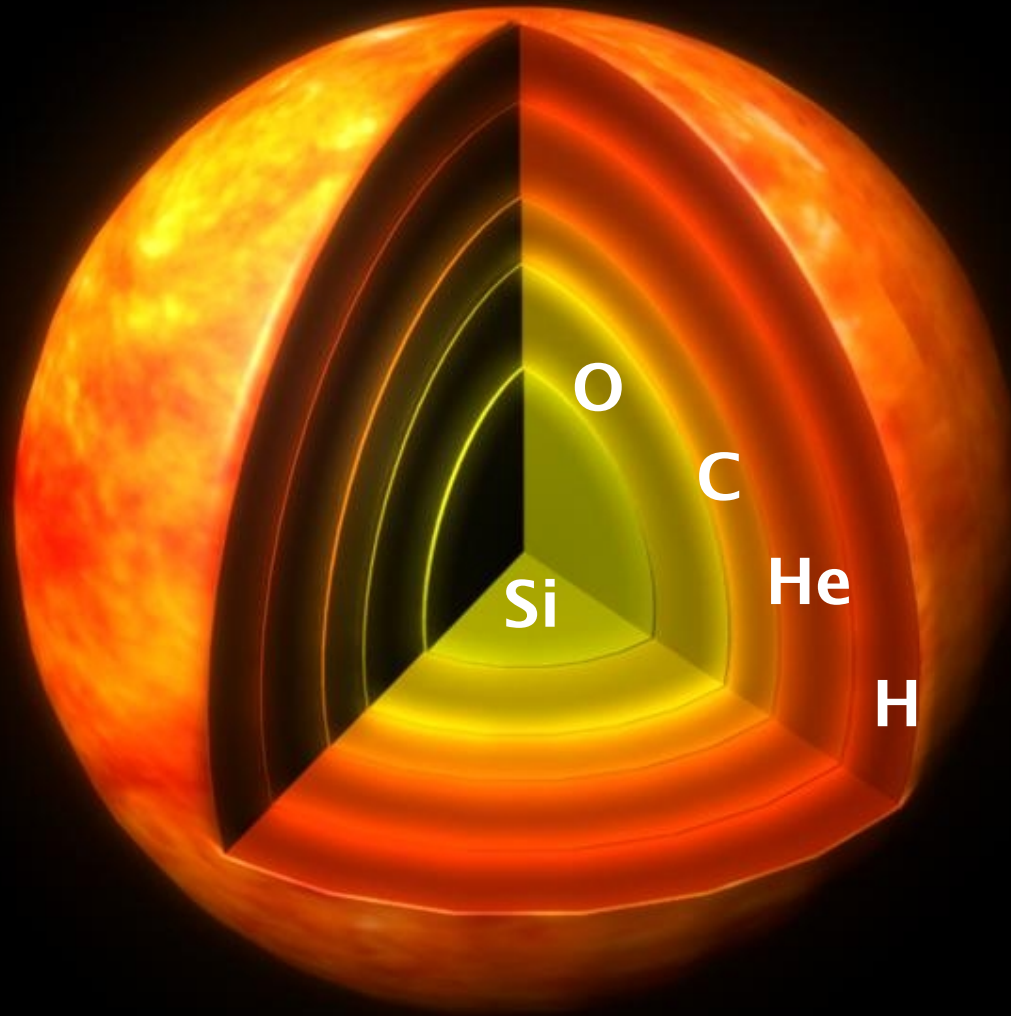
50% of liberated energy heat the star

\rightarrow **star has negative specific heat!**

Onion-like structure
of a presupernova
star several million
years after its birth:

mass: $10 \dots 10^2 M_{\text{sun}}$

radius: $50 \dots 10^3 R_{\text{sun}}$



- shells of different composition are separated by active thermonuclear burning shells

- core Si-burning leads to formation of central iron core

Note: figure not drawn to scale!

Energy sources for a core collapse supernova

Gravitational binding energy

Formation of a compact object of ~ 1 solar mass
with a radius ~ 10 km

$$\rightarrow E_b \sim 3 \times 10^{53} (M/M_{\text{sun}})^2 (R/10\text{km})^{-1} \text{ erg}$$

Fe-Ni core: $\rho \sim 10^{10} \text{ g/cm}^3$, $T \sim 10^{10} \text{ K}$

$\rightarrow P \sim P_e$ (relativistic degenerate Fermi gas)

\rightarrow maximum mass (Chandrasekhar)

Core becomes unstable due to:

- a) electron captures
- b) photo-disintegrations

Blue Giant (Red Giant: $\times 100$)

30 000 000 km

Fe-Ni core

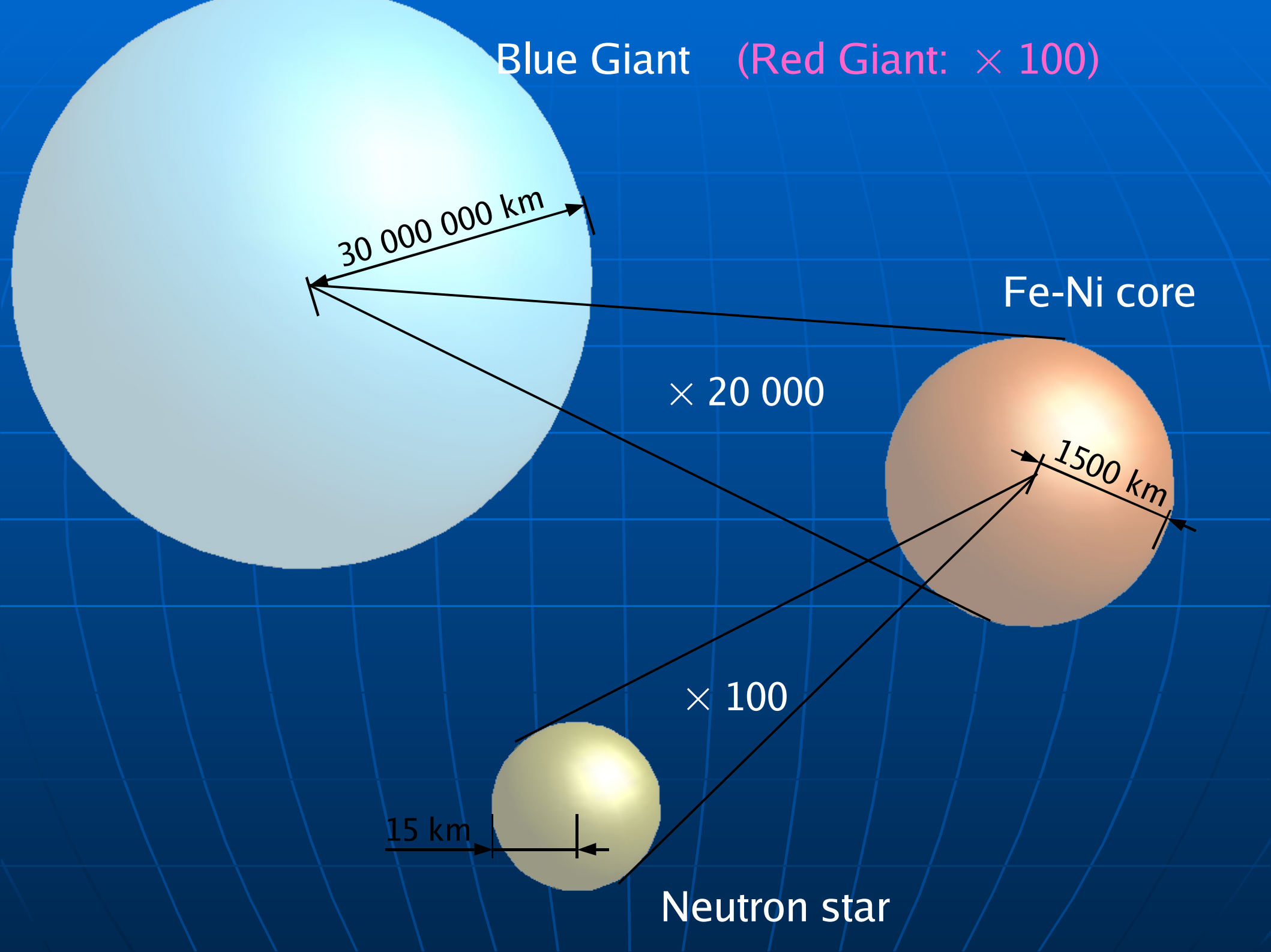
$\times 20\ 000$

1500 km

$\times 100$

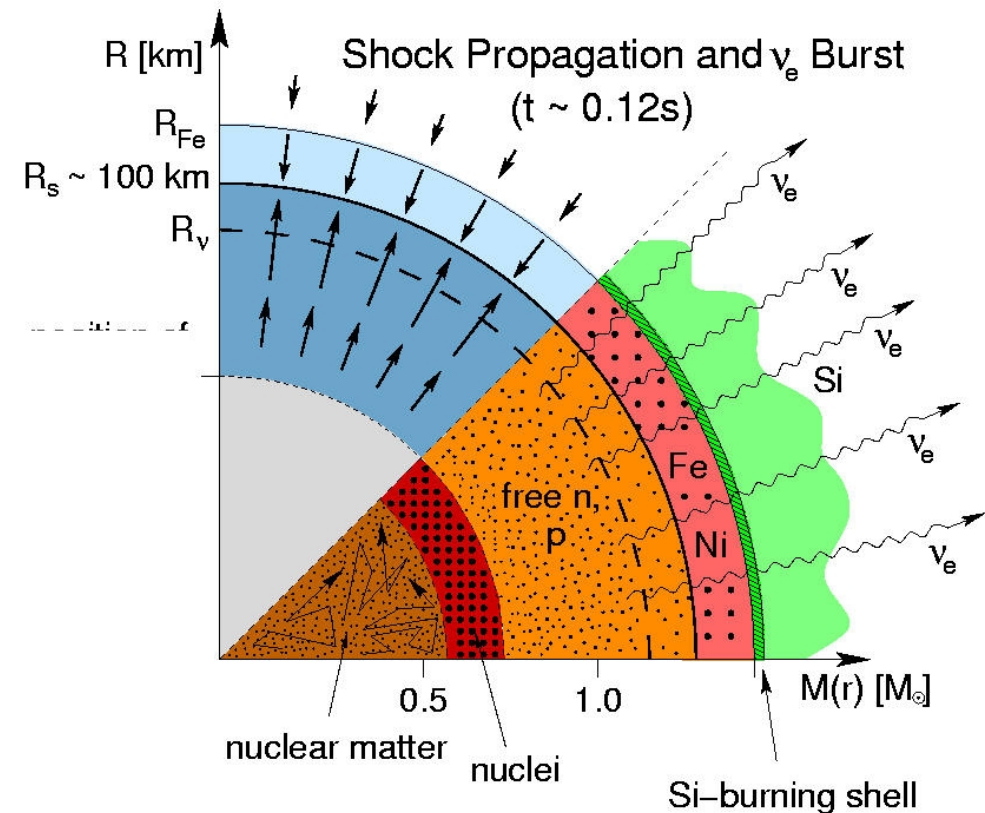
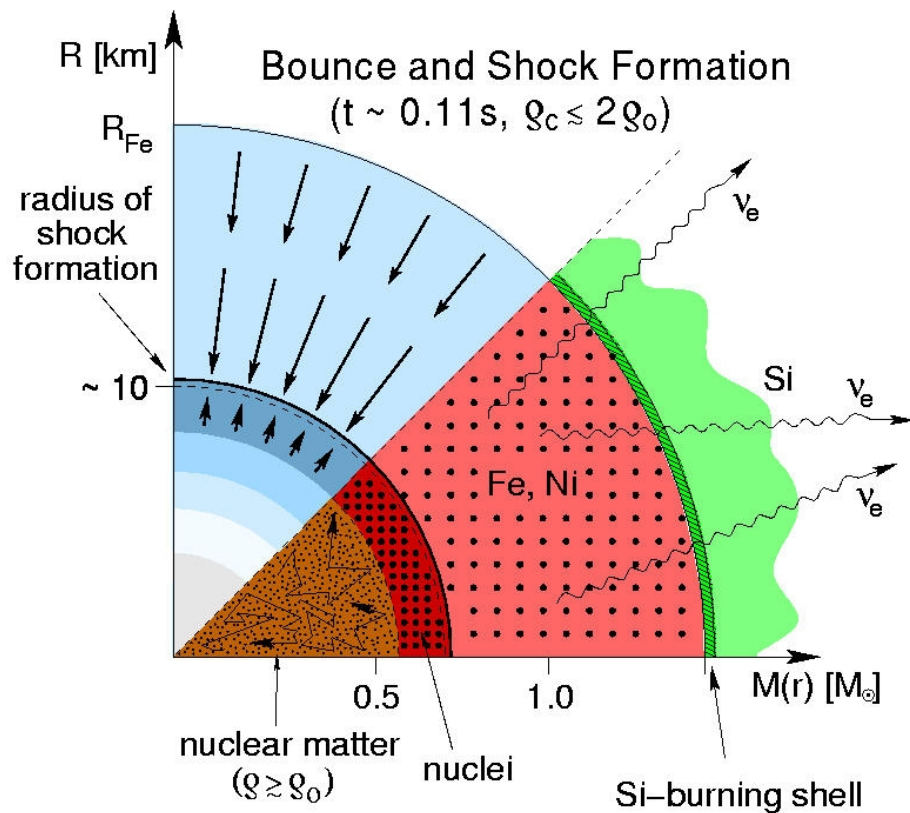
15 km

Neutron star



Core collapse supernovae:

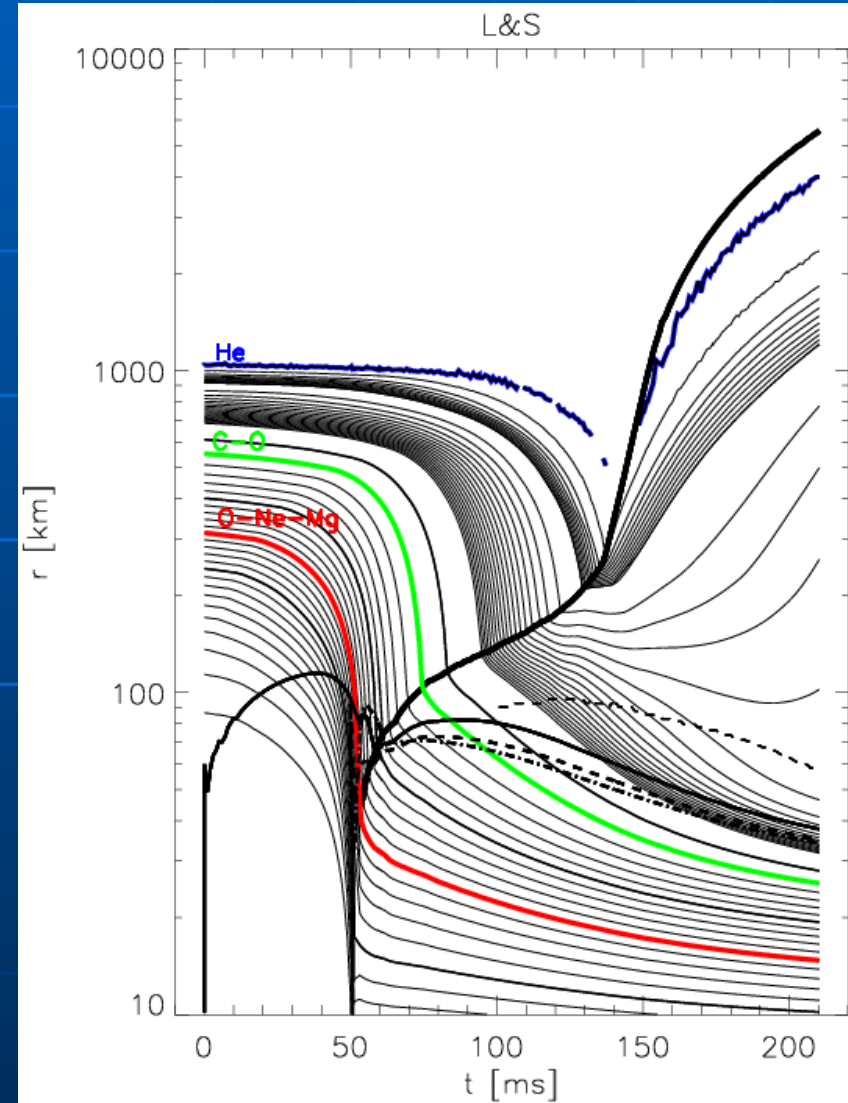
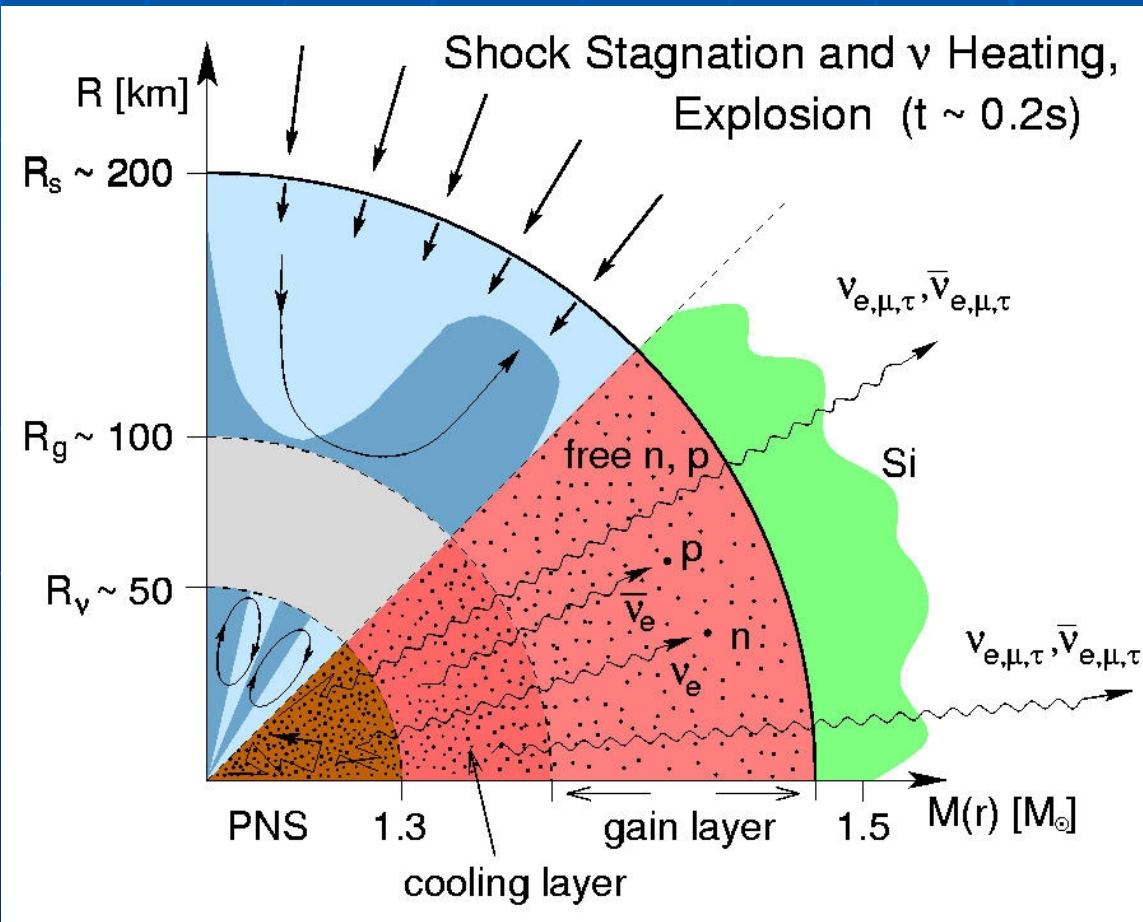
- Prompt explosion mechanism does not work (explored during the 1970's and 1980's; commonly accepted early 1990's)



- Shock wave forms close to sonic point ($M \sim 0.7 M_{\text{sun}}$)
initial energy: $(5 \dots 8) \times 10^{51}$

- Severe energy losses during shock propagation (8 MeV/nucleon or 1.6×10^{51} erg/ $0.1M_{\text{sun}}$)

- Current paradigm: neutrino-driven delayed explosions
 (discovered through computer simulations by Wilson '82, and
 first analyzed by Wilson & Bethe '95)



In its simplest form: Seems to work for low-mass core only! (Kitaura et al '05)

Seems to work for the low-mass end!

(F. Kitaura, Th.-J. Janka & WH, 2006)

Initial model $\sim 9 M_{\odot}$
(Nomoto 1984)

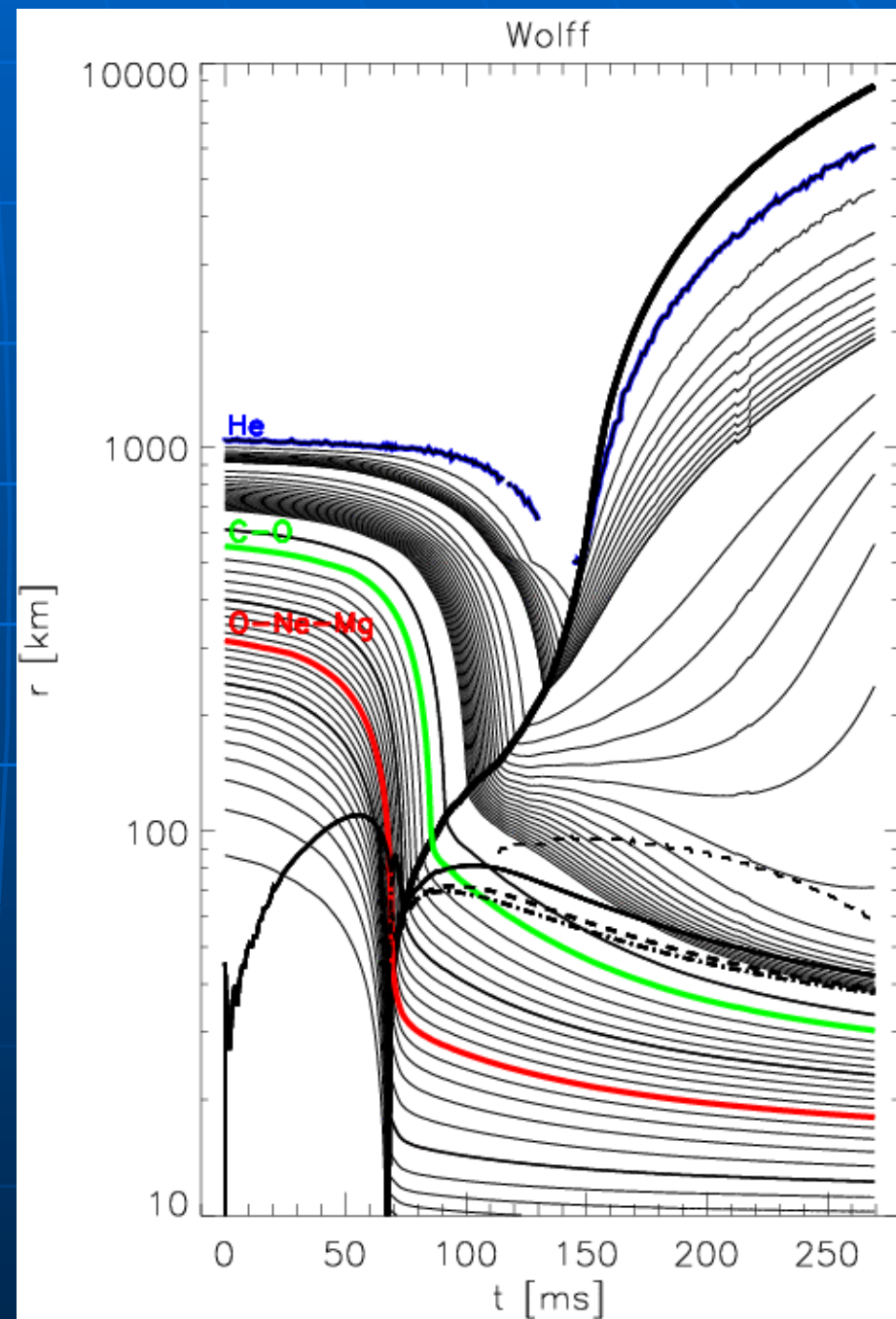
O-Ne-Mg core

1D, full neutrino transport

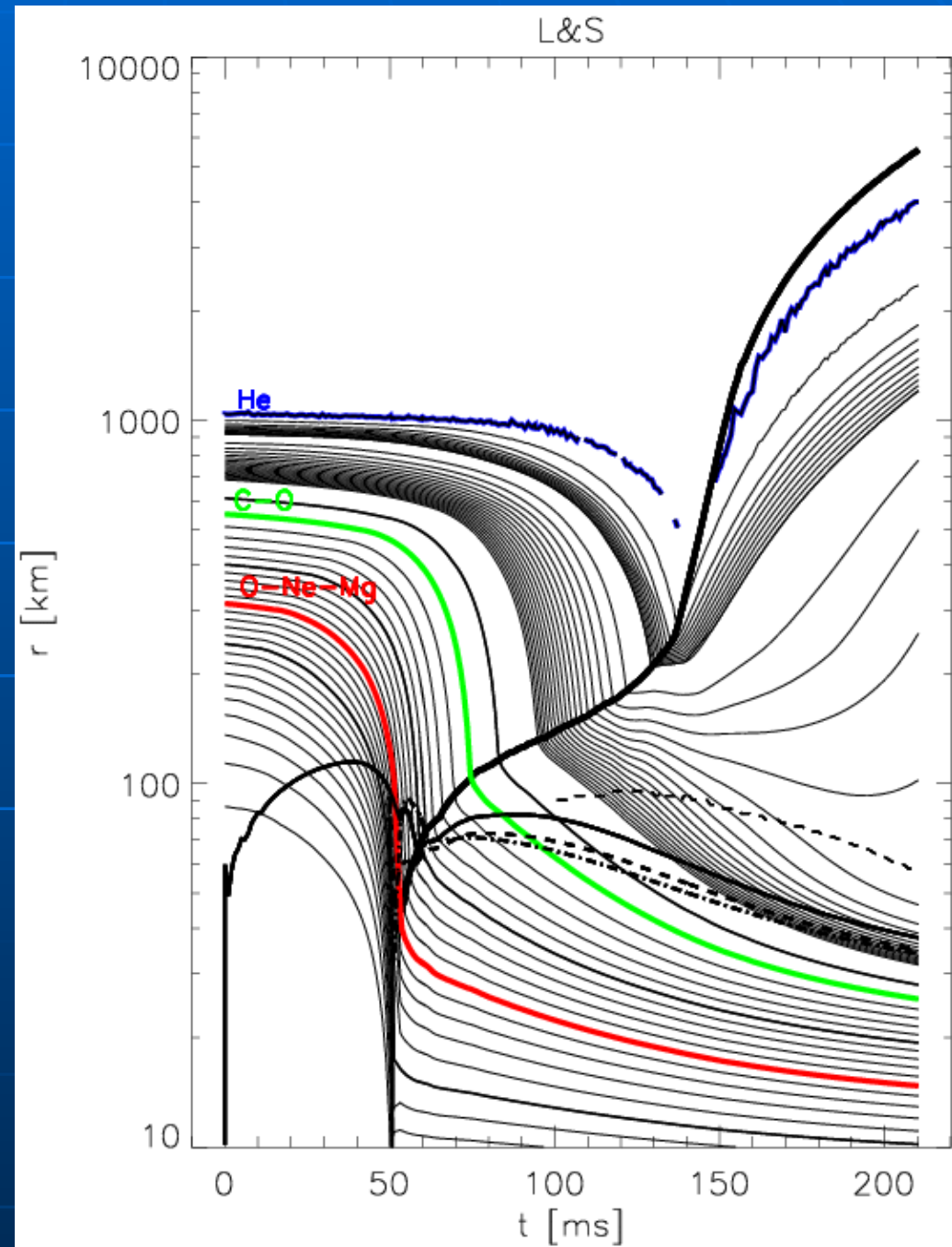
“Wolff-EOS”
(WH & Wolff 1984)

Explosion in 1D !!!

(ν -wind driven)

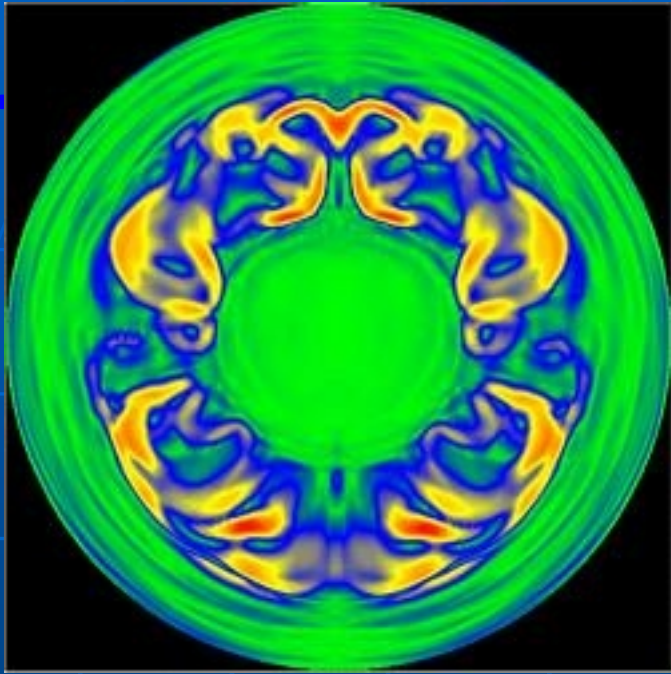


“Lattimer-Swesty-EOS”



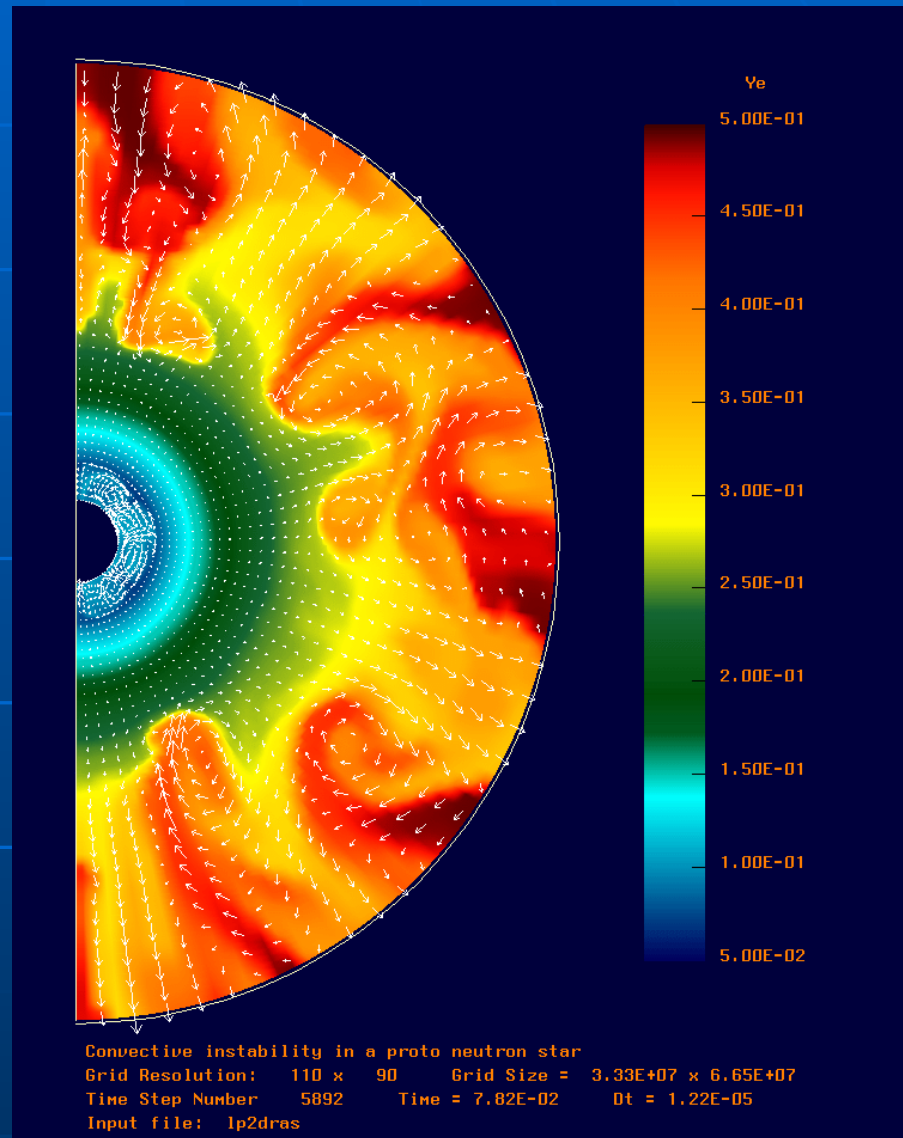
- Observations imply: non-radial flow and mixing are common in core collapse supernovae
- Theoretical models based on delayed explosion mechanism predict non-radial flow and mixing due to:
 - Ledoux convection inside the proto-neutron star (deleptonization and neutrino diffusion)
 - convection inside neutrino heated hot bubble (neutrino energy deposition behind the shock)

Core collapse supernovae need multidimensional modeling !



Ledoux convection inside proto-neutron star due to negative lepton and entropy gradients (Keil, Janka & Müller '96)

- asymmetric ν -emission (few sec) and flow (~ 100 s?)



Convection in the surface layers of the proto-neutron star and in the hot bubble 78 ms after core bounce (Janka & Müller '96)

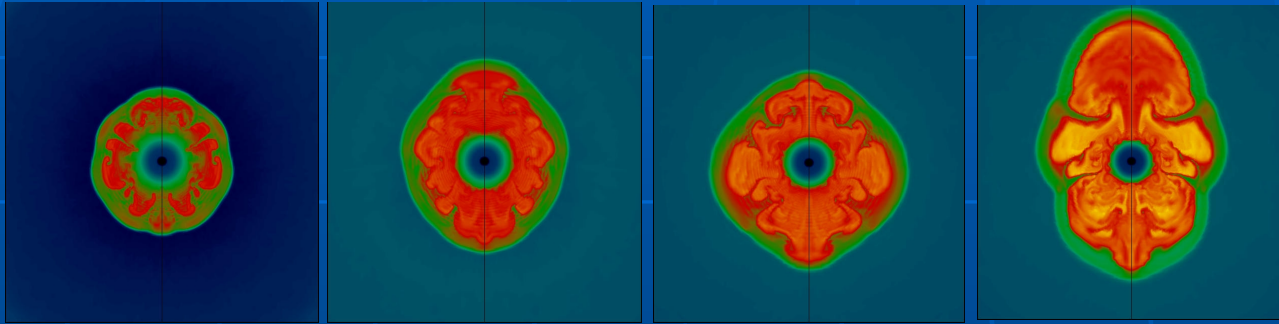
Simulations of core collapse supernovae challenging, because of:

- a) **neutrino transport** (fermions, multi-flavor)
(semi-transparent region: Boltzmann solver)

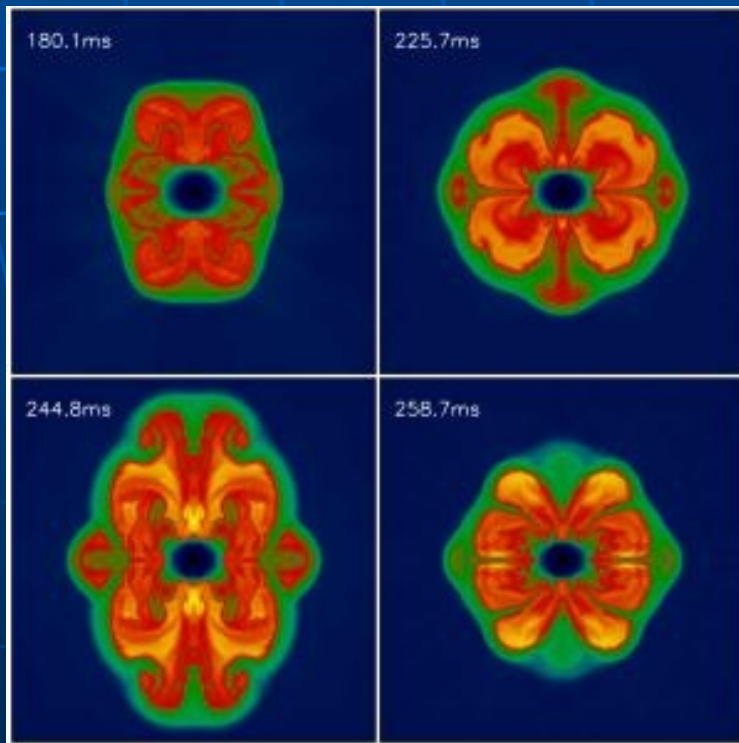
- b) very **different time and length scales**
→ adaptive mesh refinement (AMR)

- c) **multi-dimensional** flow problem

• State-of-the-art hydrodynamic simulations with Boltzmann ν -transport, realistic EOS, relativistic gravity, and realistic progenitors



Snapshots, 2D run of a non-rotating axisymmetric $11.2 M_{\text{sun}}$ progenitor
(Buras, Rampp & Janka 2003)

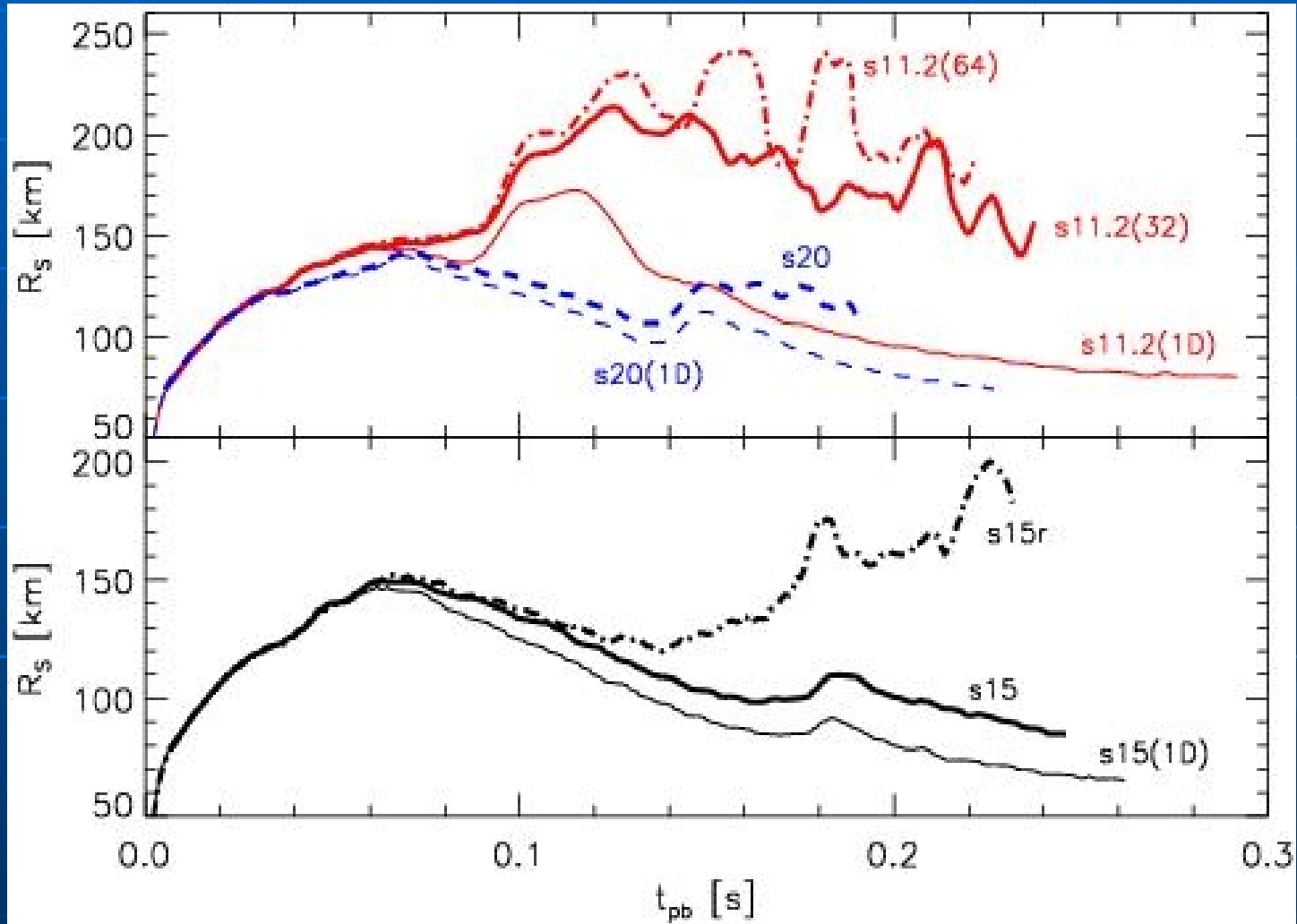


Snapshots, 2D run of a rotating
($b_{\text{initial}} = 0.05\%$, $\omega_{i,c} = 0.5 \text{ s}^{-1}$; Heger et al 2003)

axisymmetric $15 M_{\text{sun}}$ progenitor

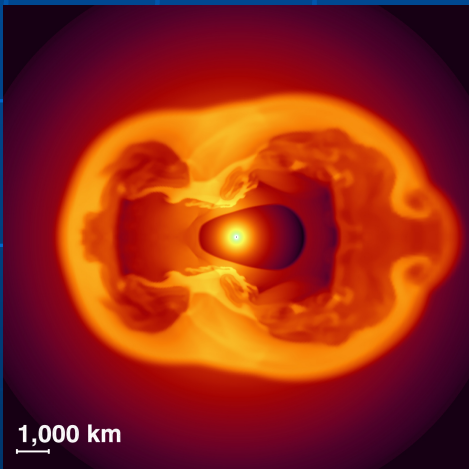
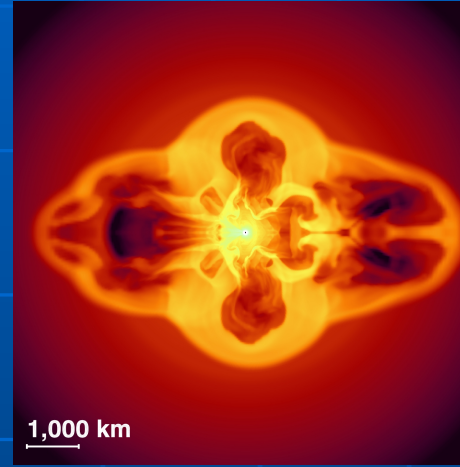
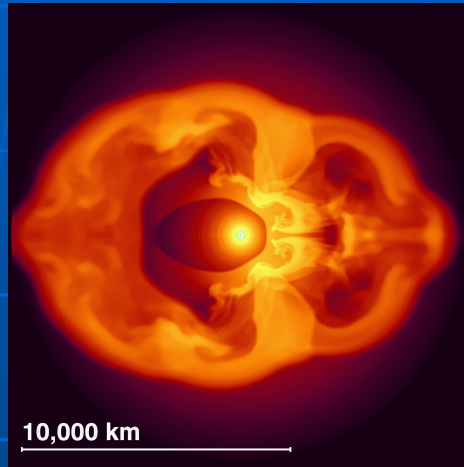
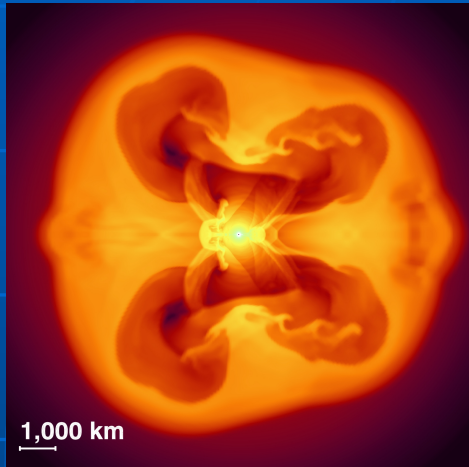
(Buras, Rampp, Janka & Kifonidis 2003)

Core Collapse Supernovae: Massive stars ($M \geq 10M_{\odot}$)

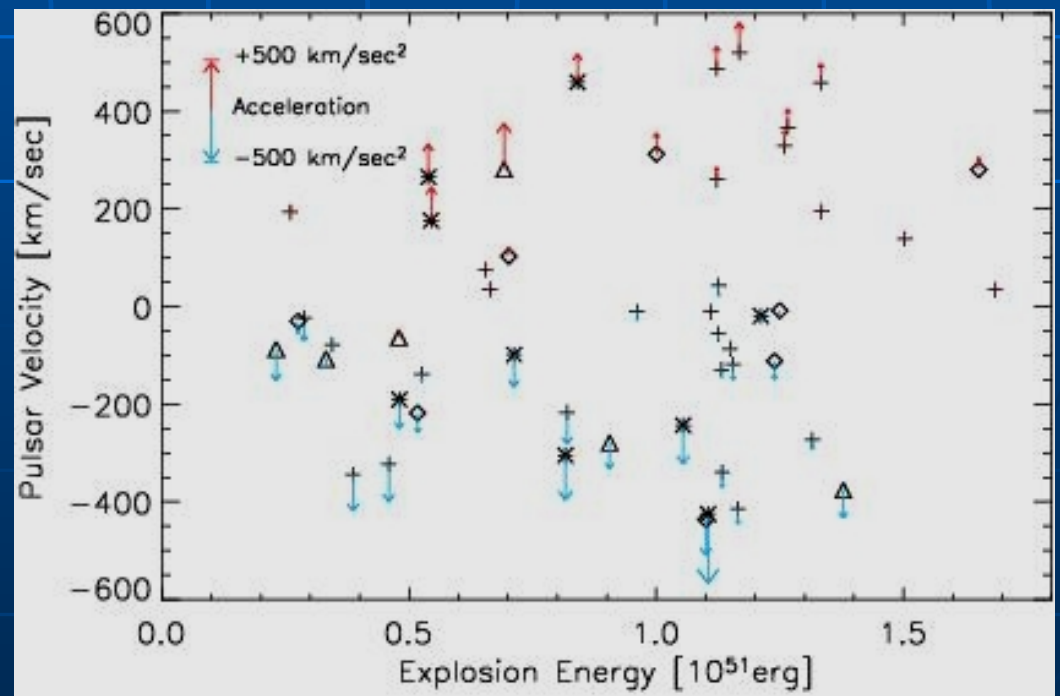


Buras et al. (2003): *No explosions (in 2-D)!!!*

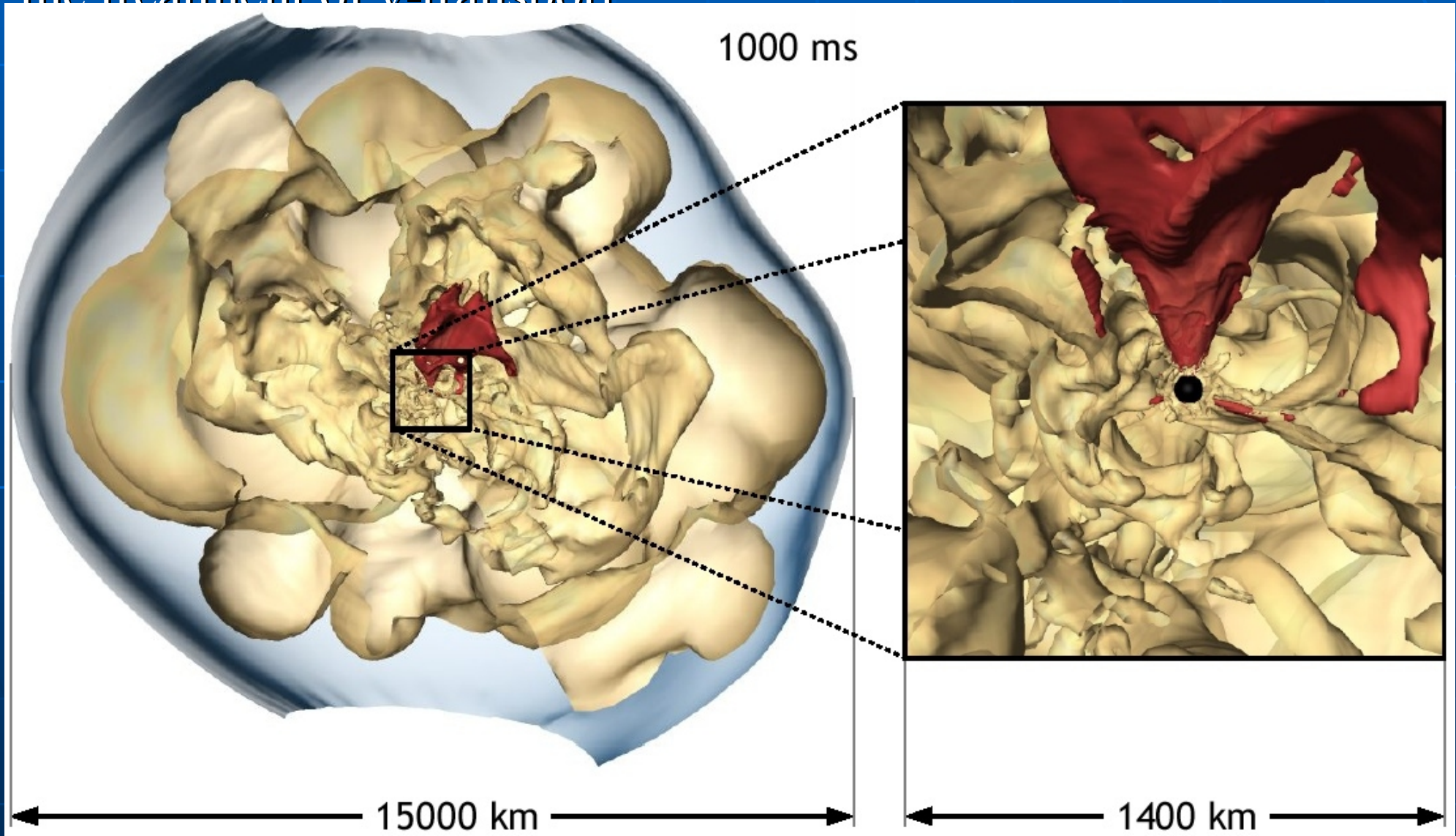
- 2D simulations show growth of dipolar ($l=1$) mode in post-shock layer \rightarrow neutron star kicks (Scheck et al. 2003)



Density distribution 1 sec
after core bounce



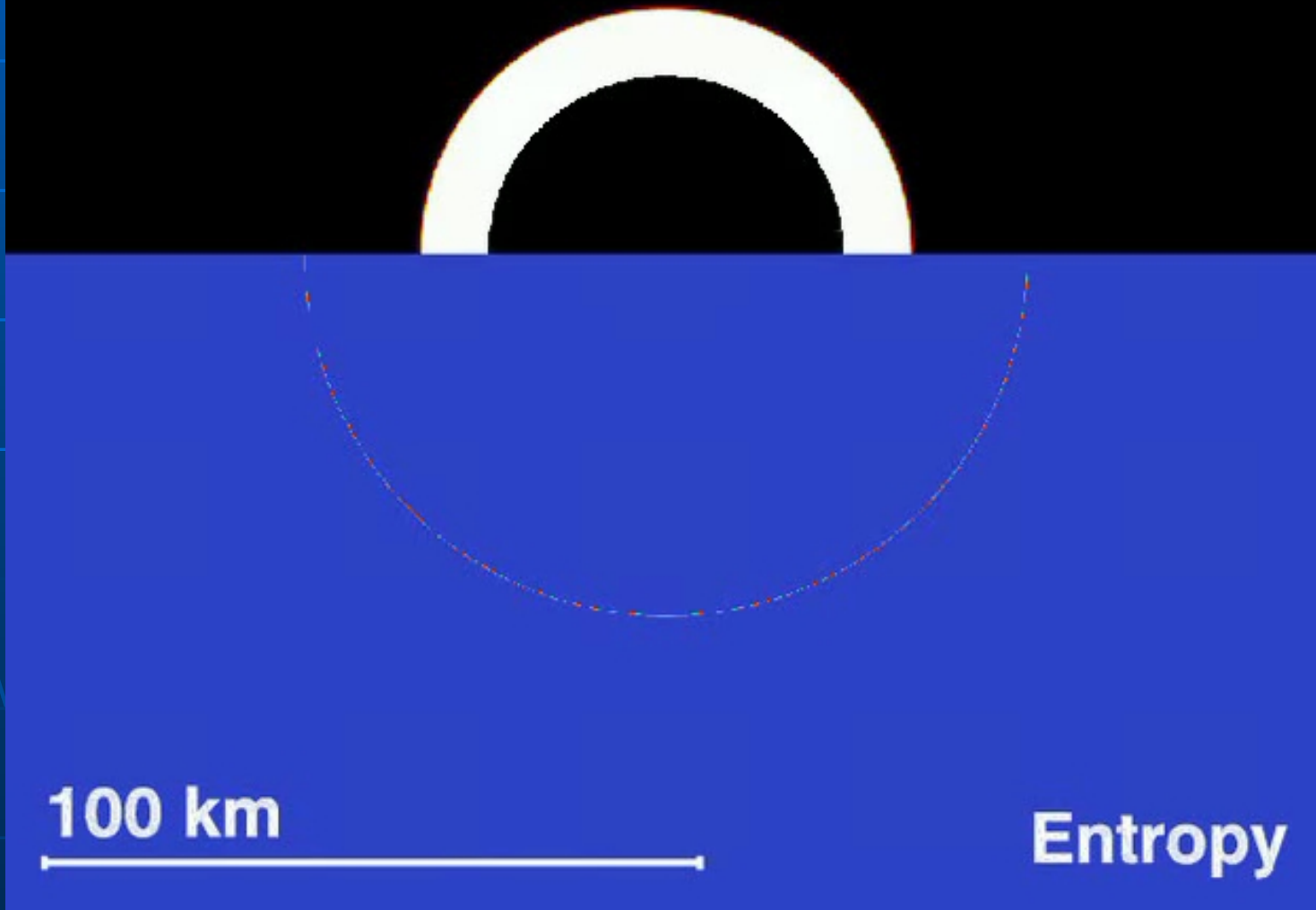
- Global dipolar oscillations of the post-shock layer also seen in recent 3D simulations neglecting (Blondin et al. '03) or simplifying (Scheck et al. '04) the treatment of ν -transport



