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Recent observations of the binary system LS 5039 with the High Energy Stereoscopic System (H.E.S.S.) revealed that its Very High Energy (VHE) γ -ray emission is modulated at the 3.9 days orbital period of the system. The bulk of the emission is largely confined to half of the orbit, peaking around the inferior conjunction epoch of the compact object. The flux modulation provides the first indication of γ -ray absorption by pair production on the intense stellar photon field. This implies that the production region size must be not significantly greater than the gamma-gamma photosphere size (~ 1 AU), thus excluding the large scale collimated outflows or jets (extending out to ~ 1000 AU). A hardening of the spectrum is also observed at the same epoch between 0.2 and a few TeV which is unexpected under a pure absorption scenario and could rather arise from variation with phase in the maximum electron energy and/or the dominant VHE γ -ray production mechanism. This first-time observation of modulated γ -ray emission allows precise tests of the acceleration and emission models in binary systems.

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Figure 1: Orbital geometry of the binary system LS 5039 viewed from above and using the orbital parameters derived by Casares et al.[8]. Shown are: phases (ϕ) of minimum (*periastron*) and maximum (*apastron*) binary separation; epoch of superior and inferior conjunctions occurring when the compact object and the star are aligned along the observer light-of-sight.

Figure 2: H.E.S.S. excess sky map around LS 5039, smoothed by the instrument point spread function. The blue star denotes the position of the VLBA source. The yellow contours correspond to the 68%, 95% and 99% confidence level region of the EGRET source 3EG J1824-1514. The extended source HESS J1825-137 observed in the same field of view can serve as a cross-check for timing-analysis.

1. Introduction

In the commonly accepted paradigm, microquasars consist of a stellar mass black hole fed by a massive star. They can exhibit superluminous radio jets[15], and hints for the presence of an accretion disk. LS 5039, identified in 1997[16] as a massive X-ray binary system with faint radio emission[14], was resolved by Paredes et al.[17] into bipolar mildly relativistic radio jets ($v \sim 0.2 c$) emanating from a central core, thus placing it into the *microquasar class*. The detection of radio and variable X-ray emission[7] and its possible association with the EGRET source 3EG J1824-1514 suggested the presence of multi-GeV particles accelerated in jets. This binary system (Fig 1) consists of a massive O6.5V star in a ~ 3.9 day mildly eccentric orbit (e = 0.35)[8] around a compact object whose exact nature (black hole or neutron star) is still under debate.

2. H.E.S.S. Observations

The High Energy Stereoscopic System (H.E.S.S.) is an array of four identical Atmospheric Cherenkov Telescopes (ACT)[4] located in the Southern Hemisphere (Namibia, 1800 m a.s.l.) and sensitive to γ rays above 100 GeV. LS 5039 was serendipitously detected in 2004 during the H.E.S.S. galactic scan[2] The 2004 observations have been followed up by a deeper observation campaign[3] in 2005, leading to a total dataset of 69.2 hours of observation after data quality selection. Data were analysed using two separate calibrations[1] and analysis pipelines. The results

presented here are based on the log-likelihood comparison of the shower images with a precalculated semi-analytical model[9].

After selection cuts, a total of 1969 γ -ray events were found within 0.1° of the VLBA radio position of LS 5039, leading to a statistical significance of 40 σ (Fig. 2). The best fit position is, in Galactic Coordinates, $l = 16.879^\circ$, $b = -1.285^\circ$ with statistical and systematic uncertainties of respectively $\pm 12''$ and $\pm 20''$. It is compatible with the VLBA position (denoted as a blue star in Fig. 2). We obtain an upper limit of 28'' (at 1 σ) on the VHE source extension.

2.1 Timing Analysis

The runwise VHE γ -ray flux at energies ≥ 1 TeV was decomposed into its frequency components using the Lomb-Scargle periodogram[20] (Fig. 3) which is appropriate for unevenly sampled datasets such as those collected by H.E.S.S. A very significant peak (chance probability of $\sim 10^{-20}$ before trials) occurs in the Lomb-Scargle periodogram at the period 3.9078 ± 0.0015 days, consistent with the most recent optical orbital period[8] (3.90603 ± 0.00017). The effect of subtracting a pure sinusoid at the orbital period is shown in Fig. 3, middle panel. The orbital peak disappears as expected, but also the numerous satellite peaks with chance probabilities less than 10^{-7} - 10^{-8} that were present in the original periodogram. These peaks are beat periods of the orbital period). The bottom panel of the same figure shows the result obtained on the neighbouring source HESS J1825-137 observed in the same field of view as LS 5039, which doesn't show any statistically significant peak, thus demonstrating that the observed periodicity is genuinely associated with LS 5039.

2.2 Flux Modulation

The runwise Phasogram (Fig 4) of integral flux at energies ≥ 1 TeV vs. orbital phase (ϕ) shows an almost sinusoidal behaviour, with the bulk of the emission largely confined in a phase interval $\phi \sim 0.45$ to 0.9, covering about half of the orbital period. The emission maximum ($\phi \sim 0.7$) appear to lag behind the apastron epoch and to align better with the *inferior conjunction* ($\phi = 0.716$), when the compact object lies in front of the massive star (see Fig. 1). The VHE flux minimum occurs at phase ($\phi \sim 0.2$), slightly further along the orbit than *superior conjunction* ($\phi = 0.058$). Neither evidence for long-term secular variations in the VHE flux independent of the orbital modulation nor any other modulation period are found in the presented H.E.S.S. data.

2.3 Spectral Modulation

Due to changing environment with orbital phase (magnetic field strength, stellar photon field, relative position of compact object and star with respect to observer, ...), the VHE γ -ray emission spectrum is expected to vary along the orbit. We first define two broad phase interval: **INFC** centered on the inferior conjunction (0.45 < $\phi \le 0.9$) and its complementary **SUPC** centered on the superior conjunction, corresponding respectively to high and low flux states. The high state VHE spectral energy distribution (Fig 5) is consistent with a hard power law with index $\Gamma = 1.85 \pm 0.06_{\text{stat}} \pm 0.1_{\text{syst}}$ and exponential cutoff at $E_0 = 8.7 \pm 2.0$ TeV. In contrast, the spectrum for low state is compatible with a relatively steep ($\Gamma = 2.53 \pm 0.06_{\text{stat}} \pm 0.1_{\text{syst}}$) pure power law extending from 200 GeV to ~ 20 TeV. Interestingly, the flux appears to be almost unmodulated at 200 GeV as well as around 20 TeV, whereas the modulation is maximum around a few (~ 5) TeV.





Figure 3: Lomb-Scargle (LS) periodogram of the VHE runwise flux of LS 5039 above 1 TeV (Chance probability to obtain the LS power vs. frequency). From [3]. Zoom: inset around the highest peak, which corresponds to a period of 3.9078 ± 0.0015 days. This period is found to be compatible with the optical orbital period[8] denoted as a red line on the inset. Middle: LS periodogram of the same data after subtraction of a pure sinusoidal component at the orbital period of 3.90603 days (see text). Bottom: LS periodogram