3D core collapse simulation: shock, $Y_e = \text{const}$ & downflow to NS (Scheck 2004)

t = 0.002 sec

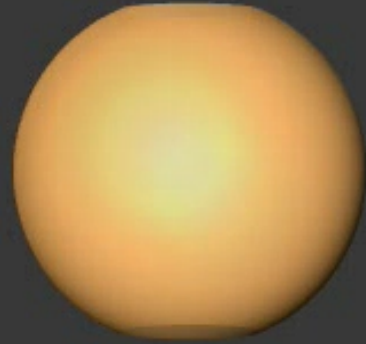
Density



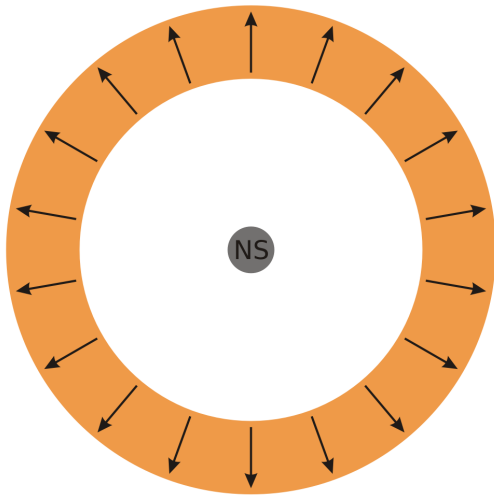
L. Scheck

100 km

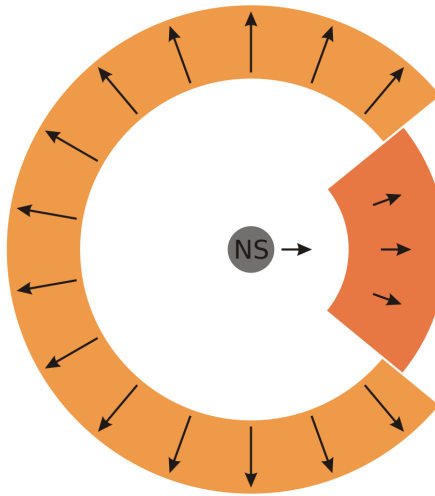
Entropy



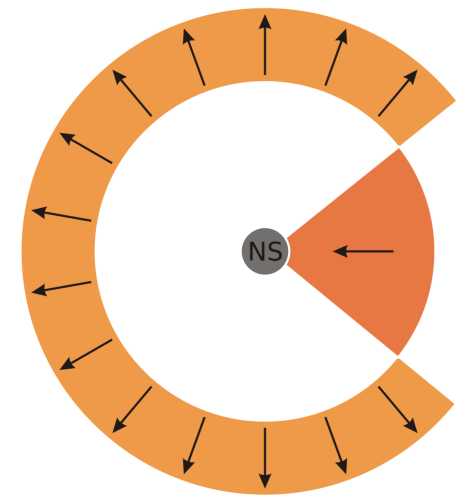
$T = 0$ msec



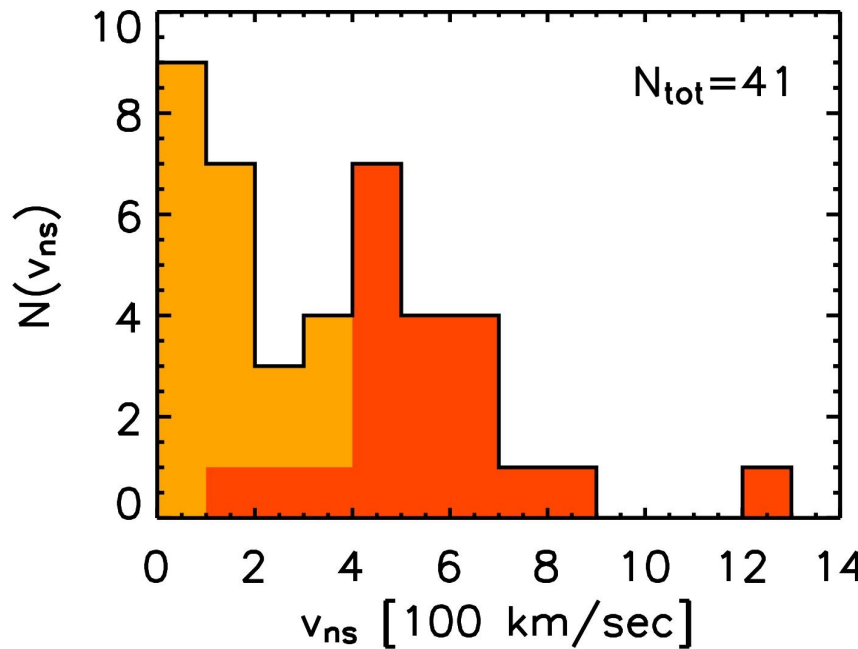
spherical explosion:
no kick



anisotropic explosion: kick due
to **gravitational acceleration**



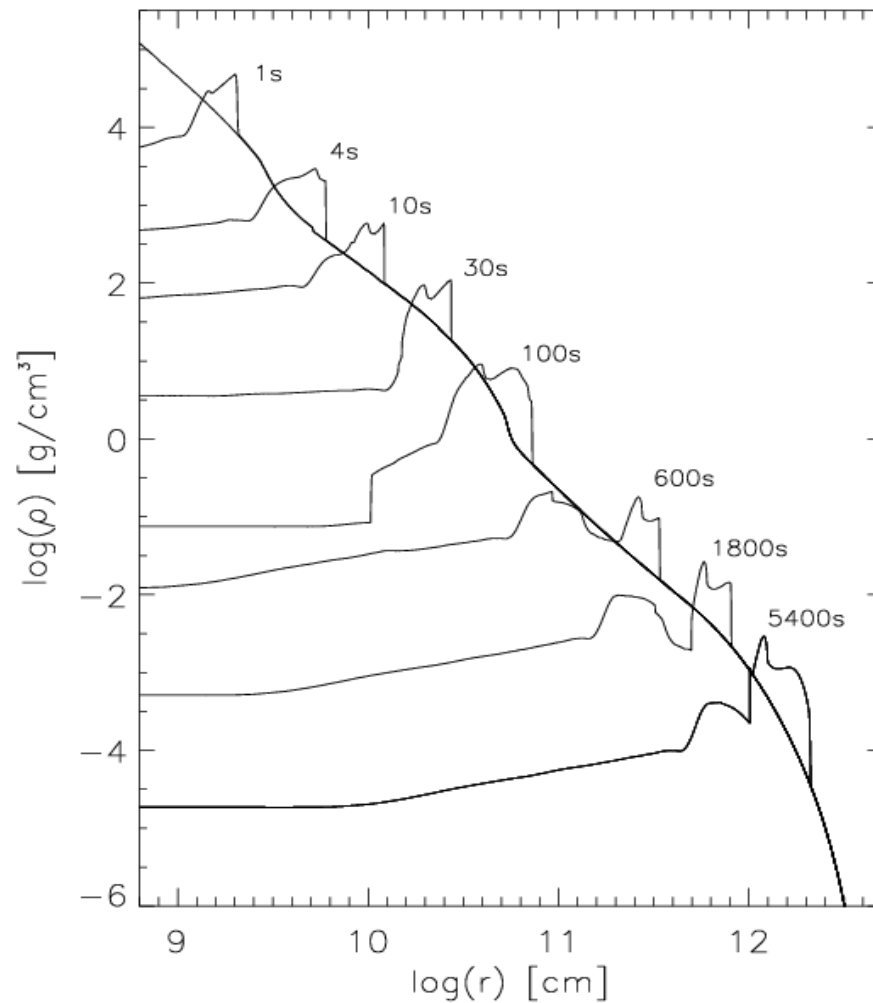
anisotropic explosion: kick due
to **anisotropic accretion**



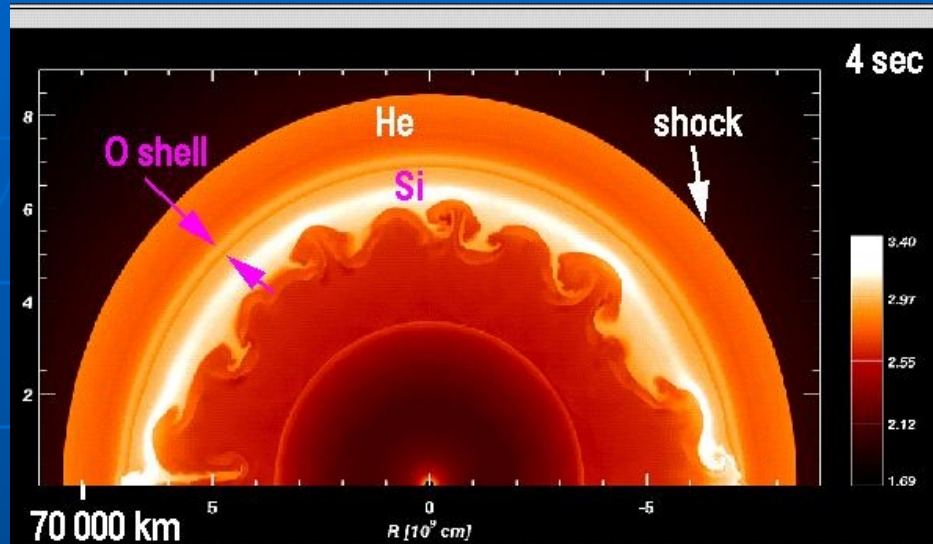
Large set of
simulations shows
bi-modal kick
velocity distribution
(Scheck 2005)

Core Collapse Supernovae: Further evolution

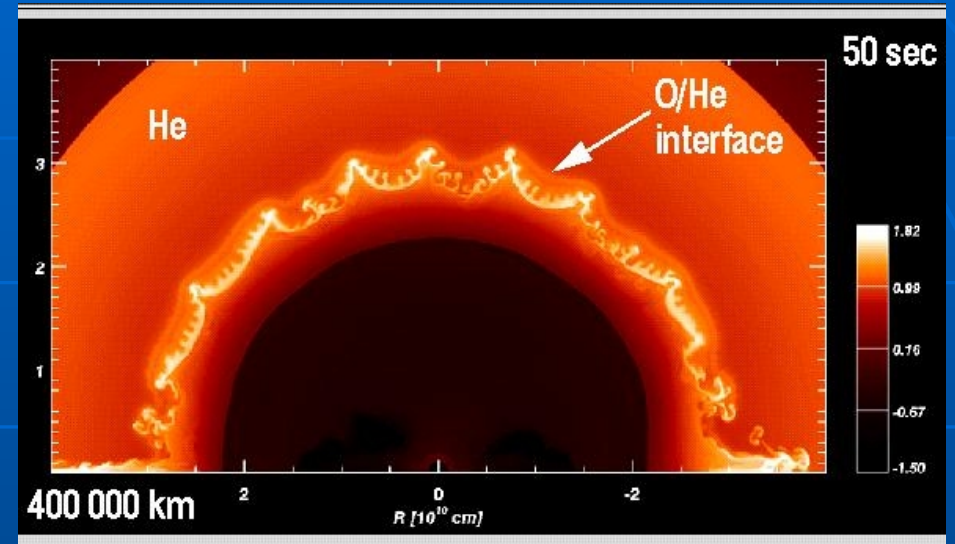
Unsteady shock propagation through stellar envelope --> **Rayleigh Taylor unstable regions**



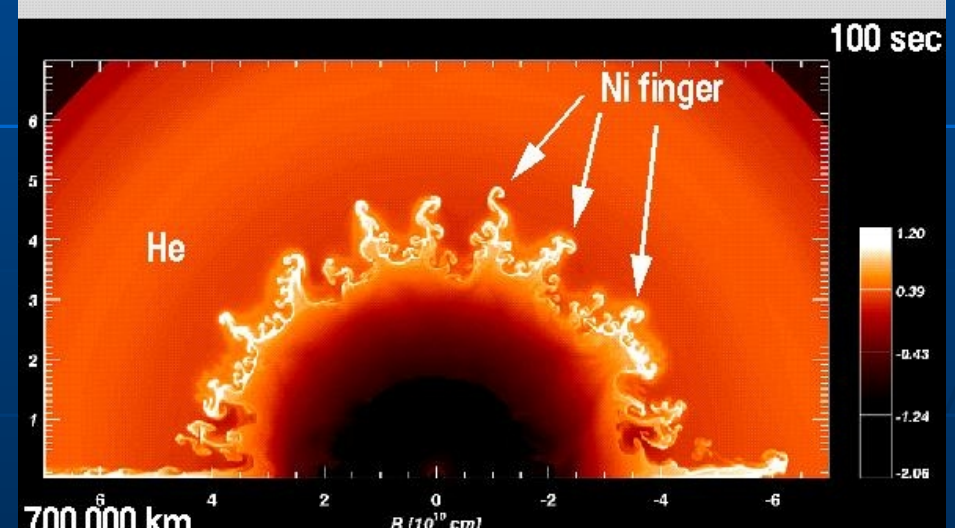
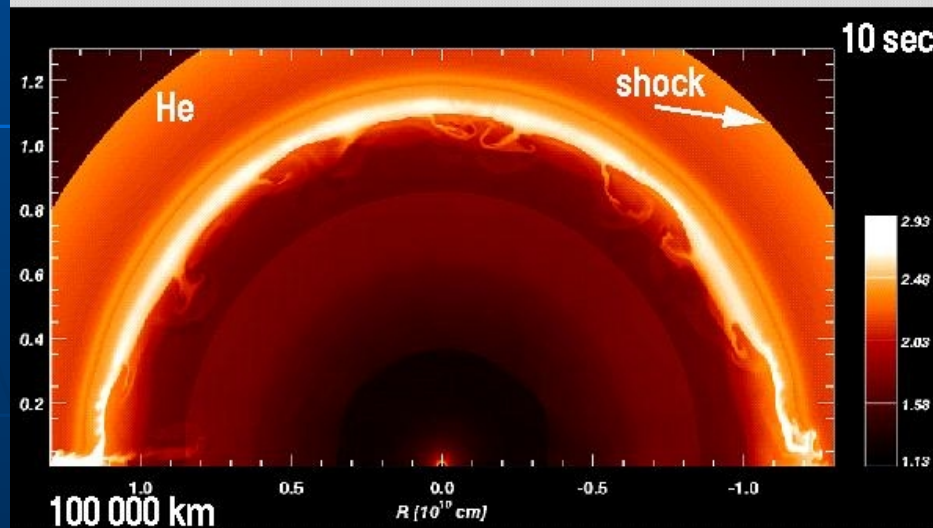
Instabilities, mixing and nucleosynthesis



density



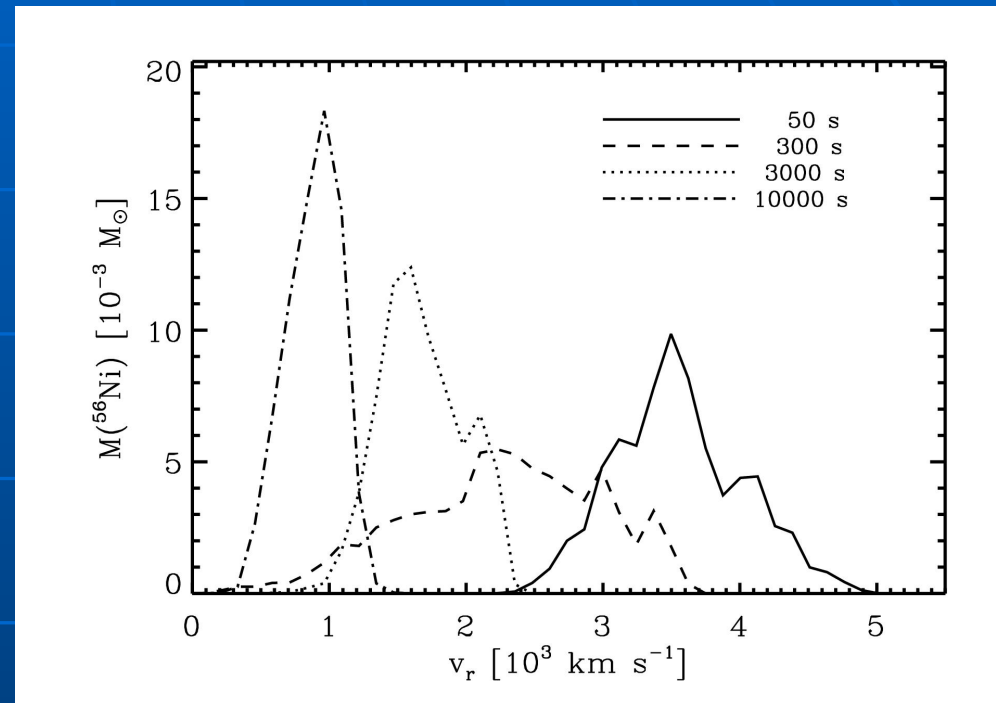
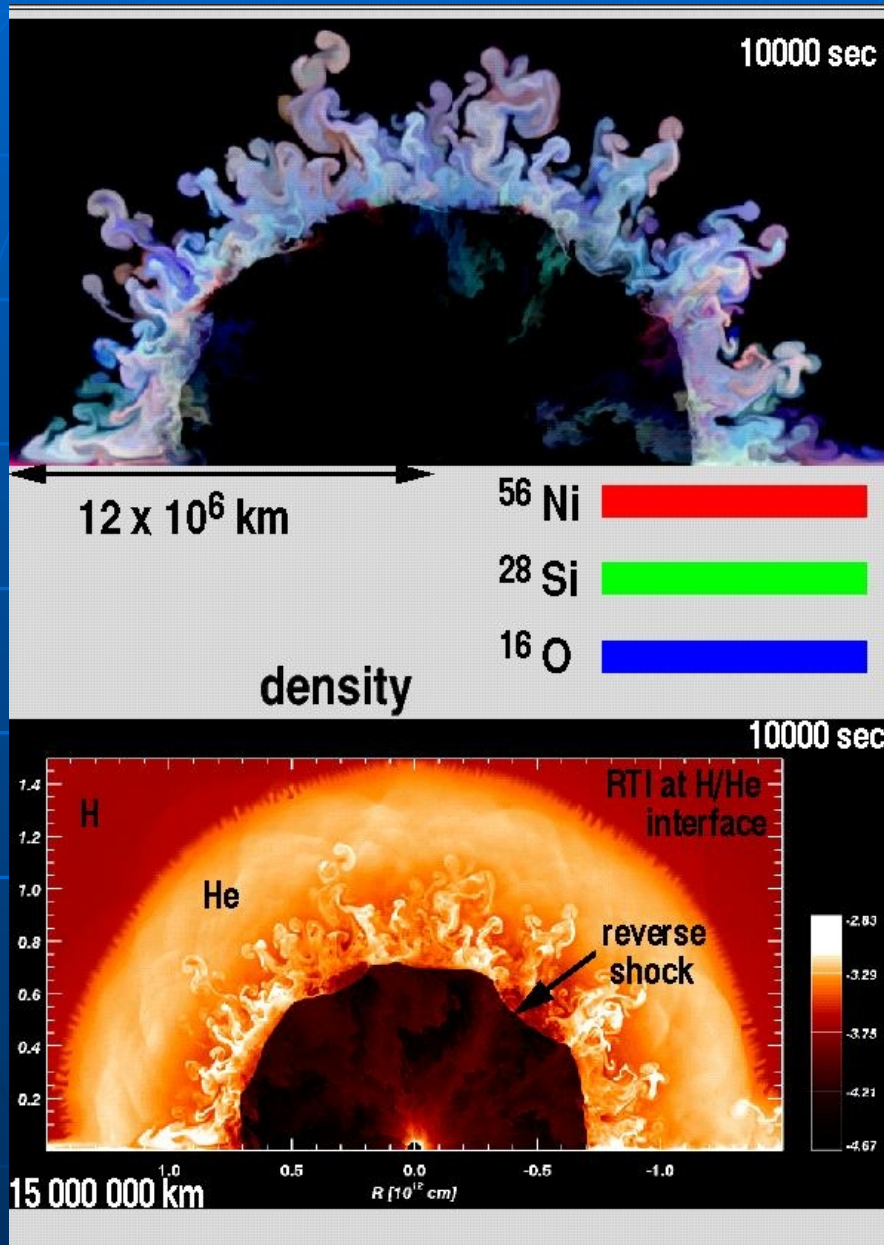
density



1 sec



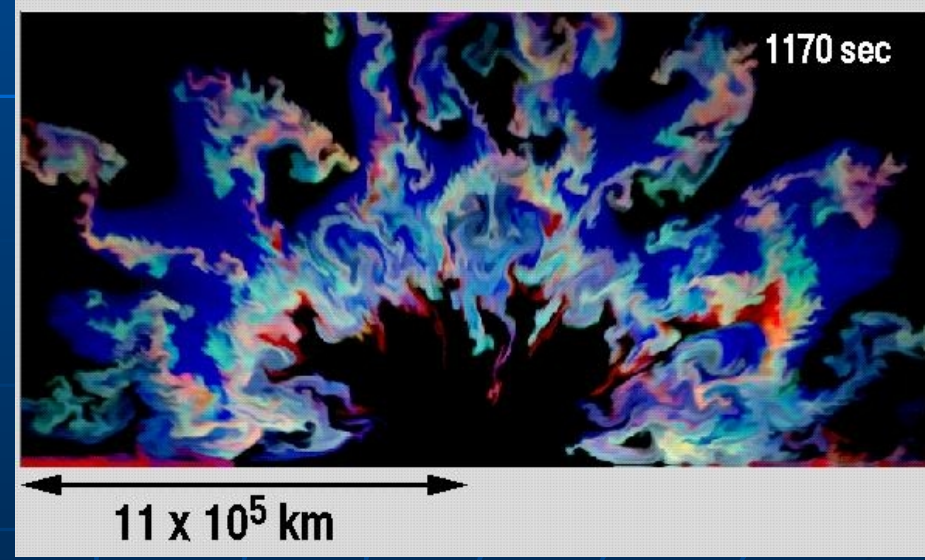
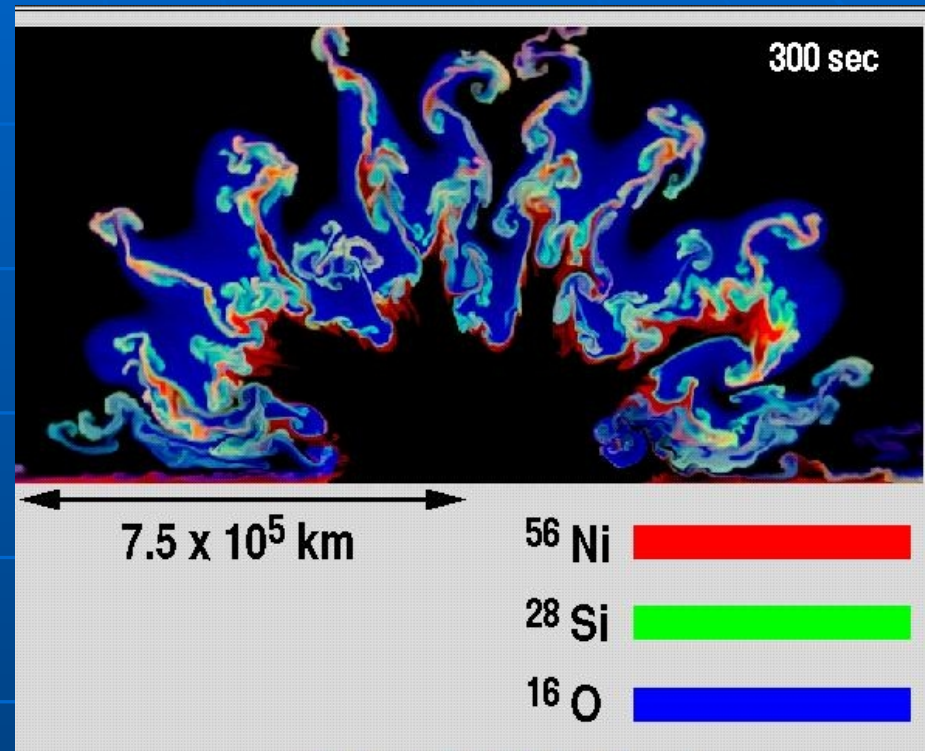
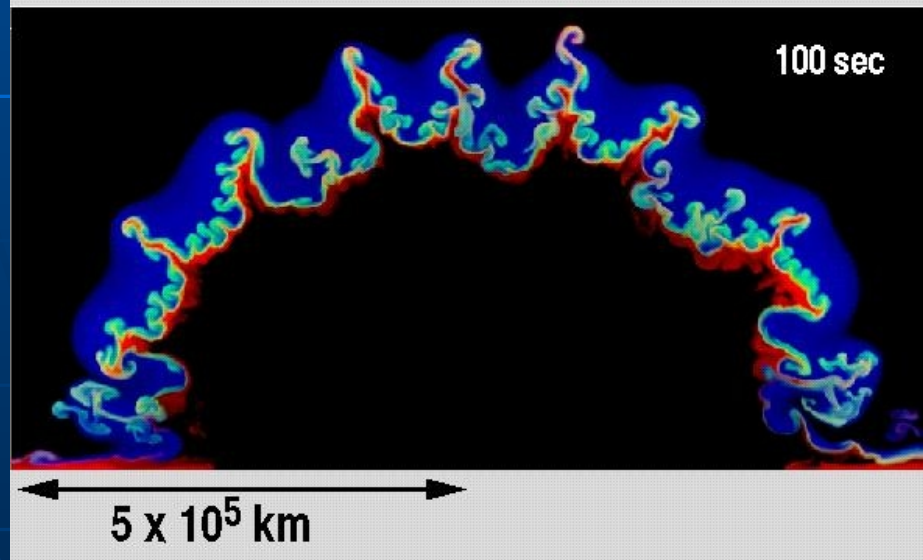
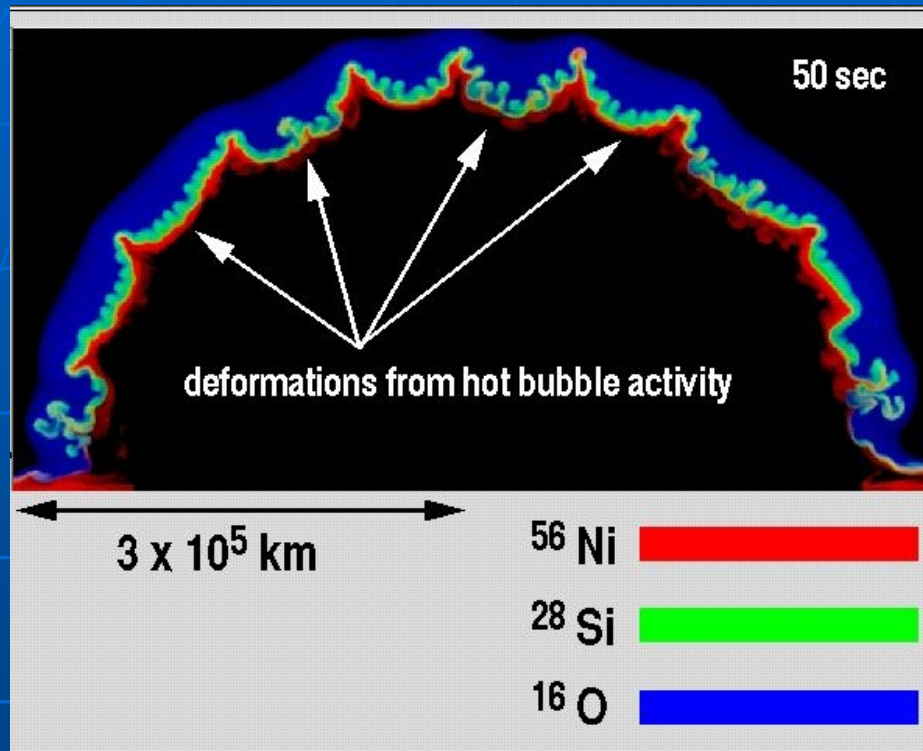
Instabilities, mixing and nucleosynthesis



- results of simulations in accordance with observations of SNe Ib/Ic
- simulations do not reproduce large velocities of Fe/Ni observed in SN 1987A

AMR simulation of shock propagation through stellar envelope (Kifonidis, Plewa, Janka & Müller 2003)

Instabilities, mixing and nucleosynthesis (cont.)



Summary (Part I)

- ❑ Core-collapse supernova explosions are triggered by neutrino interactions with matter and hydrodynamic instabilities and/or rotation (magnetic fields?).
- ❑ Even the best models available predict weak explosions only.
- ❑ What is the missing physics?
- ❑ “Artificially” triggered explosion models predict nuclear abundances in fair agreement with observations.

Thermonuclear (Type Ia) Supernovae



Example:

SN 2002bo in NGC 3190;

Discovered: March 9, 2002

B-maximum: March 22, 2002

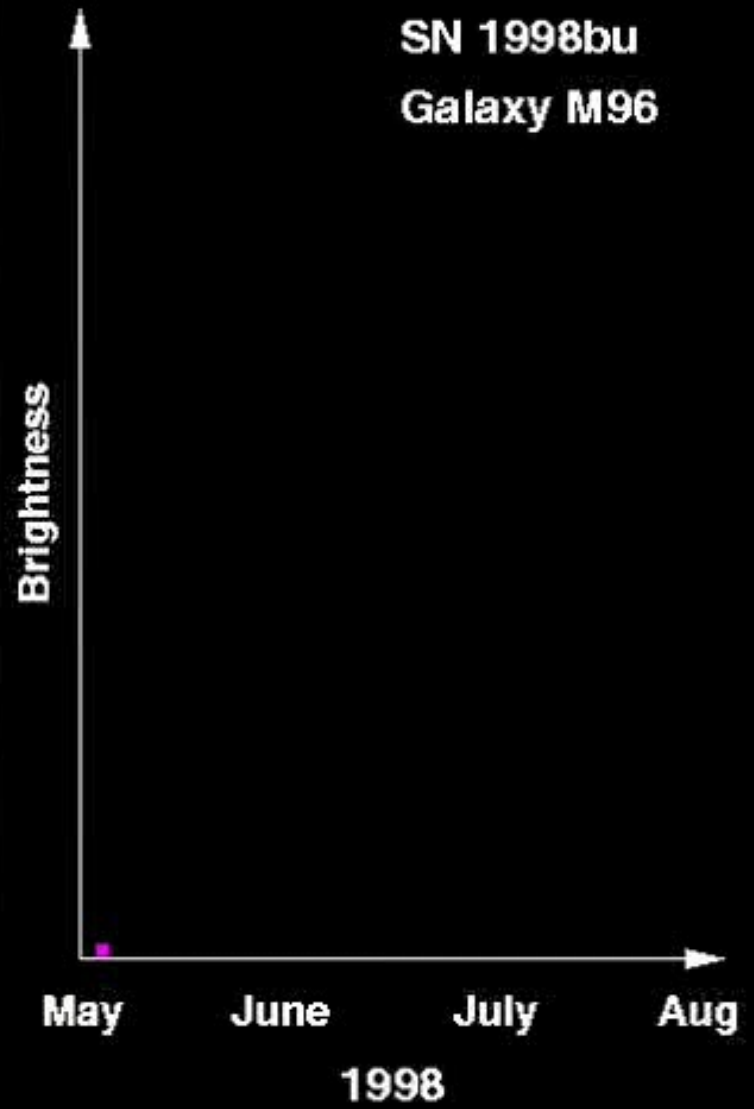
(RTN/ESC)



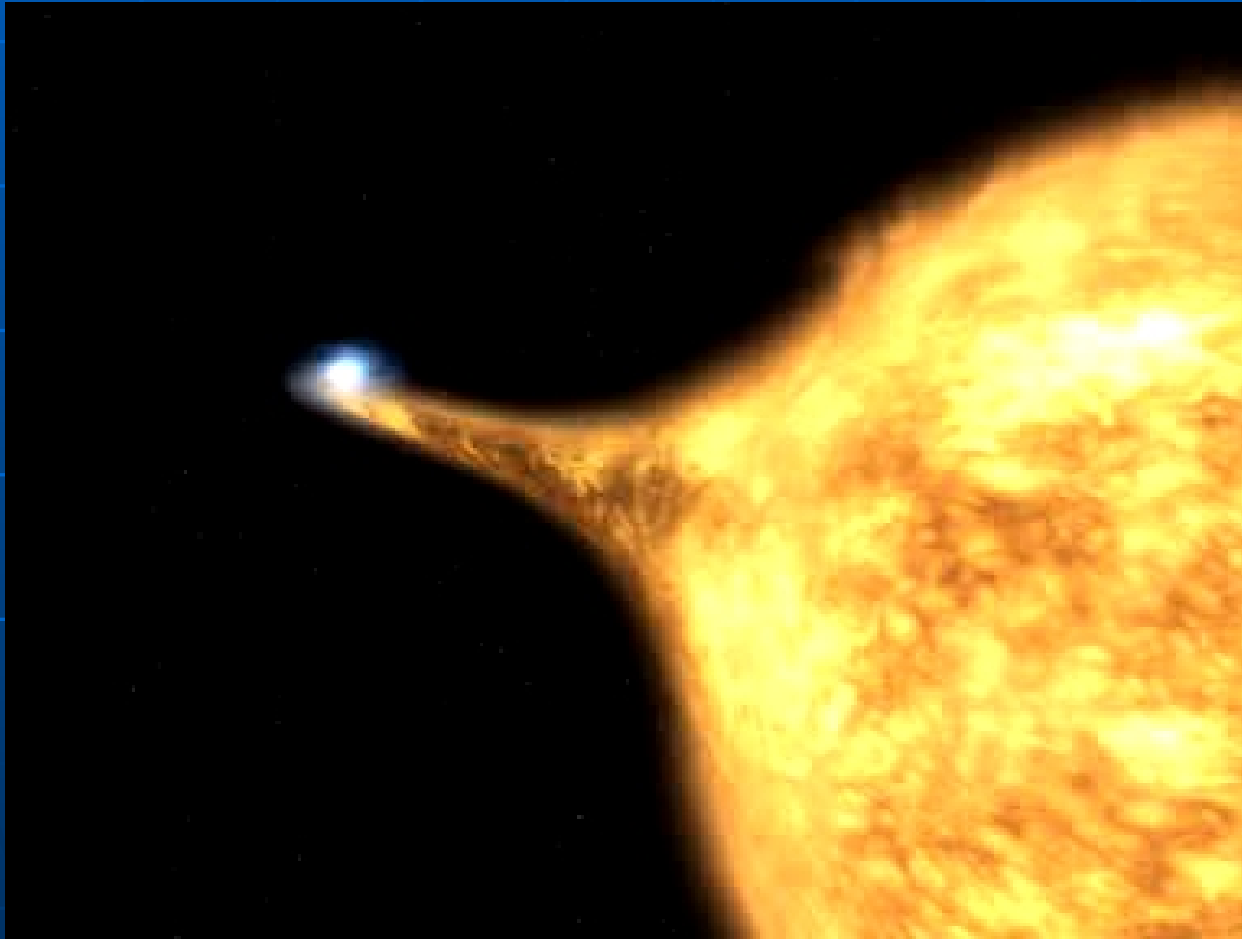
The last (observed) galactic SN (Ia?):

SN 1604 (“Kepler’s Supernova”)





The “standard model”

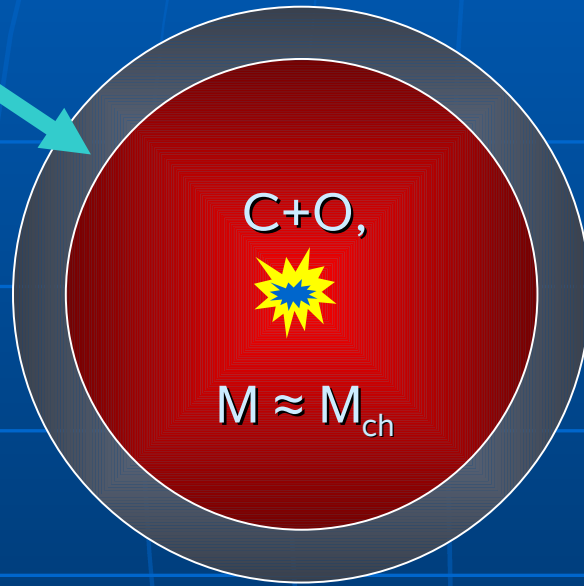


- White dwarf in a binary system
- Growing to the Chandrasekhar mass by mass transfer



How does the model work?

He (+H)
from binary
companion



Density $\sim 10^9 - 10^{10}$ g/cm

Temperature: a few 10^9 K

Radii: a few 1000 km

Explosion energy:

*Fusion C+C,
C+O, O+O \rightarrow "Fe"*

Laminar burning
velocity:

$U_L \sim 100$ km/s $\ll U_S$

Too little is burned!

The physics of turbulent combustion

- Everyday experience:
Turbulence increases the burning velocity.
- In a star:
Reynoldsnumber $\sim 10^{14}$!
- In the limit of strong turbulence: $U_B \sim V_T$!
- Physics of thermonuclear burning is very similar to premixed chemical flames.



Relevant length scales in simulations of SN Ia explosions

beginning of the explosion:

flamelet regime

Kolmogorov scale

Gibson scale

ignition radius



flame width

resolution in
3D models

WD radius



resolved flame
simulations
(Woosley et al.)



complementary small-scale
studies (Röpke et al.,
Schmidt et al.)

SGS
turbulence
model

Large-scale supernova
simulations

A couple of definitions:

Kolmogorov (length) scale

$$\eta := (v^3/\varepsilon)^{1/4}$$

(Turbulent) Reynolds number

$$Re := v'/s_L \cdot l/l_F$$

(Turbulent) Damköhler number

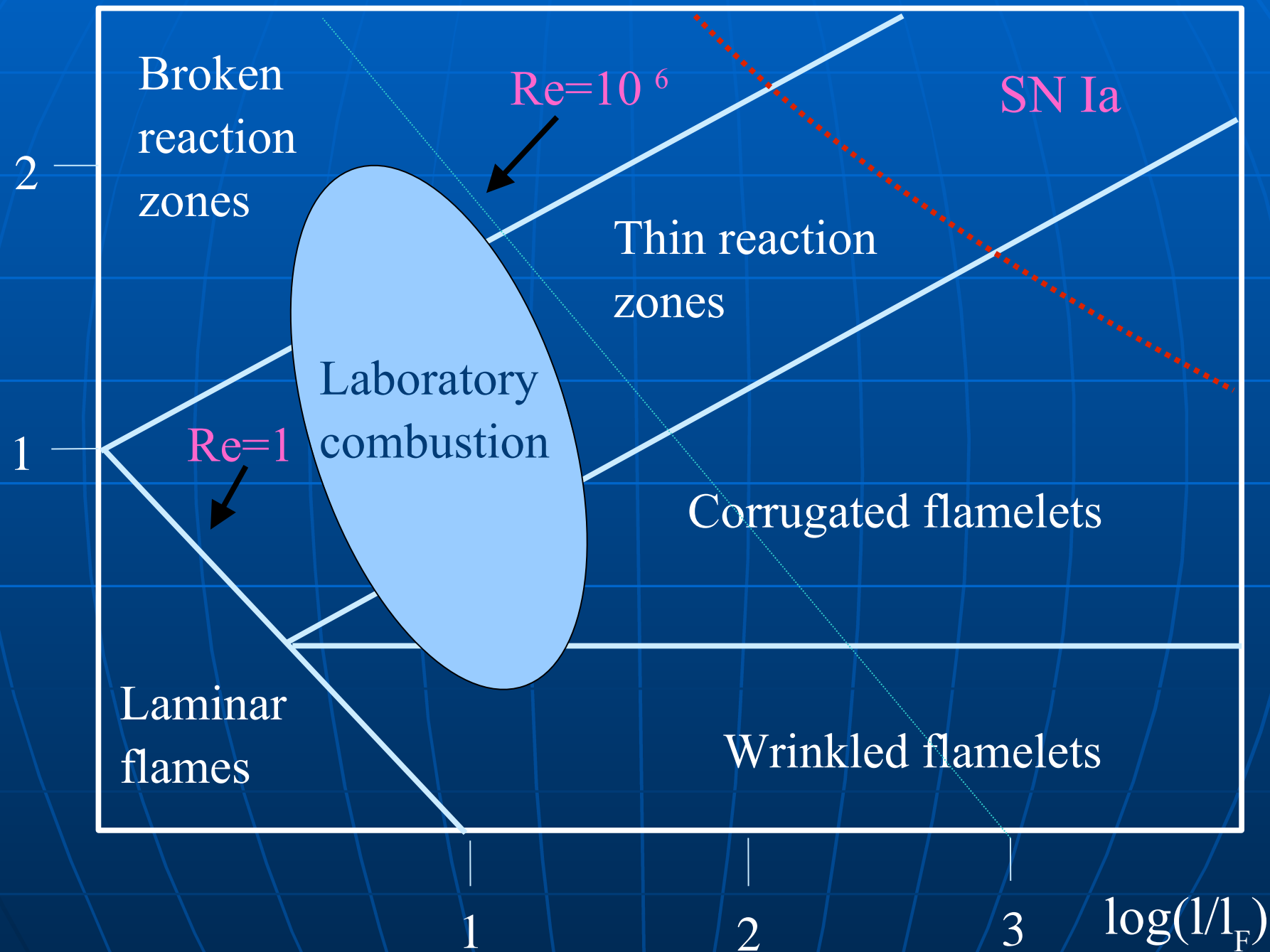
$$Da := s_L/v' \cdot l/l_F$$

(Turbulent) Karlovitz number

$$Ka := l_F^2/\eta^2$$

$$\Rightarrow Re = Da^2 \cdot Ka^2$$

$\log(v'/s_L)$



Burning regimes of pre-mixed flames

1. Cellular burning, wrinkled flamelets

$$u_{\text{cell}} = s_L [1 + \varepsilon(\mu)] ; \mu = \rho_b / \rho_u ,$$

$$\varepsilon(\mu) \approx 0.41 (1 - \mu)^2$$

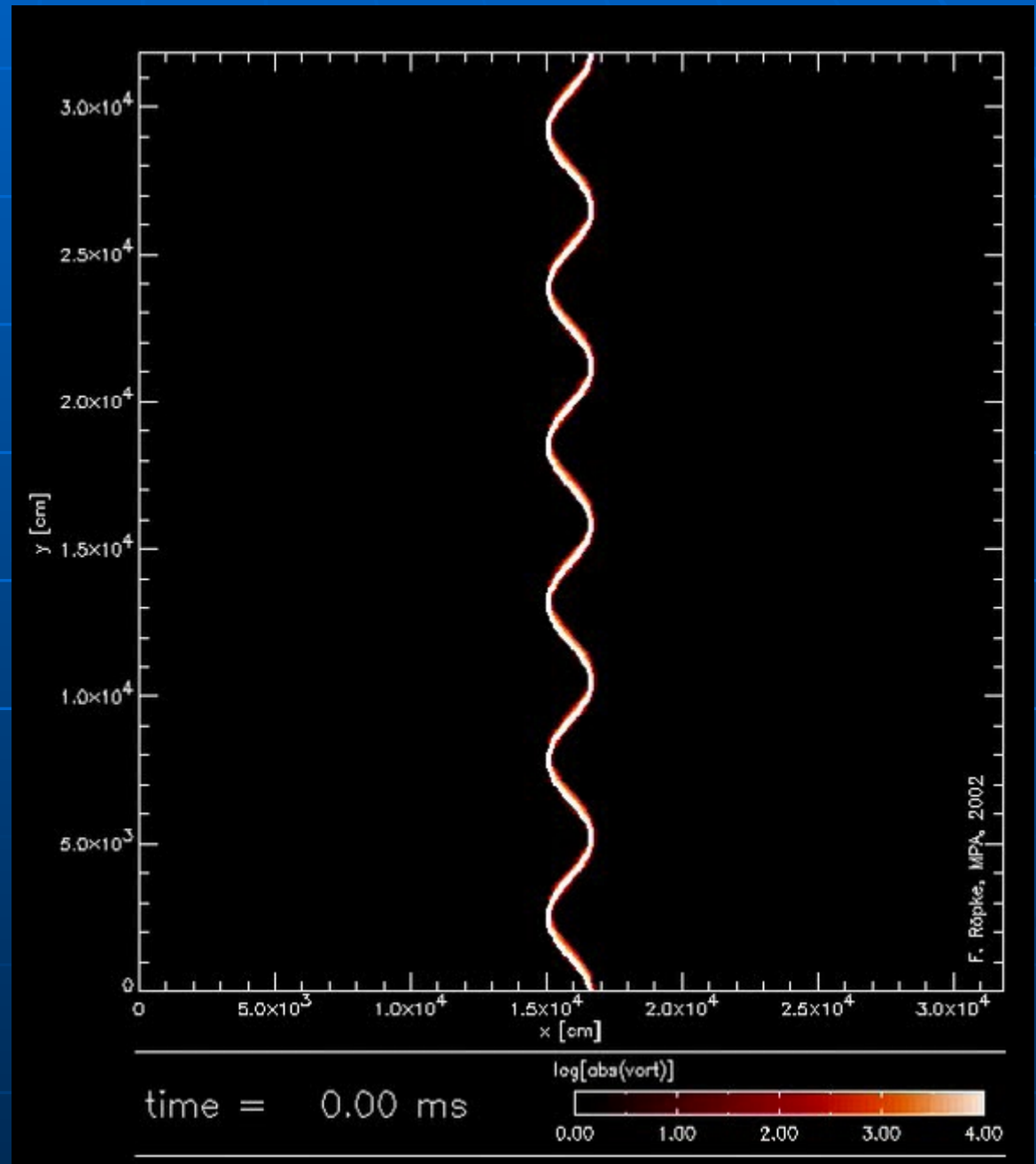
Or: “Fractal model”

$$u_{\text{cell}}(l) = s_L (l/l_{\text{crit}})^{D-1}$$

The Landau-Darrieus instability and its interaction with turbulence:

Quiescent fuel

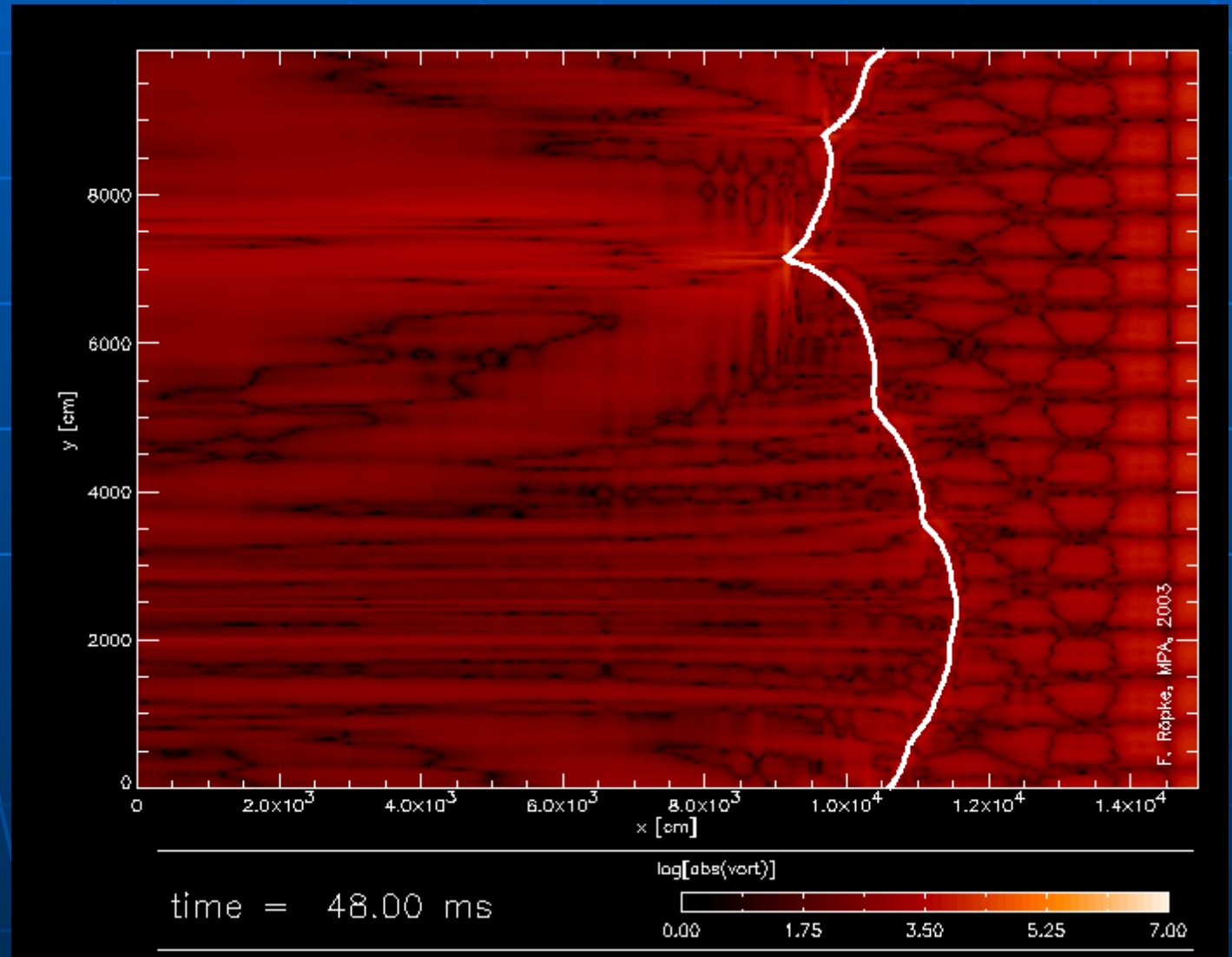
(Röpke et al., 2003a)



The Landau-Darrieus instability and its interaction with turbulence:

*Weak vortical
flow*

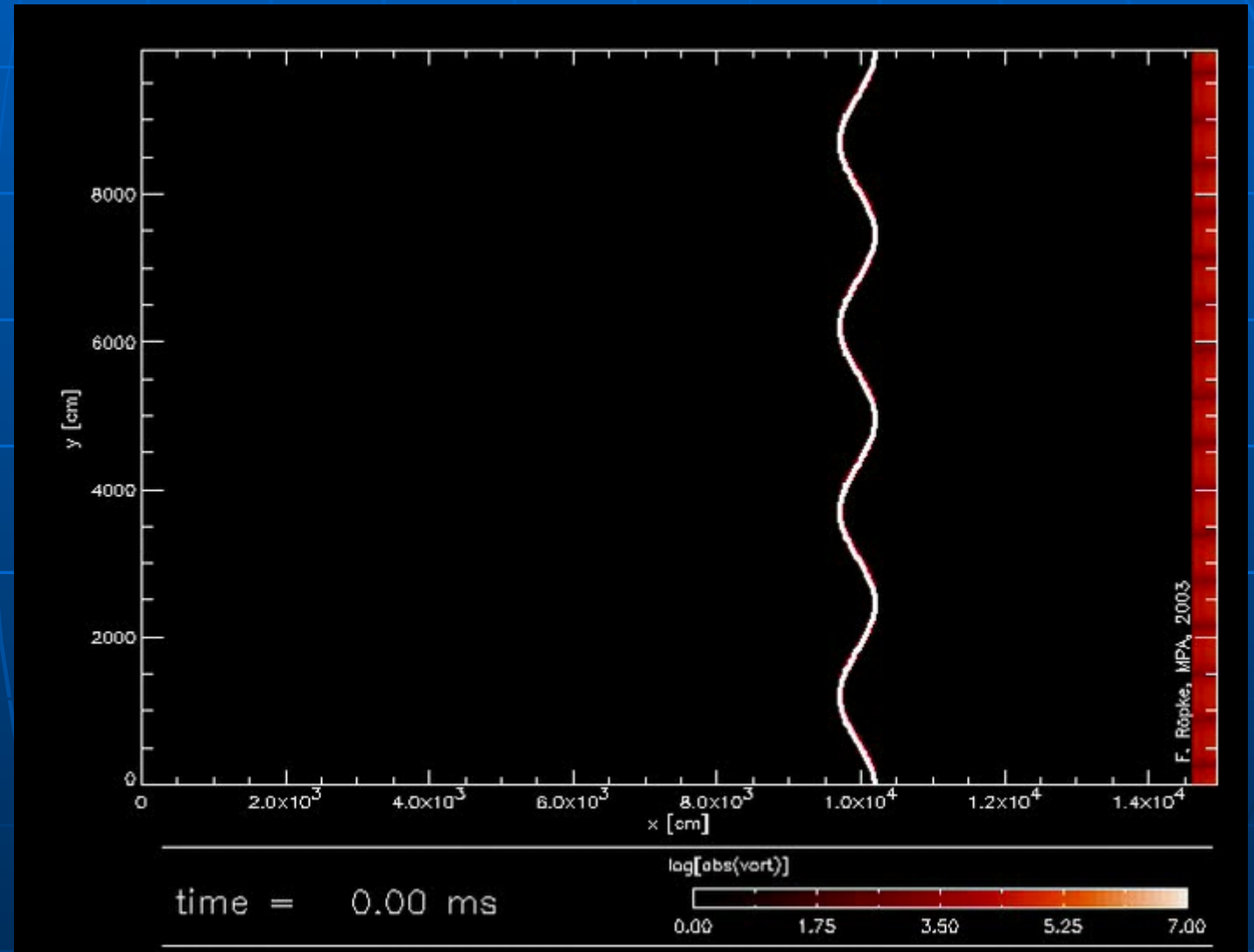
(Röpke et al.,
2003b)



The Landau-Darrieus instability and its interaction with turbulence:

Strong vortical flow

(Röpke et al.,
2003b)



Burning regimes of pre-mixed flames

2. The corrugated flamelet regime

Transition at the “Gibson scale”:

$$v(l_{\text{Gibs}}) = u_{\text{cell}}(l_{\text{Gibs}})$$

In the limit of strong turbulence:

$$s_{\text{turb}}(l) \approx v'(l), \quad l > l_{\text{Gibs}} \quad (\text{independent of } s_L!!!)$$

$$d_{\text{turb}} \approx l \quad (\text{“turbulent flame brush”})$$

Fully developed turbulence?

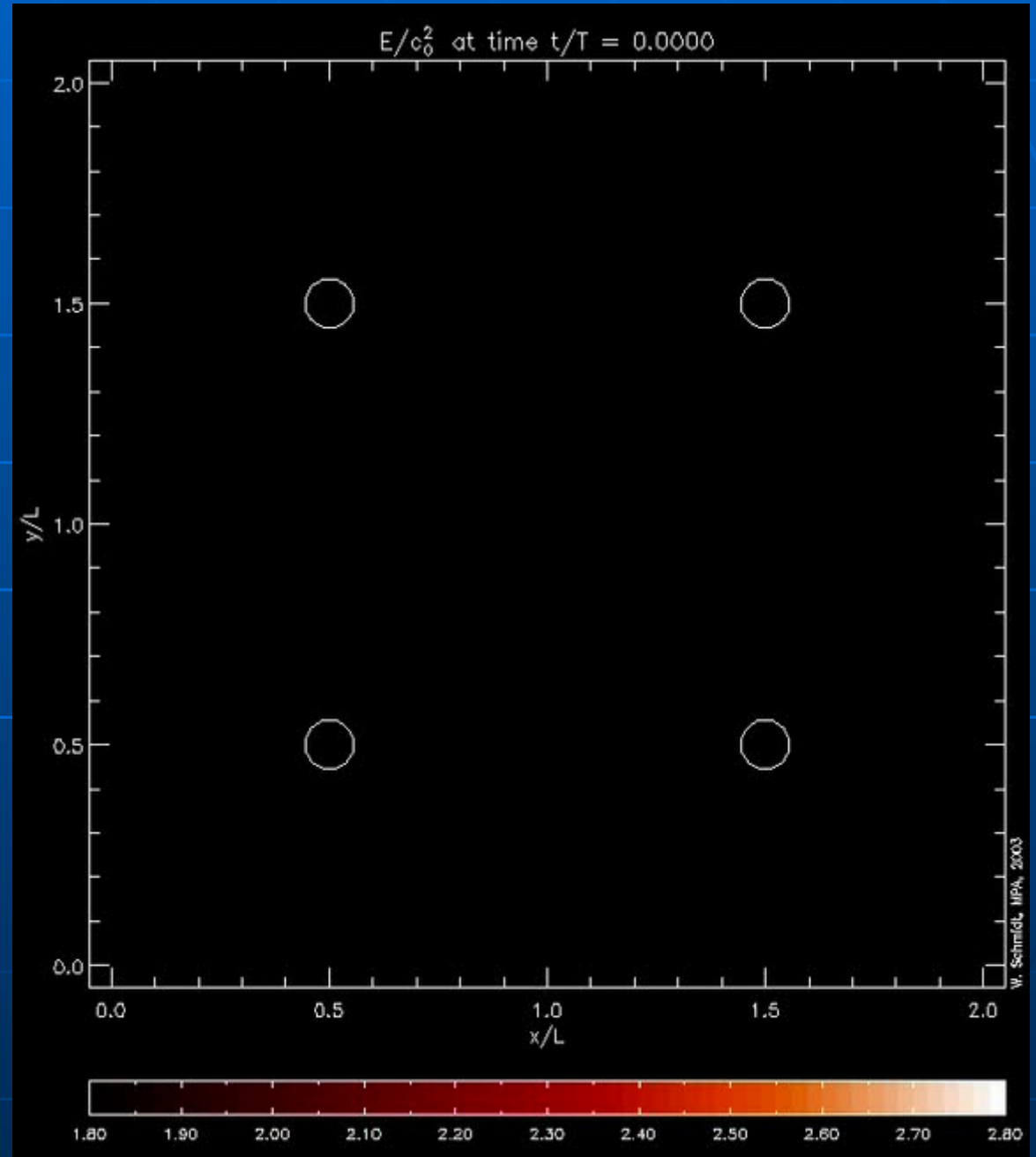
3-D “direct”
numerical simulations
of flames moving in
white dwarf matter:
Energy

$$\rho = 2.9 \cdot 10^9 \text{ gcm}^{-3}$$

$$V/s_{lam} = 4$$

$$V/c_0 = 0.043$$

(Schmidt et al., 2004)



Fully developed turbulence?

3-D “direct”
numerical simulations
of flames moving in
white dwarf matter:

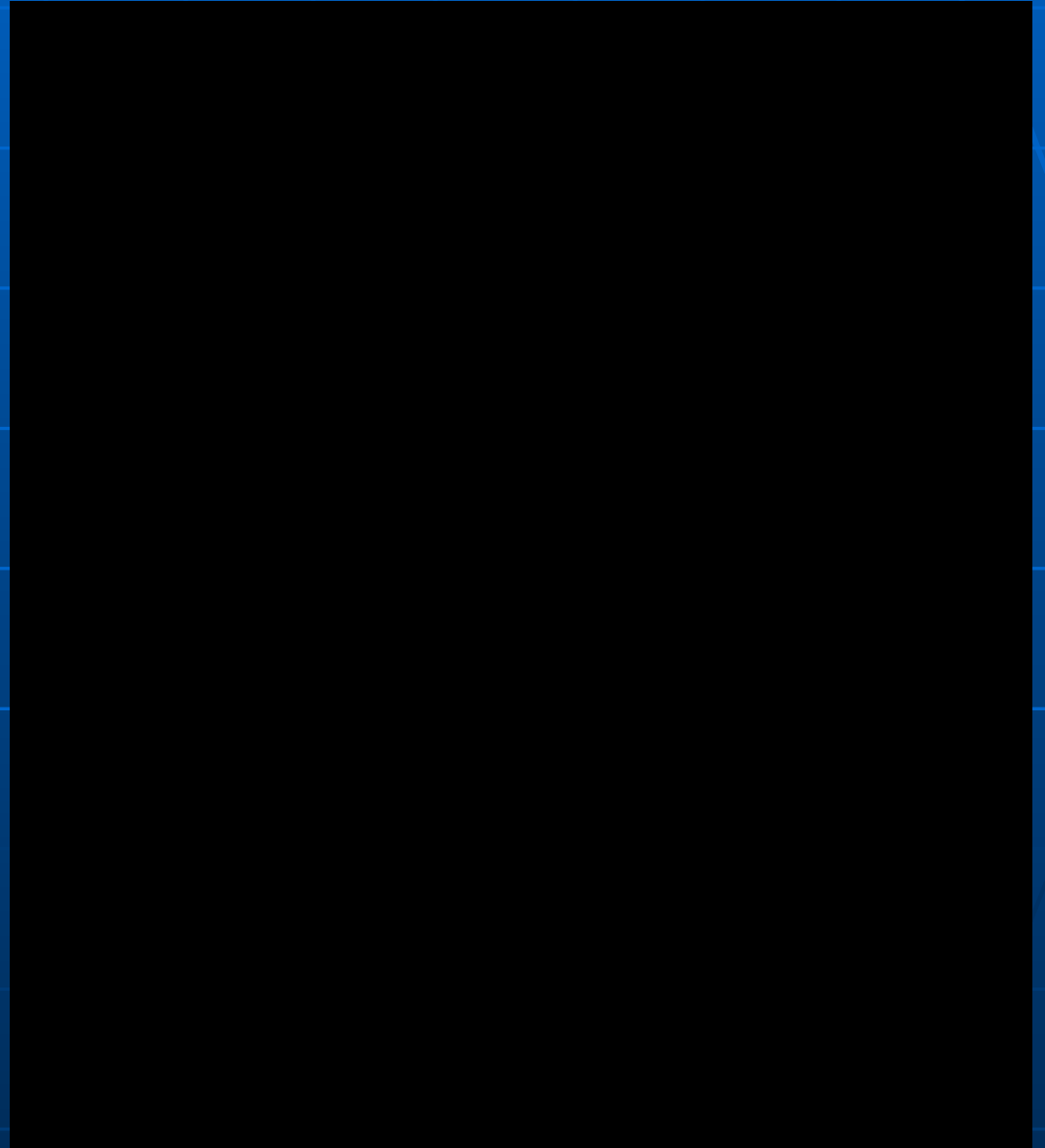
Vorticity

$$\rho = 2.9 \cdot 10^9 \text{ gcm}^{-3}$$

$$V/s_{\text{lam}} = 4$$

$$V/c_0 = 0.043$$

(Schmidt et al., 2004)



Fully developed turbulence?

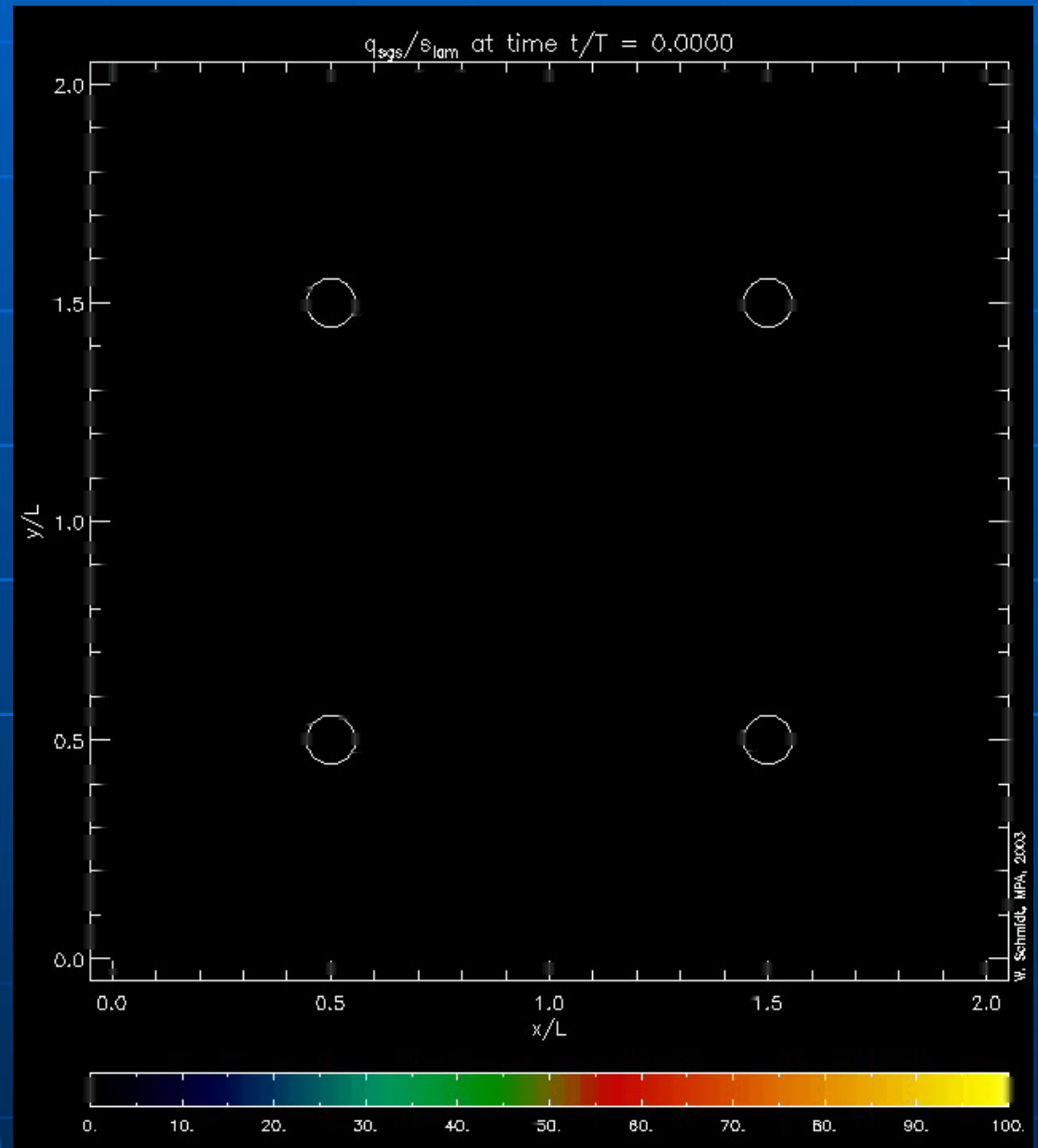
3-D “direct”
numerical simulations
of flames moving in
white dwarf matter:
Subgrid-scale energy

$$\rho = 2.9 \cdot 10^9 \text{ g cm}^{-3}$$

$$V/s_{lam} = 4$$

$$V/c_0 = 0.043$$

(Schmidt et al., 2004)



Burning regimes of pre-mixed flames

3. The distributed-burning

Turbulent eddies interact with the flame:

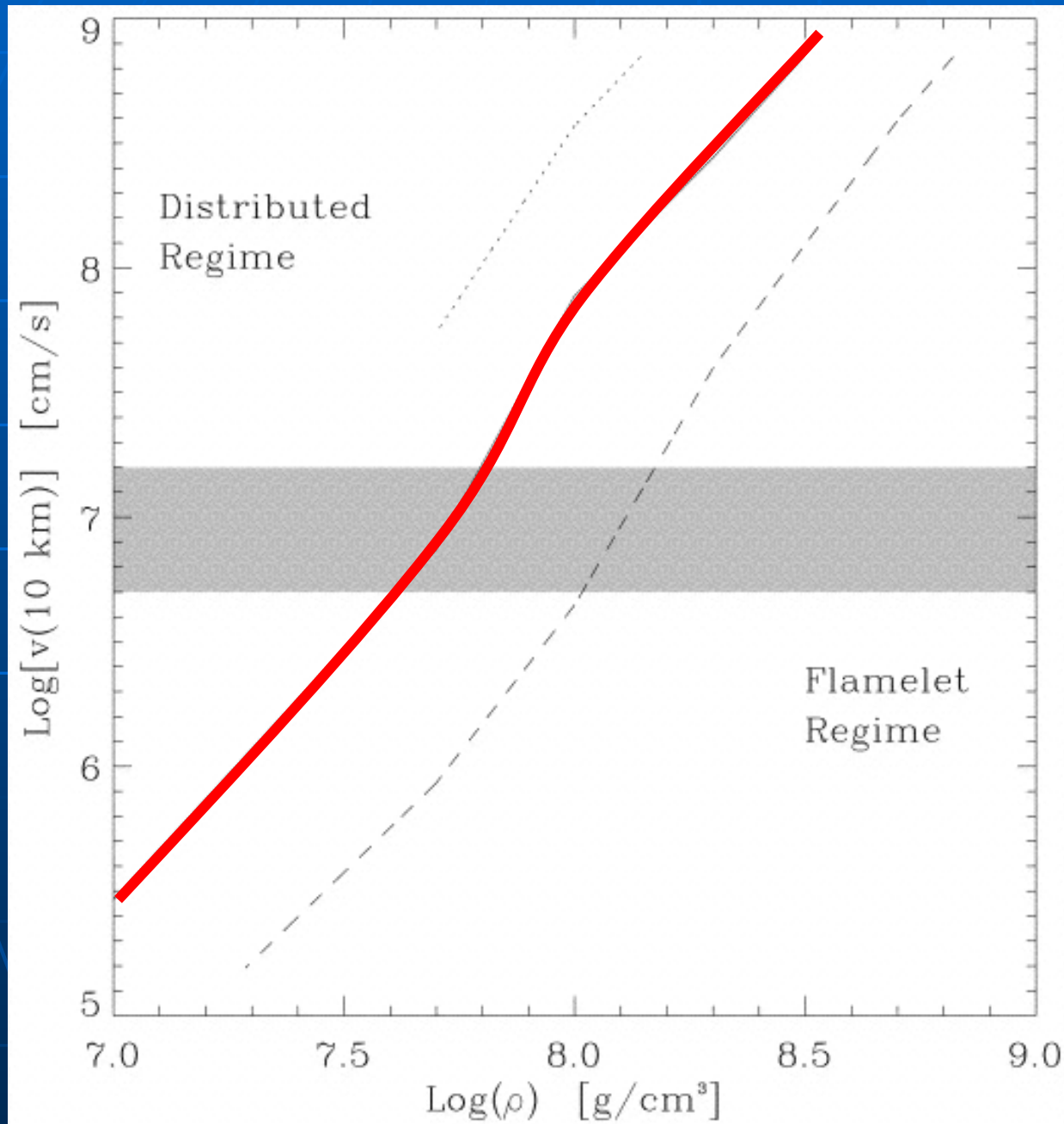
$$l_F \geq l_{\text{Gibs}}$$

Rough estimate (“Damköhler scaling”):

$$s_{\text{turb}}/s_L \approx \text{const} (D_t/D)^{1/2} \text{ (dependent on } s_L \text{ !!!)}$$

$$\text{const} = O(1)$$

Application to type Ia supernova



Niemeyer &
Woosley (1997)

Burning regimes of pre-mixed flames

4. The Rayleigh-Taylor regime

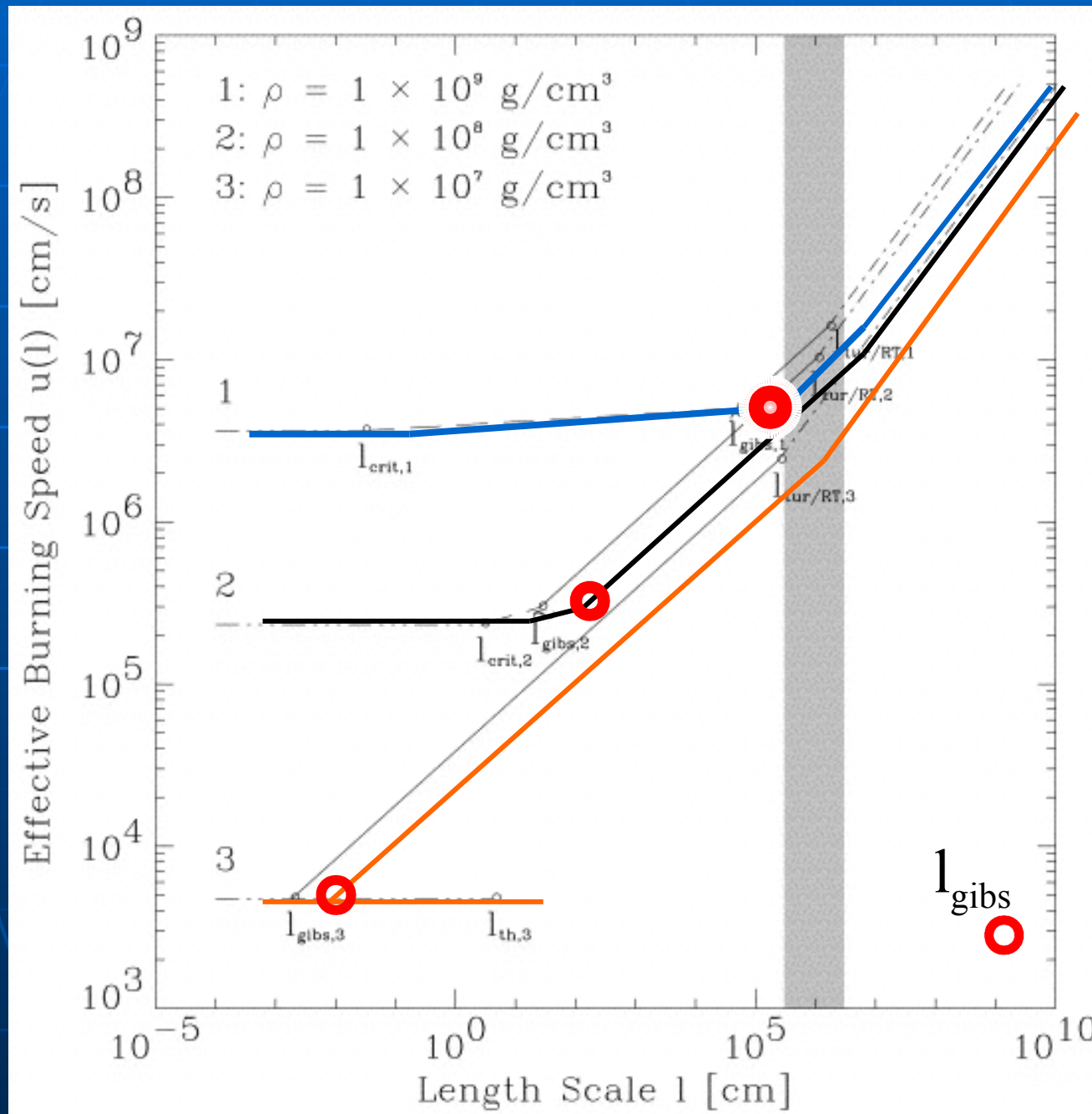
$$v_{RT} = B \sqrt{(g_{eff} l)} ; B \approx 0.5 ; g_{eff} = At \cdot g$$

Sharp-Wheeler model:

$$r_{sw} \approx 0.05 g_{eff} t^2 ; v_{sw} \approx 0.1 g_{eff} t ;$$

$$l_{tur/RT} \approx 10^6 \text{ cm}$$

Effective burning velocities in SN Ia



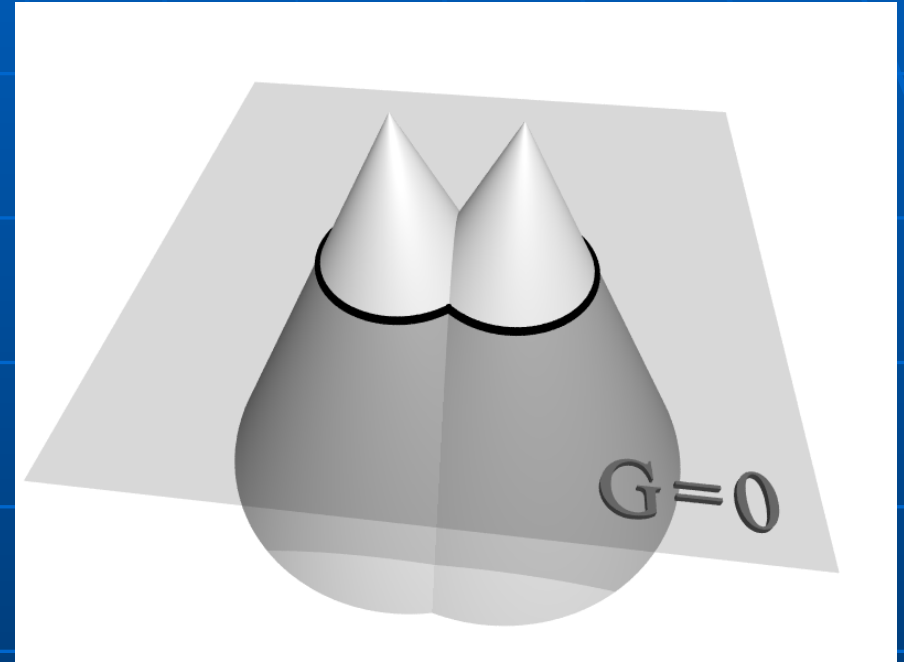
Niemeyer &
Woosley
(1997)

How to model thermonuclear flames?

- The "flames" cannot be resolved numerically.
- The amplitudes of turbulent velocity fluctuations in the length scale of the flame are determined on the integral scale.



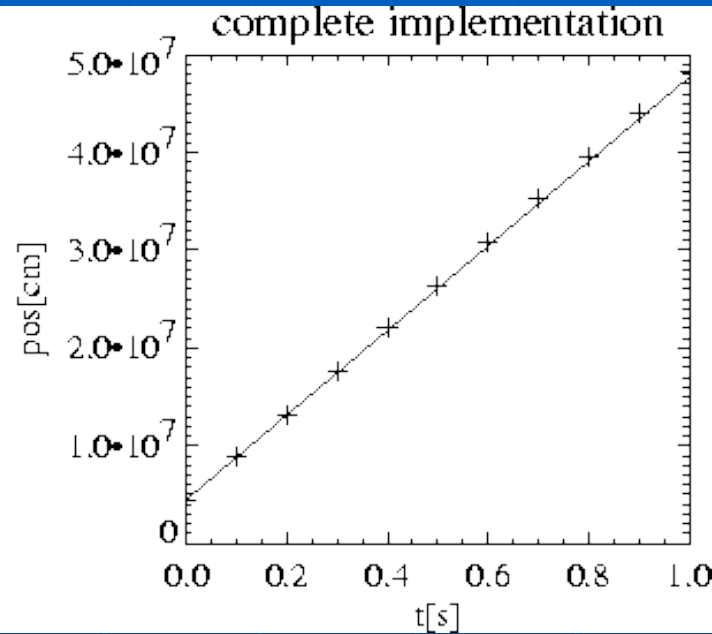
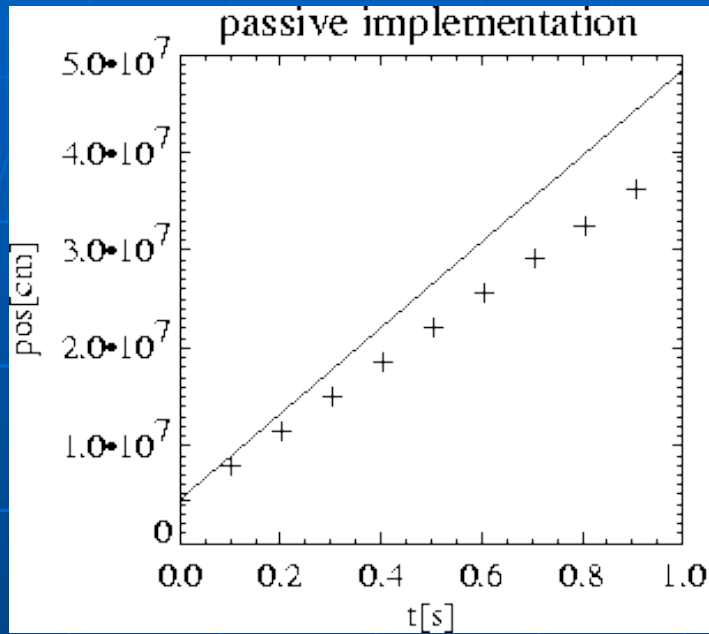
"LES" + "Level Set"



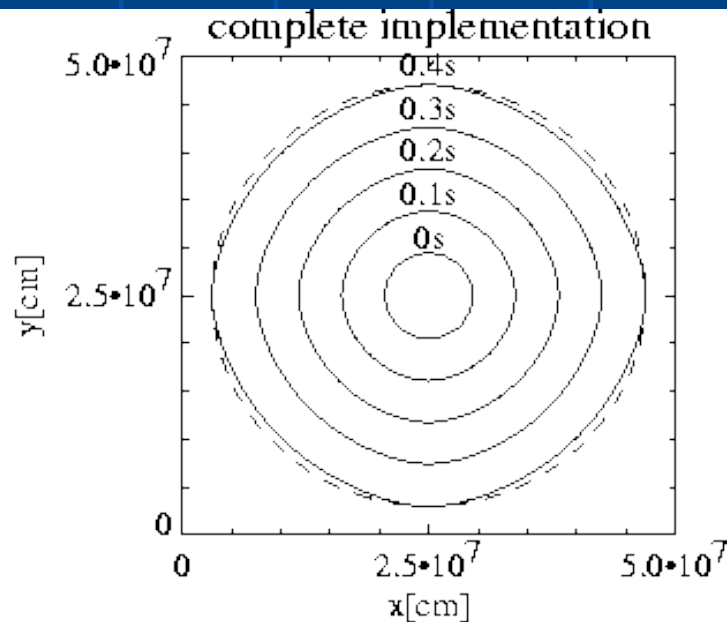
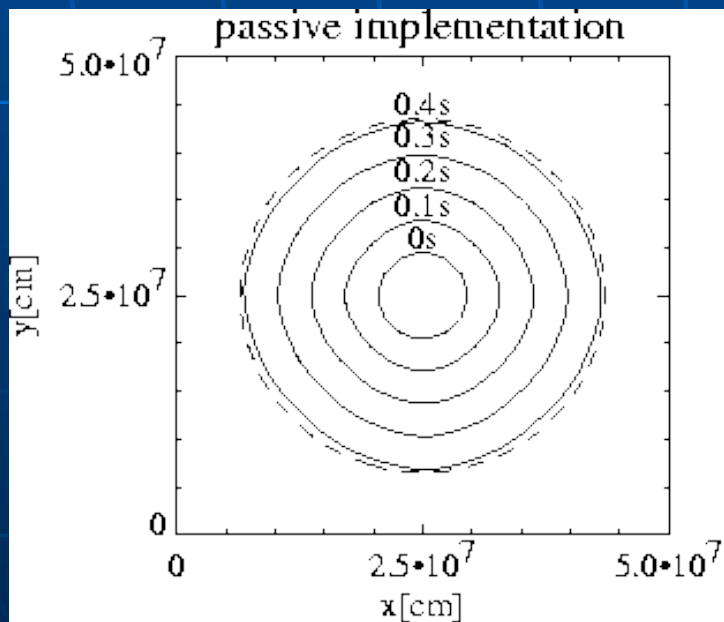
$$\partial G / \partial t = -\mathcal{D}_f \nabla G$$

$$\mathcal{D}_f = \mathbf{v}_u + s_{\text{tur}} \mathbf{n}; \quad |\nabla G| = 1$$

Some test of the code



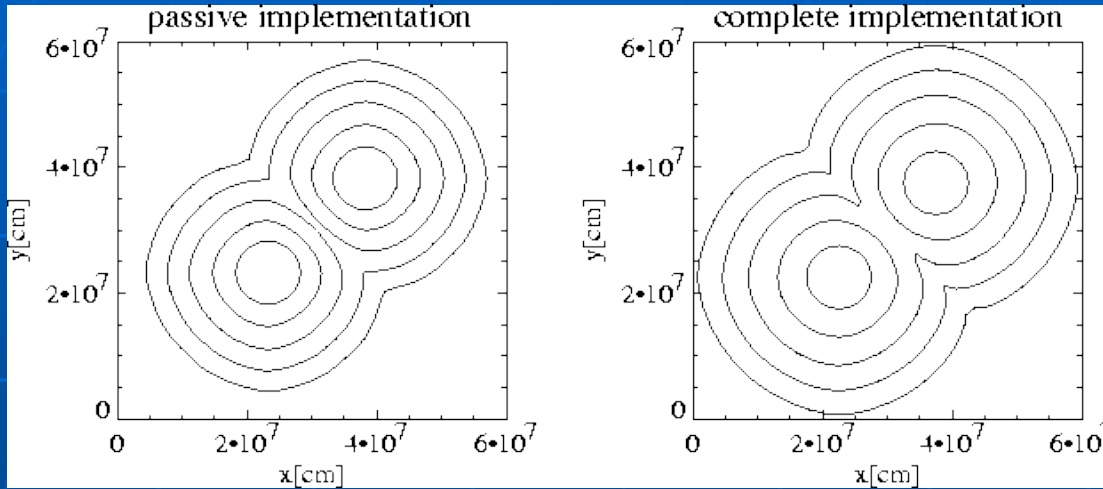
Planar flame



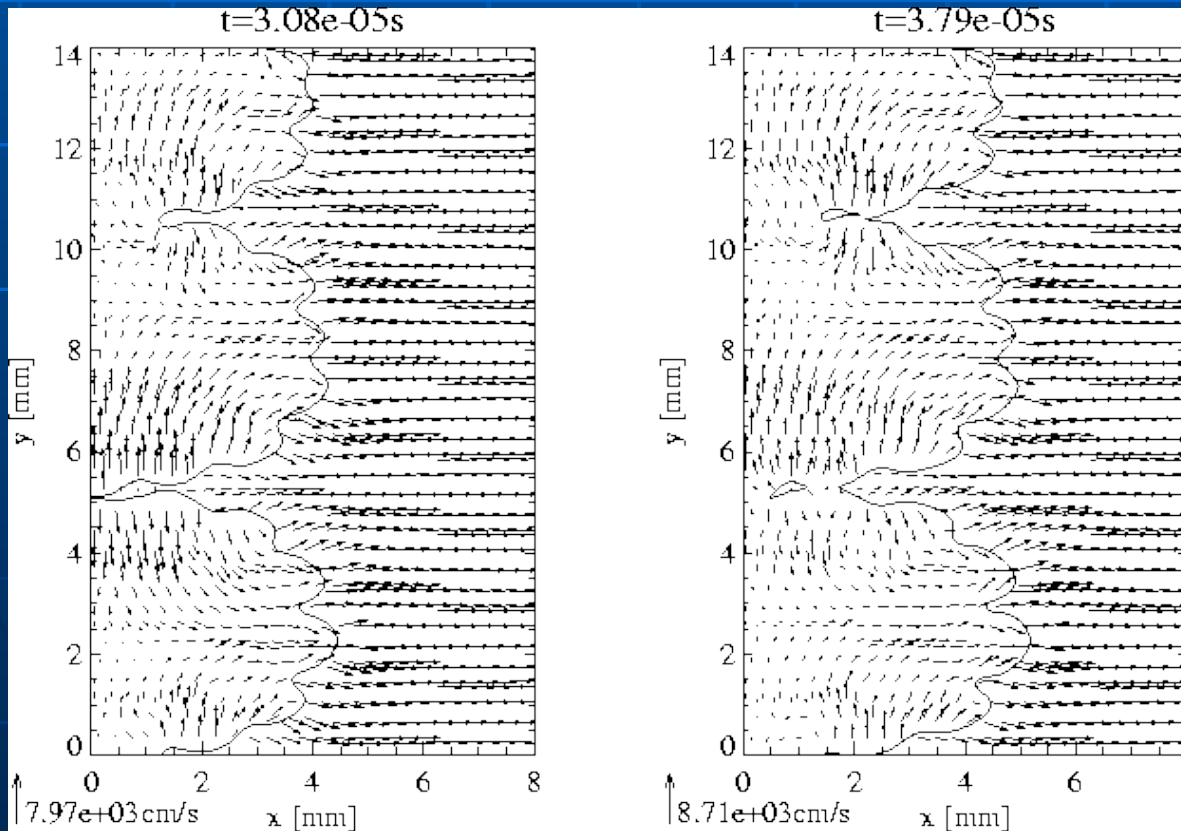
Circular flame

Reinecke et al.
(1999)

Some test of the code (ctn.)



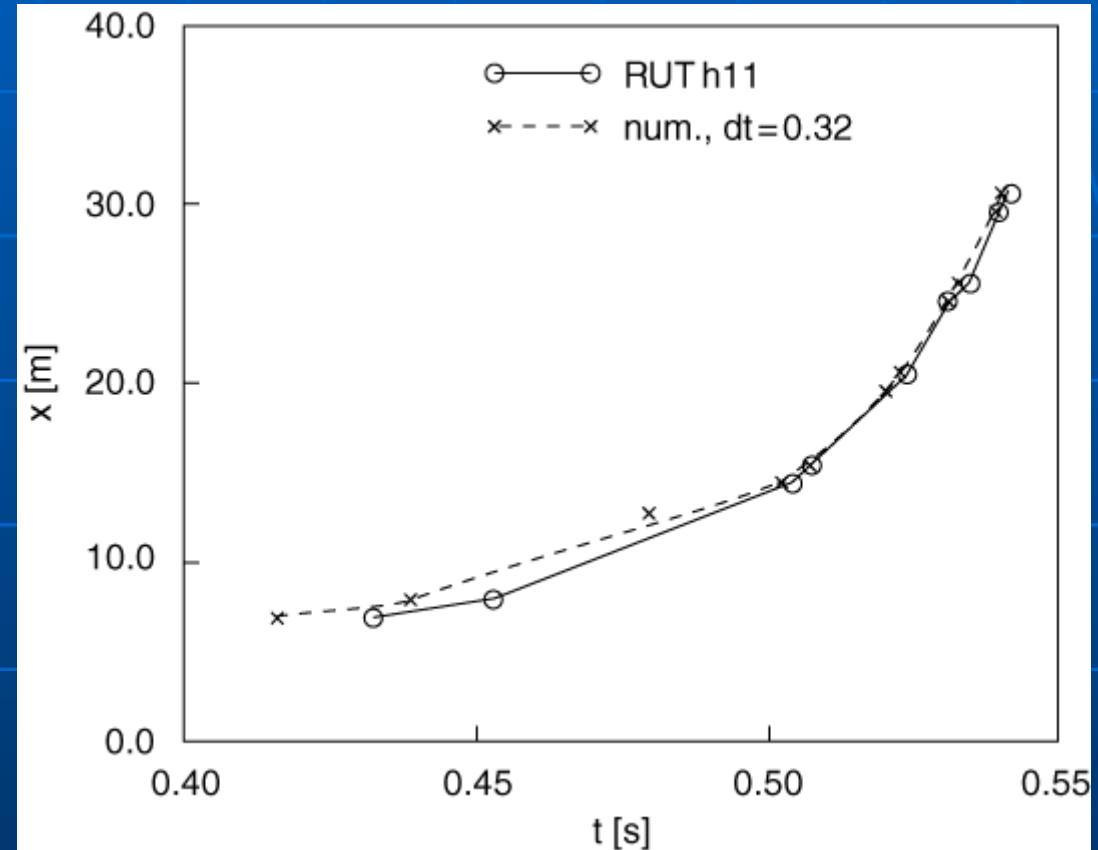
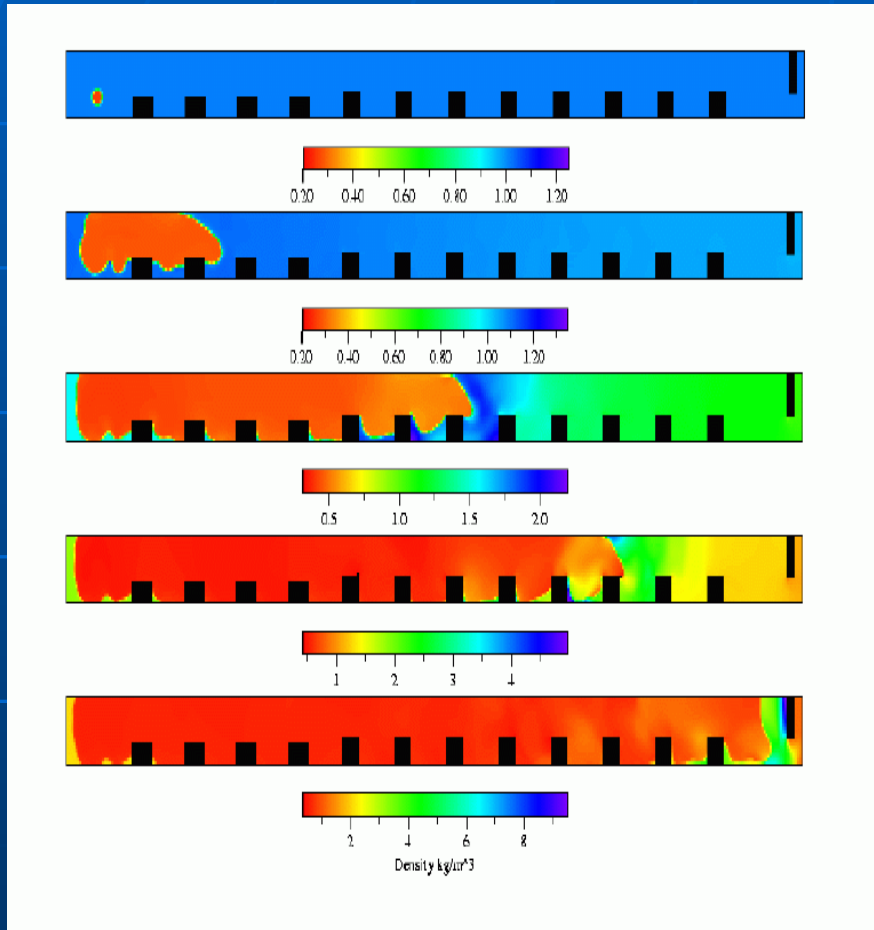
Merging circular flames



Hydrogen-in-air flames

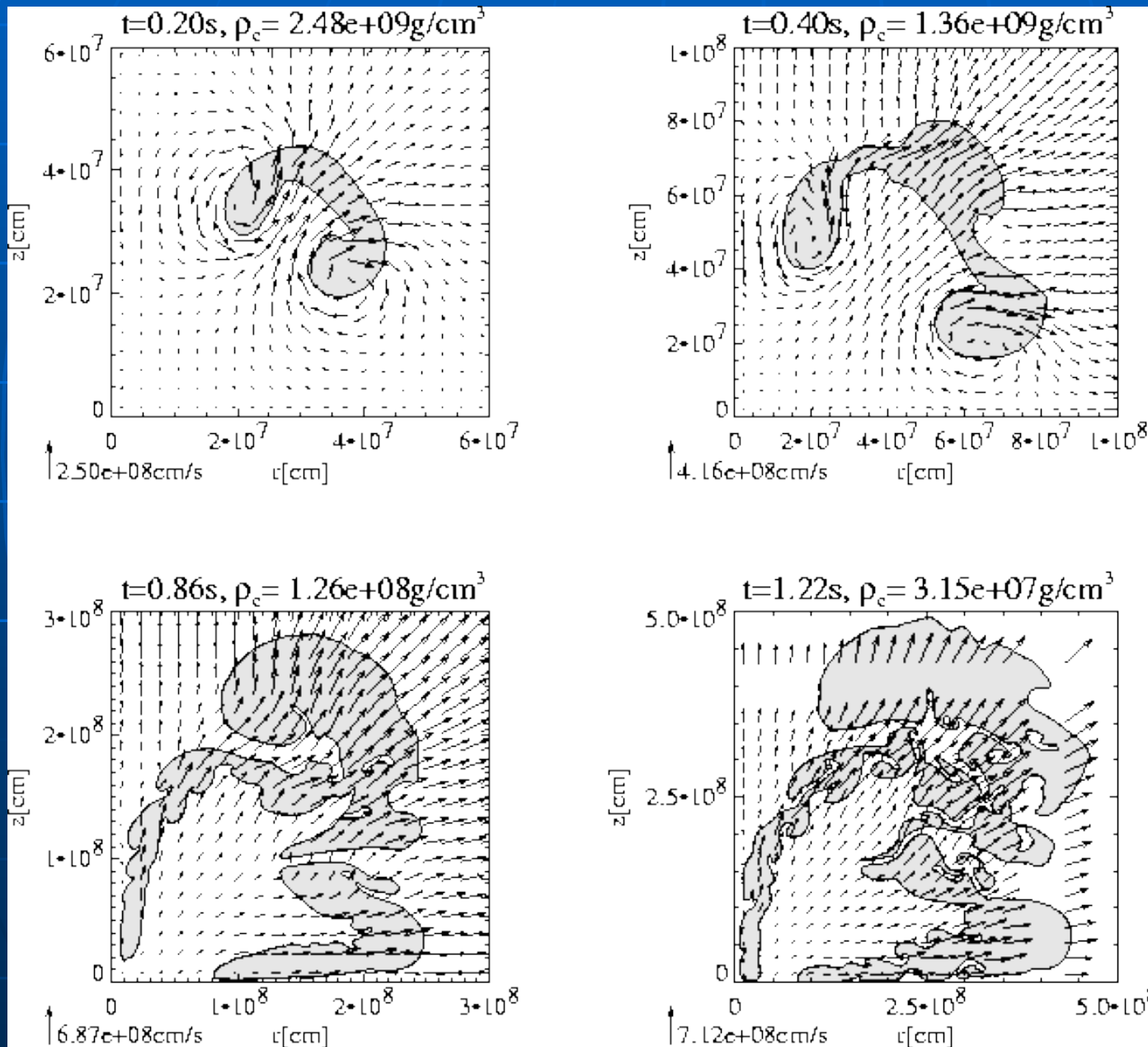
Reinecke et al. (1999)

Other laboratory flames



The method can reproduce terrestrial experiments well!
(Smiljanowski et al. 1997)

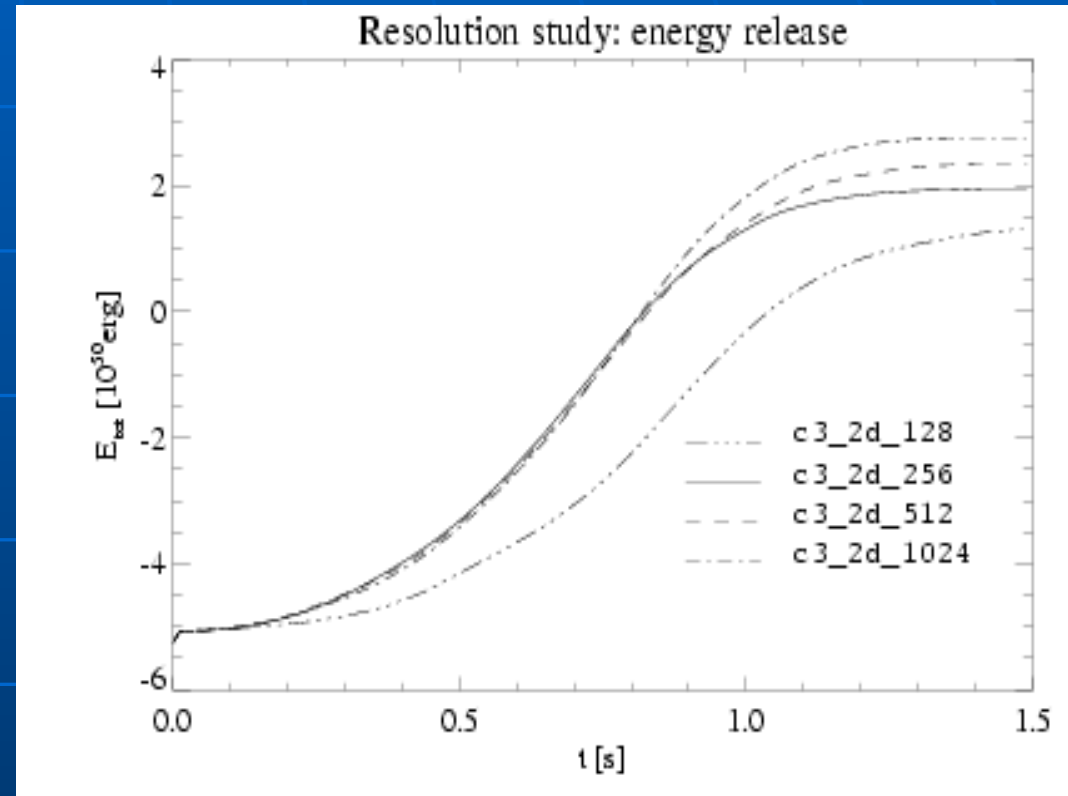
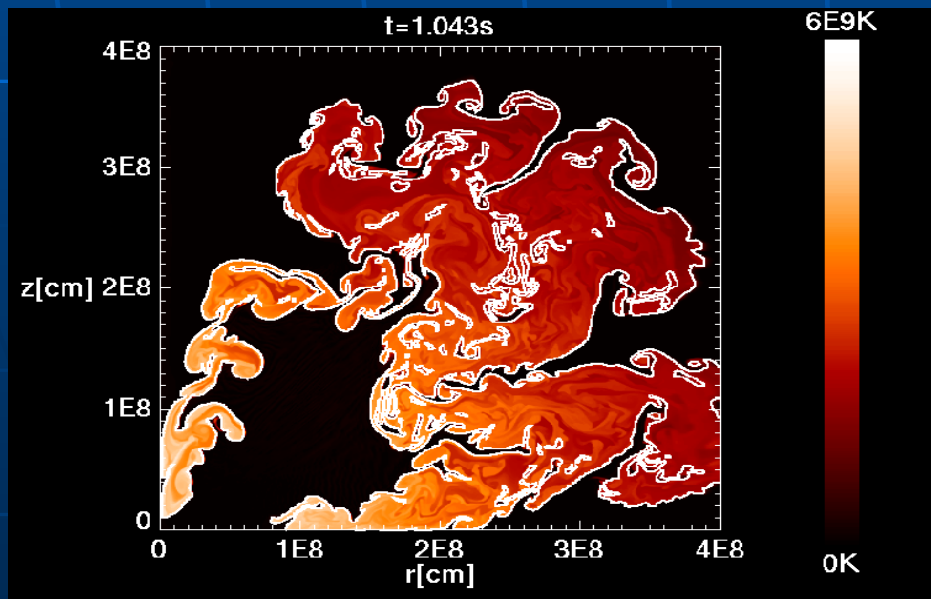
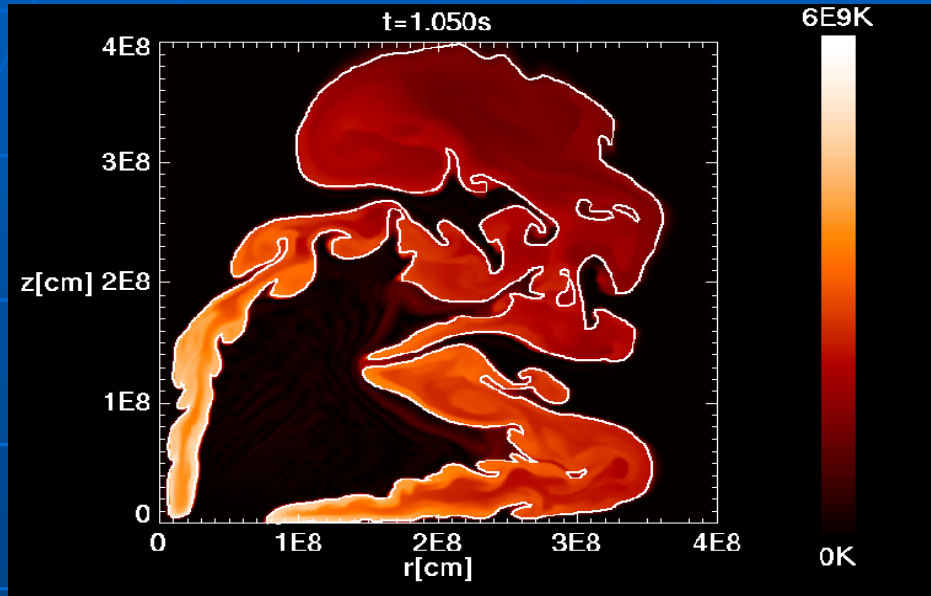
Application to the SN Ia problem



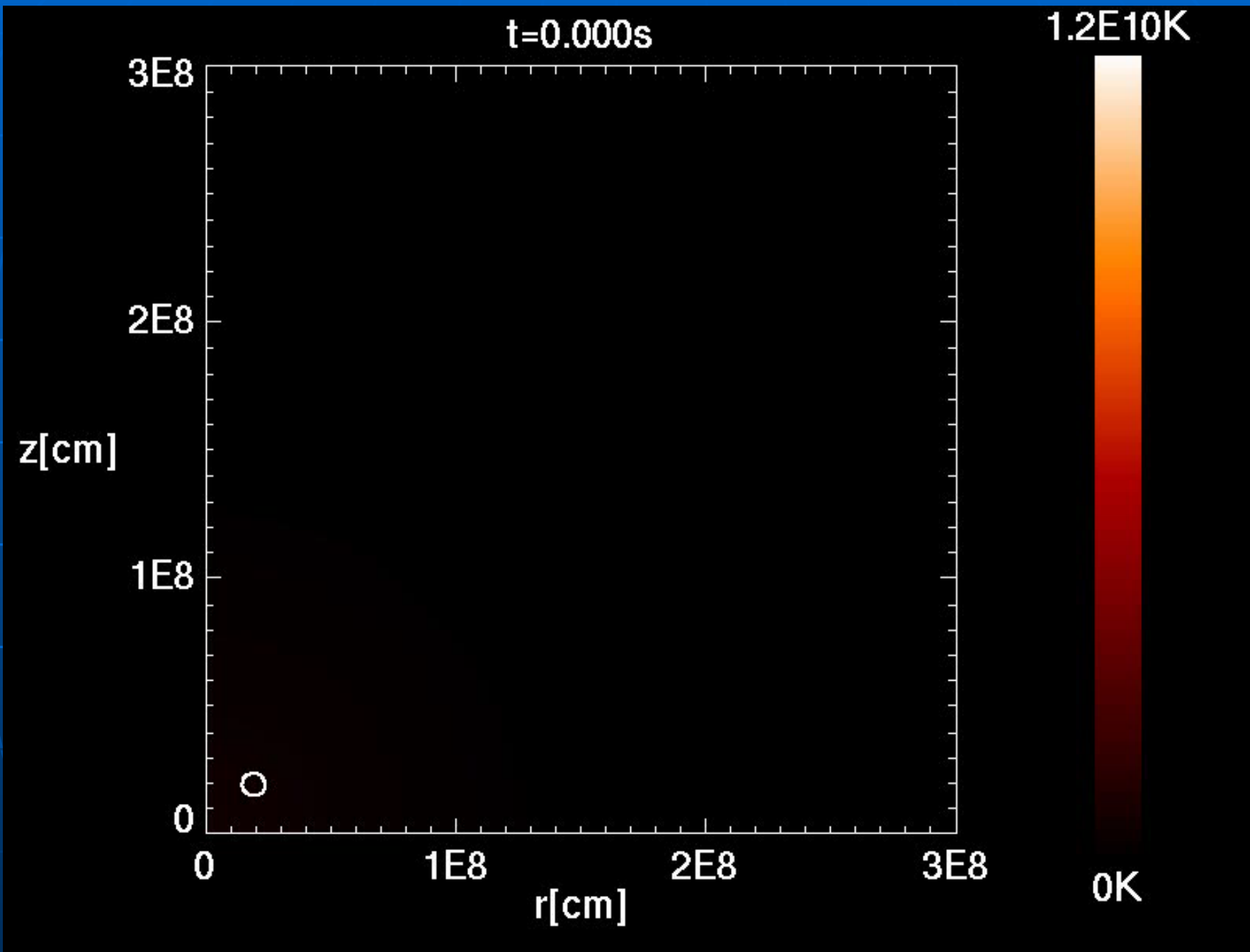
One rising
blob (in 2D)

Reinecke et al.
(1997)

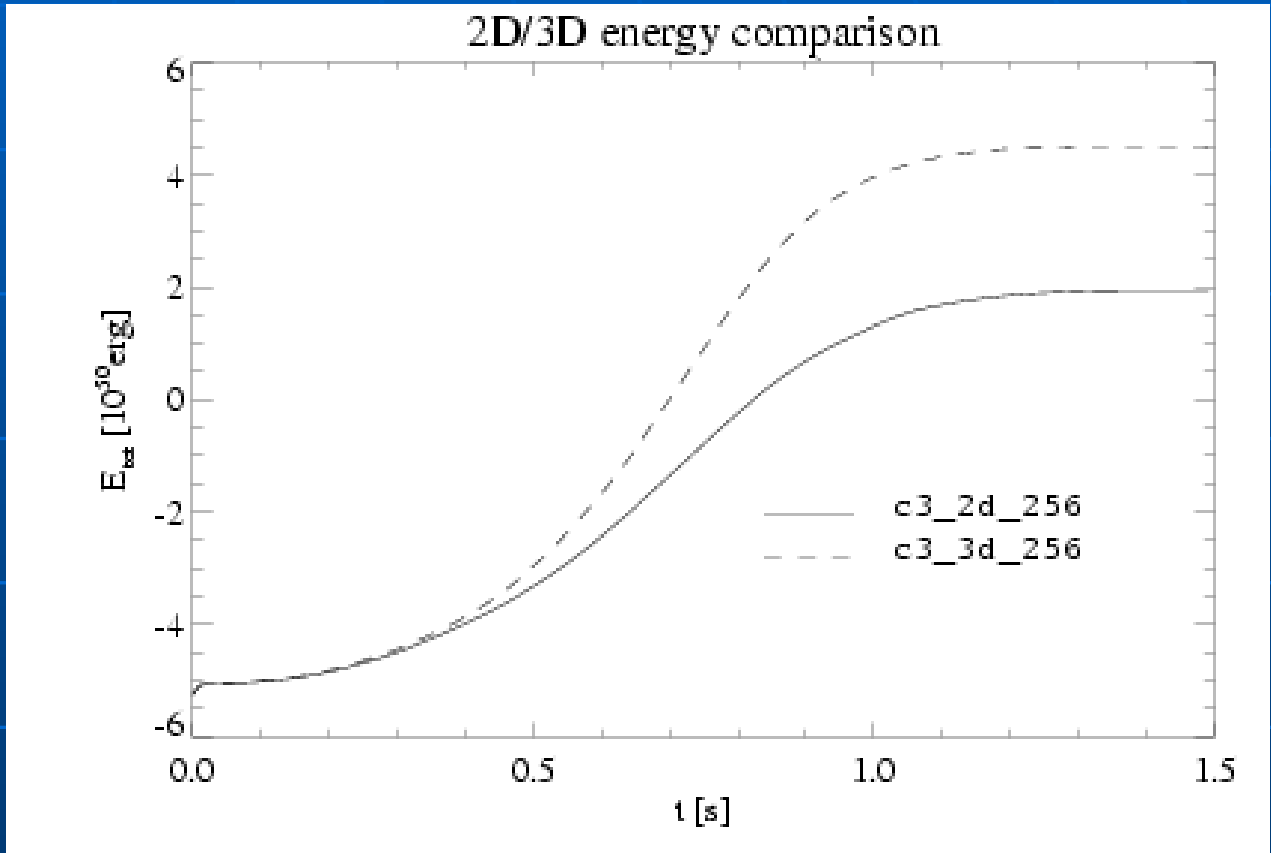
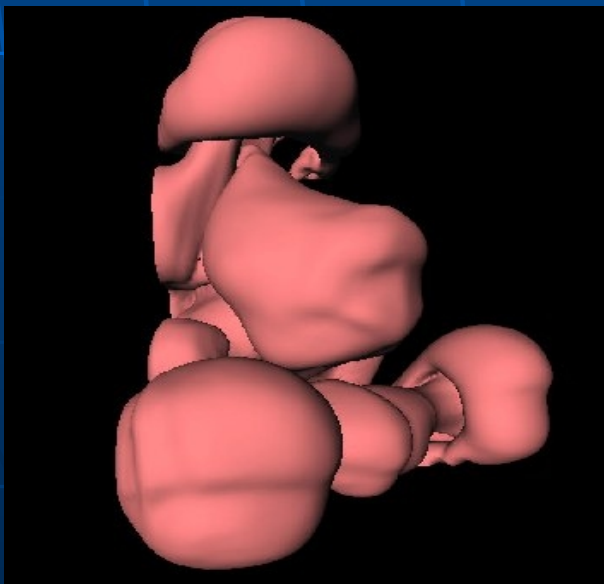
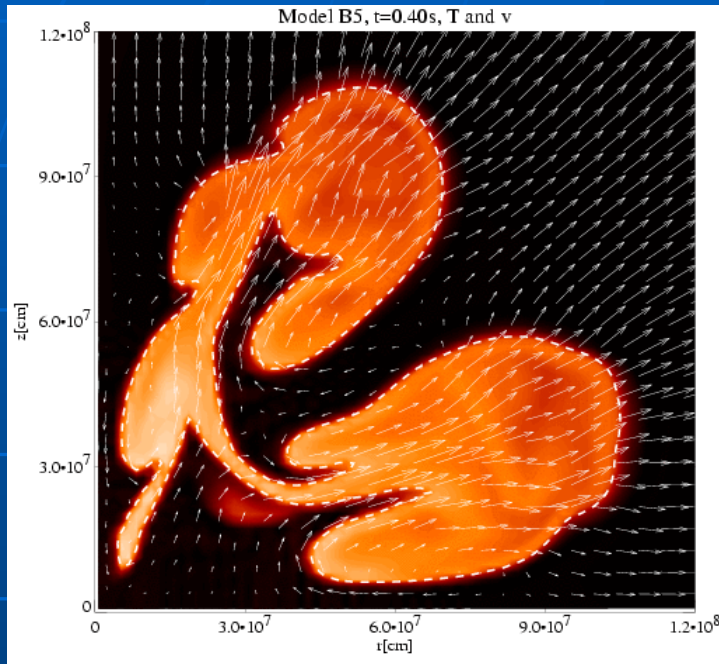
Convergence tests in 2D



Global results are
independent of the
numerical
resolution!
Reinecke et al. (1999, 2002)



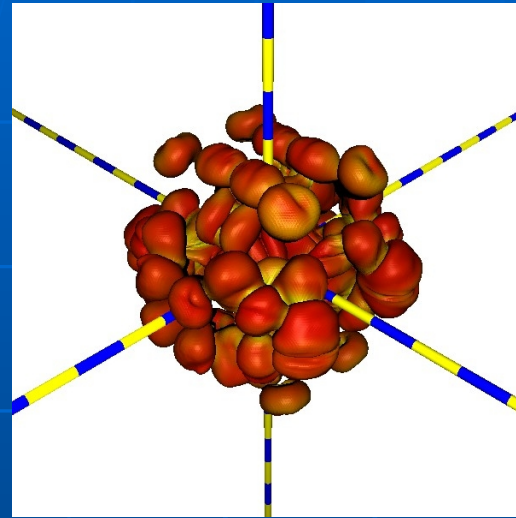
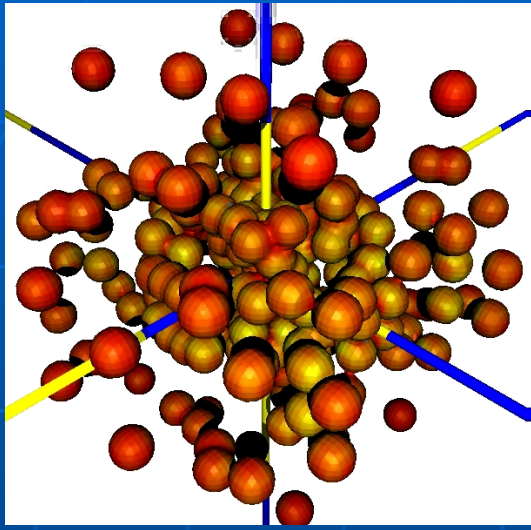
2D \Rightarrow 3D



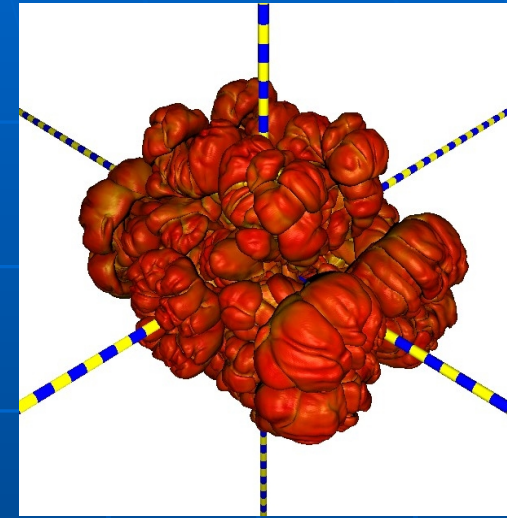
Because of larger surface area:
More energy is produced!

Reinecke et al. (2001)
(See also Gamezo et al., 2003)

3D models: The best we could do until recently

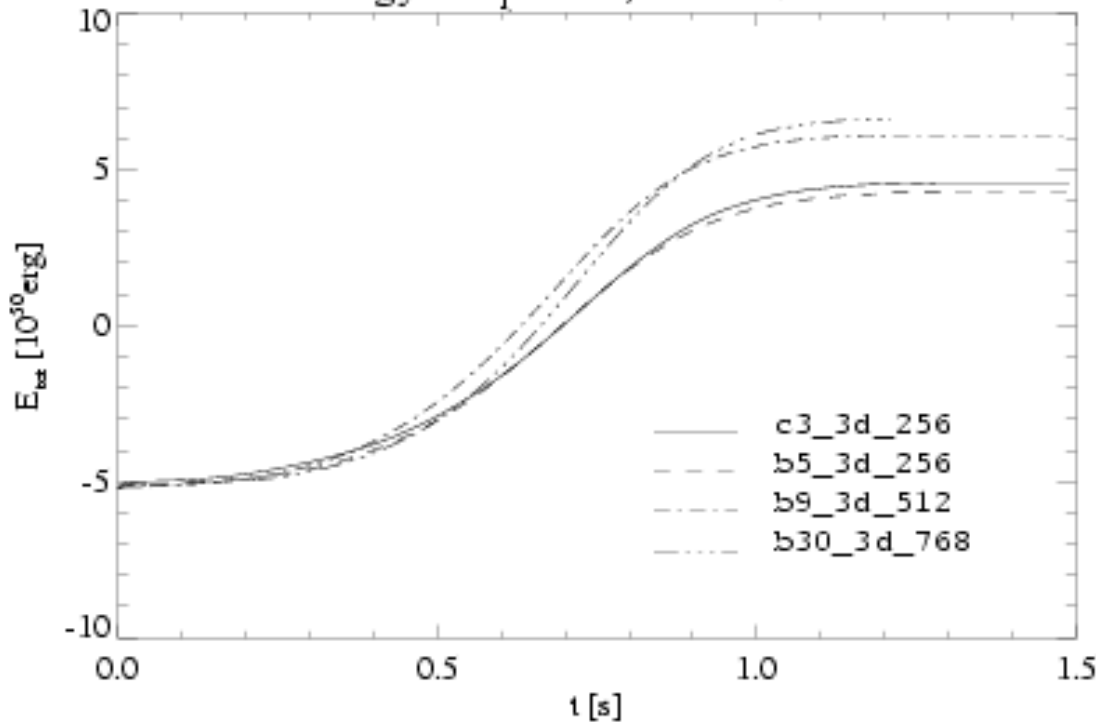


0.25s



0.6s

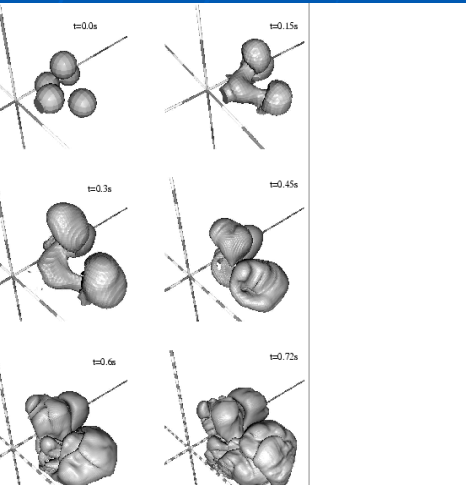
Energy comparison, 3D simulations



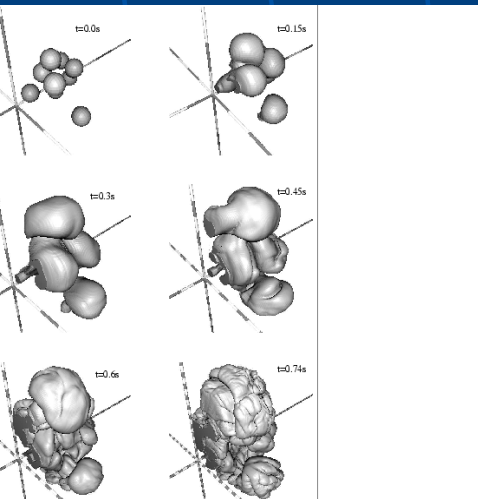
Mod b30_3d

(Reinecke et al., 2003)

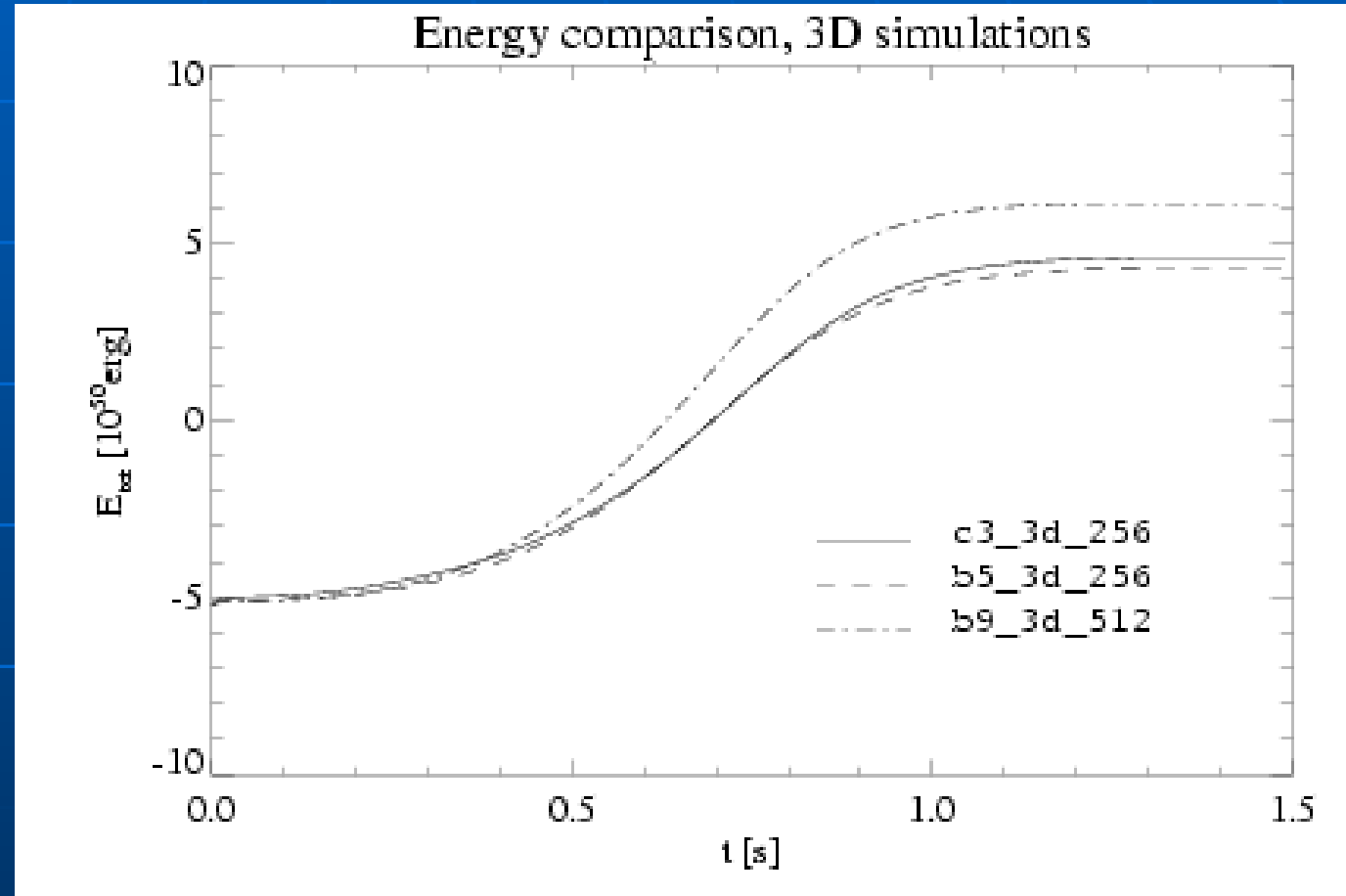
Modeling Flames in 3D: Dependence on initial conditions?



Mod.
b5_3d



Mod.
b9_3d



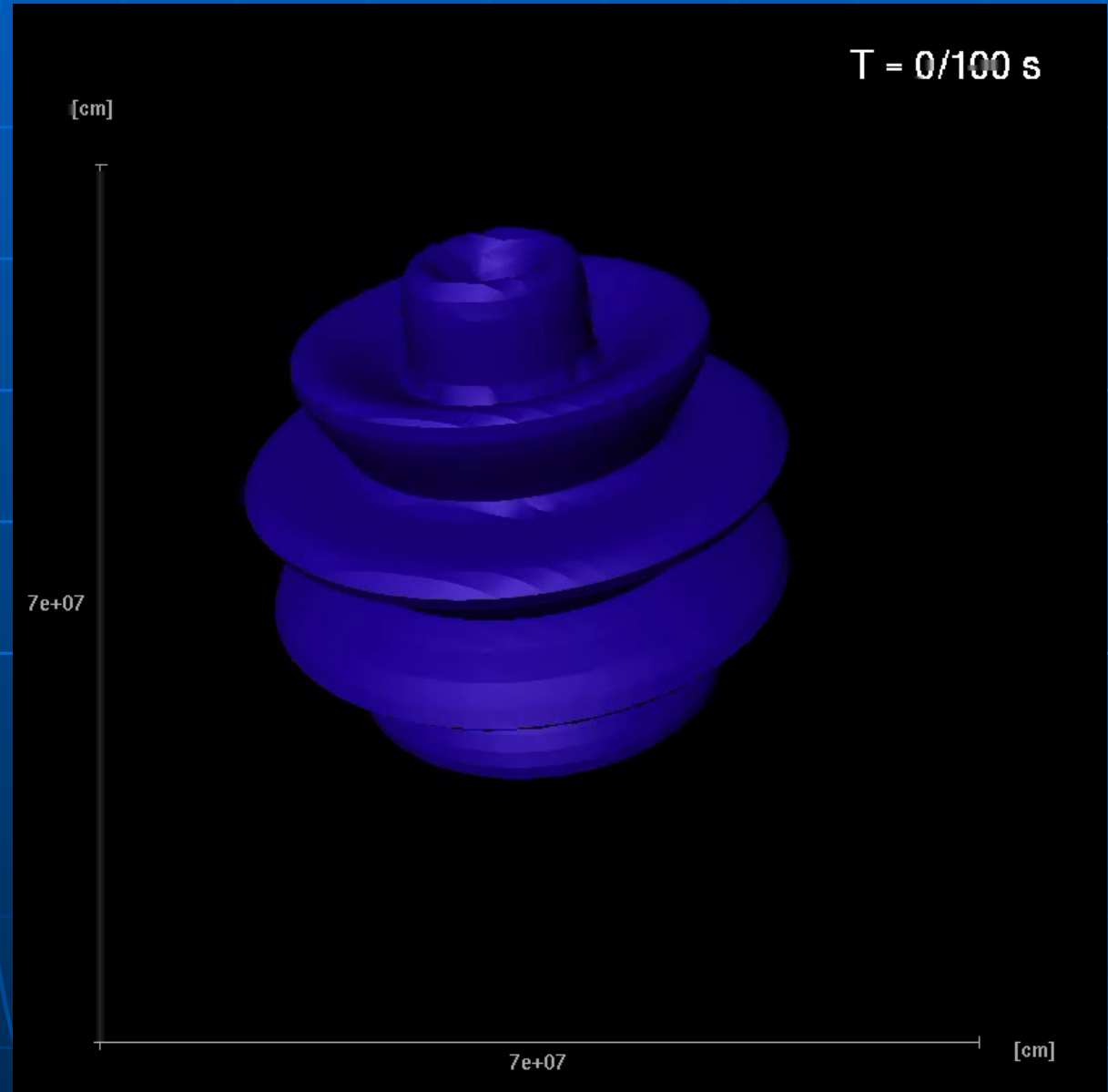
Moderate dependence
on initial conditions!

(Reinecke et al., 2002)

Recent modifications of the code:

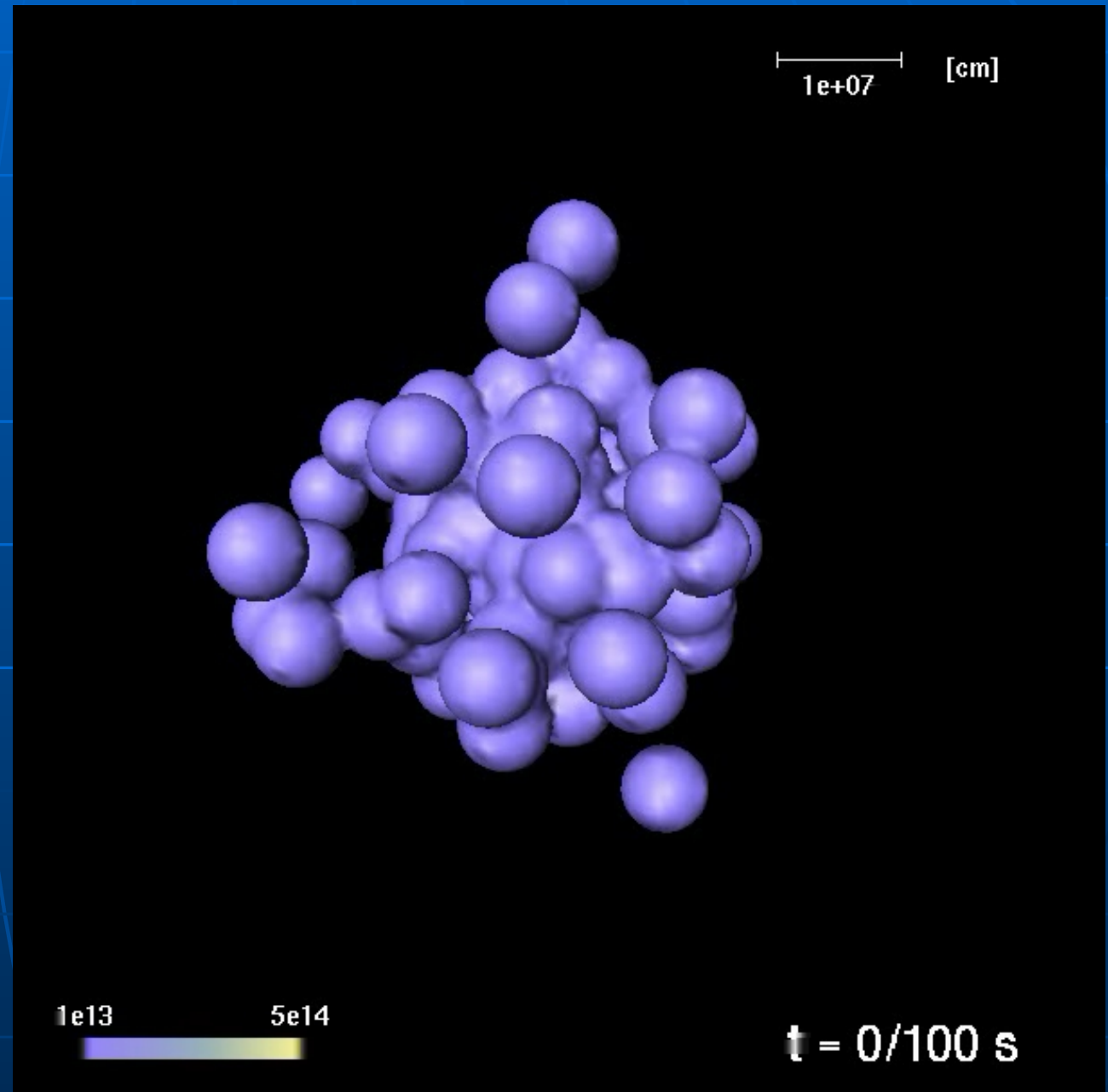
1. Moving grid

Röpke (2004)



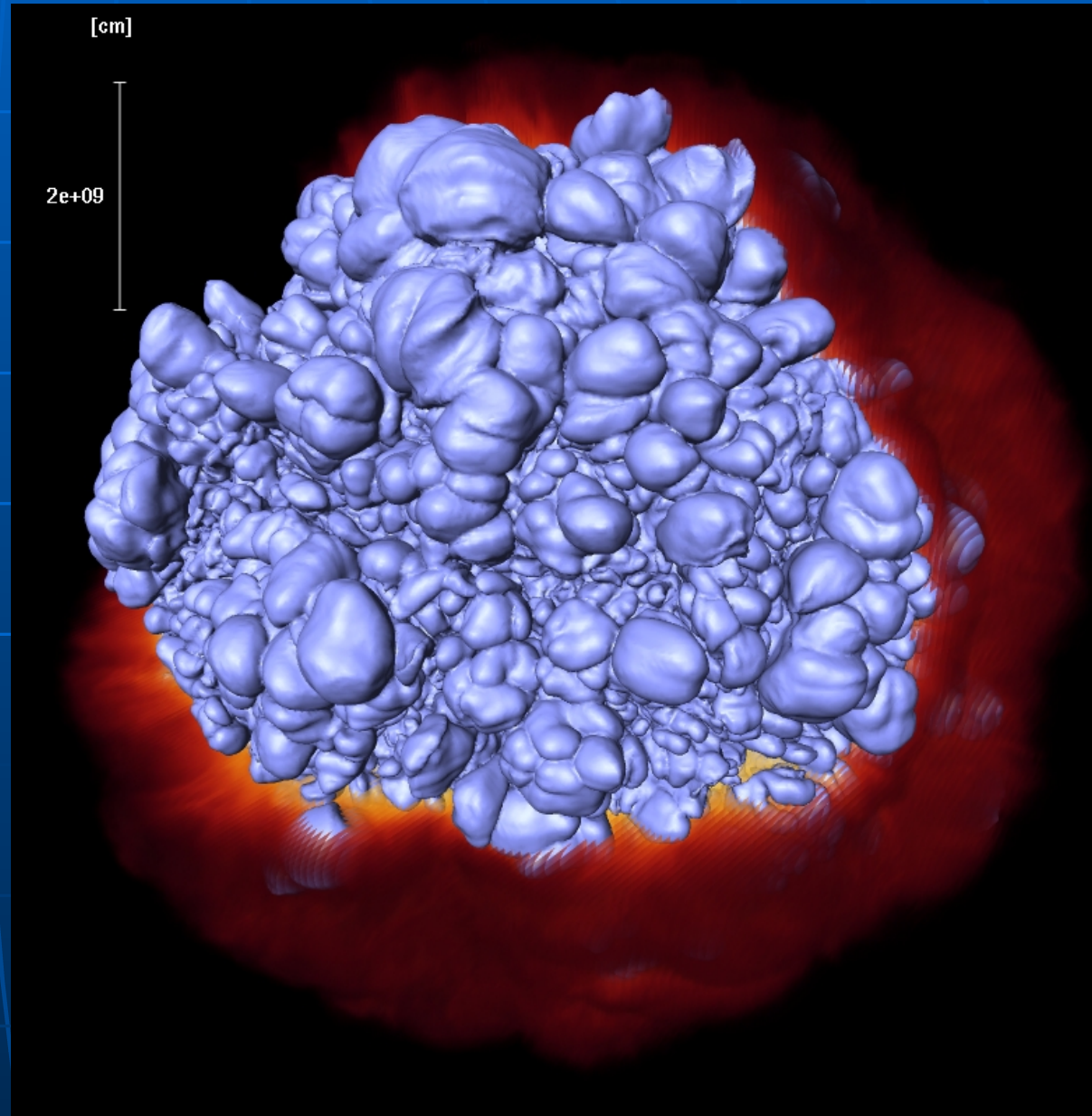
2. Full star (“ 4π ”)

Röpke & Hillebrandt
(2004)

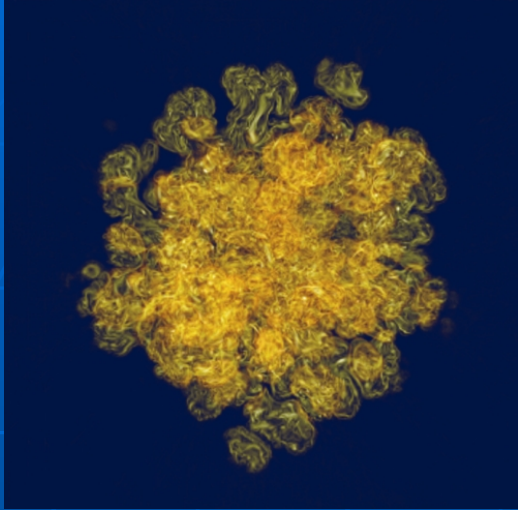


A high-resolution model (“the SNOB run”)

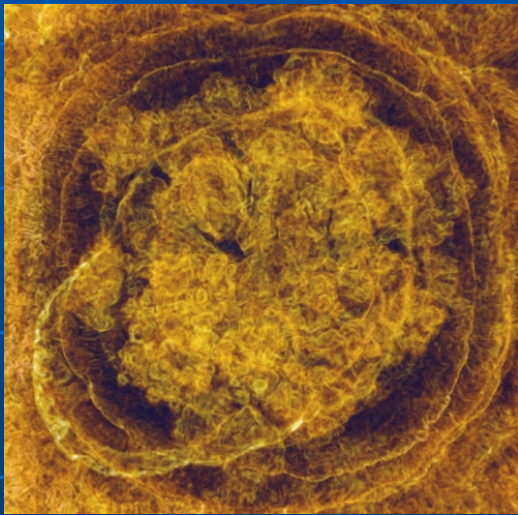
- “ 4π ”
- 1024^3 grid
- initial resolution near the center $\approx 800\text{m}$
- moving grid
- Local & dynamical sgs-model
- $\sim 100,000$ h on 512 processors, IBM/Power4, at RZG



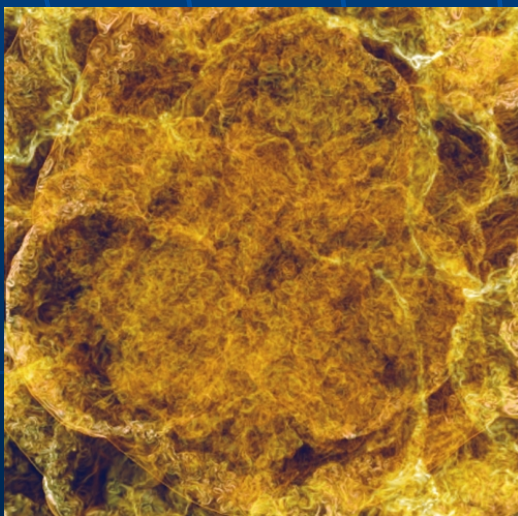
Turbulence?



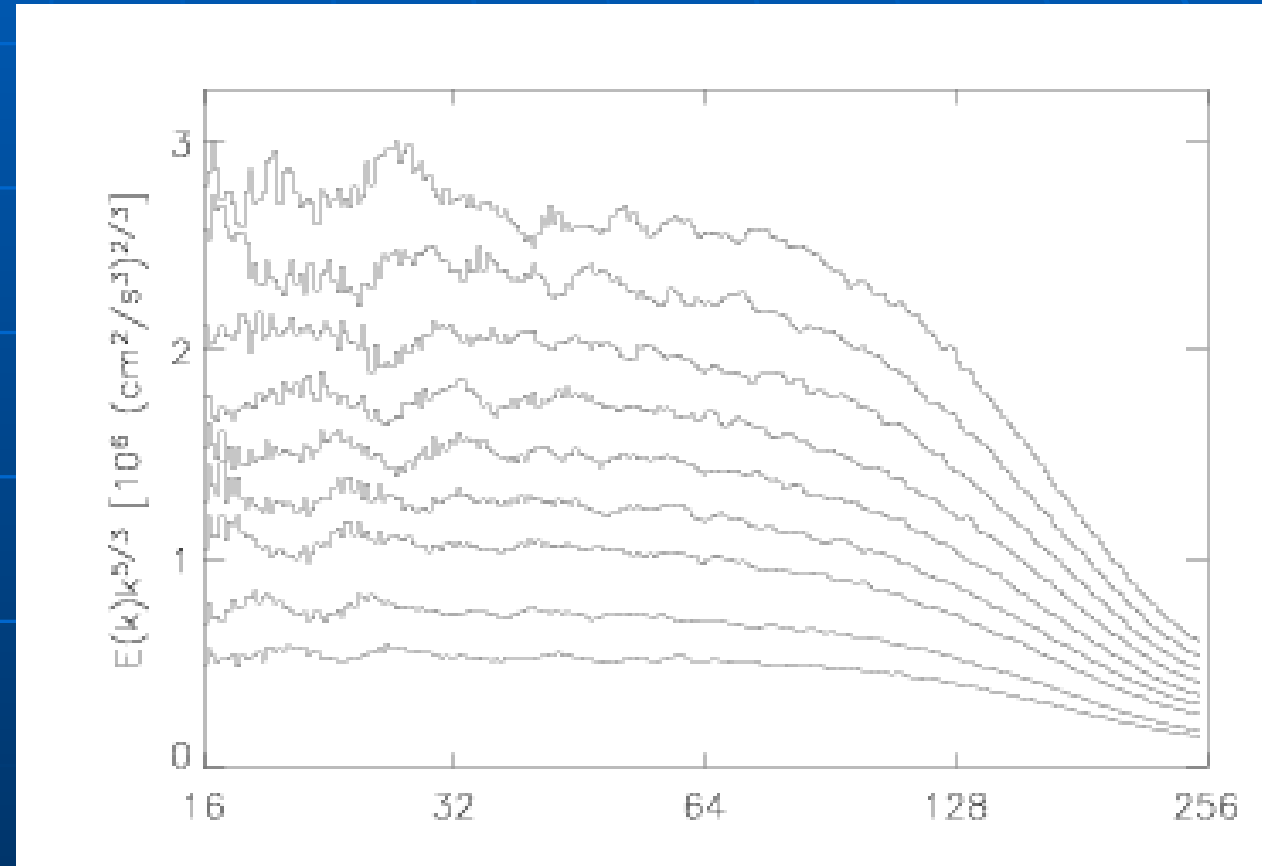
0.25s



0.50s



0.75s



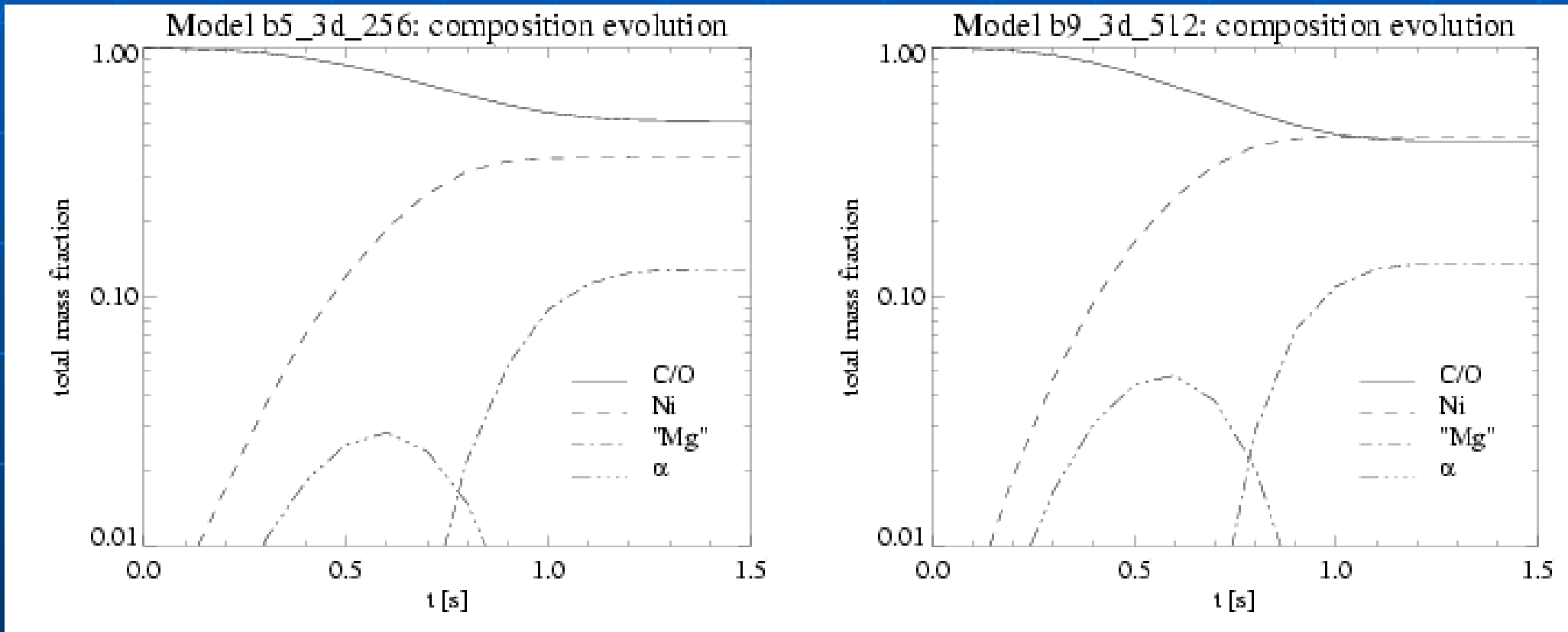
Schmidt et al. (in preparation)

Some (preliminary) results:

- $E_{\text{kin}} = 8.1 \cdot 10^{50}$ erg
- Iron-group nuclei: $0.61 M_{\text{sun}}$ ($\sim 0.41 M_{\text{sun}} {}^{56}\text{Ni}$)
- Intermediate-mass nuclei: $0.43 M_{\text{sun}}$ (from hydro)
- Unburnt C+O: $0.37 M_{\text{sun}}$ (from hydro)
(less than $0.08 M_{\text{sun}}$ at $v < 8000 \text{ km/s}$)
- $V_{\text{max}} \approx 17,000 \text{ km/s}$

Good agreement with observations!

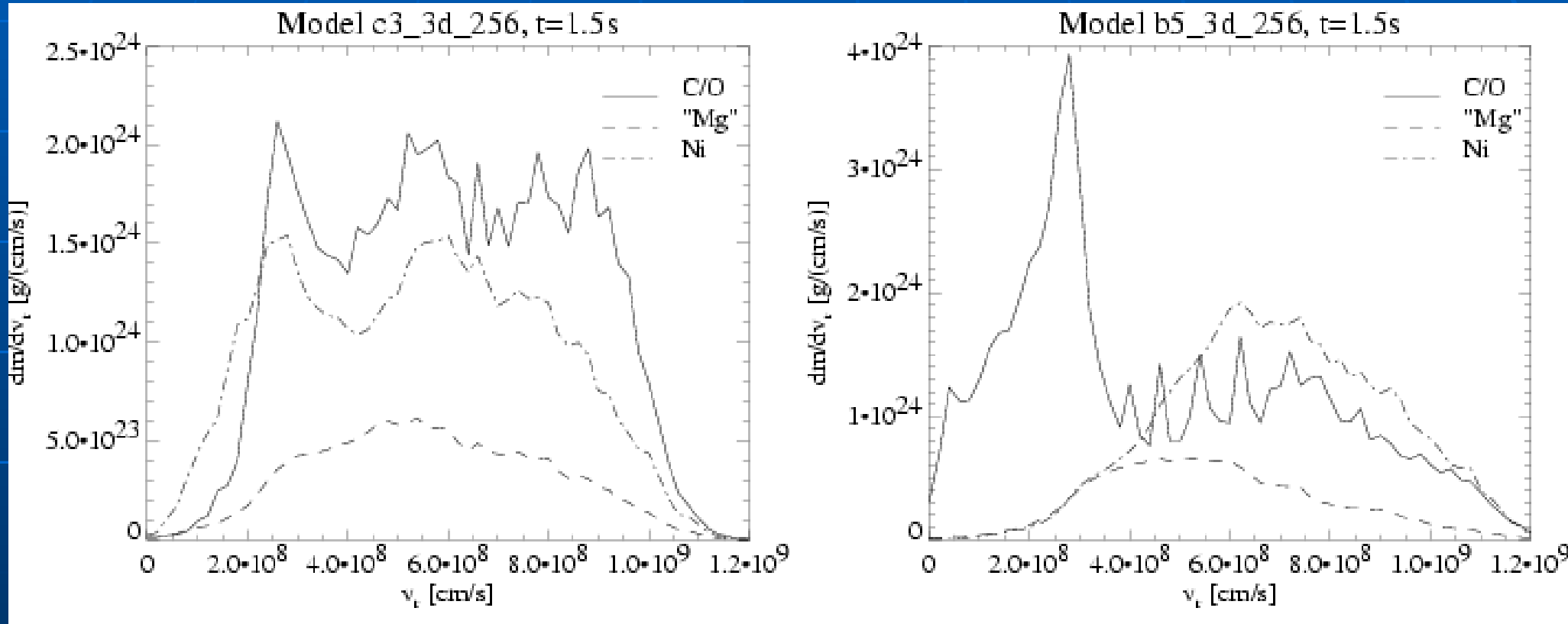
Observable Predictions: Chemical composition?



Significant amounts of unburned C and O!

(Reinecke et al., 2002)

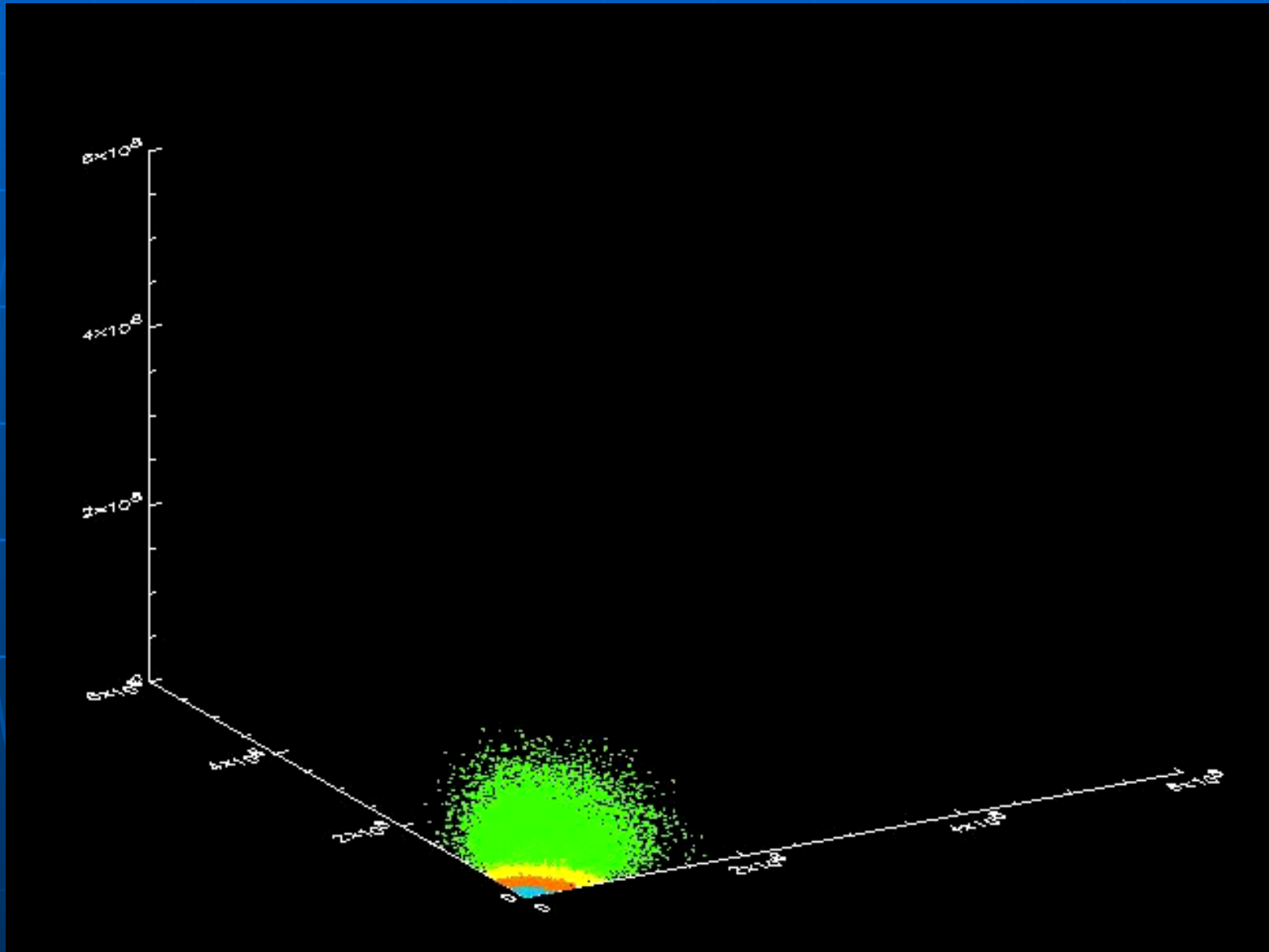
Observable Predictions: Chemical composition in velocity space?



Velocity distribution sensitive to ignition conditions!

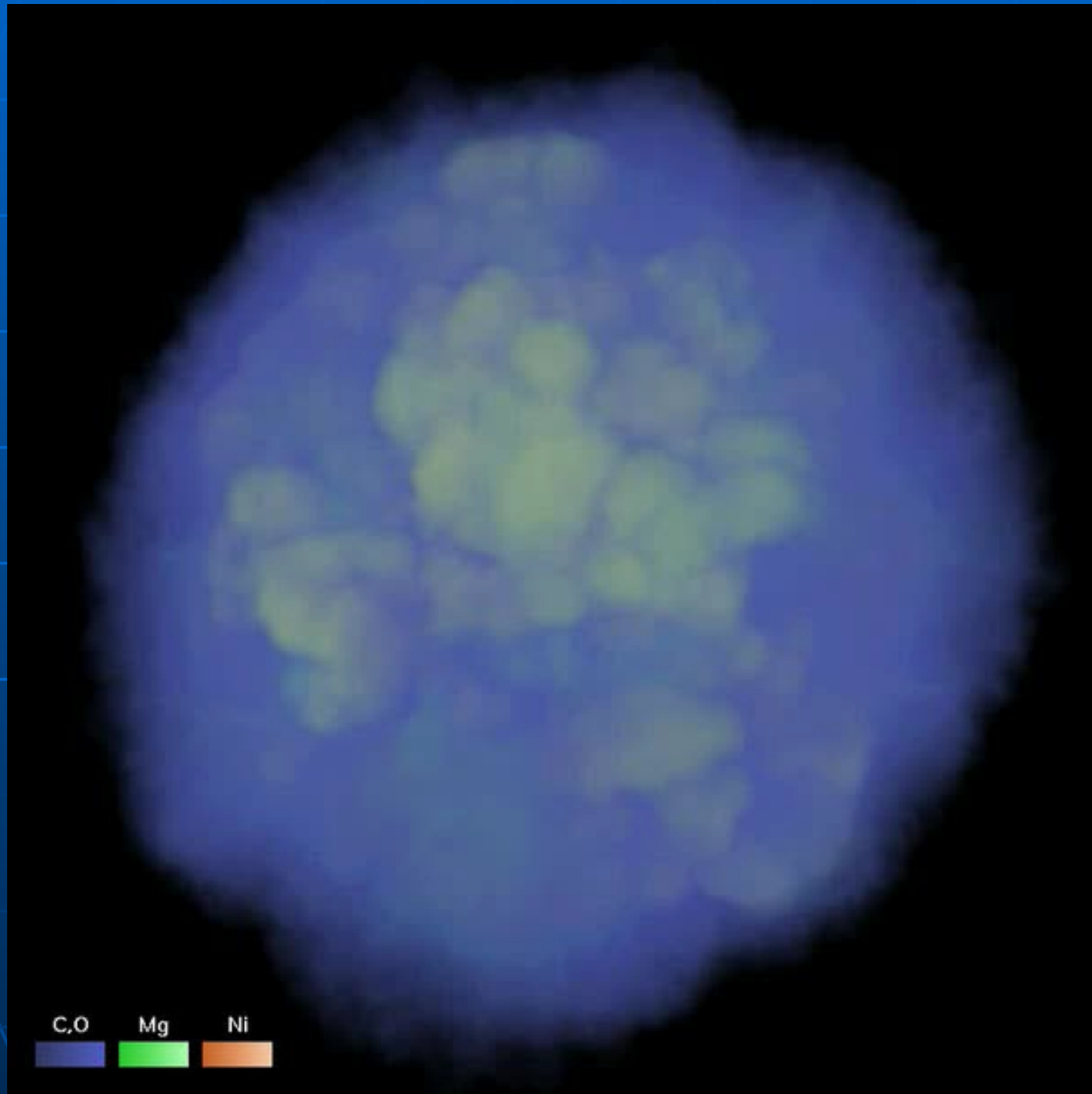
(Reinecke et al., 2002)

Nucleosynthesis (in 'post-processing' mode)

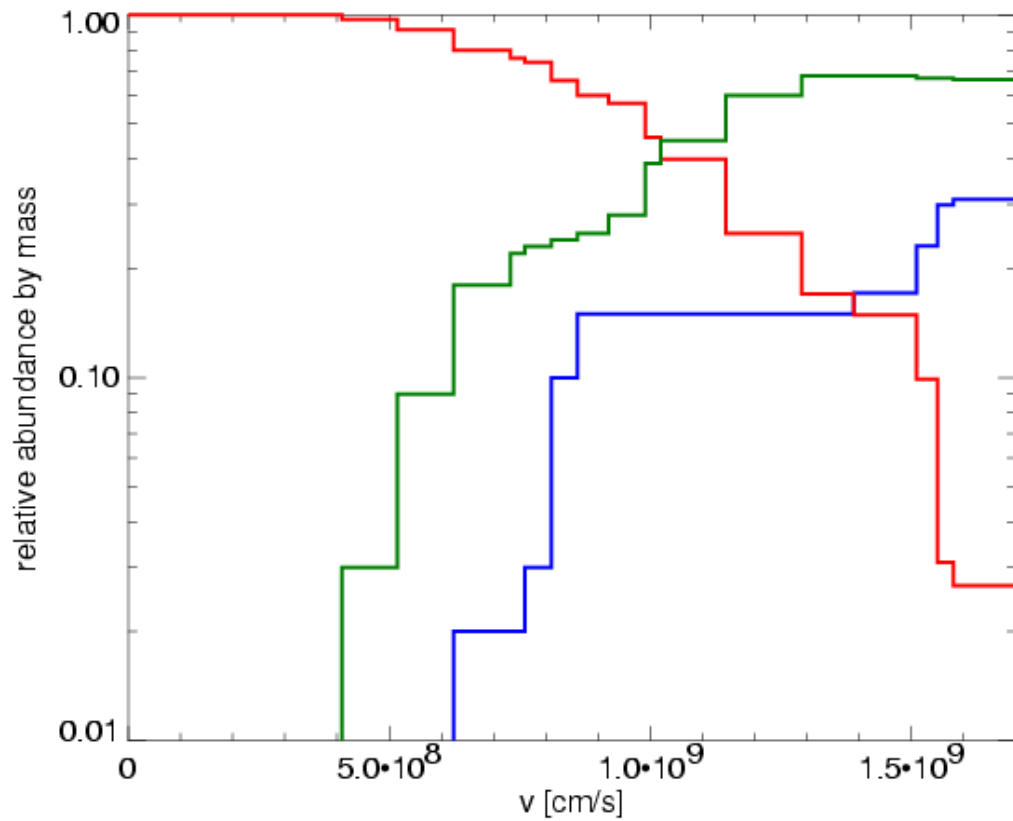


(Travaglio et al., 2004)

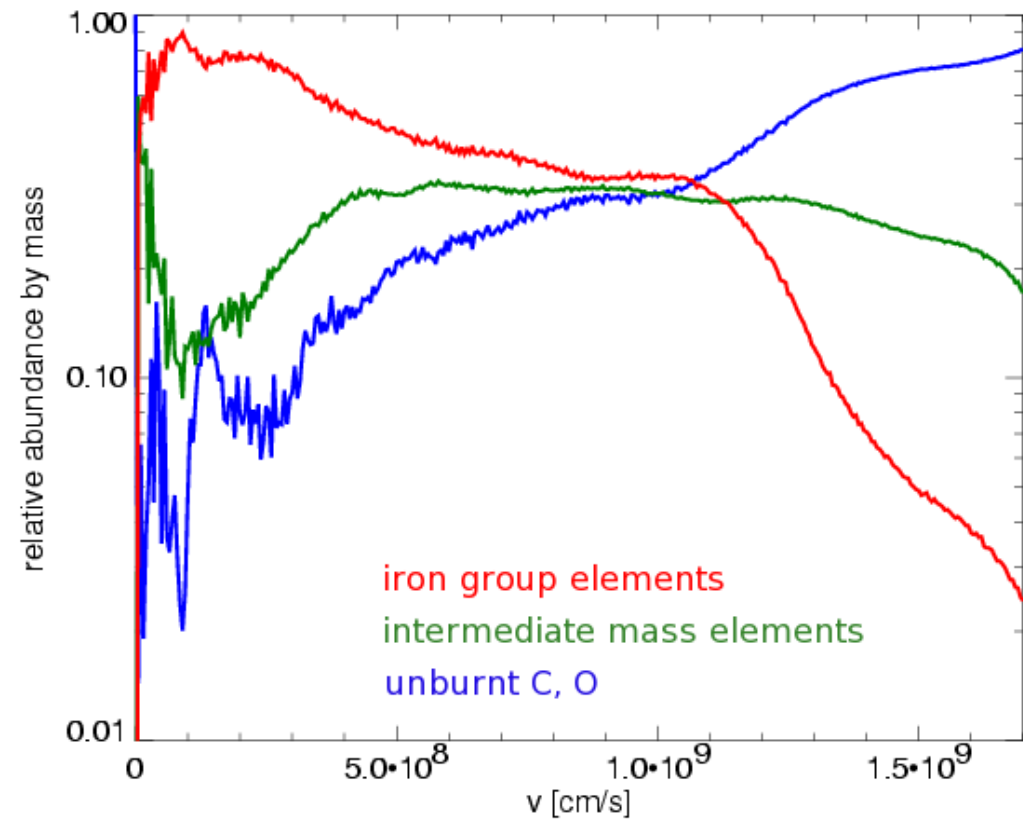
Example: Abundances of the SNOB run...



.... and “abundance tomography”

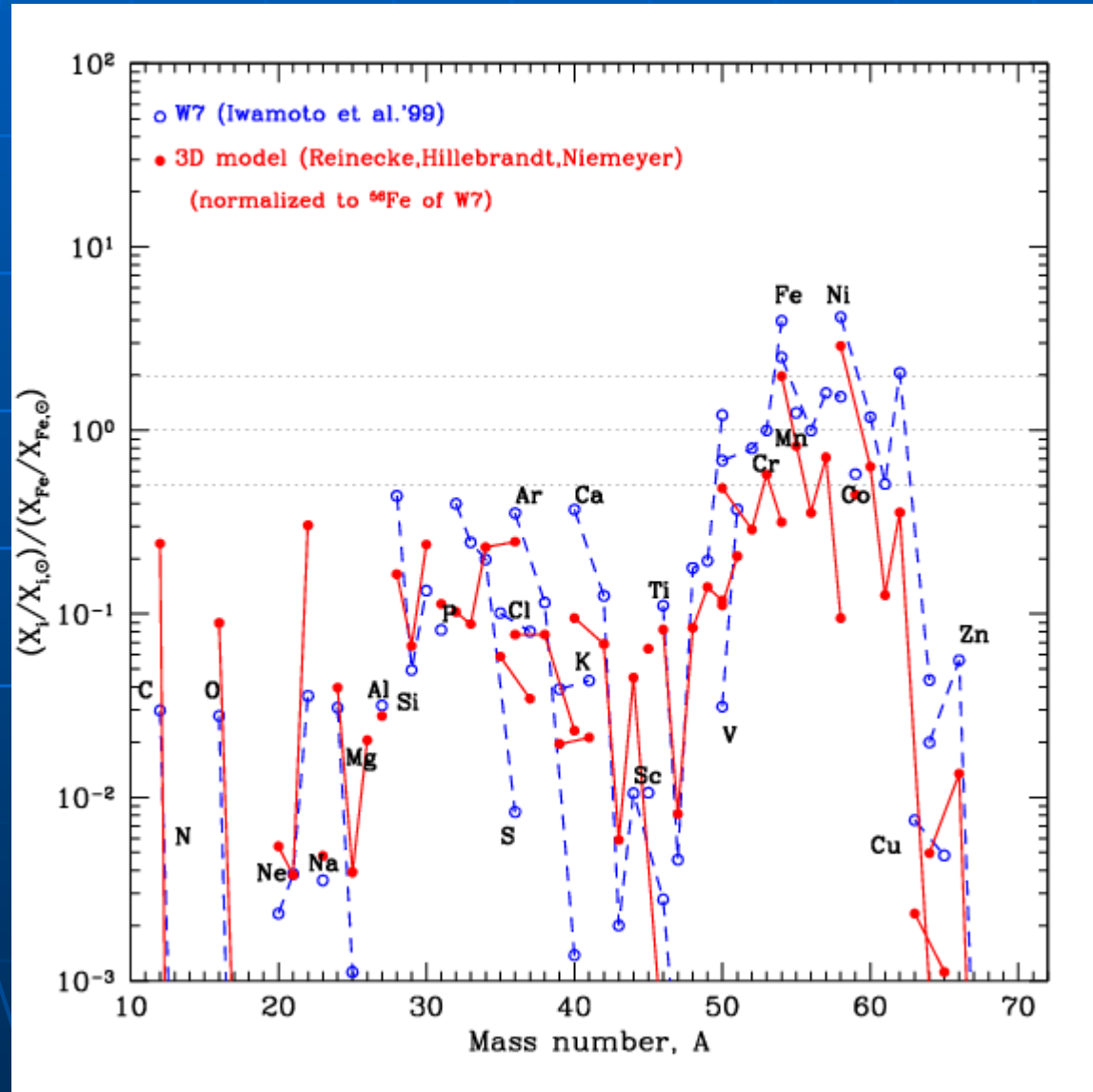


SN 2002bo



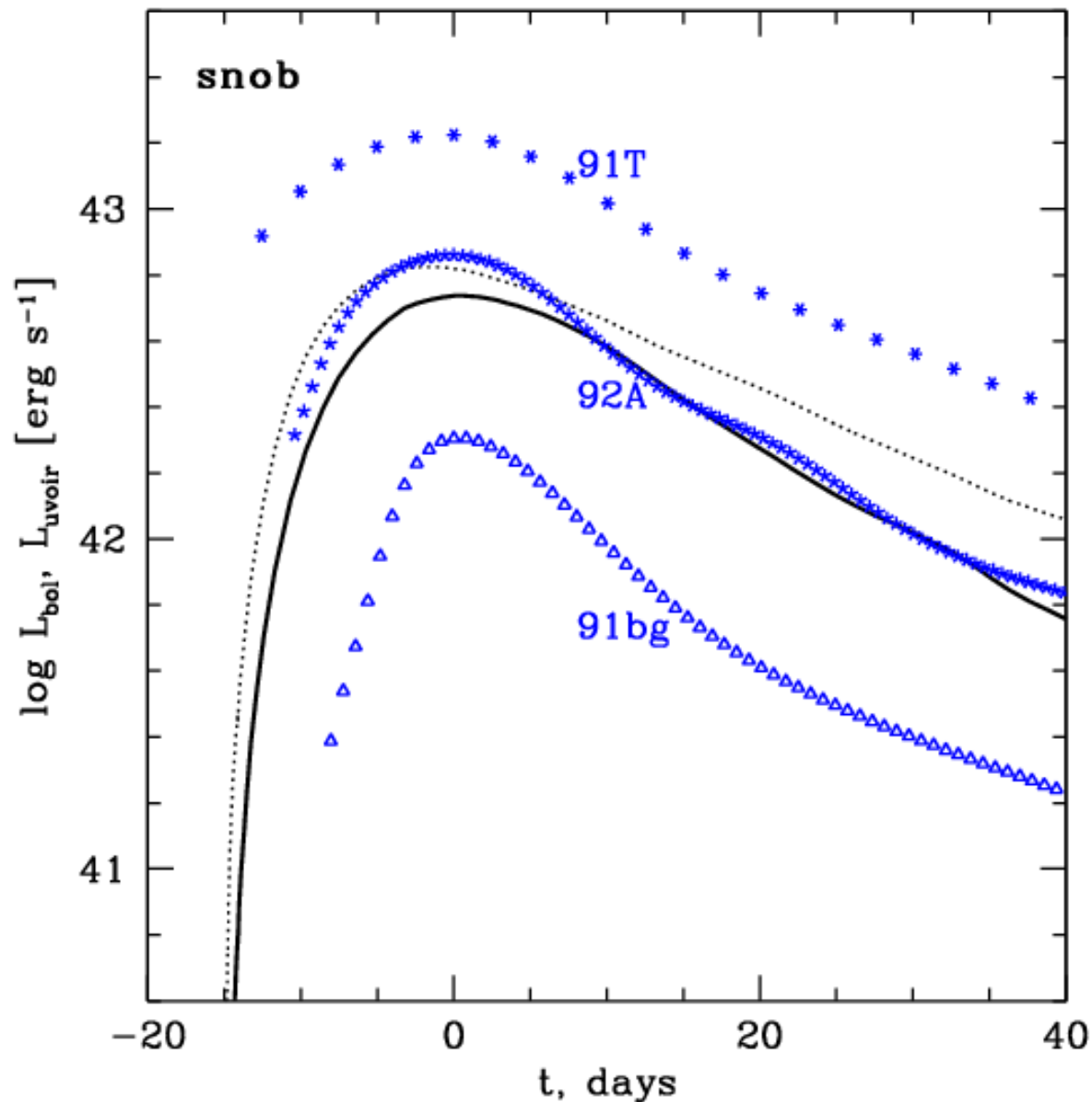
SNOB-run

Chemical composition: What is different from W7?



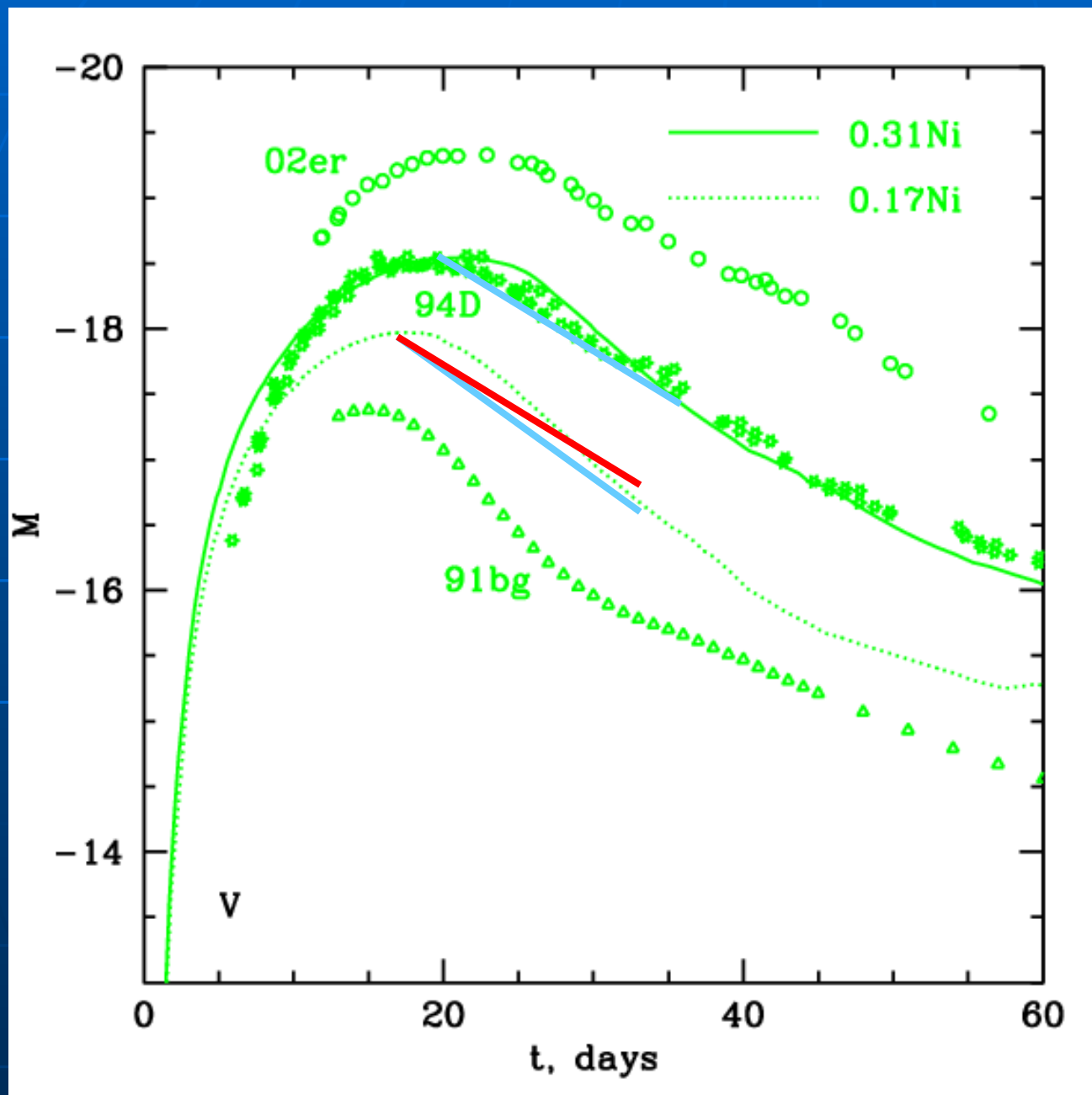
(Off-center ignited 3D model, Travaglio et al., 2003)

Example: Bolometric light curves from SNOB



Note:

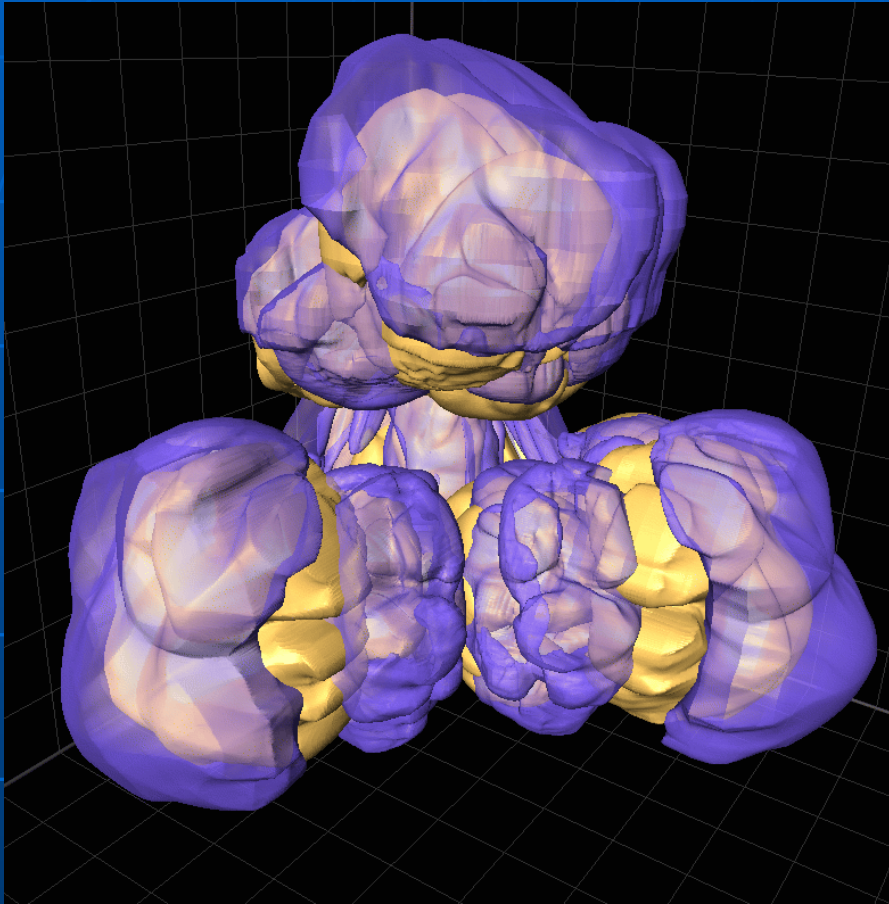
These are predictions, not fits!



Prediction from Theory :

Light-curve shape / luminosity correlation?

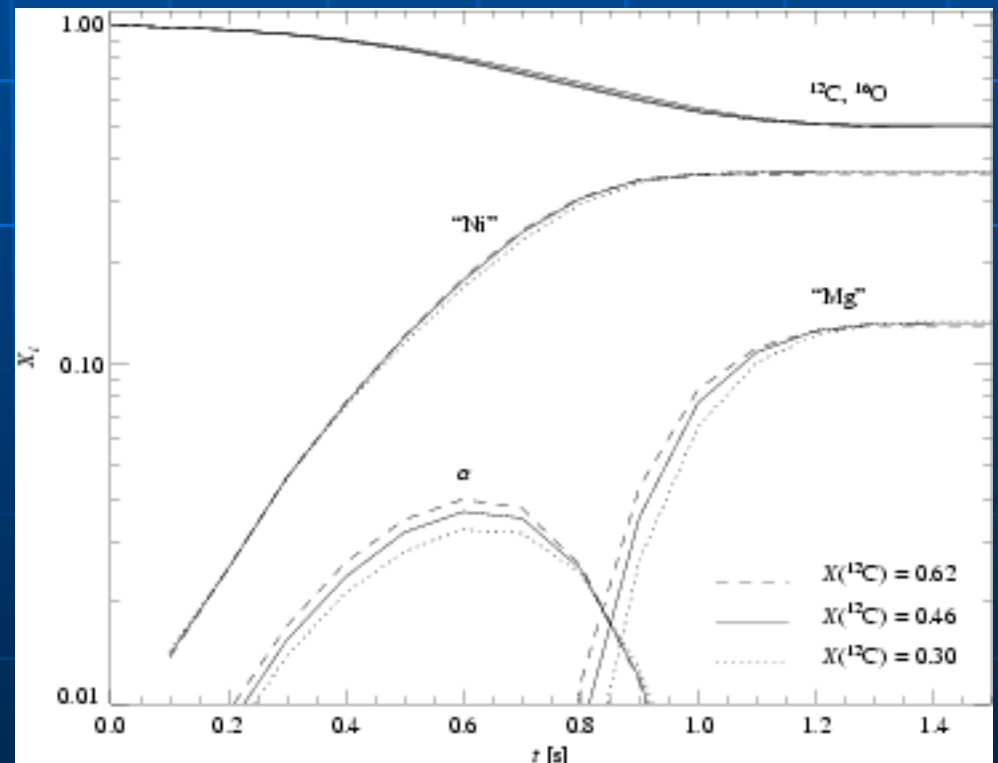
Dependence on the initial C/O ratio?



$X(^{12}\text{C})$	E_{nuc} (10^{50}erg)	$M(\text{Ni}) (M_{\odot})$	M_{α}^{max} (M_{\odot})
0.30	8.85	0.5178	0.0458
0.46	9.46	0.5165	0.0518
0.62	9.97	0.5104	0.0564

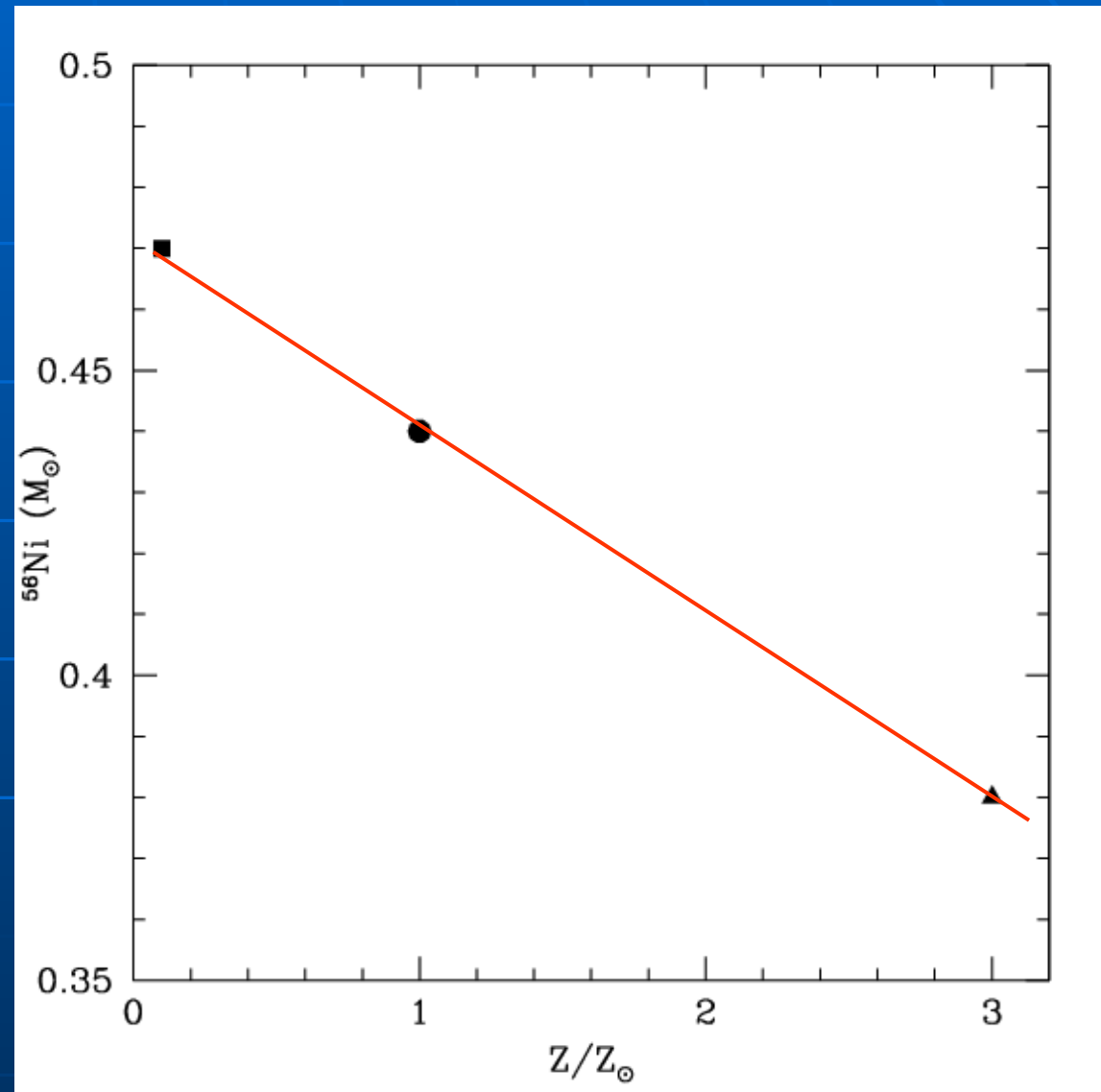
Ni-mass (luminosity)
independent of initial C/O!

(Röpke & Hillebrandt, 2004)



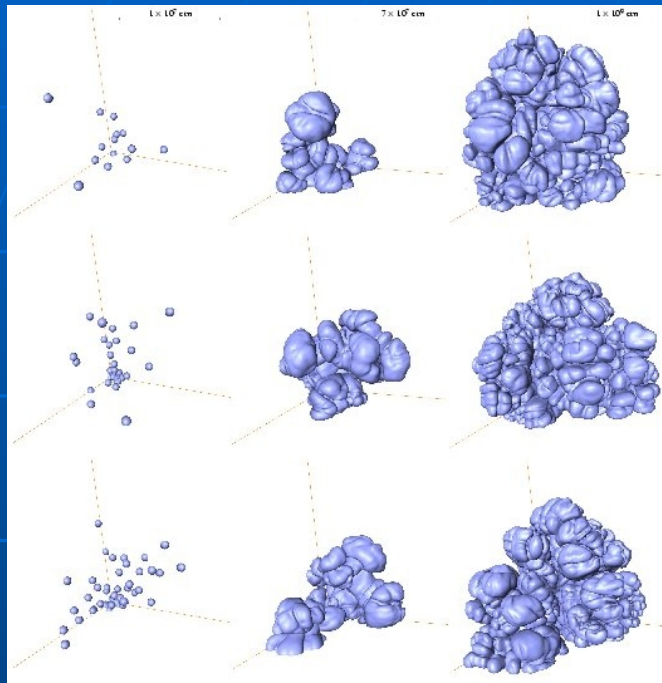
Metallicity dependence (Travaglio et al. 2005)

(Model “b30_3d”)

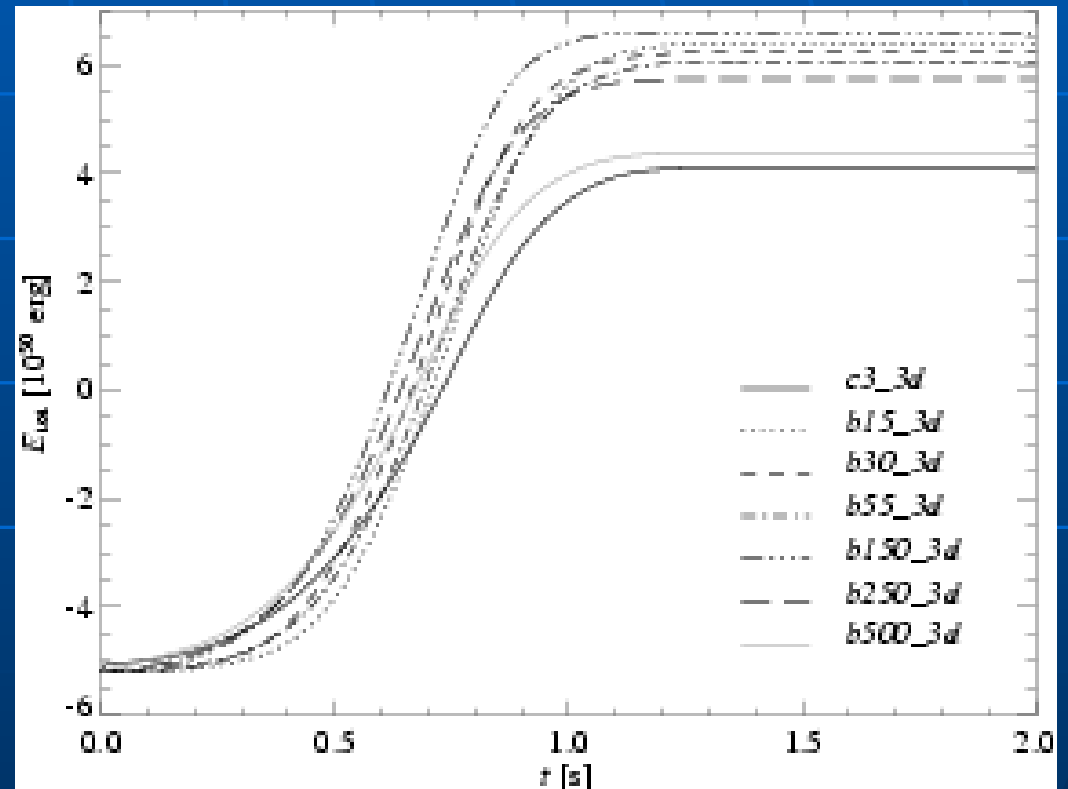
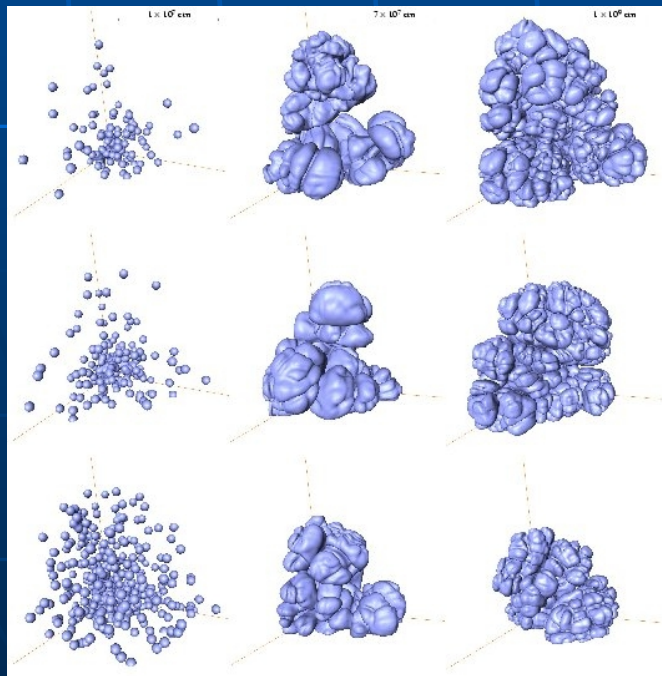


Weak metallicity dependence
(in agreement with Timmes
et al. 2003)

Ignition conditions: Reason for the diversity?

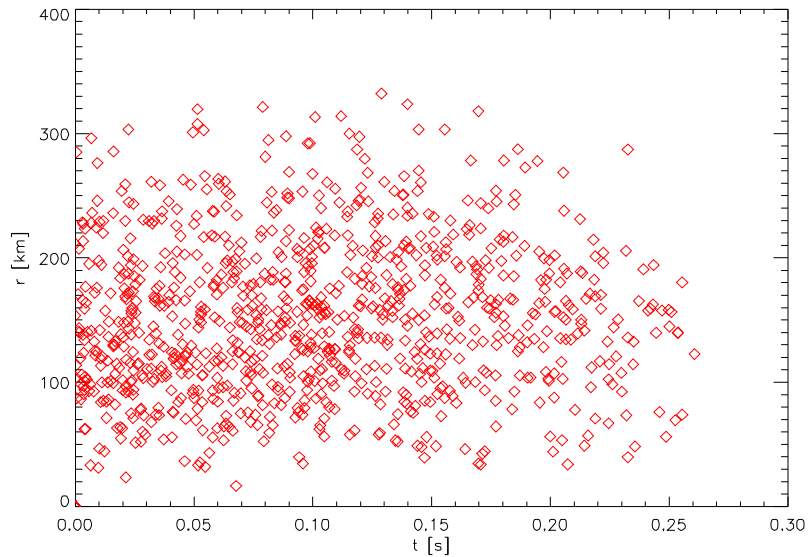


"Multi-spot"

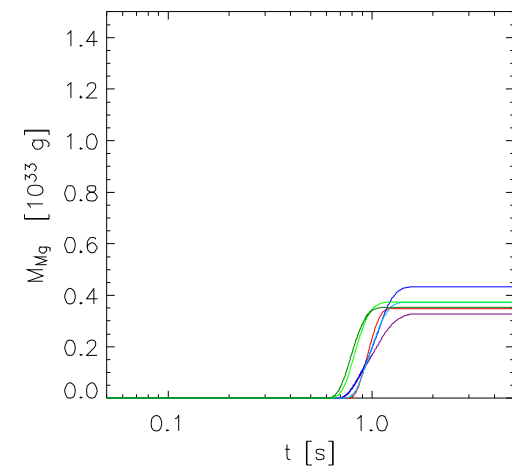
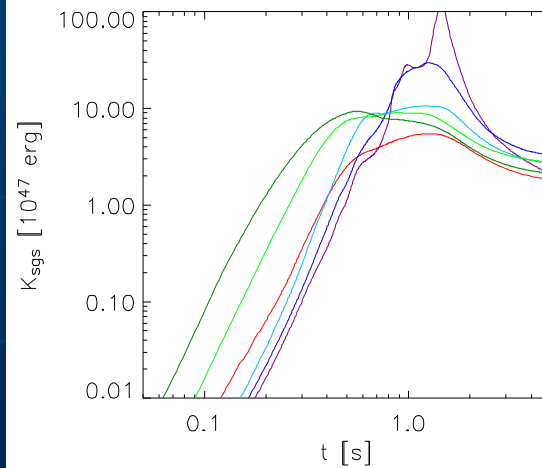
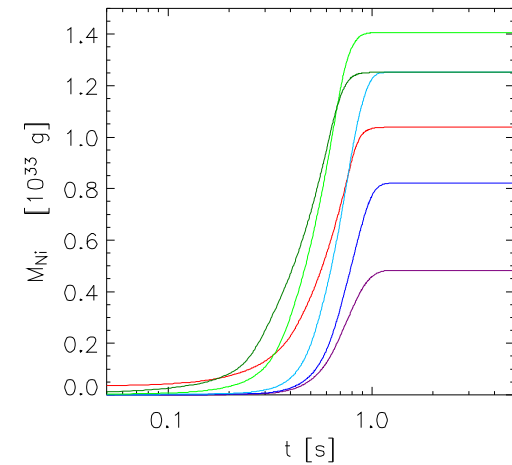
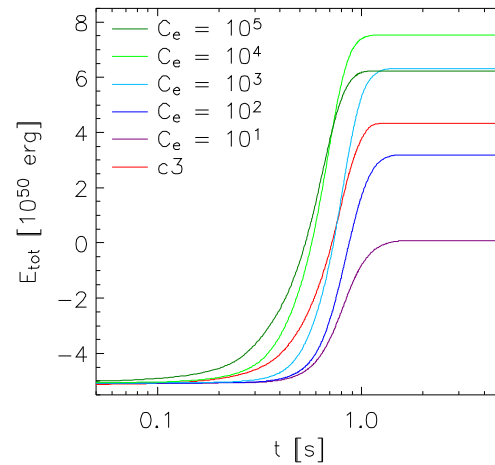


Röpke et al. (2005)

Ignition conditions (cont.):



"Stochastic ignition"



Schmidt & Niemeyer. (2005)

Summary (Part II)

- "Parameter-free" thermonuclear models of type Ia supernovae, based on Chandrasekhar-mass C+O white dwarfs explode with about the right energy.
- They allow to predict light curves and spectra, depending on physical parameters!
- They can explain (most of ?) the observed properties well.
- The diversity may be due to randomness in the ignition conditions, (C/O), and

Questions and challenges

➤ Ignition conditions:

How do WDs reach M_{Ch} ? Center/off-center ignition?
One/multiple “points”?

➤ Combustion modeling:

Interaction of nuclear flames with turbulence;
“distributed burning”; “active turbulent combustion” ?
Deflagration/detonation transition: Does it happen? Is it
“needed”?

➤ New generation of “full-star” models:

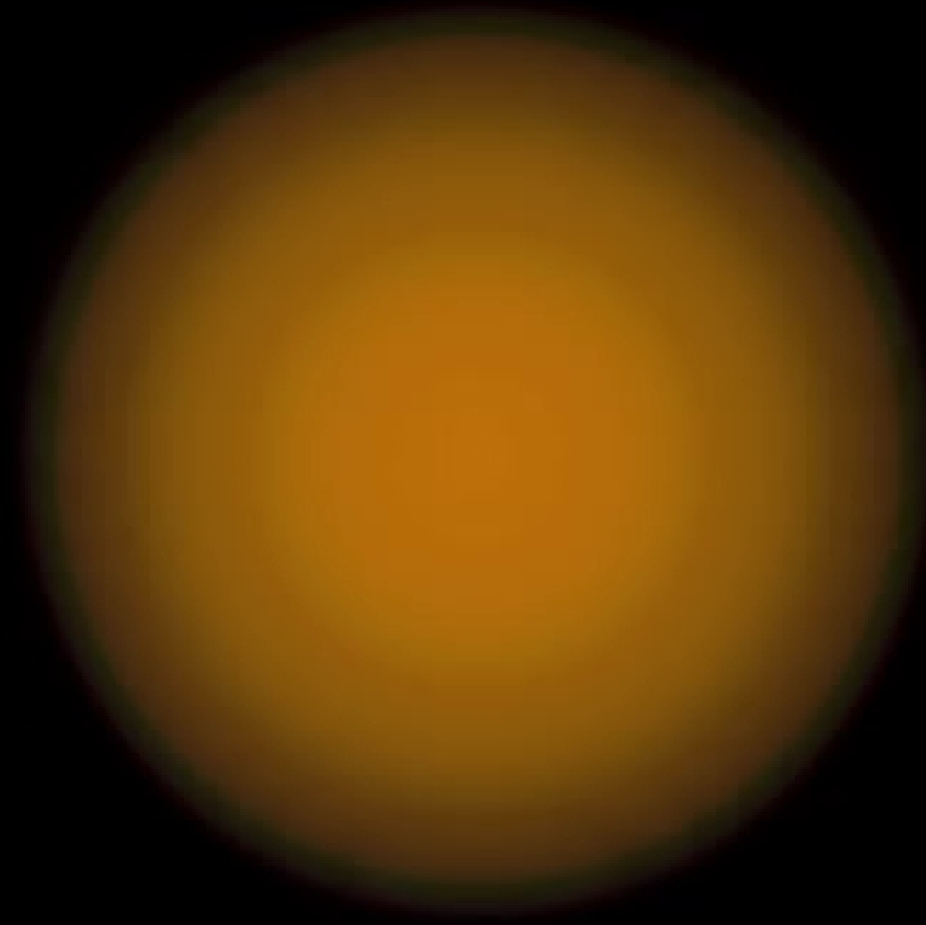
Light curves? Spectra?

➤ Other progenitors:

Mergers? Sub-Chandrasekhars?

Size (km) : 5342.16

Time [s] : 0.0120128

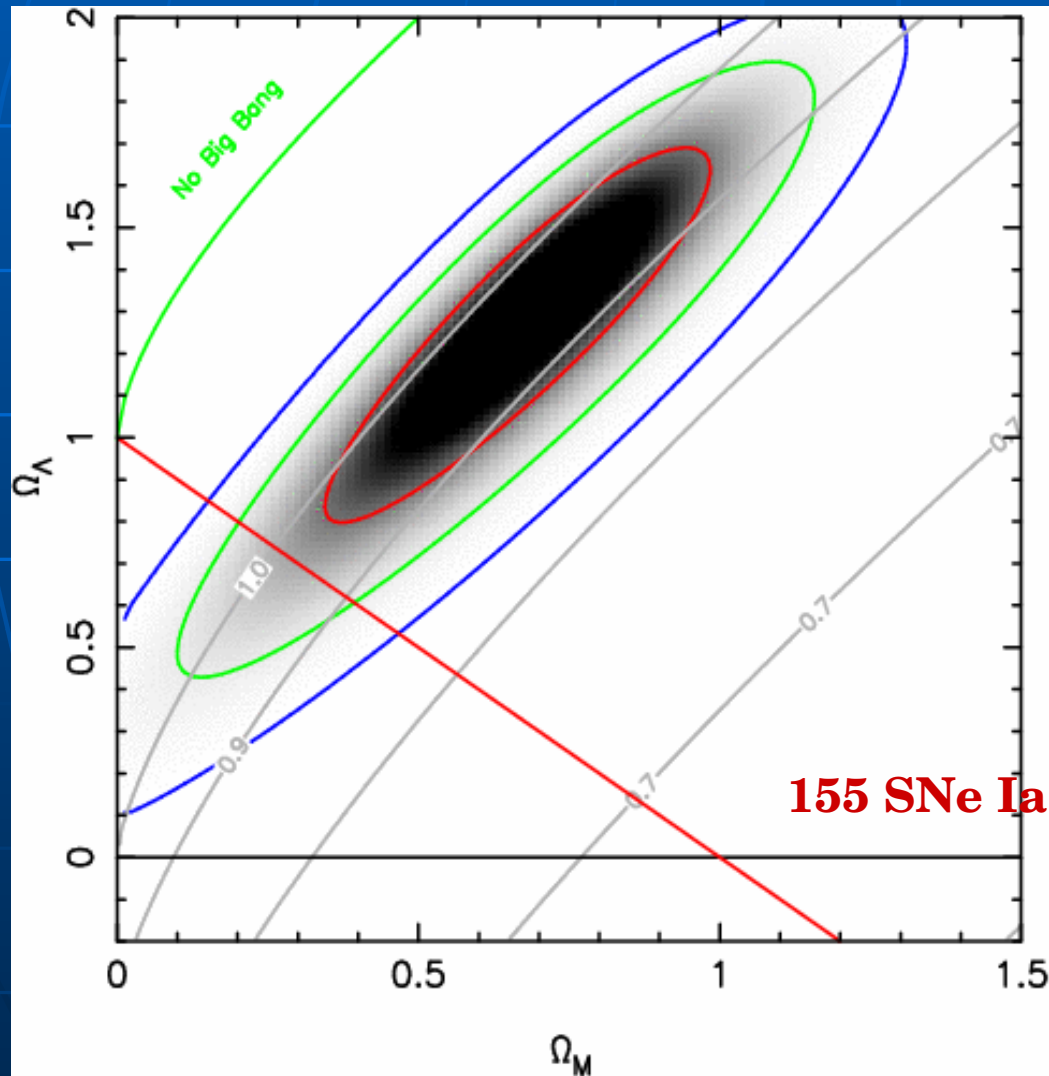


Web address for movies:

<http://www.mpa-garching.mpg.de/~wfh/>

Supernovae and Cosmology: The Quest for Precise Luminosity Distances

(In part “borrowed” from a lecture by Bruno Leibundgut)



(Tonry et al. 2003)

Distances in the local universe

- Assume a linear expansion (“Hubble law”)

$$v = cz = H_0 \cdot D$$

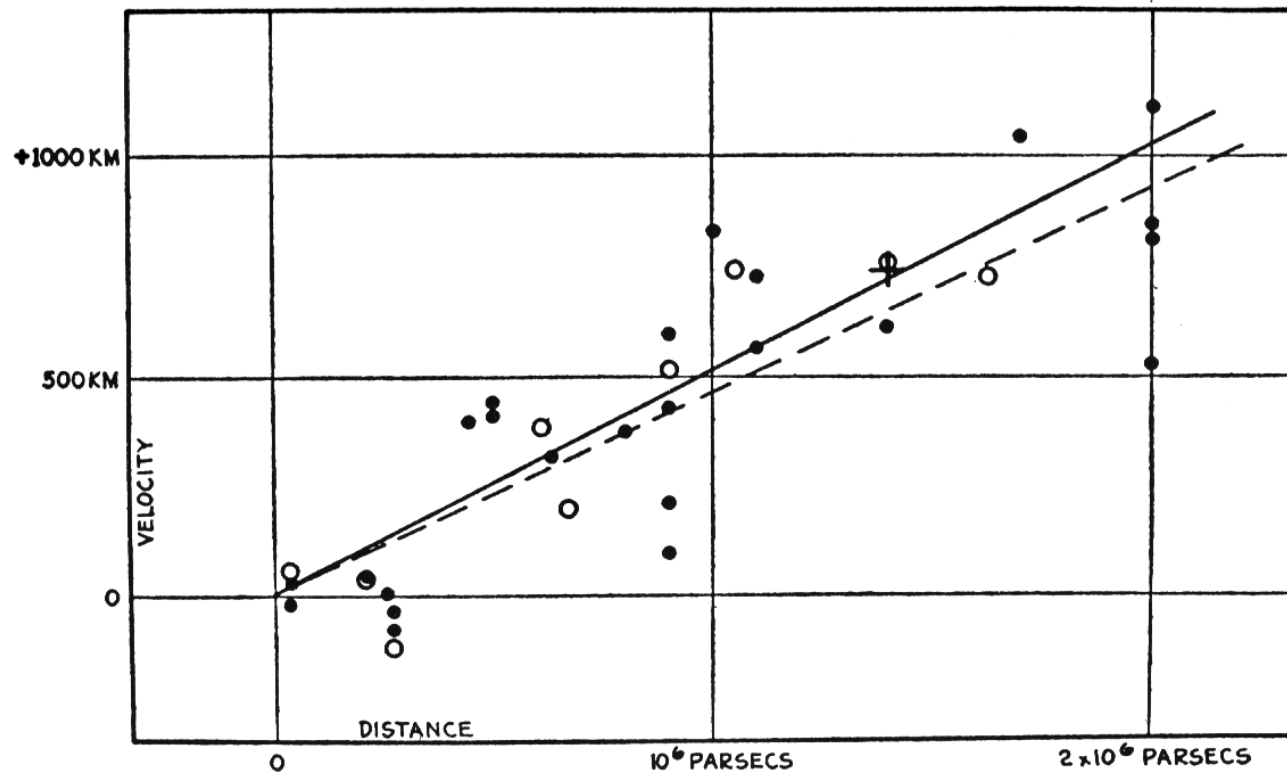
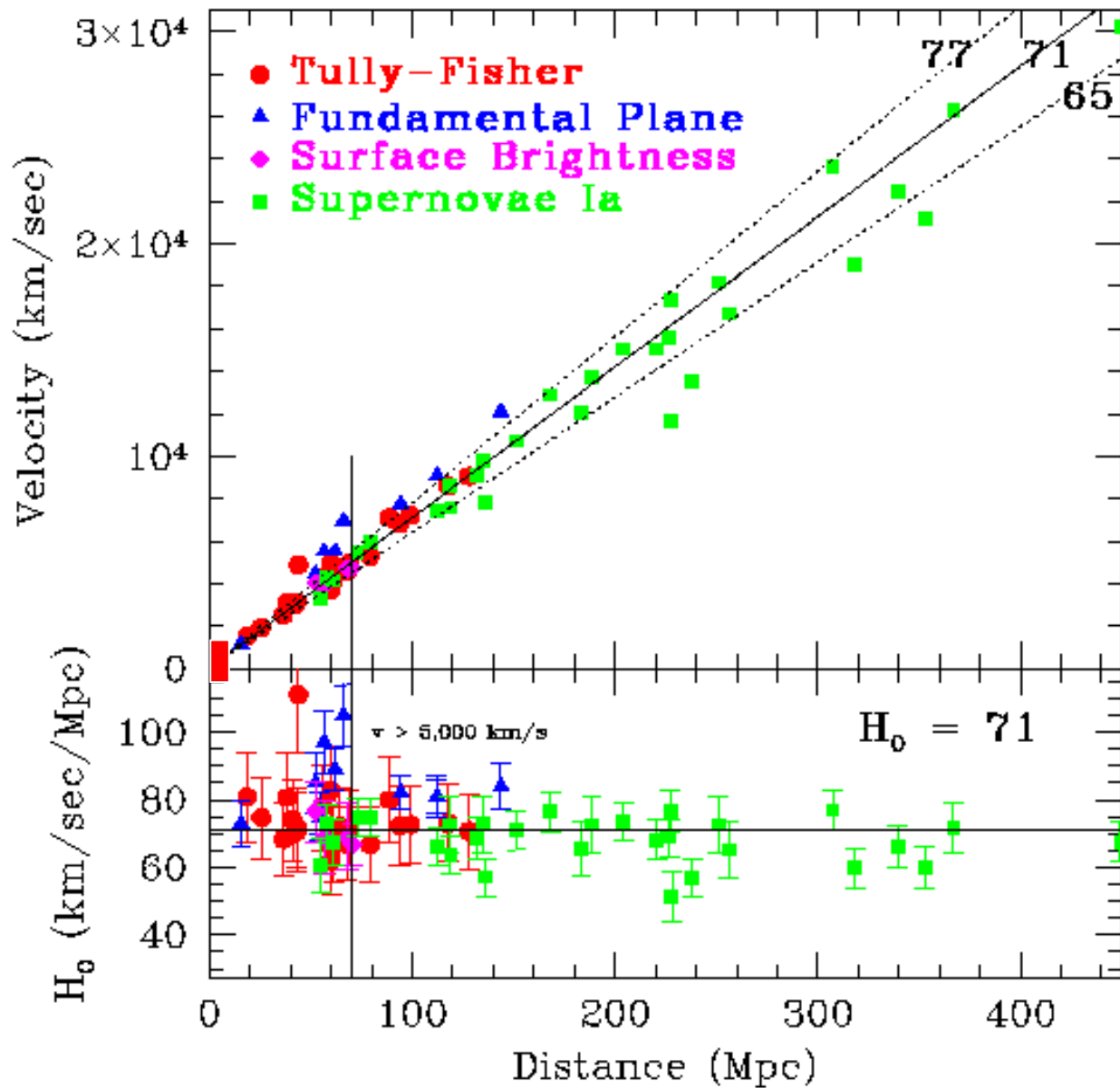
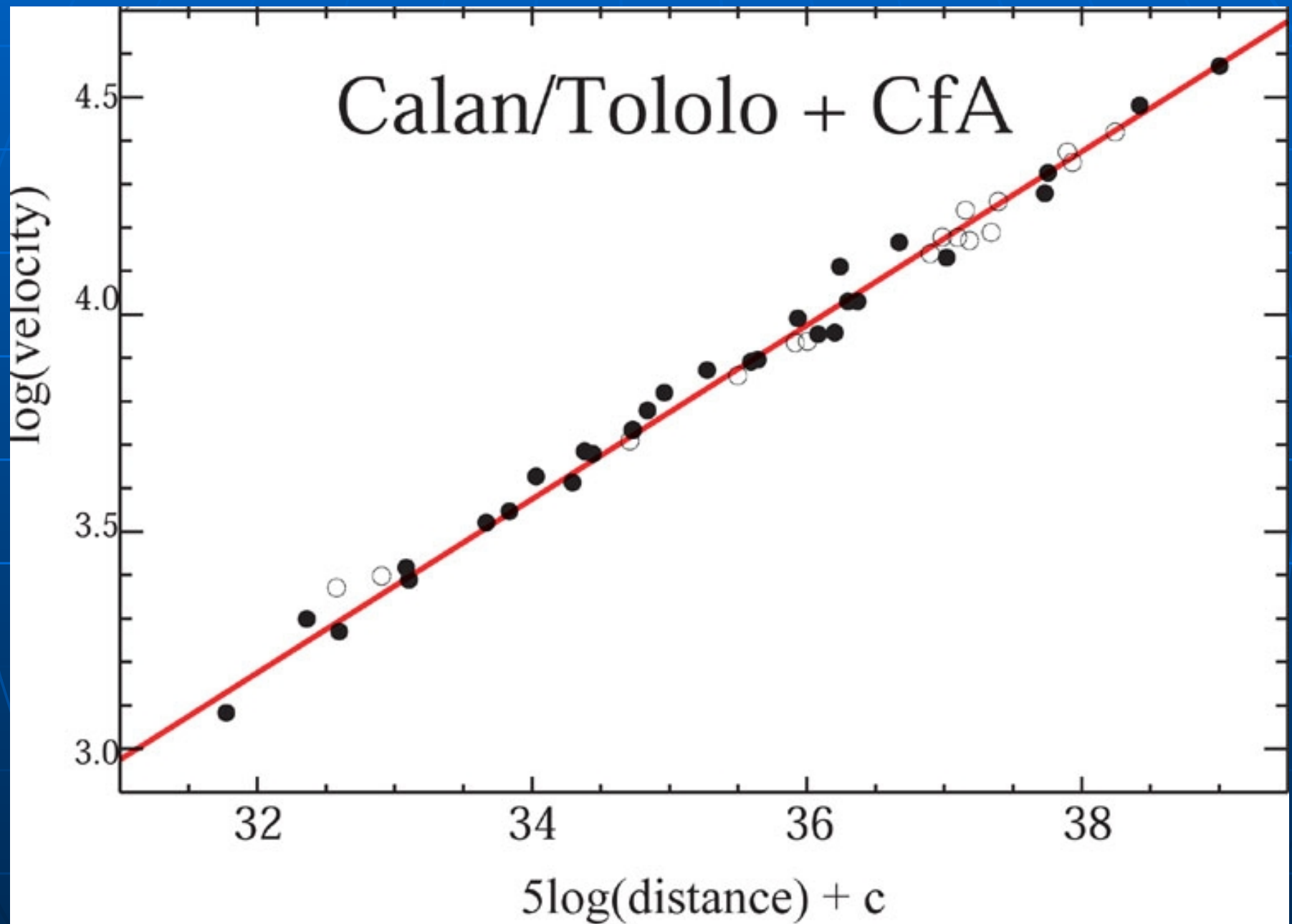


FIG. 9. *The Formulation of the Velocity-Distance Relation.*

A “modern” Hubble diagram



Universal expansion



Distances in the local universe

- Assume a linear expansion (Hubble law):

$$v = cz = H_0 \cdot D$$

- Use the distance modulus

$$m - M = 5 \log(D/10 \text{ pc}) - 5$$

- Distances of a 'standard candle' ($M = \text{const.}$)

$$m = 5 \log(z) + b$$

$$b = M + 25 + 5 \log(c) - 5 \log(H_0)$$

The Hubble constant

- Sets the absolute scale of cosmology
 - (“replaces these annoying h ’s in all the theorists talks” ; B. Leibundgut)
- Measure redshifts and distances in the nearby universe
 - Supernovae can do this in two ways:
 - Expanding photosphere method of core-collapse SNe
 - accurate (relative) distances from SN Ia

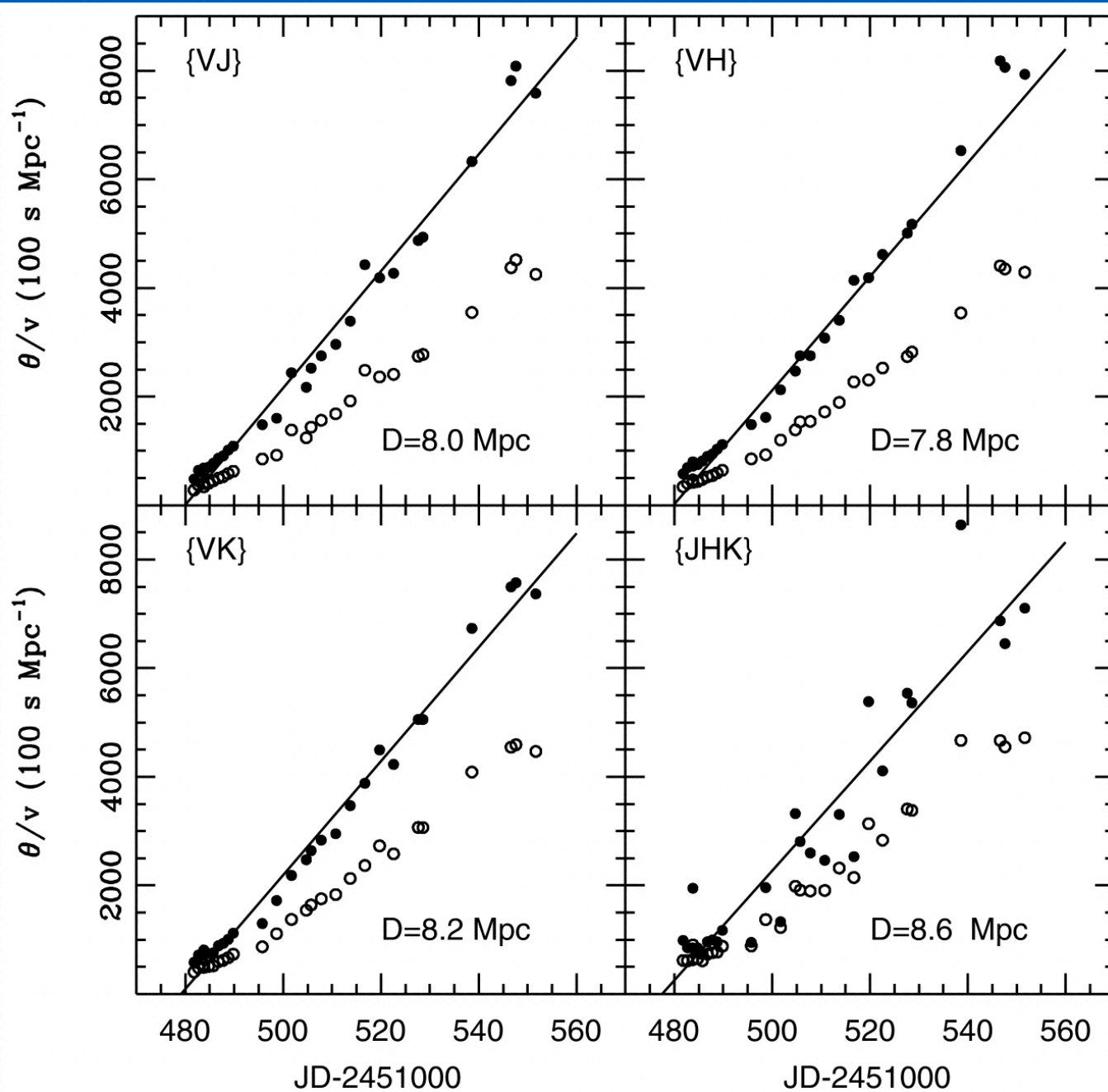
Expanding Photosphere Method

- Baade (1926), Schmidt et al. (1993), Eastman et al. (1996), Hamuy et al. (2001)
- Assume homologous expansion: $R(t) = R_0 + v(t - t_0)$
- Photometric angular diameter

$$\Theta = \frac{R}{D} = \sqrt{\frac{f_{\lambda}}{\zeta_{\lambda}^2 \pi B_{\lambda}(T) 10^{-0.4A(\lambda)}}$$

Distances from EPM

(SN 1999em, Hamuy et al. 2001)



$$\frac{\Theta_i}{v_i} \approx \frac{t_i - t_0}{D}$$

Slope gives the distance

Intercept the size of the progenitor and/or time of explosion

Distances from EPM

- Note that this distance measurement is completely **independent** of any other astronomical object!
 - no distance ladder
- Assumption:
 - massive envelope that creates a photosphere
 - spherical symmetry
 - not true for many core collapse supernovae
 - correction factors for deviation from black body spectrum
 - model dependent

EPM so far

■ Limitations

- needs large and extensive data sets
- difficulties to get into the Hubble flow
- distances only to galaxies with supernovae
 - difficult to build large sample

■ Promise

- completely independent distance measurements
 - checks on the Cepheid distance scale

Distances with Type Ia Supernovae

- Use the Hubble diagram ($m-M$ vs. $\log z$)
 - $m-M=5\log(z)+25+5\log(c)-5\log(H_0)$
- Note that the slope is given here.
- Hubble constant can be derived when the absolute luminosity M is known
 - $\log H_0 = \log(z) + 5 + \log(c) - 0.2(m-M)$

Hubble constant from SNe Ia

- Calibrate the absolute luminosity
 - through Cepheids
 - ‘classical distance ladder’
 - depends on the accuracy of the previous rungs on the ladder
 - LMC distance, P-L(-C) relation, metallicities
 - HST program (Sandage, Tammann)
 - HST Key Programme (Freedman, Kennicutt, Mould, Madore)
 - through models
 - extremely difficult (but possible!)

Absolute Magnitudes of SNe Ia

SN	Galaxy	m-M	M_B	M_V	M_I	Δm_{15}
1937C	IC 4182	28.36 (12)	-19.56 (15)	-19.54 (17)	-	0.87 (10)
1960F	NGC 4496	31.03 (10)	-19.56 (18)	-19.62 (22)	-	1.06 (12)
1972E	NGC 5253	28.00 (07)	-19.64 (16)	-19.61 (17)	-19.27 (20)	0.87 (10)
1974G	NGC 4414	31.46 (17)	-19.67 (34)	-19.69 (27)	-	1.11 (06)
1981B	NGC 4536	31.10 (12)	-19.50 (18)	-19.50 (16)	-	1.10 (07)
1989B	NGC 3627	30.22 (12)	-19.47 (18)	-19.42 (16)	-19.21 (14)	1.31 (07)
1990N	NGC 4639	32.03 (22)	-19.39 (26)	-19.41 (24)	-19.14 (23)	1.05 (05)
1998bu	NGC 3368	30.37 (16)	-19.76 (31)	-19.69 (26)	-19.43 (21)	1.08 (05)
1998aq	NGC 3982	31.72 (14)	-19.56 (21)	-19.48 (20)	-	1.12 (03)
Straight mean			-19.57 (04)	-19.55 (04)	-19.26 (0 6)	
Weighted mean			-19.56 (07)	-19.53 (06)	-19.25 (0 9)	

(Saha et al. 1999)

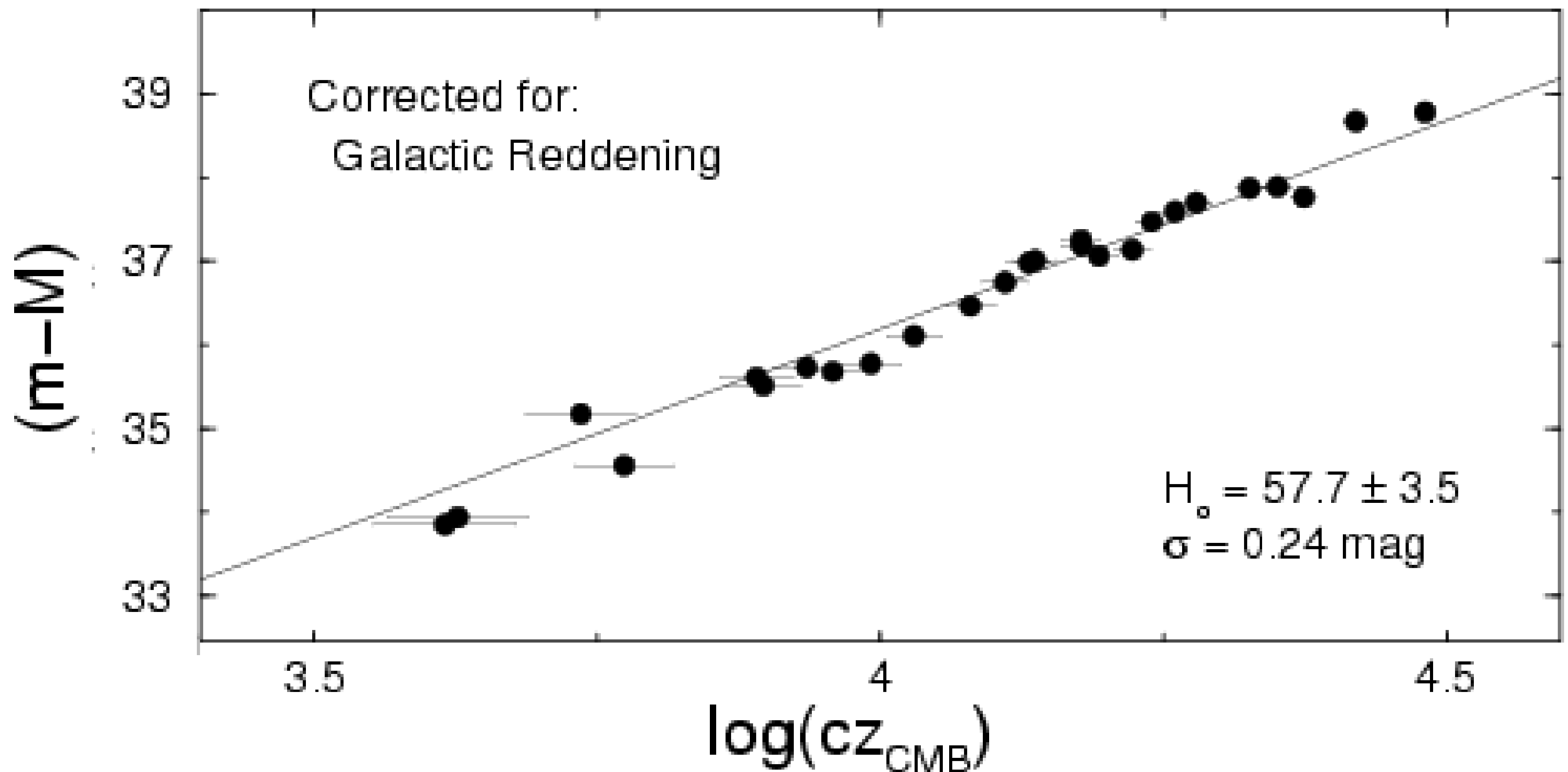
Testing the SNe Ia as distance indicators

- Hubble diagram of SNe Ia in the local, linear expansion, Hubble flow
- Calibration through “primary” distance indicators
- Theoretical models

Nearby SNe Ia

Phillips et al. (1999)

Calan/Tololo "Low Extinction" Sample



Light curve shape – luminosity

- Δm_{15} relation

Phillips (1993), Hamuy et al. (1996), Phillips et al. (1999)

- MLCS

Riess et al. (1996, 1998), Jha et al. (2003)

- stretch

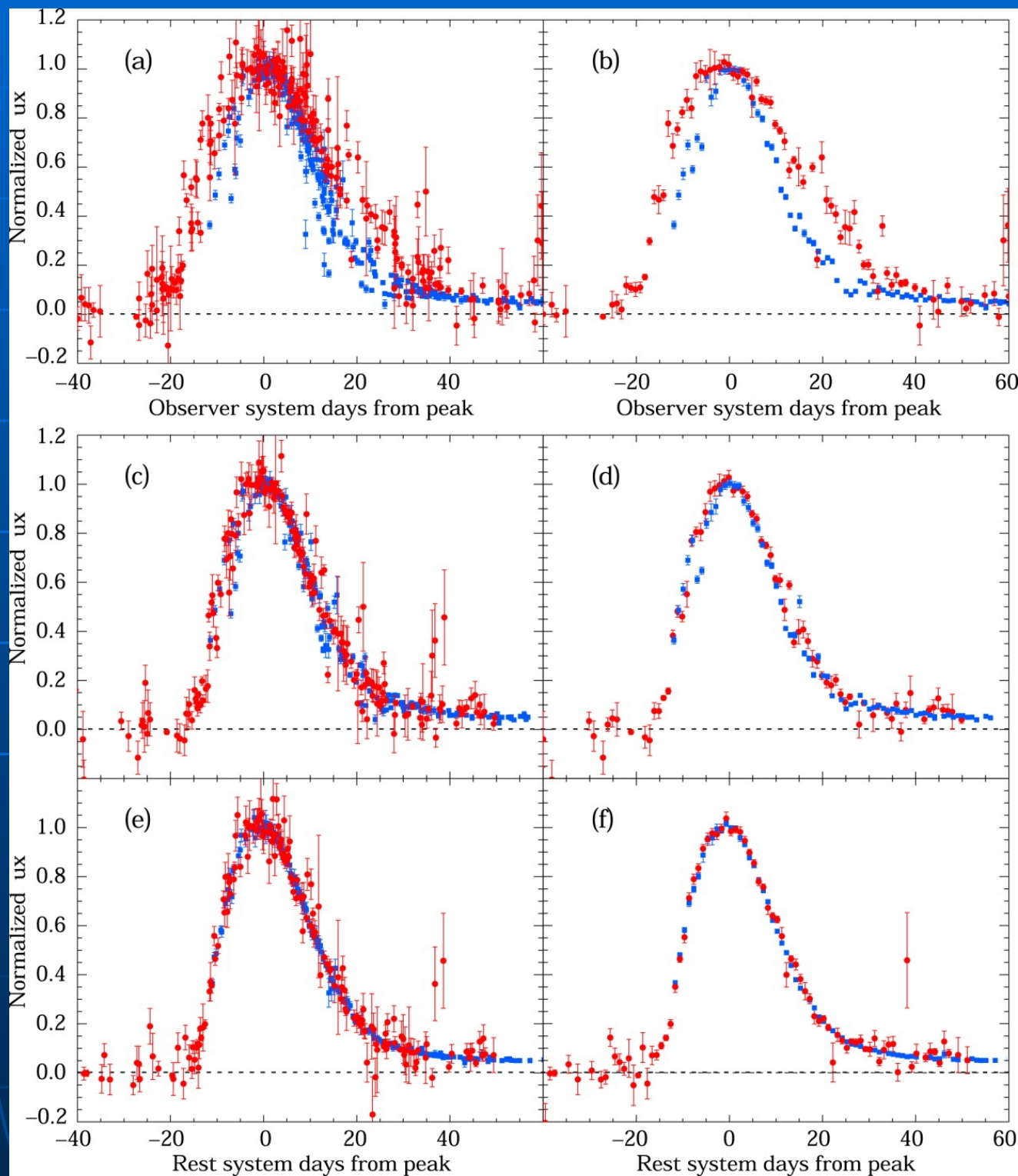
Perlmutter et al. (1997, 1999), Goldhaber et al. (2001)

- MAGIC

Wang et al. (2003)

The principles of light-curve calibrations

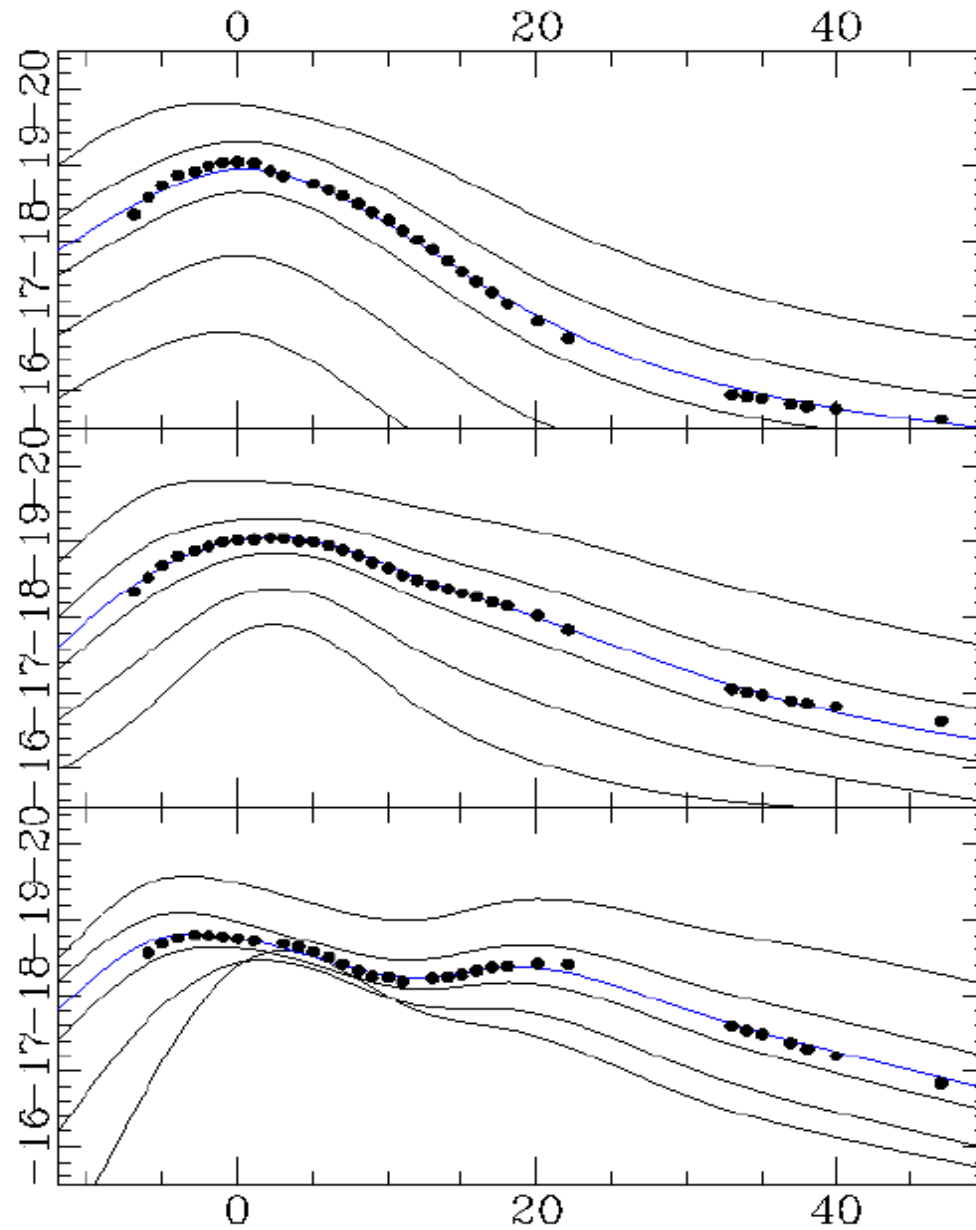
(Goldhaber et al. 2001)



B

V

I



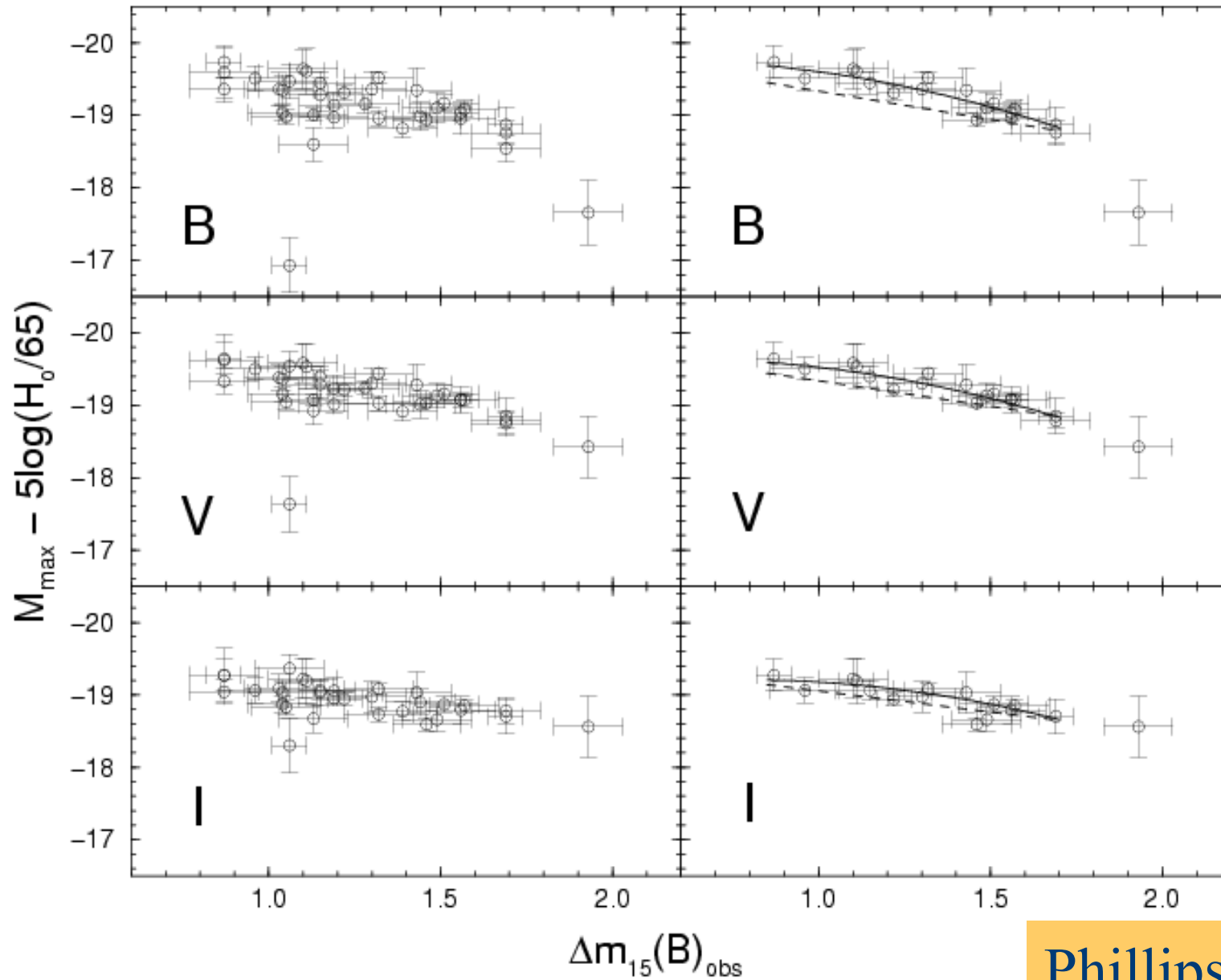
The SN Ia luminosity
can be normalised:

Bright = slow

Dim = fast

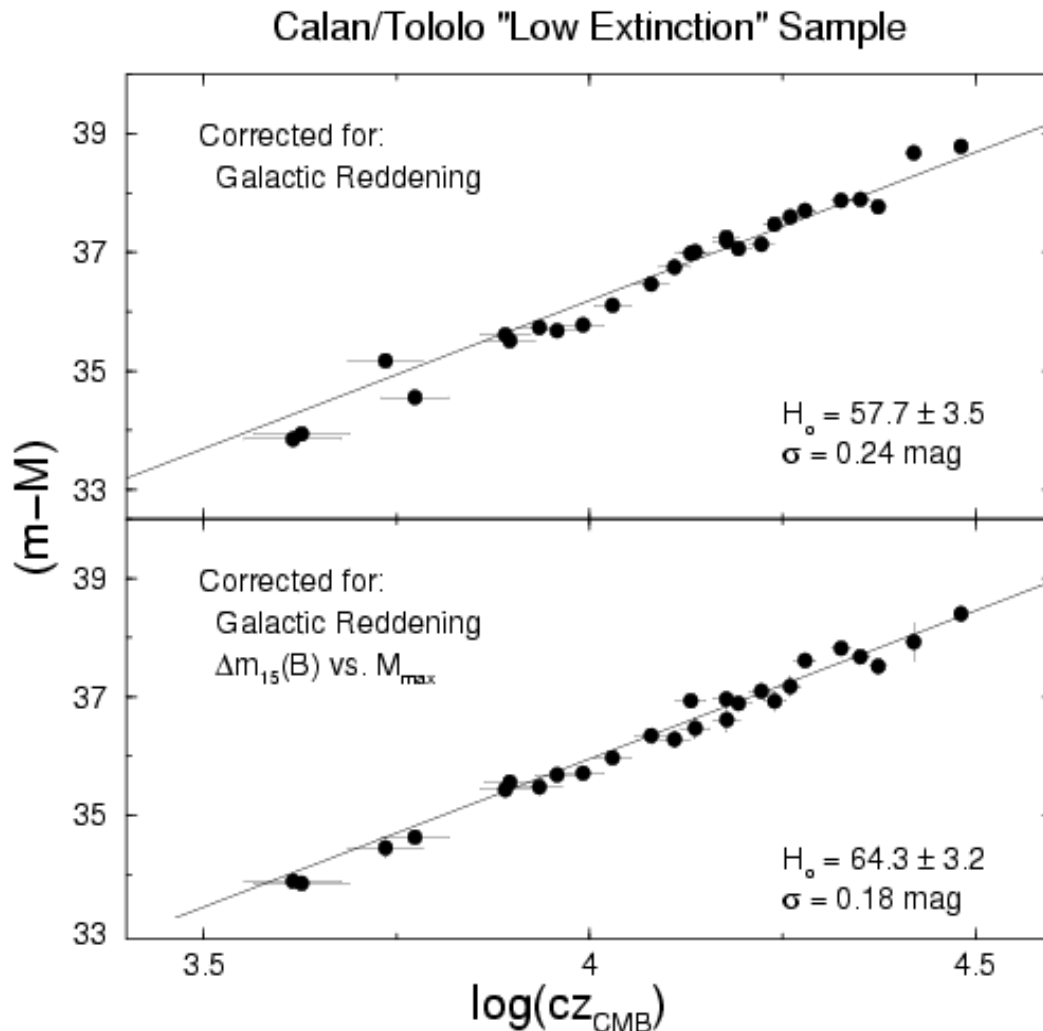
(Riess et al. 1996)

Correlations



Normalisation of the peak luminosity

Phillips et al. 1999



■ Using the luminosity-decline rate relation one can normalise the peak luminosity of SNe Ia

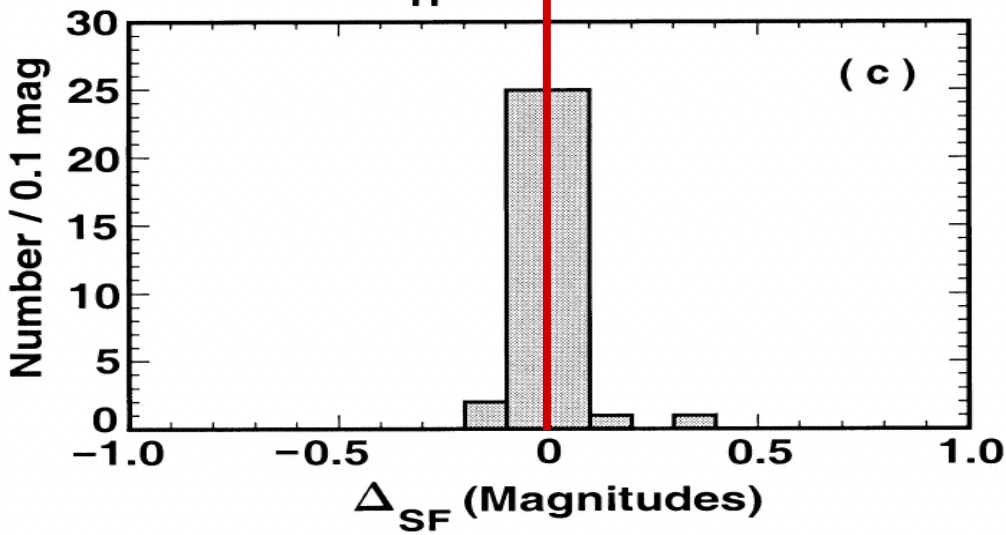
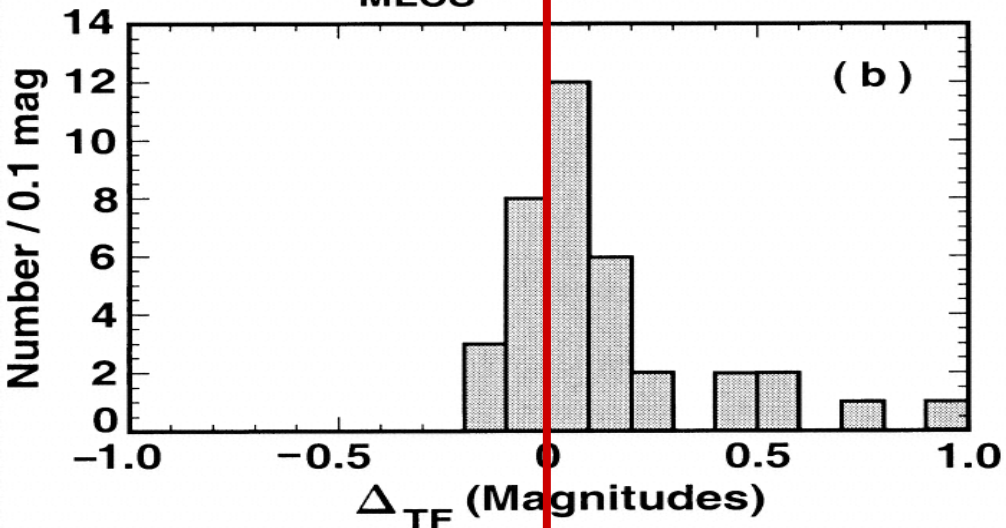
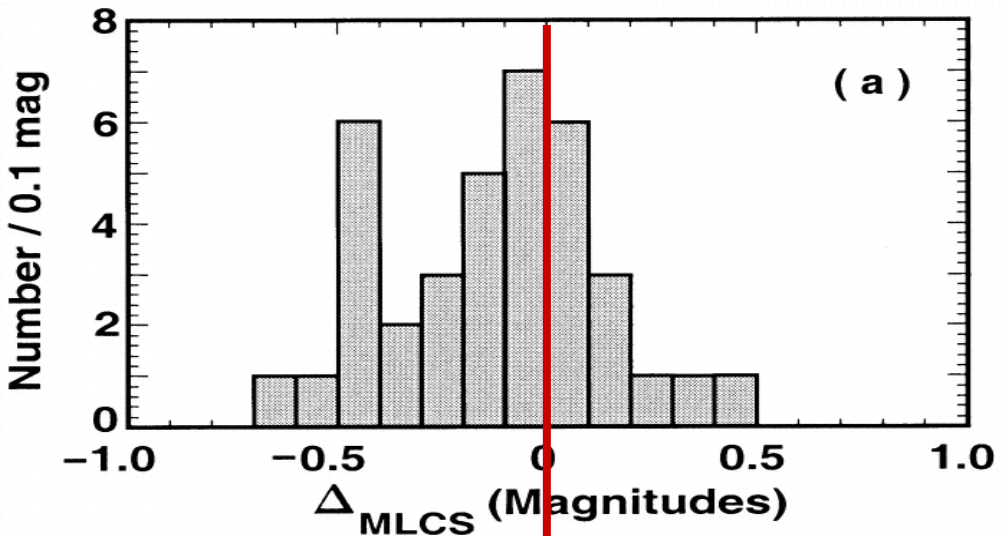


Reduces the scatter!

SN Ia Correlations

- **Luminosity vs. decline rate**
 - Phillips 1993, Hamuy et al. 1996, Riess et al. 1996, 1998, Perlmutter et al. 1997, Goldhaber et al. 2001
- **Luminosity vs. rise time**
 - Riess et al. 1999
- **Luminosity vs. color at maximum**
 - Riess et al. 1996, Tripp 1998, Phillips et al. 1999
- **Luminosity vs. line strengths and line widths**
 - Nugent et al. 1995, Riess et al. 1998, Mazzali et al. 1998
- **Luminosity vs. host galaxy morphology**
 - Filippenko 1989, Hamuy et al. 1995, 1996, Schmidt et al. 1998, Branch et al. 1996

SN Ia Correlations



Riess et al. 1998



Phillips et al. 1999

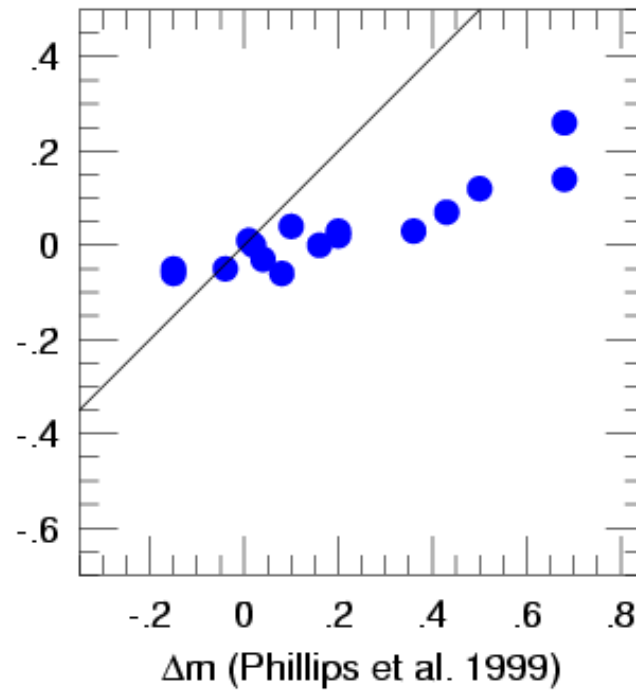
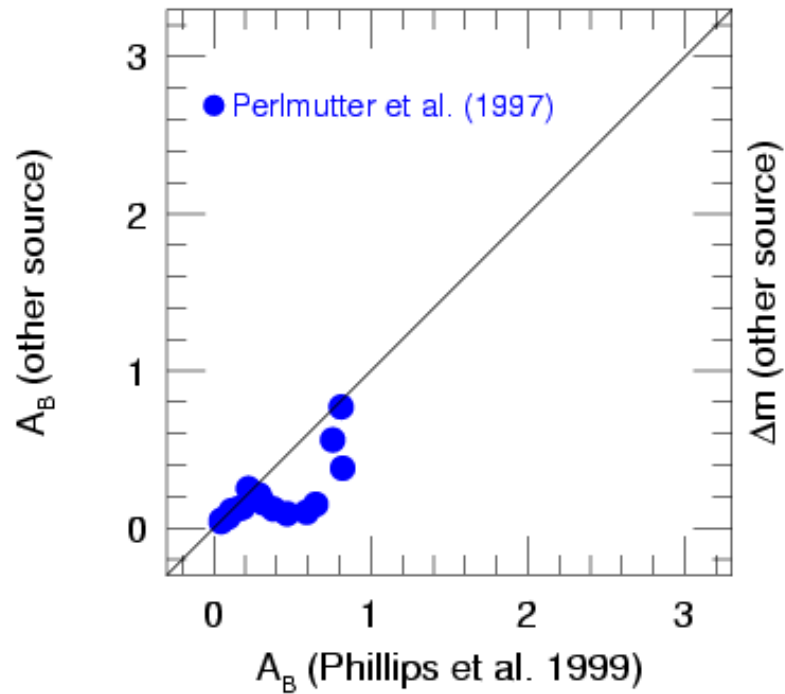
Perlmutter et al. 1997



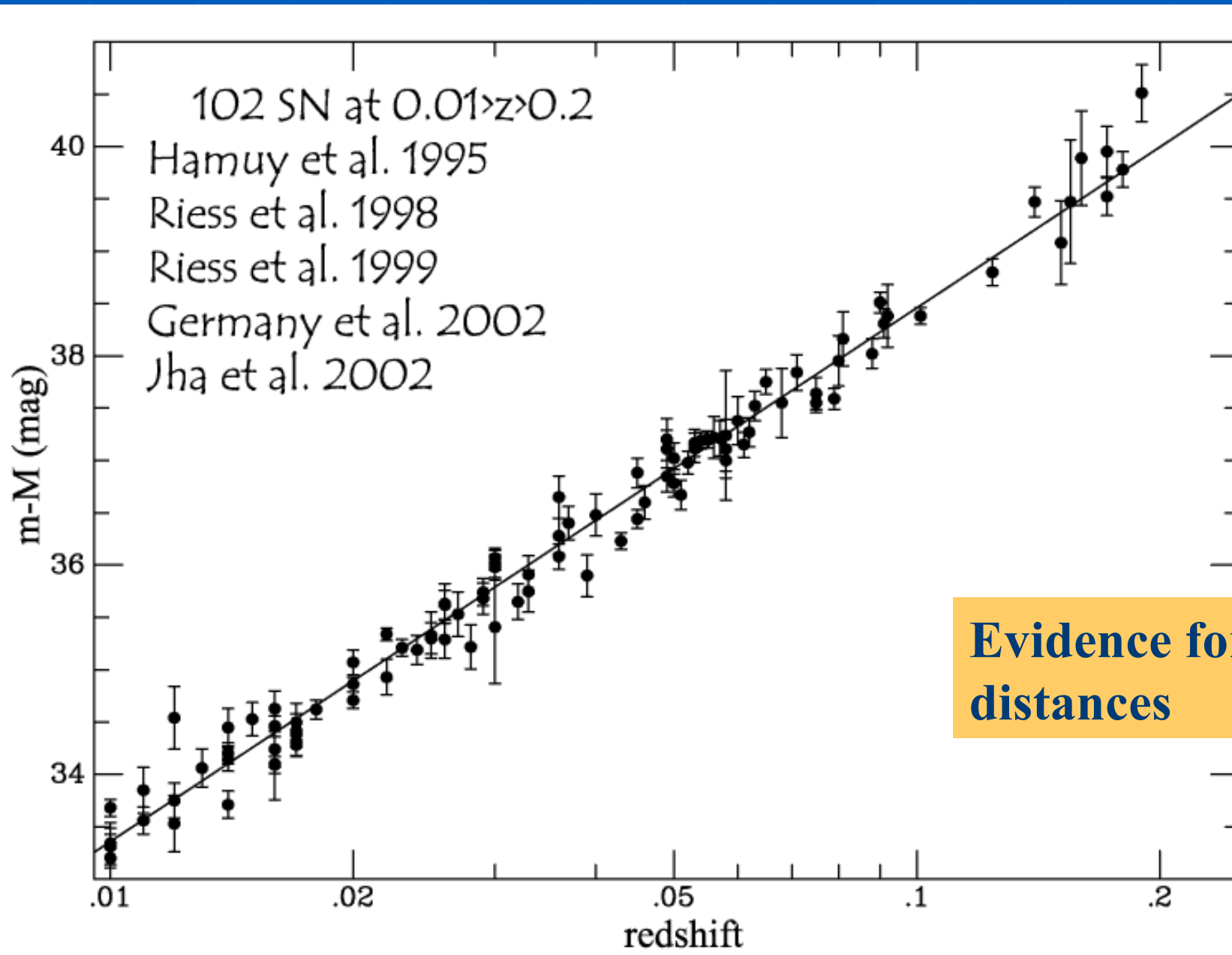
(Drell et al. 2000)

SN Ia Correlations

Leibundgut 2000



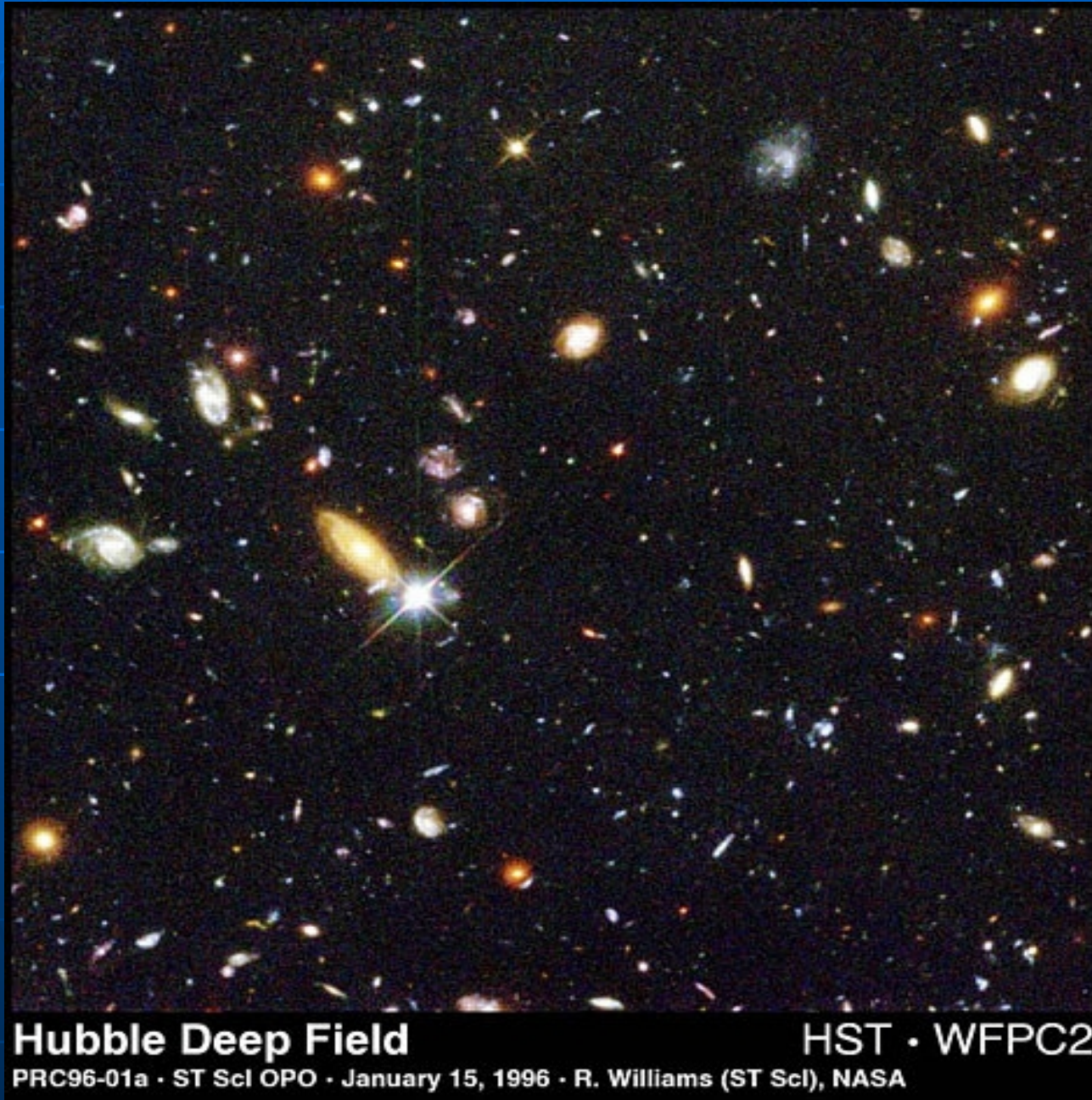
The nearby SN Ia sample



Hubble constant from SNe Ia

- Extremely good (relative) distance indicators
 - distance accuracy around 10%
- Uncertainty in H_0 mostly from the LMC and the Cepheid P-L relation

Very distant supernovae

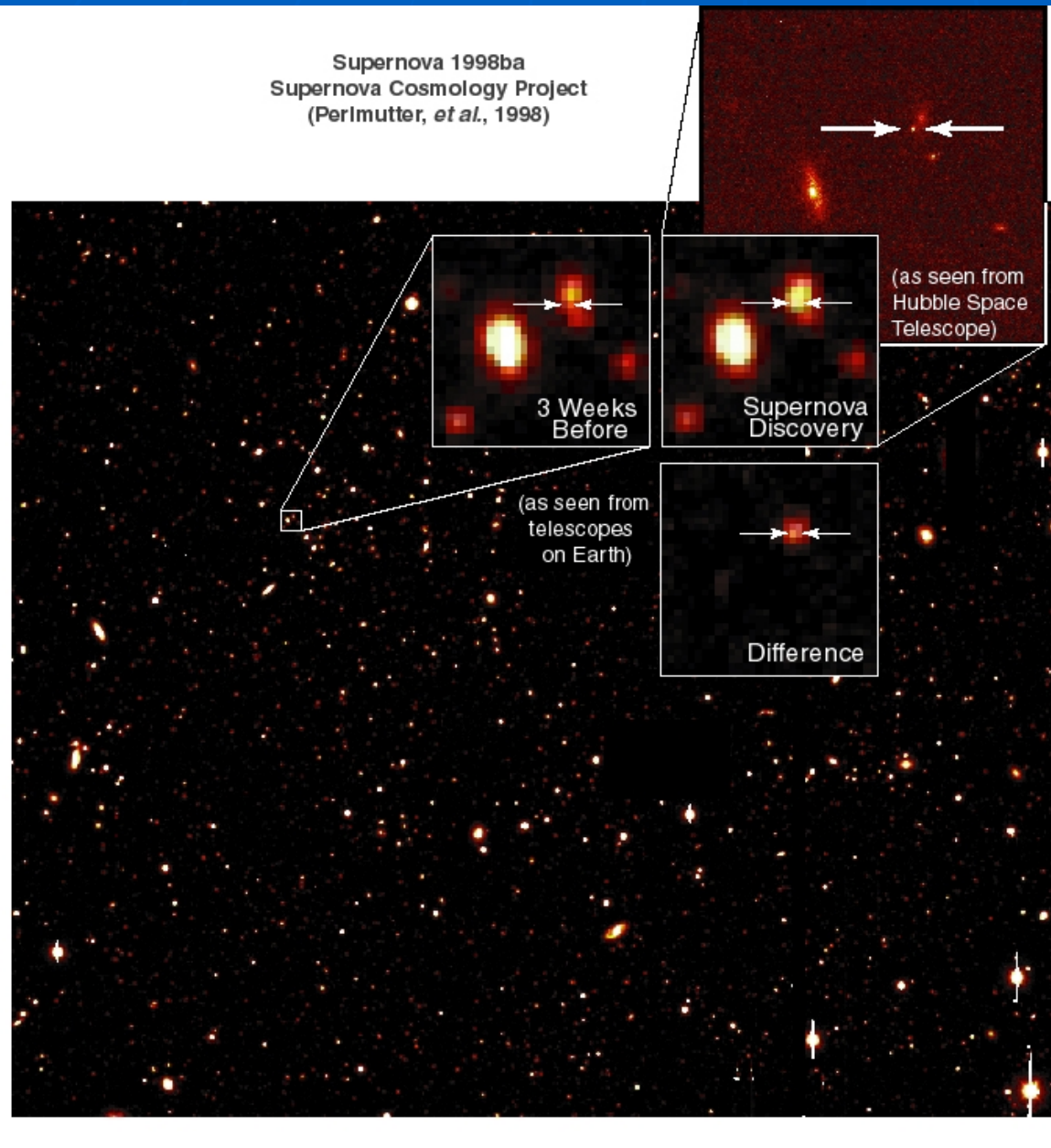


Supernovae are very rare, ~ 1 SN per 100 years and galaxy.

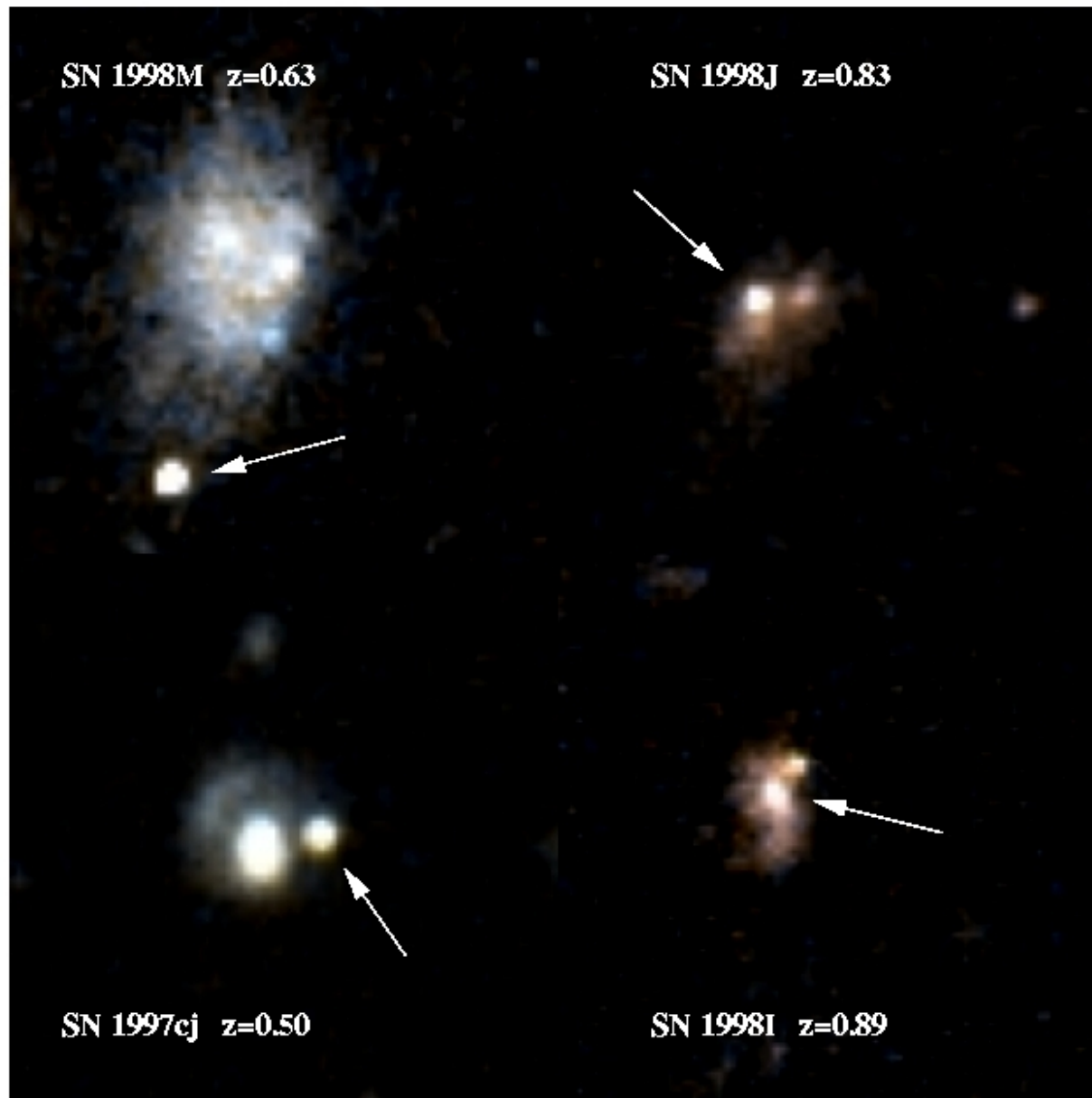


One has to observe very many galaxies!

Search strategy:



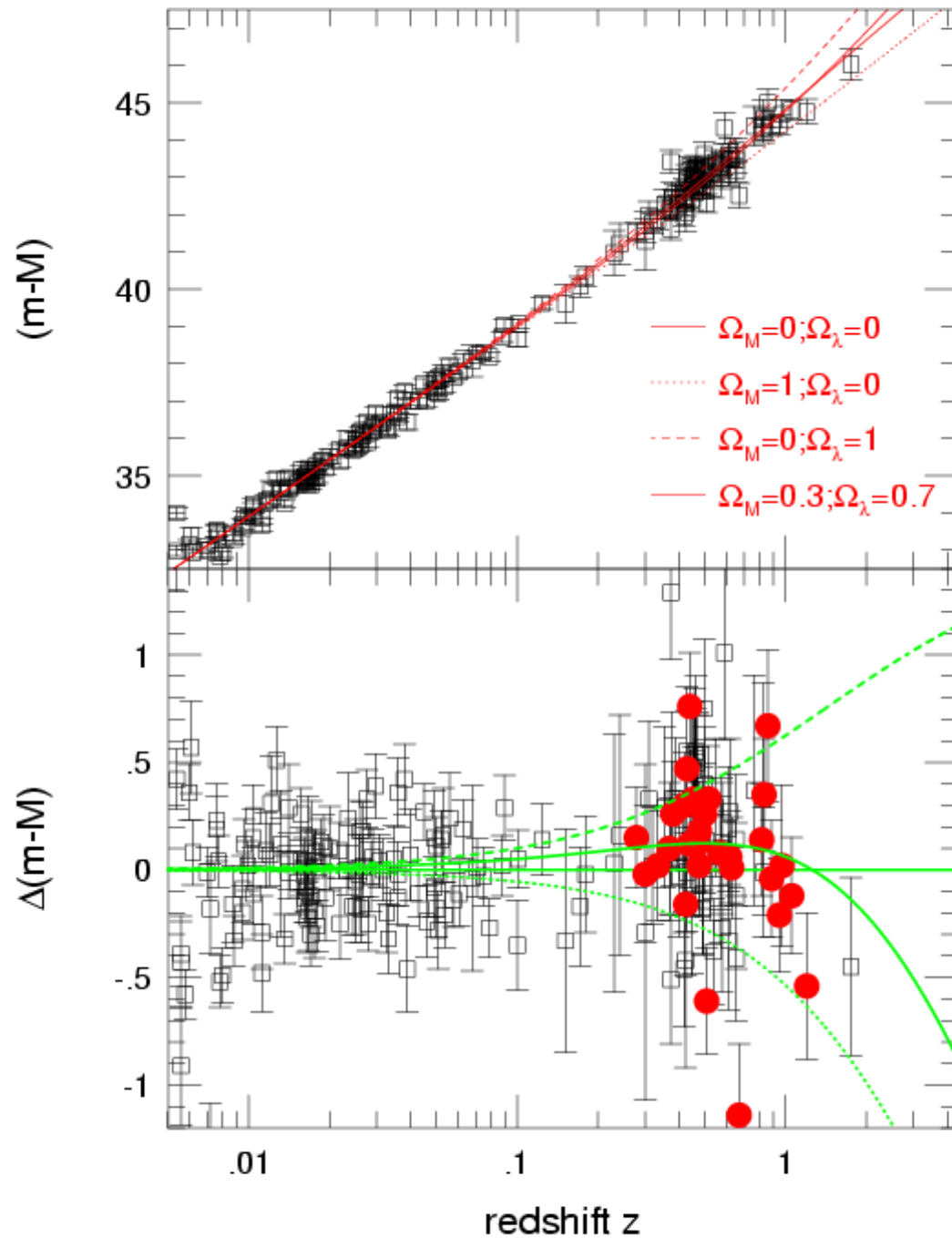
1. Repeated scanning of a certain field.
2. Electronic readout of the data.
3. Follow-up observations, e.g., HST, VLT, ...



Supernovae are routinely detected at redshifts $Z > 0.4$:

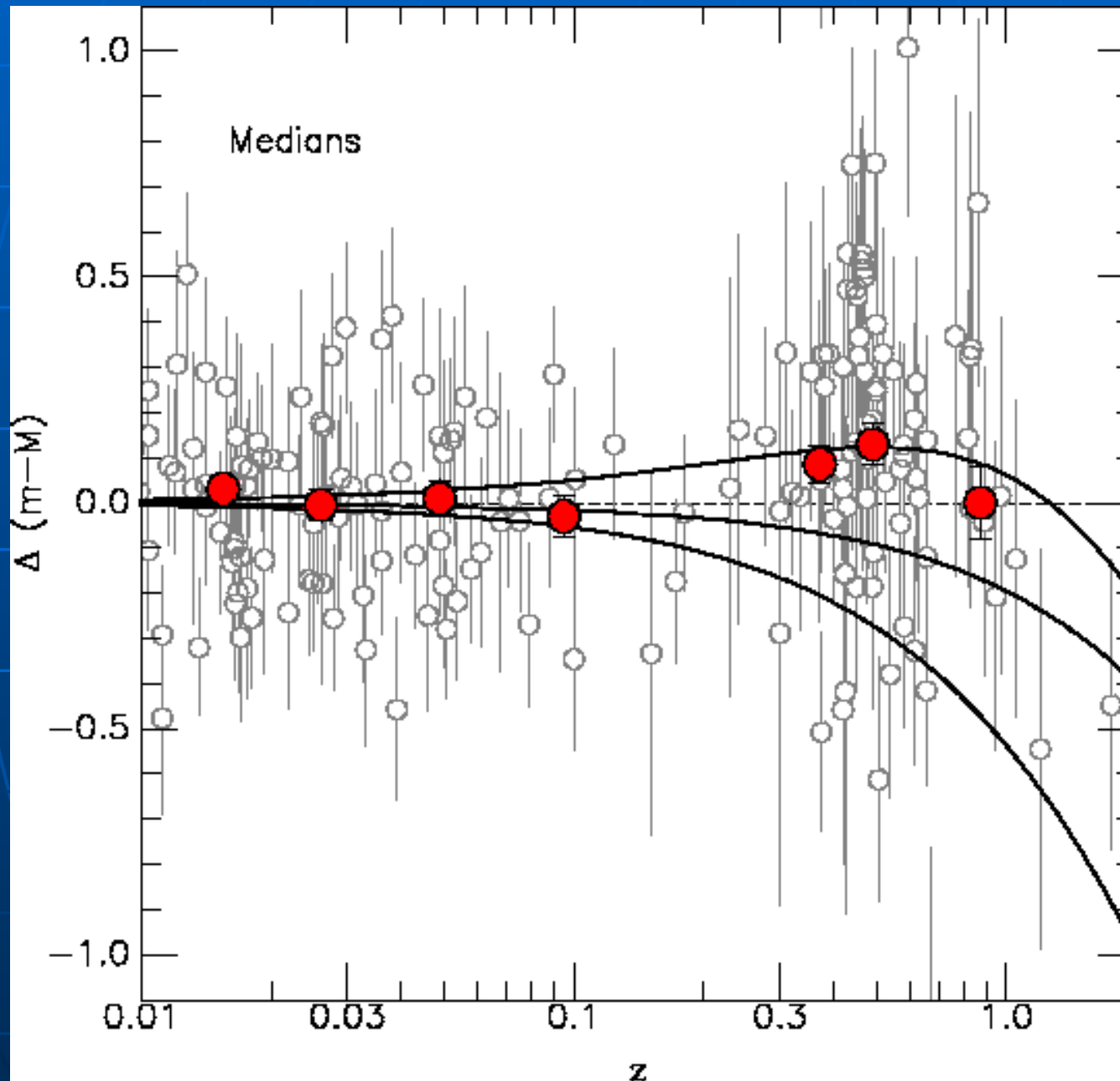
- What is the intrinsic scatter in luminosities?
- Are they different from the local sample?
- Do we understand the differences?

Supernovae at high redshifts

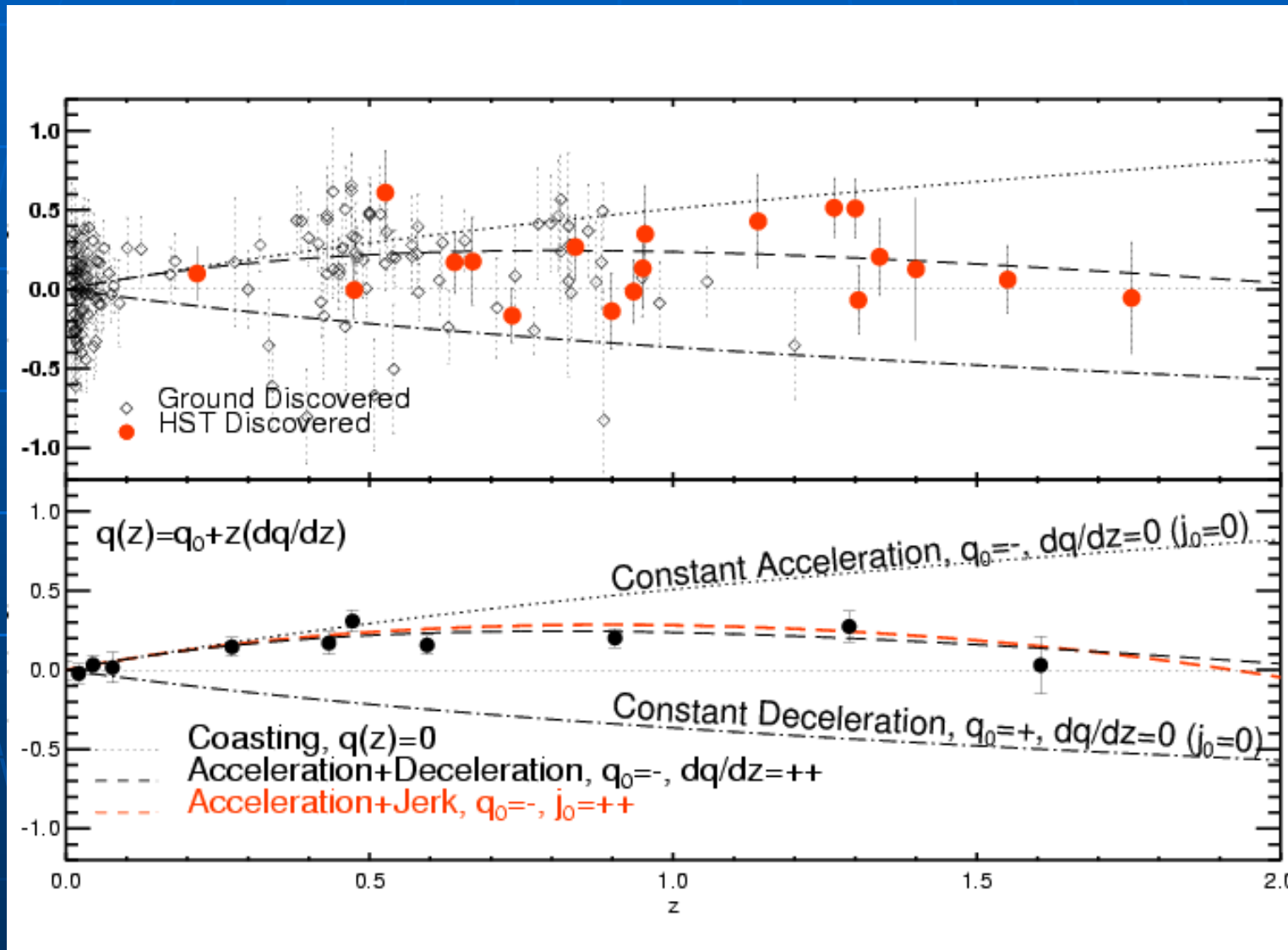


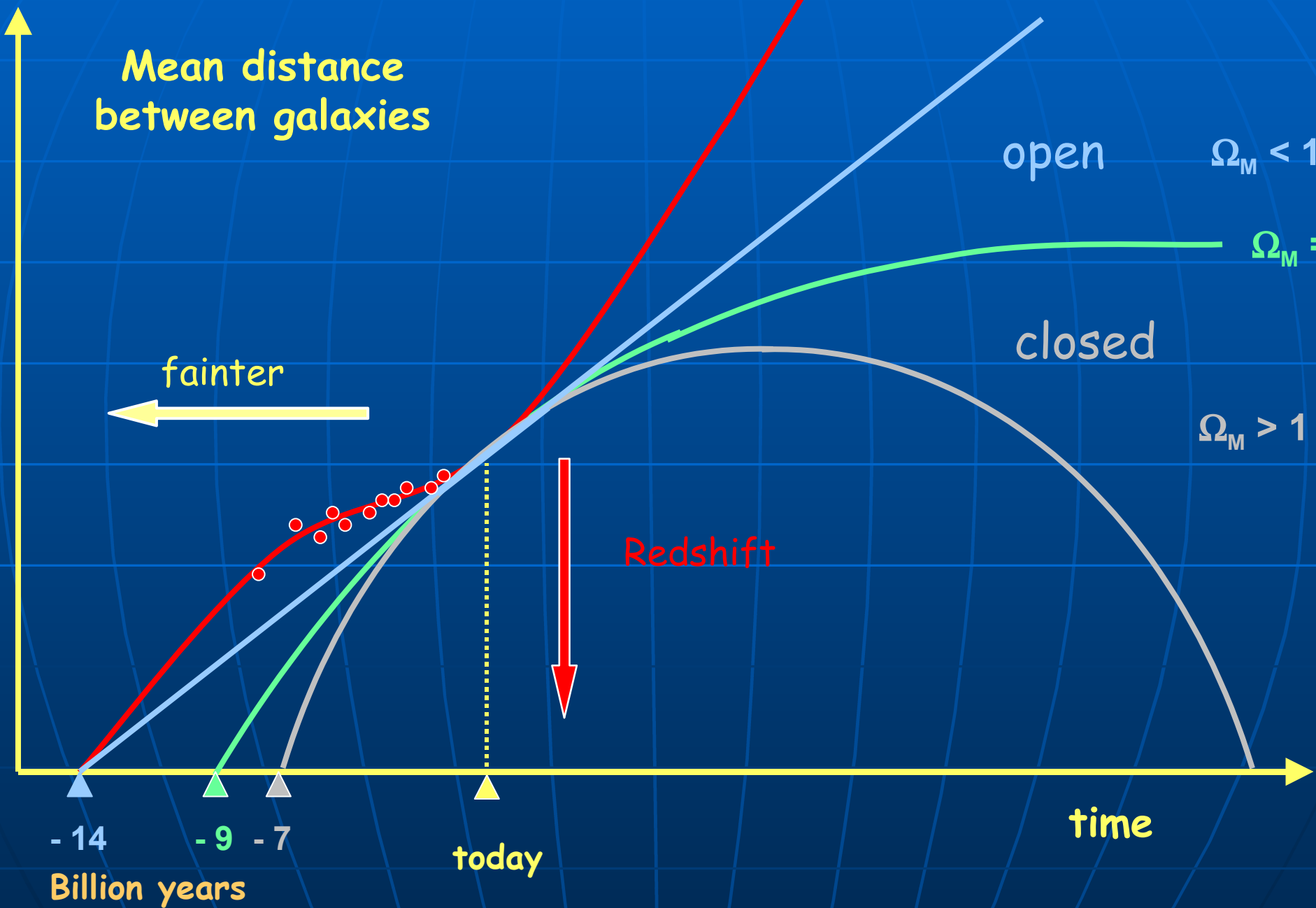
Tonry et al. 2003

209 SNe Ia and medians



Very high redshift SNe Ia





General luminosity distance

$$D_L = \frac{(1+z)c}{H_0 \sqrt{|\Omega_k|}} S \left\{ \sqrt{|\Omega_k|} \int_0^z \left[\Omega_k (1+z')^2 + \sum_i \Omega_i (1+z')^{3(1+w_i)} \right]^{-1/2} dz' \right\}$$

- with $\Omega_k = 1 - \sum_i \Omega_i$ and $w_i = \frac{p_i}{\rho_i c^2}$

$w_M = 0$ (matter)

$w_R = 1/3$ (radiation)

$w_\Lambda = -1$ (cosmological constant)

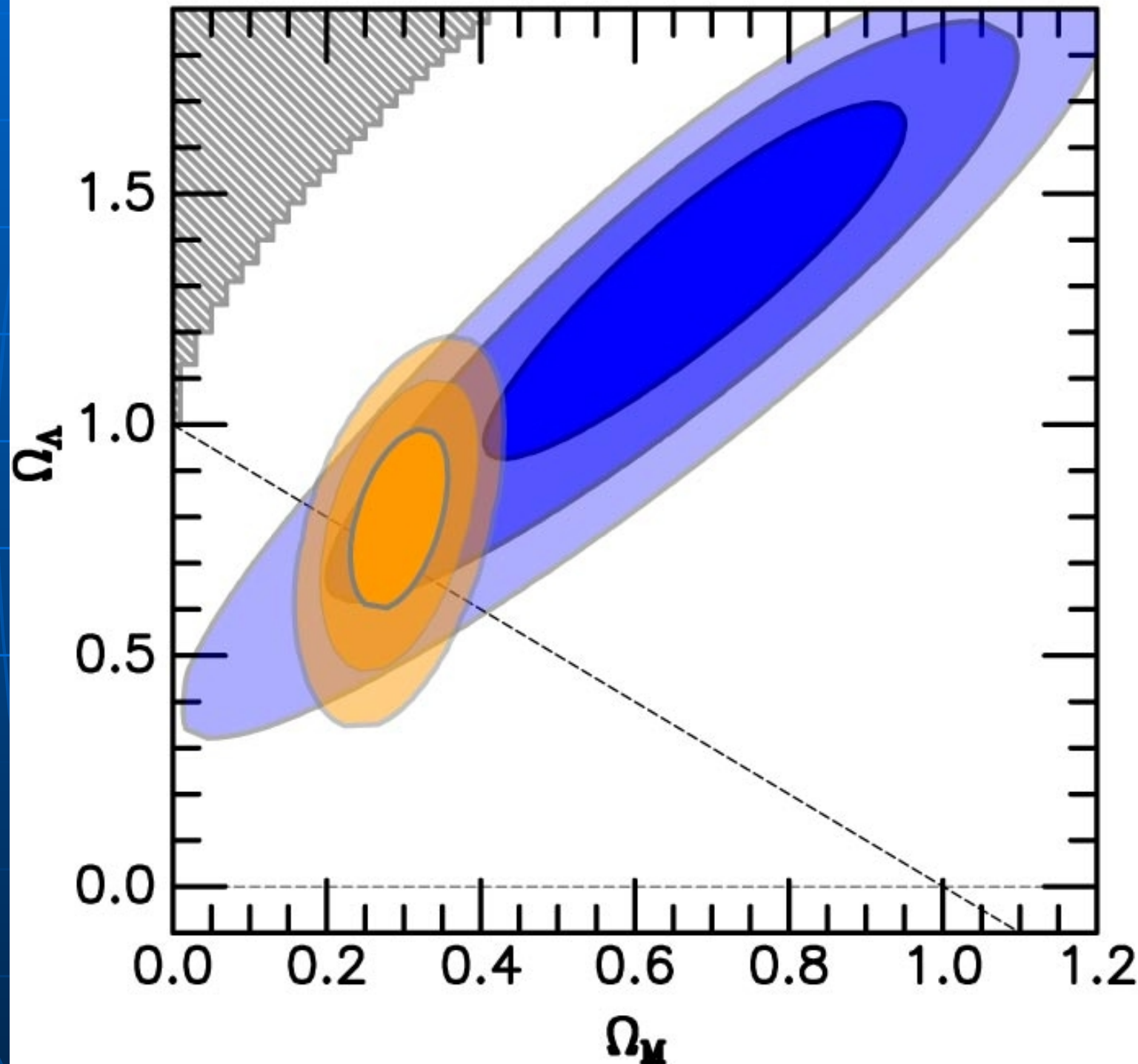
2dF:

$$\Omega_M = 0.2 \pm 0.03$$

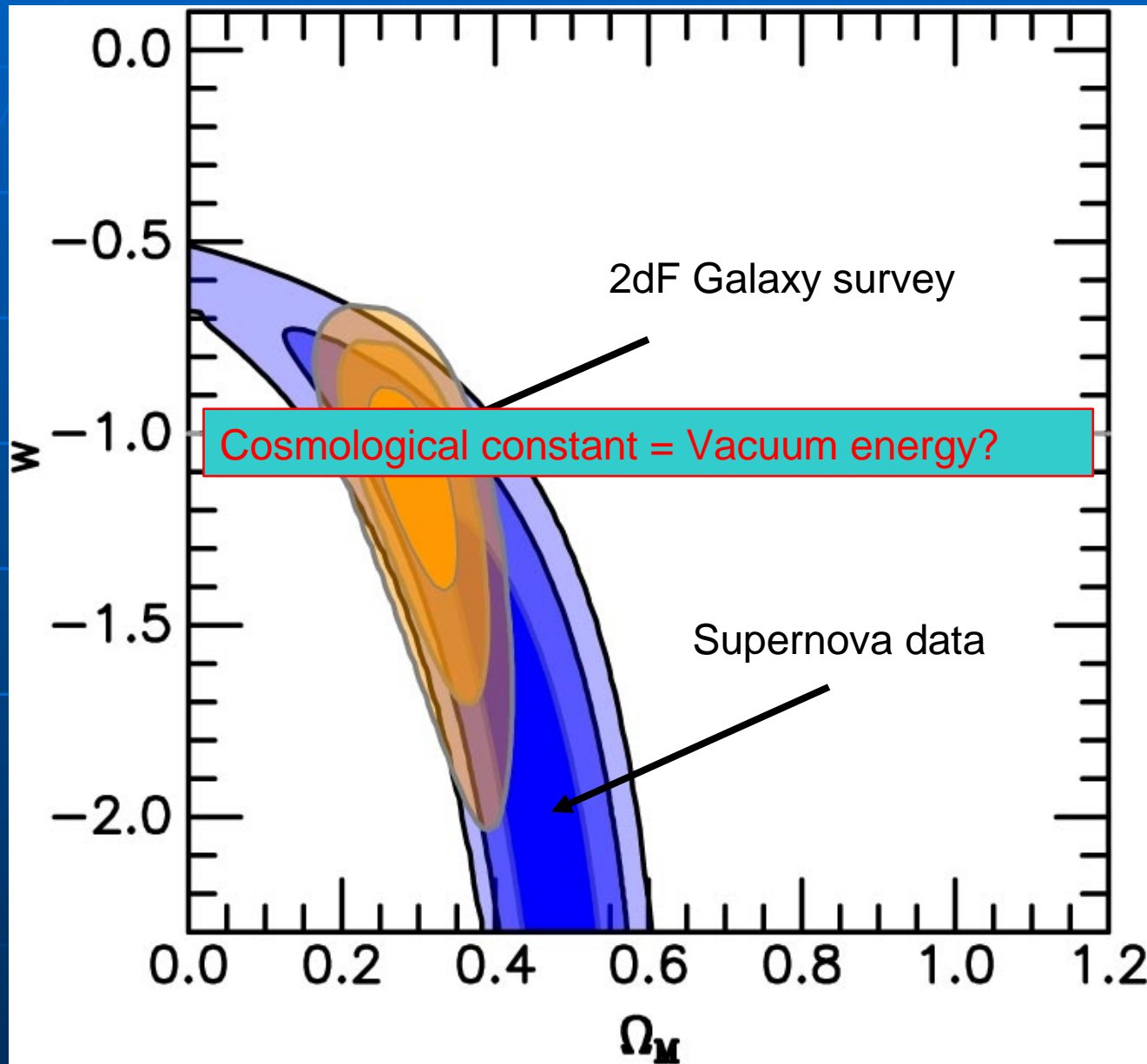
KP:

$$h = 0.72 \pm 0.08$$

Entire High-Z SN Ia Data Set



Cosmology and Typ Ia supernovae



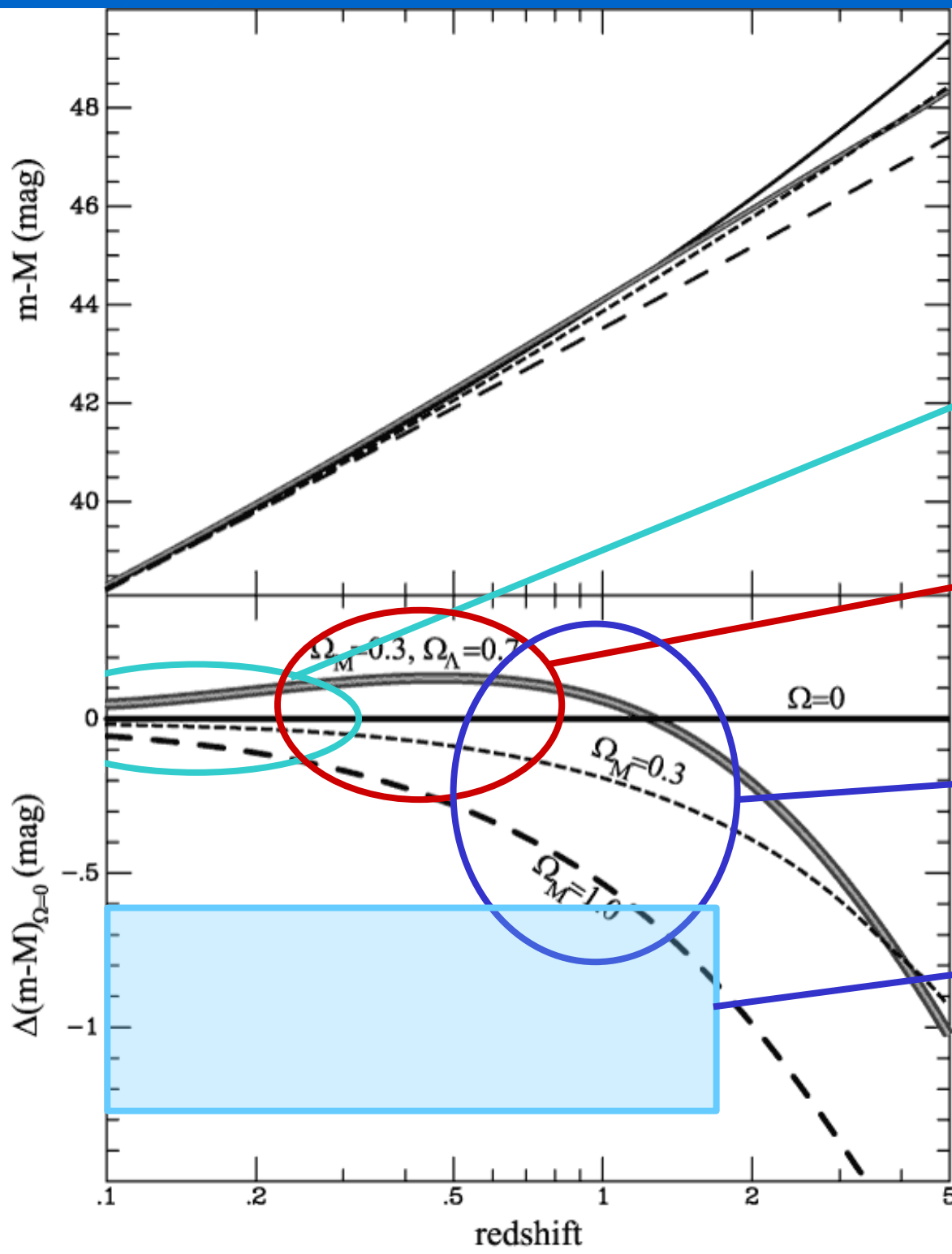
The “equation of state” of the Universe:

$$p = w\rho$$

$$\ddot{a} \sim (\rho + 3p)$$

$$w < -1/3 :$$

Acceleration!



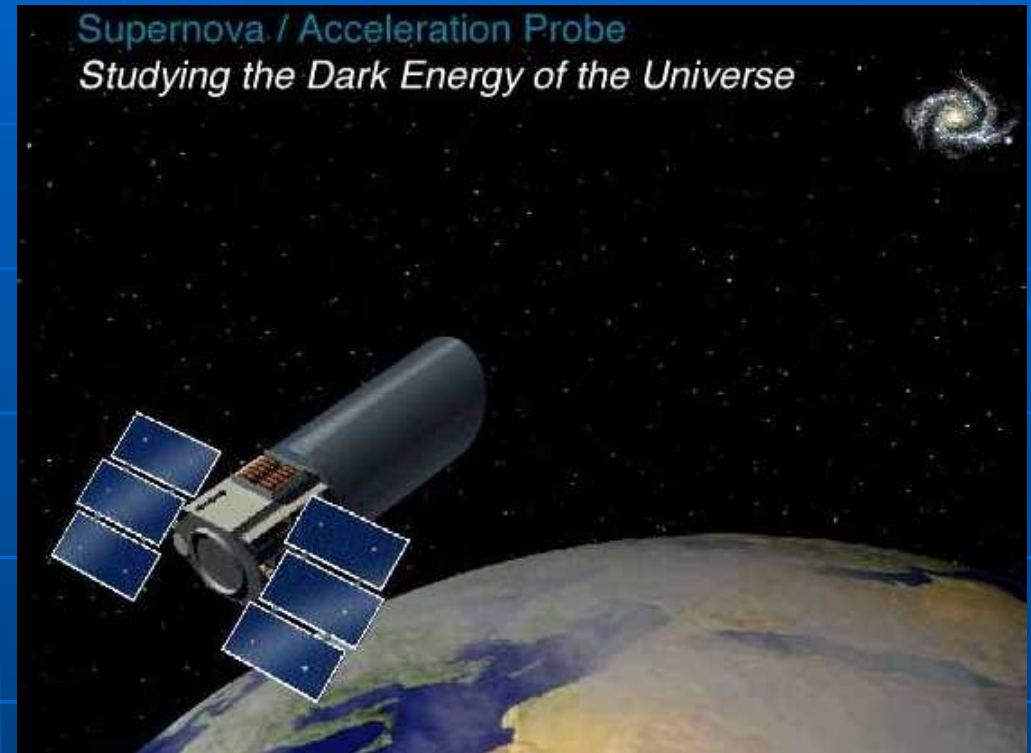
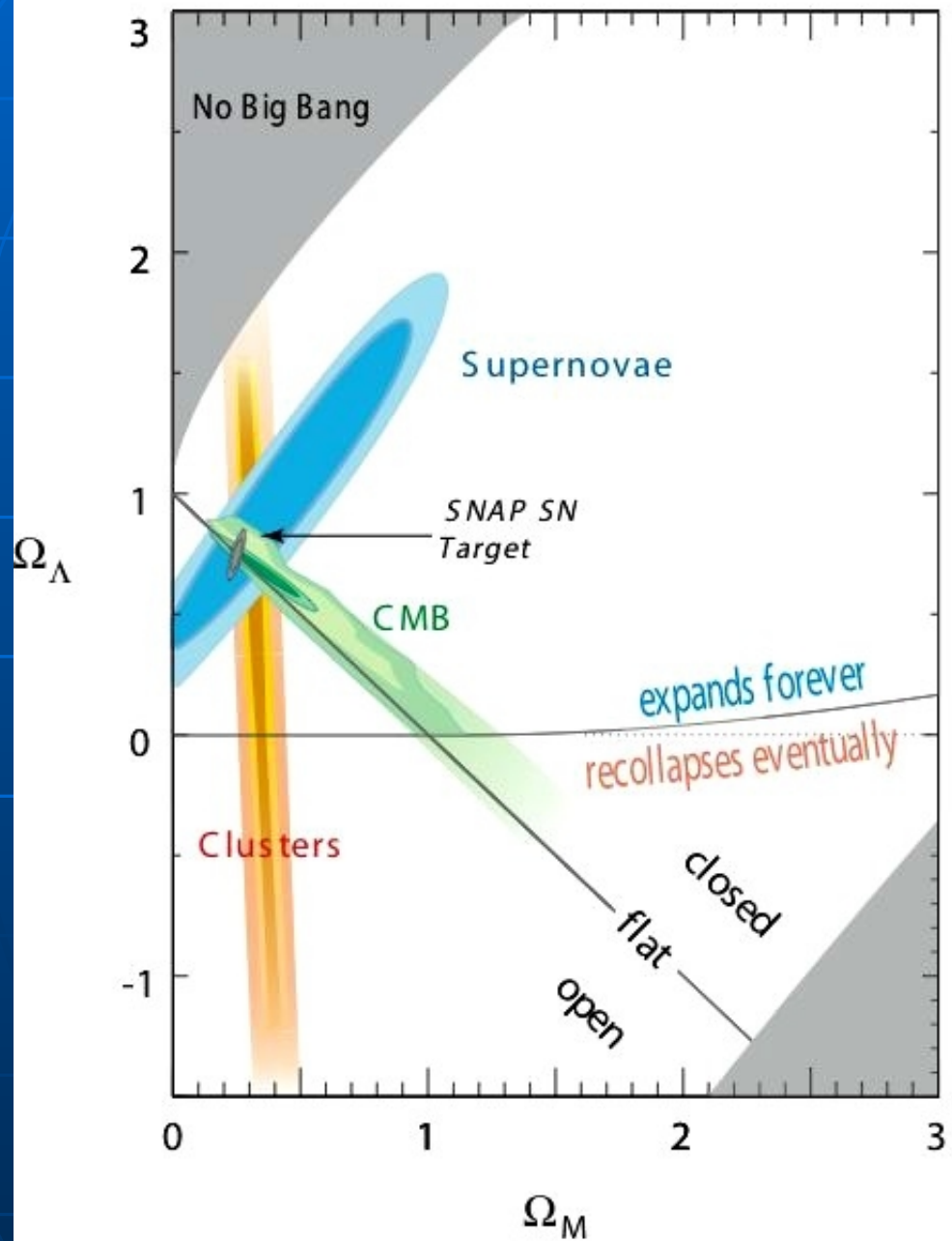
SN Projects

SN Factory
Carnegie SN Projekt

ESSENCE
CFHT Legacy Survey

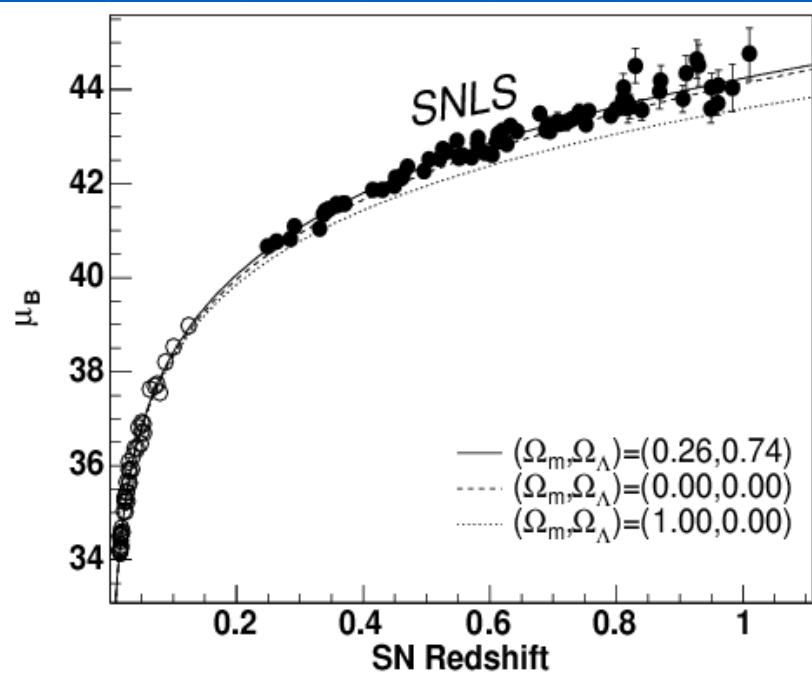
High-z SN Search
(GOODS)

SNAP
(Supernova
acceleration
Probe)



SNAP:
“Supernova/Acceleration Probe”

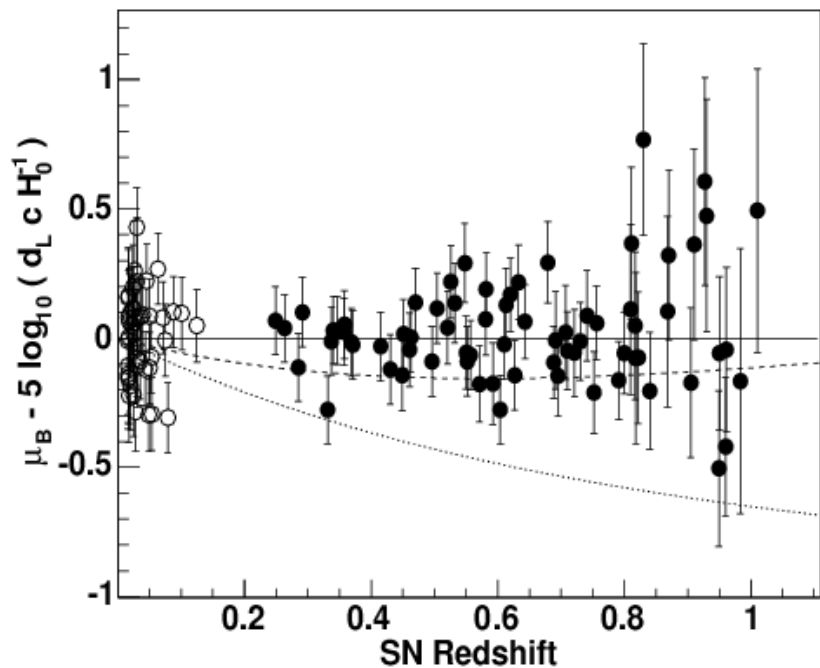
SNLS's first year Hubble diagram



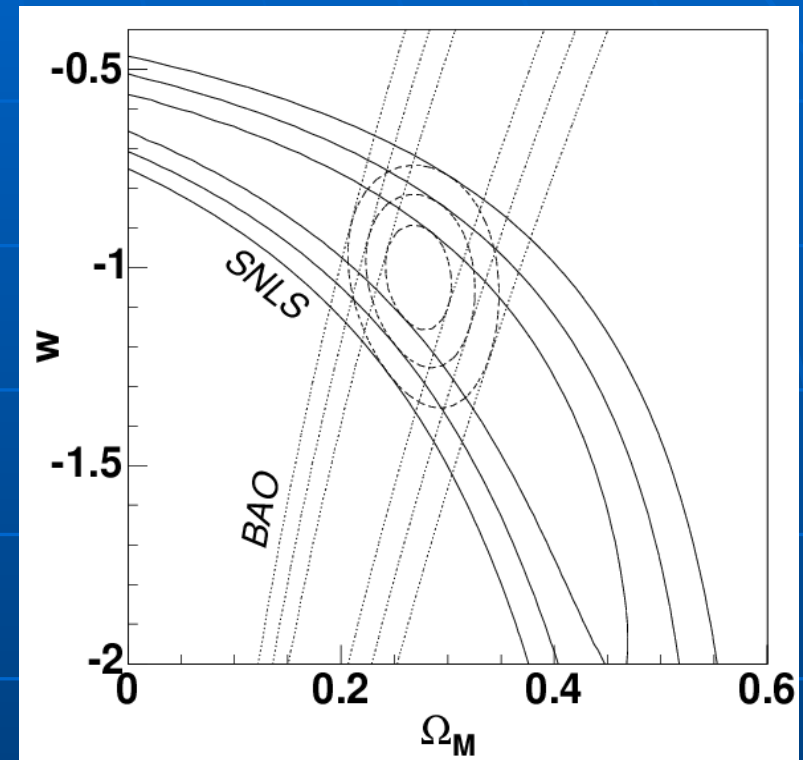
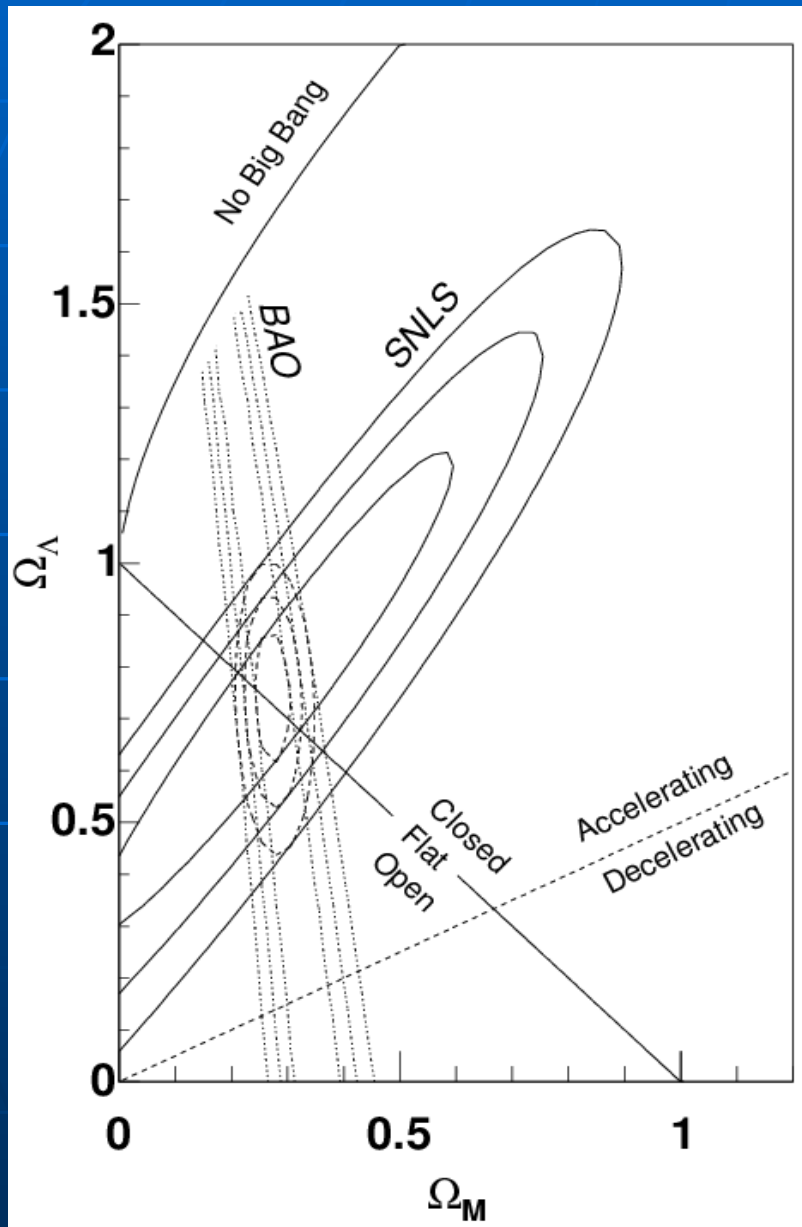
Final sample :
45 nearby SNe from literature
+71 SNLS SNe

Intrinsic scatter: (0.13 ± 0.02) mag

(Astier et al., 2006)



SNLS's Cosmological parameters



Solid contours : SNLS

Dotted contours : Baryon acoustic oscillations (BAO) (SDSS, Eisenstein et al., 2005)
(68.3, 95.5 and 99.7% CL)

(Astier et al., 2006)

SNLS 1st year results on Cosmology

For a flat Λ CDM cosmology :

$$\Omega_M = 0.264 \pm 0.042 \text{ (stat)} \pm 0.032 \text{ (sys)}$$

Combined with BAO (Eisenstein, 2005) :

$$\Omega_M = 0.271 \pm 0.021 \text{ (stat)} \pm 0.007 \text{ (sys)}$$

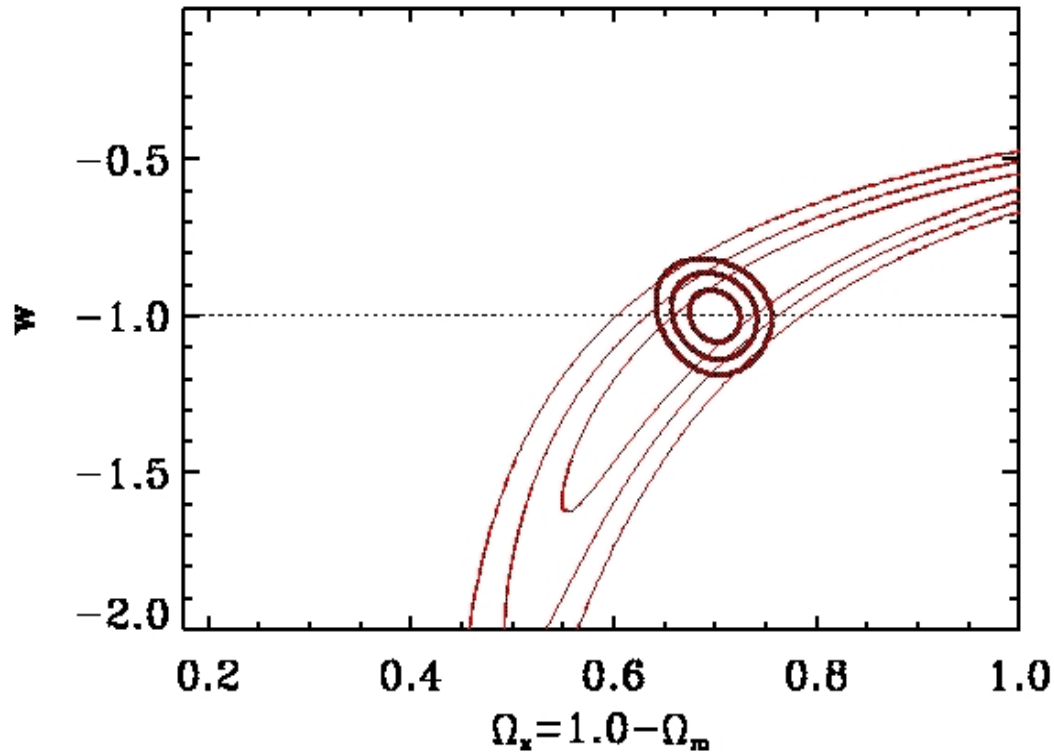
$$w = -1.02 \pm 0.09 \text{ (stat)} \pm 0.054 \text{ (sys)}$$

(Preliminary numbers!)

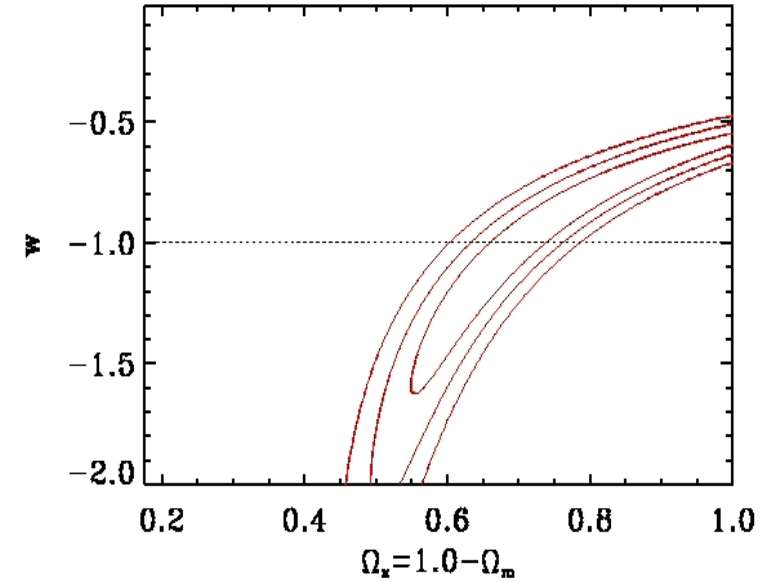
(Astier et al., 2006)

ESSENCE: Anticipated Cosmology Limits

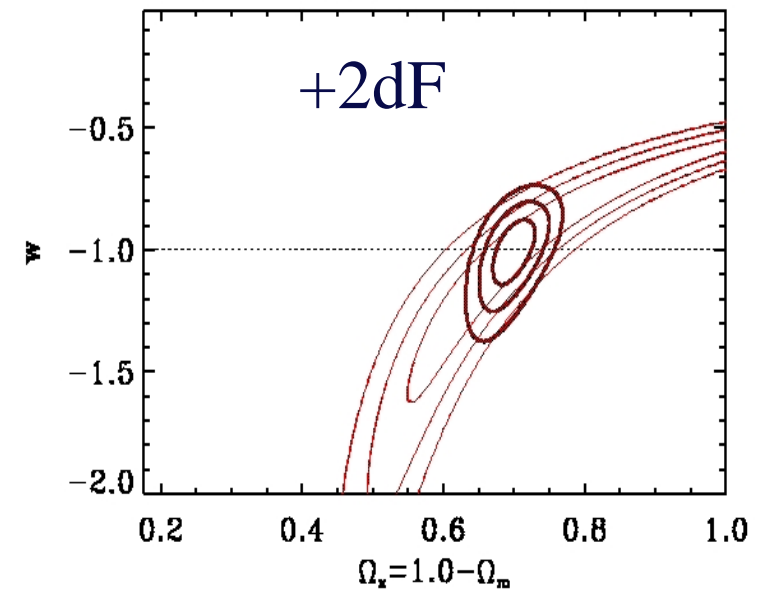
ESSENCE + WMAP



ESSENCE: 200 SN, z limit 0.8



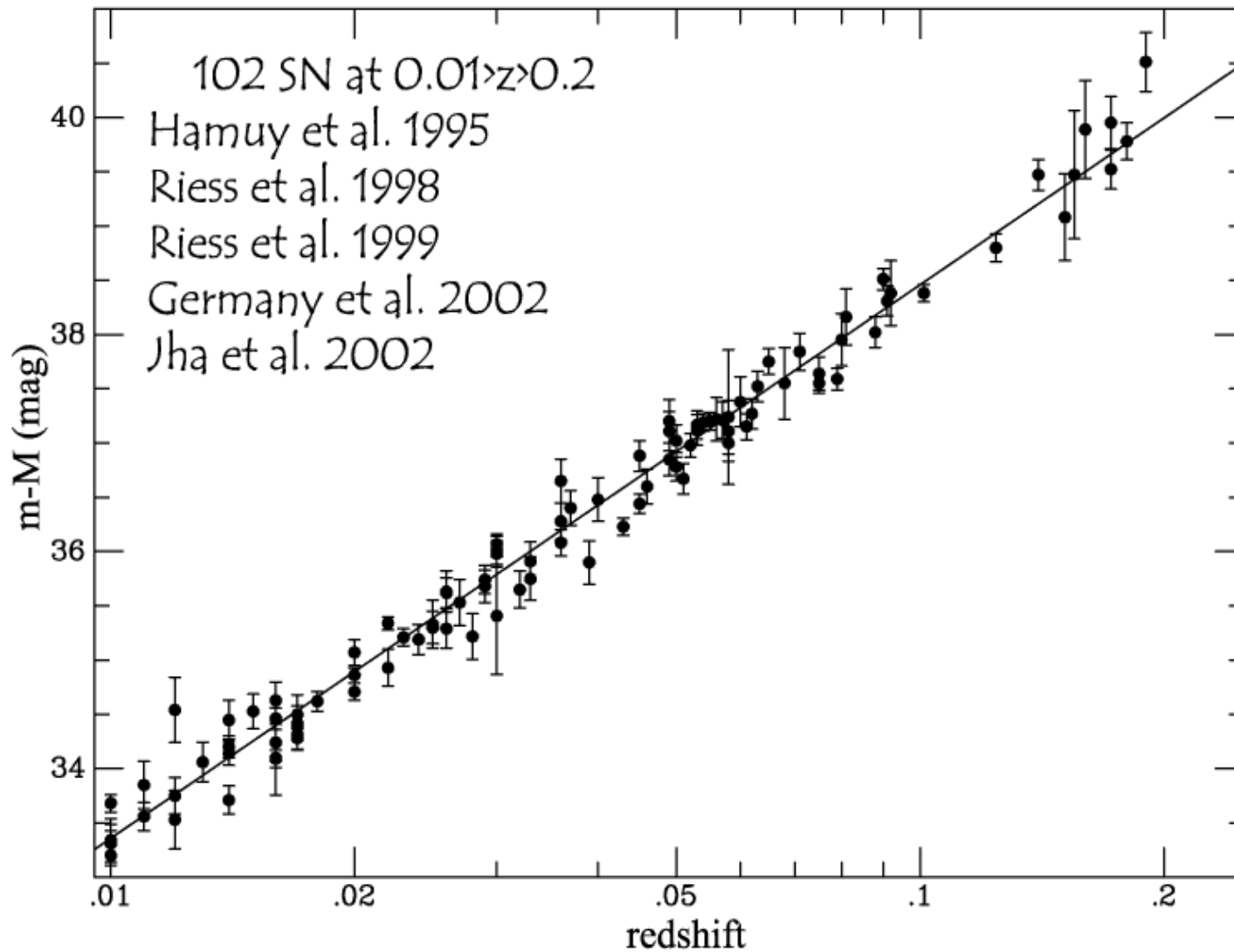
ESSENCE: 200 SN, z limit 0.8



What can still be wrong???

- Systematic errors?
- Pollution of high- Z samples?

Systematics: Nearby supernovae?



Hubble diagram:

Scatter in $(m - M)$:
 ≈ 0.15 mag

(≈ 0.13 mag, SNLS)

Required for w :

≈ 0.02 mag !

Statistics?

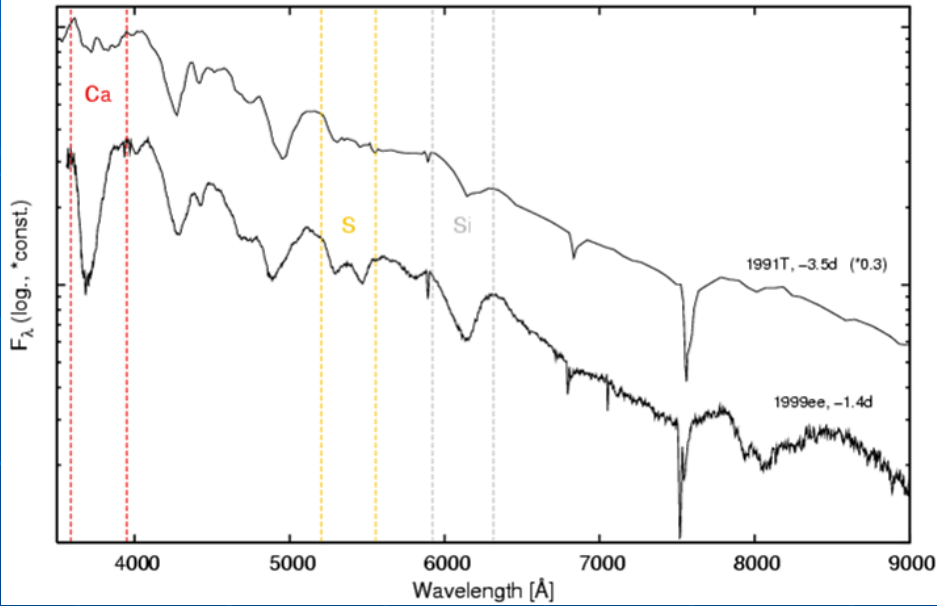
Systematics?

(Tonry et al. 2003)

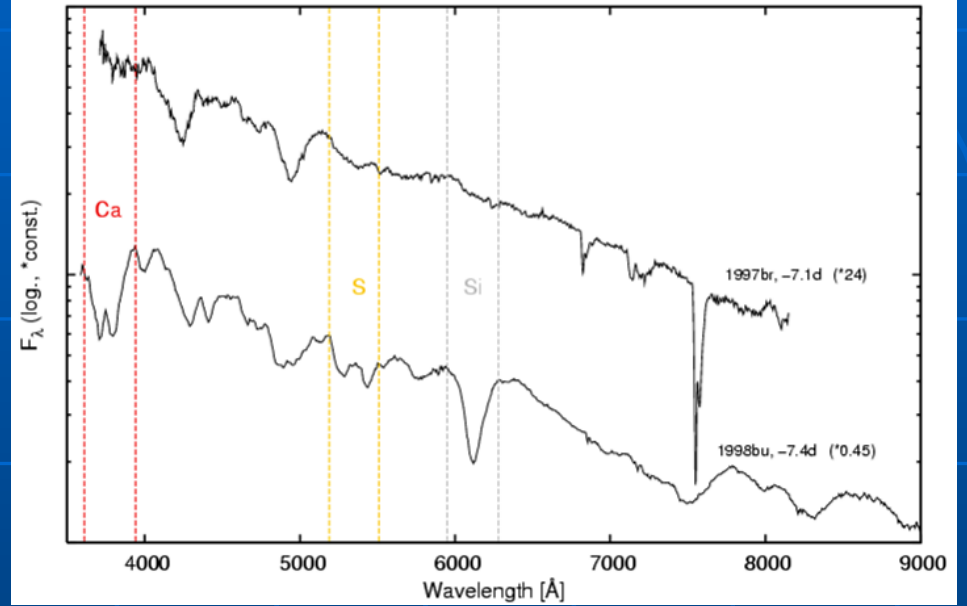
Are they different?

(Early spectra; court. Stephan Hachinger)

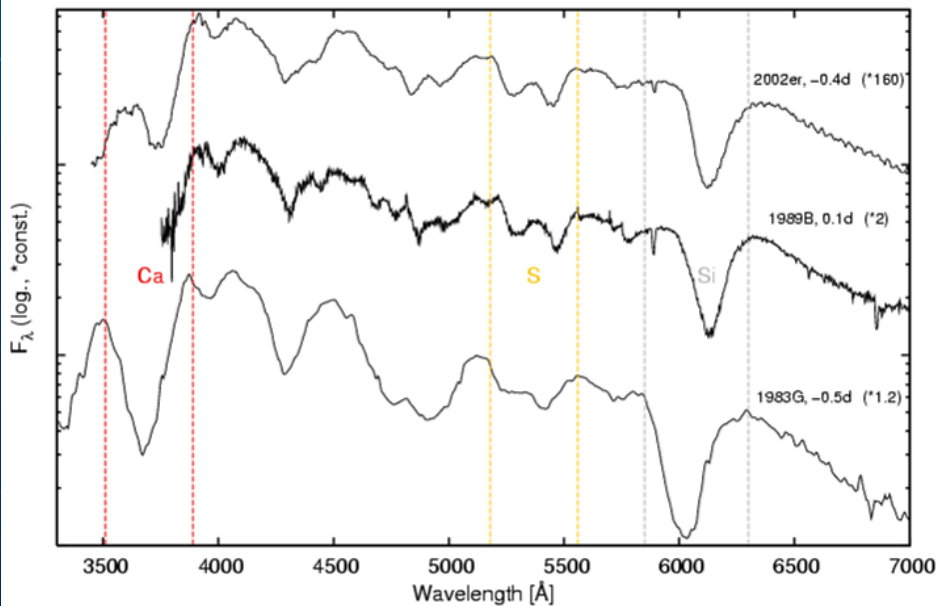
Comparing SNe Ia with $\Delta m_{15} = 0.9$



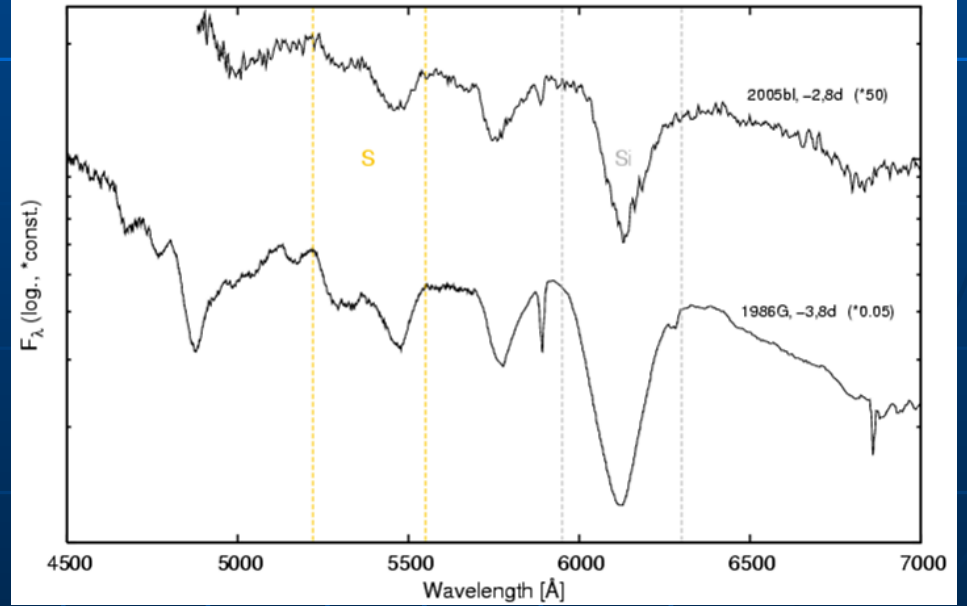
Comparing SNe Ia with $\Delta m_{15} = 1.04$



Comparing SNe Ia with $\Delta m_{15} = 1.35$



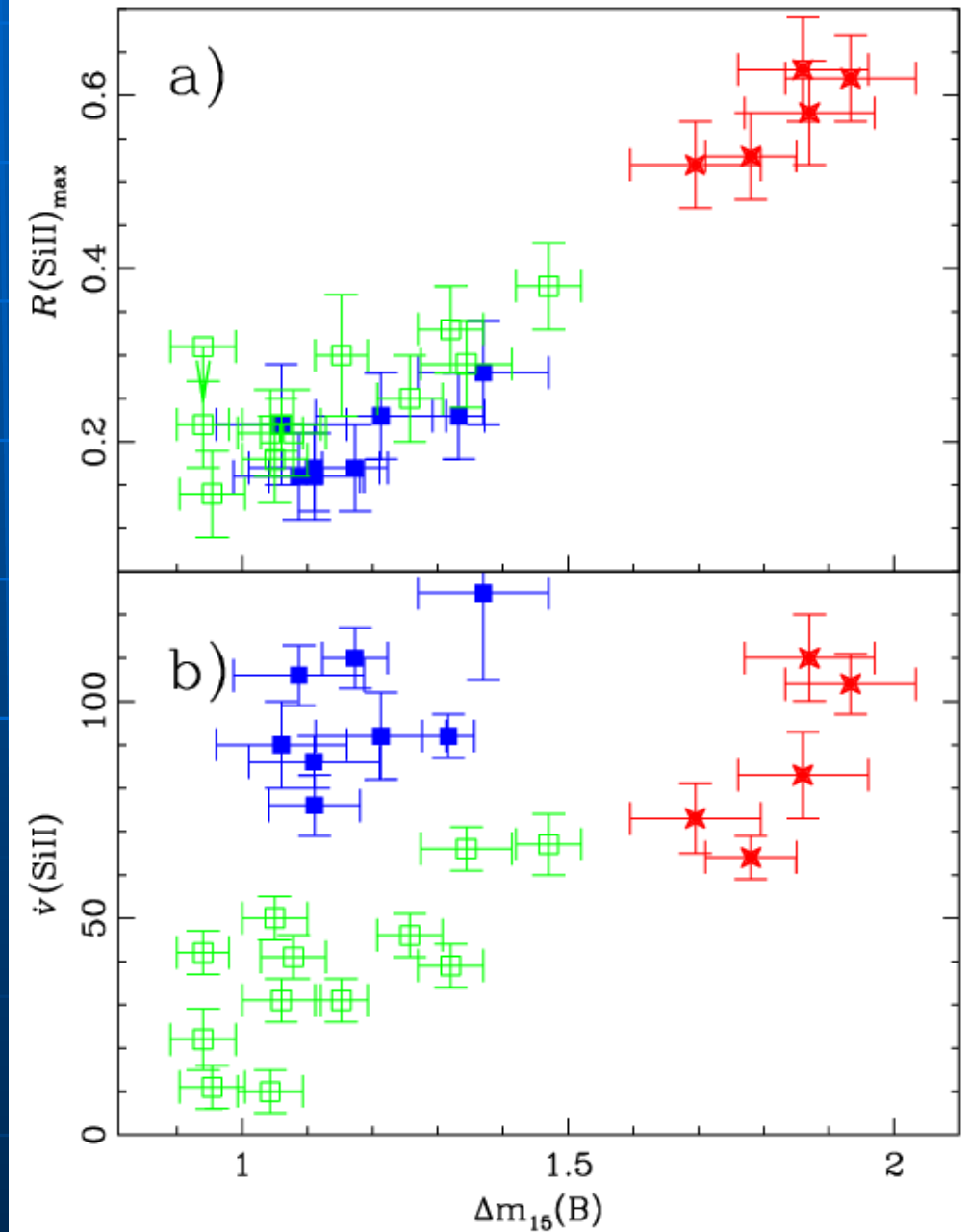
Comparing SNe Ia with $\Delta m_{15} = 1.8$



Expansion velocities?

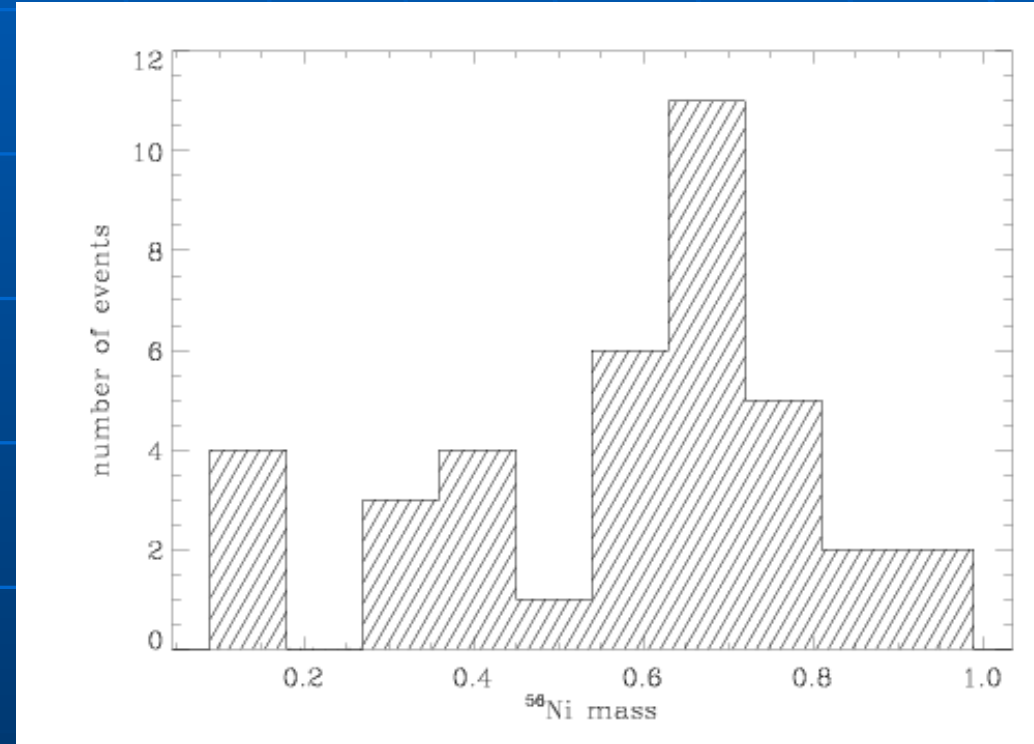
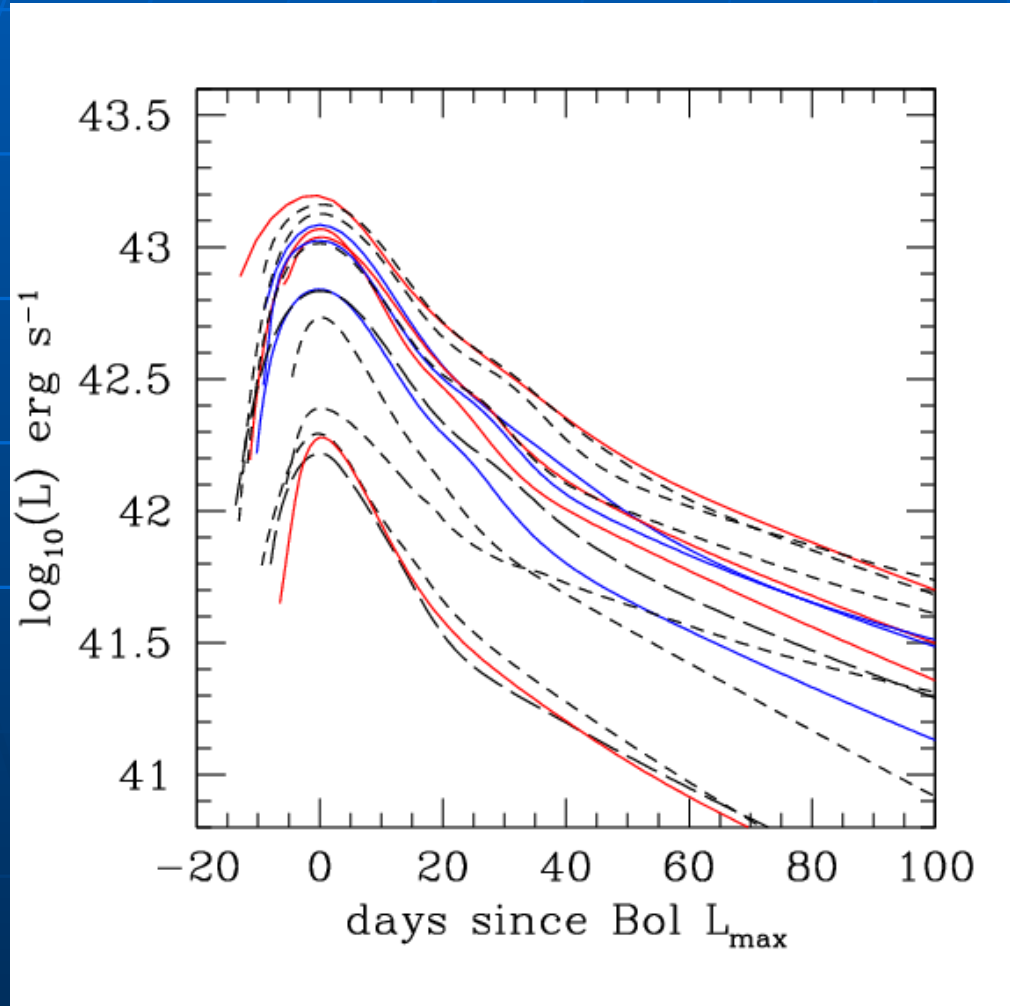
(mostly RTN/ESC data)

(Benetti et al. 2004, 2005)



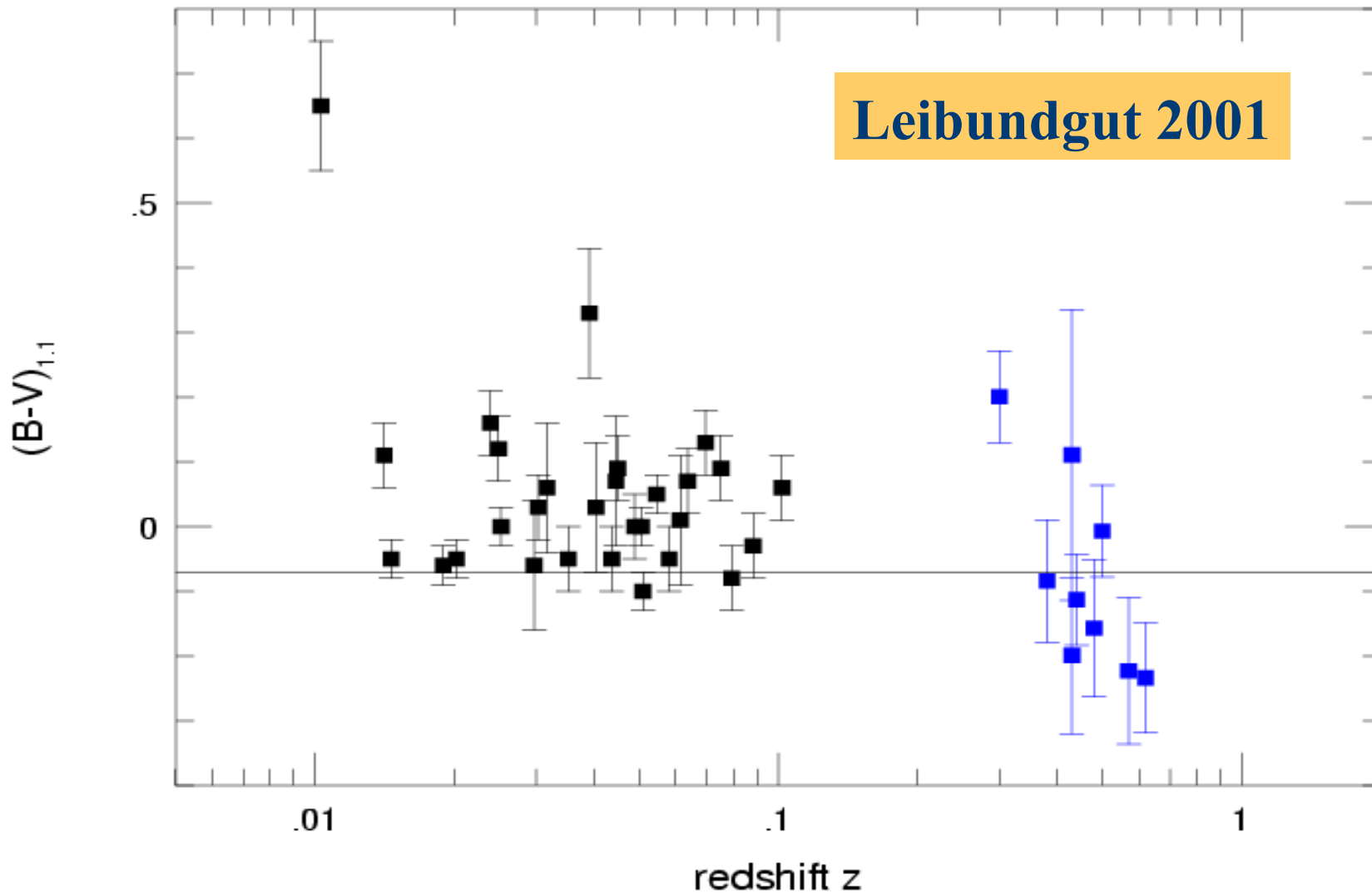
Bolometric LC's and Ni-masses?

(mostly RTN/ESC data)



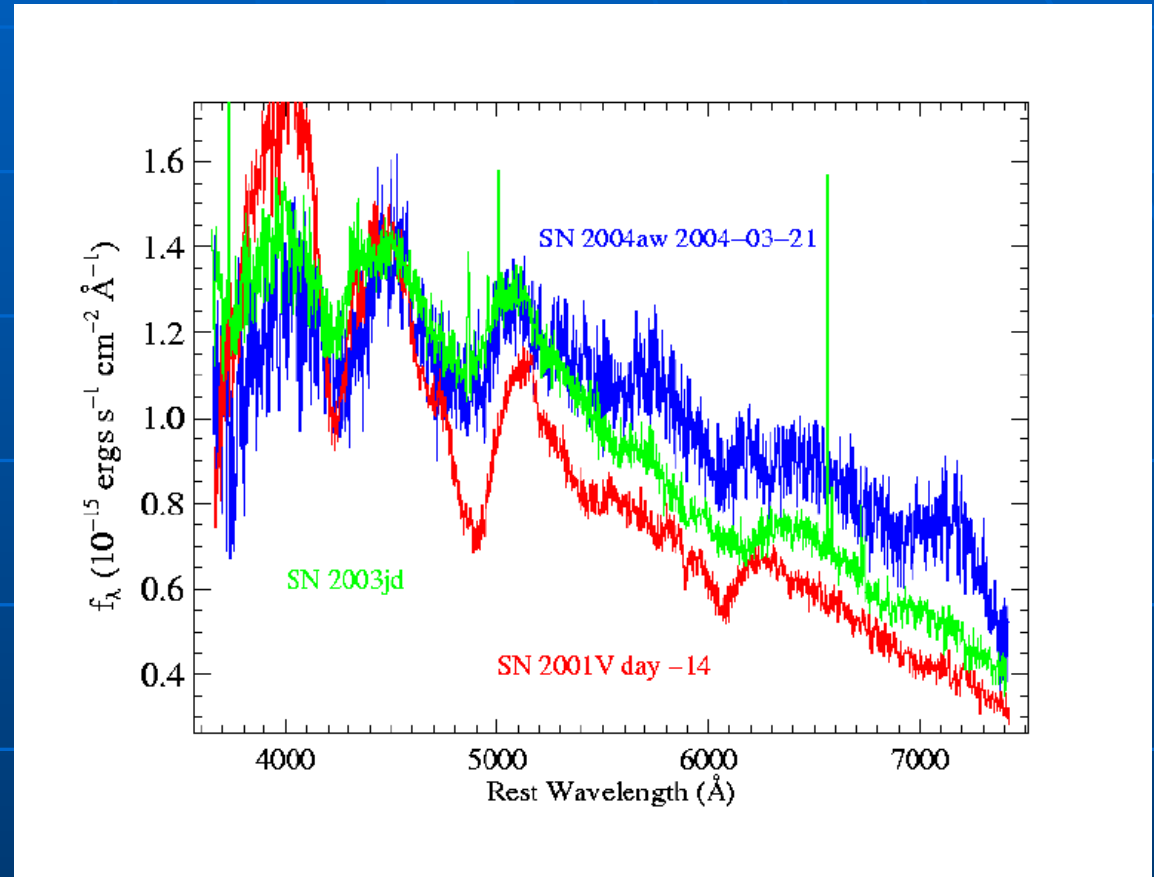
(Court. M. Stritzinger)

Is evolution a problem?



Absorption distributions

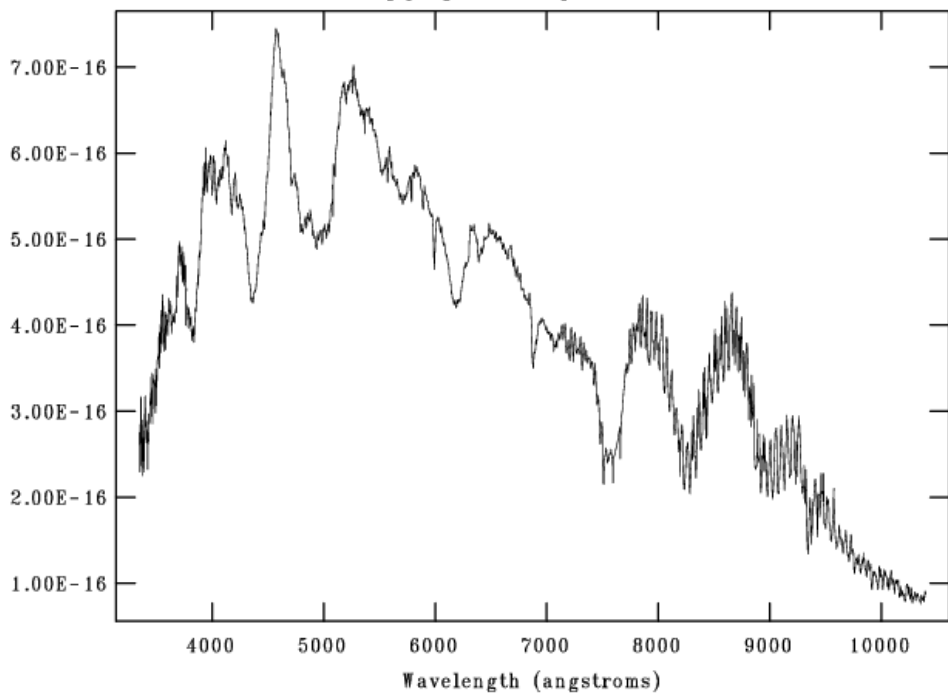
Pollution of the samples: SN 2004aw!



- Discovery: *March 19, 2004*
- Host galaxy: *NGC 3997; $V = 4771 \text{ km/s}$*
- B-Maximum: *\sim April 5, 2004*

First classification

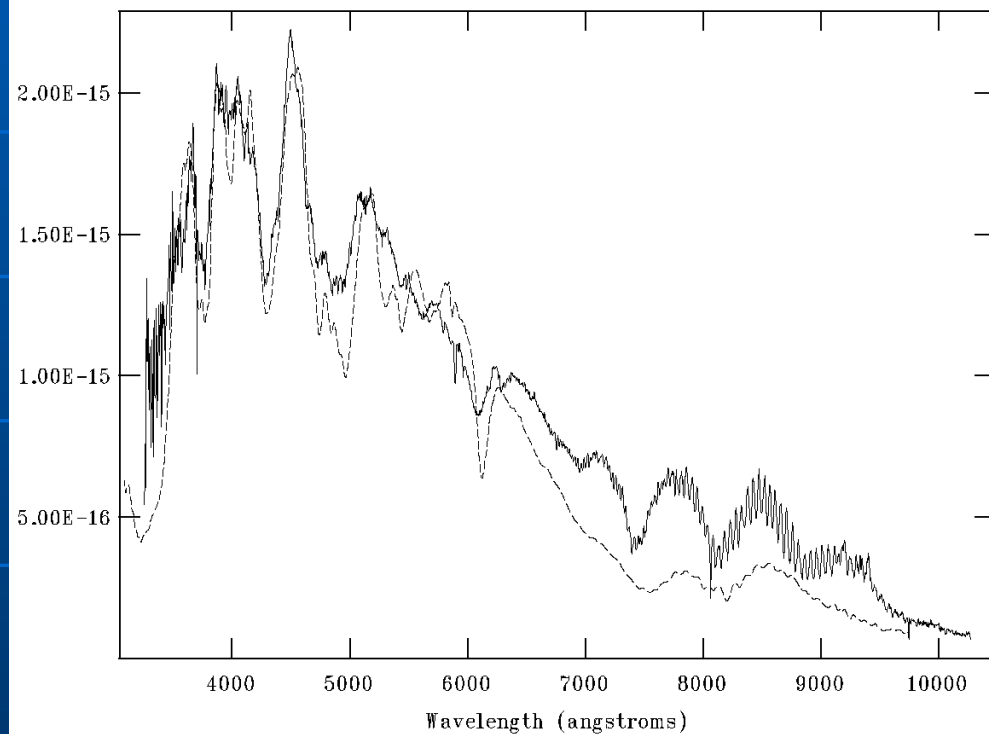
NOAO/IRAF V2.12.2-EXPORT tauben@ncf-11 Fri 16:04:42 30-Apr-2004
[spec]: 1200. ap:1 beam:1



Early spectrum of SN 2004aw

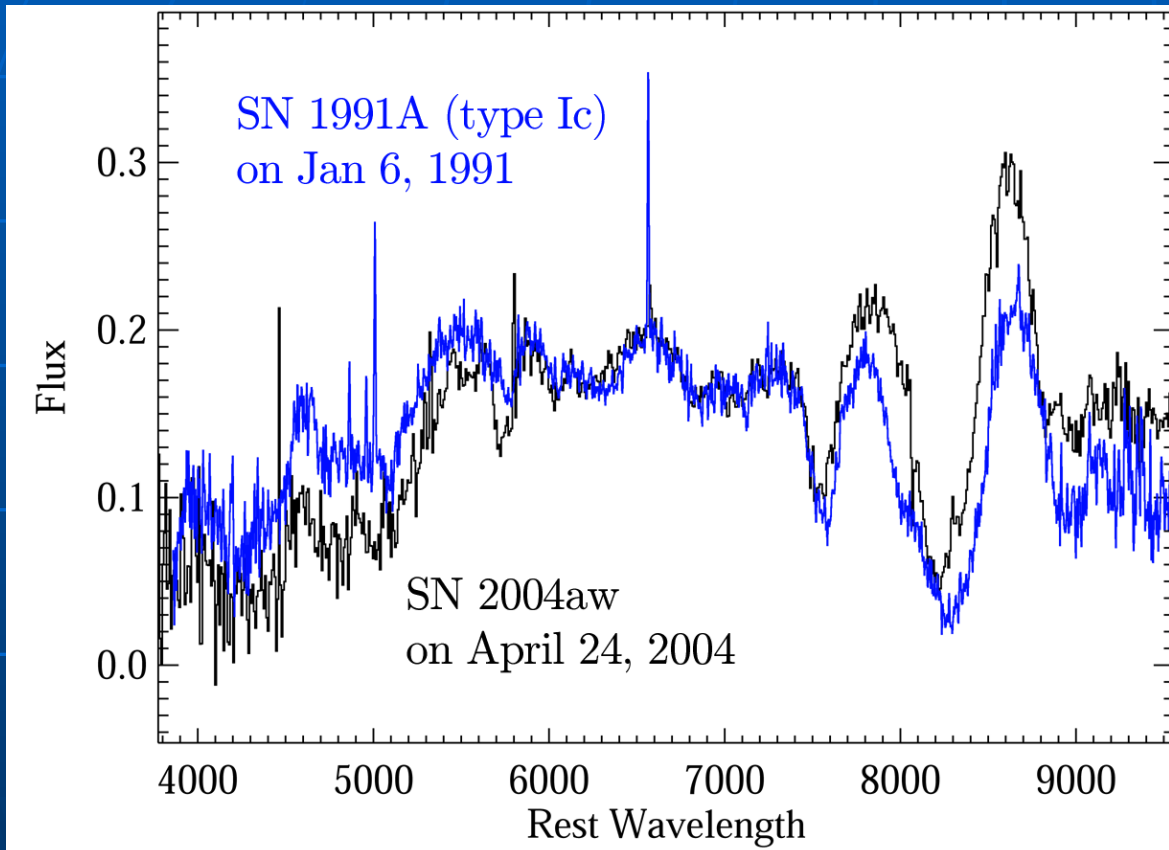
(S. Benetti)

NOAO/IRAF V2.12.1-EXPORT sbenetti@graspa.pd.astro.it Fri 09:00:58 26-Ma
[2004aw_20040324_z0_der]: 1200. ap:1 beam:1

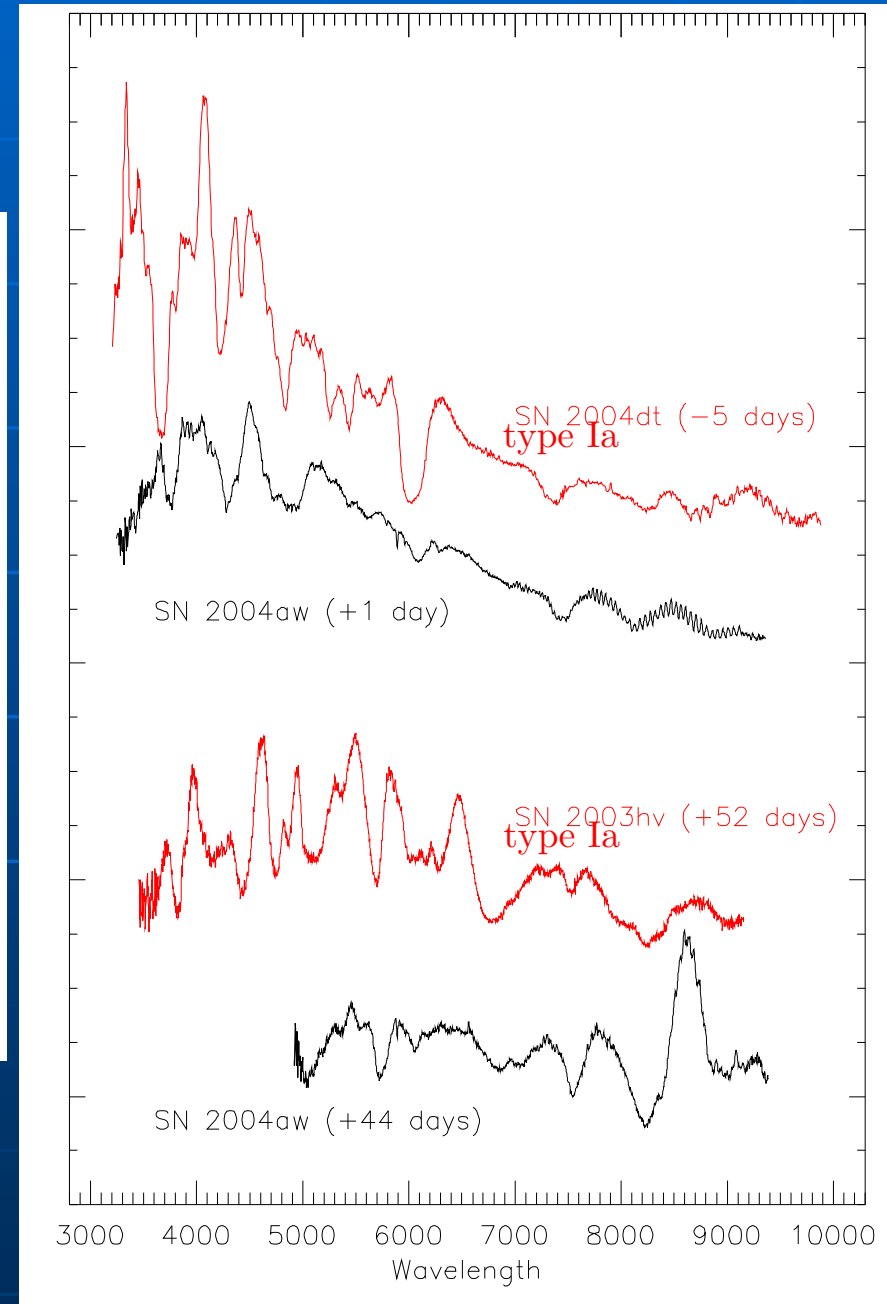


... and SN 1991T

After one month of observations: Reclassification by Alex Filippenko

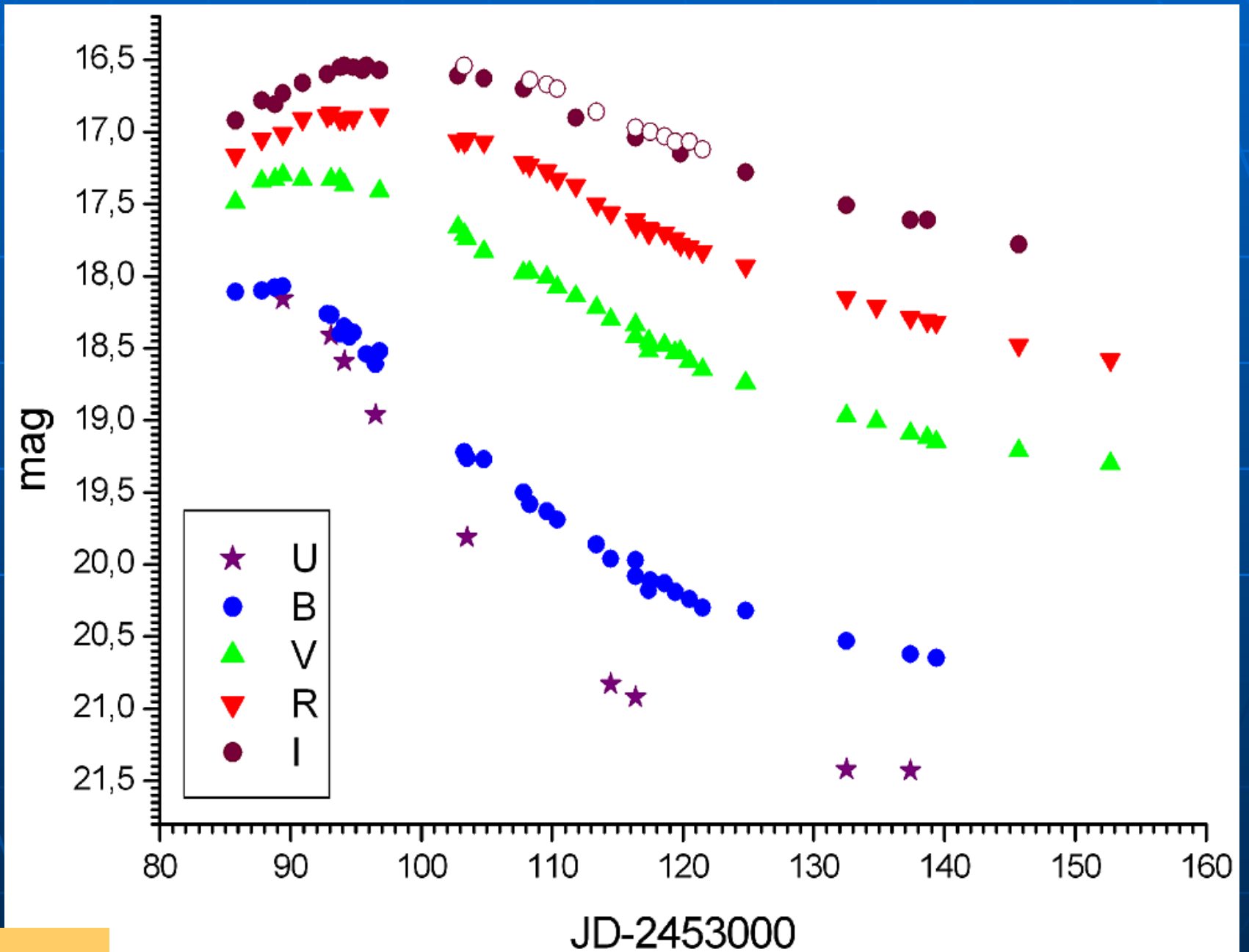


It's a Ic !!!



S. Taubenberger

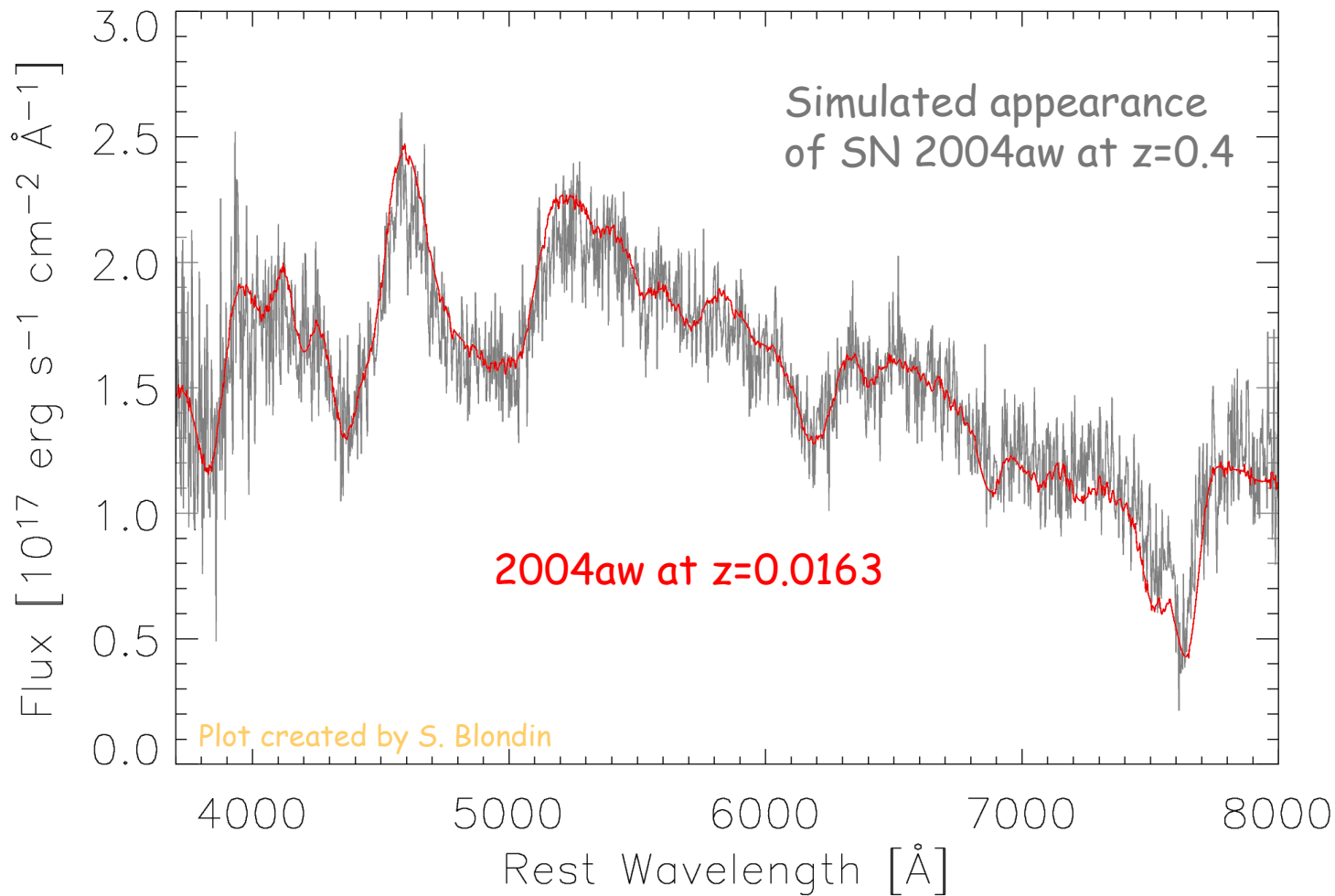
Light curves of SN 2004aw



Absolute magnitudes

- For 2004aw extinction and distance highly uncertain:
 - $A(B) = 1.80$ mag and $A(V) = 1.36$ mag from EW measurements
 - $\forall \mu = 34.23$ from host galaxy recession velocity
- Results for maximum brightness:
- $B = -17.95 \pm 0.47$ mag , $V = -18.26 \pm 0.39$ mag, ...
(Errors almost arbitrary !)
- For comparison:
 - peak luminosities of normal type Ia SNe
 $B = -19.4$ mag and $V = -19.4$ mag with scatter (± 0.3 mag)
 - For type Ic SNe:
 $B = -16.2 \dots -19.4$ mag and $V = -16.7 \dots -19.8$ mag

High-z sample contamination ?



+1.5 days
spectrum

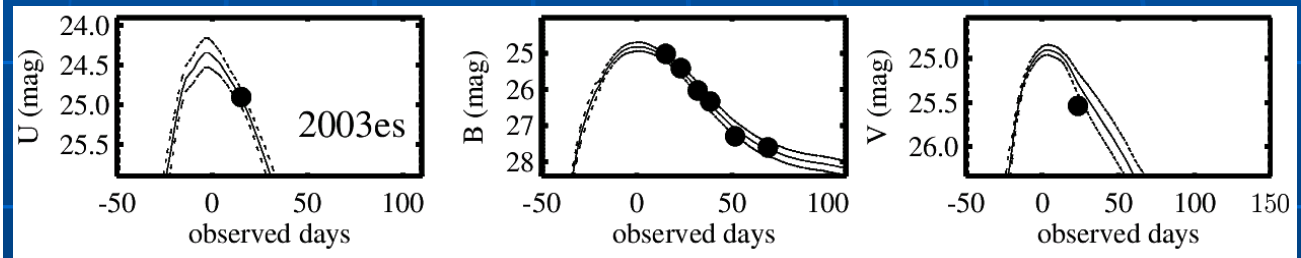
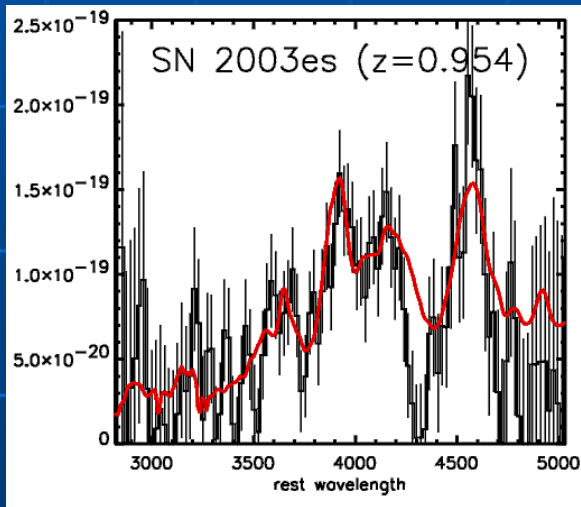
(S. Taubenberger)

Result of a classification code of S. Blondin (for the $z = 0.4$ case):

- best match: 1991T (Ia pec) @ +17.4d
- second best: 1992A (Ia) @ +9.0d
- third best: 1995D (Ia) @ +8.1d

What does this mean?

- Assume a hypothetical 2004aw at $z = 0.4$ and UNREDDENED (otherwise too faint)
- Typical dataset obtained for such SNe: one spectrum + sparse photometry in 2 or 3 filters



Figures from Riess et al. 2004

- 2004aw would be classified as Ia (pec) according to the spectrum
- B-V color determined, result: 0.28 mag (instead of -0.05 mag for a typical type Ia)

- Mis-interpreted as extinction, correction for $E(B-V) = 0.33$ applied
- This brings 2004aw to absolute magnitudes of **-19.3** in B and V
- Similar to type Ia SNe, aligns rather well in the Hubble diagram
But: Only by chance !
- Of course the procedure is much more complex and sophisticated in reality.
- Nevertheless: the danger of contamination remains!

Summary (Part III)

- Type II supernovae are good distance indicators out to a few Mpc.
- They measure absolute distances without any calibration!
- Type Ia supernovae are very good distance indicators in the local Universe .
- They allow to measure relative distances very accurately (after calibration).
- They provide the best distance indicators for cosmological distances if systematic errors can be controlled.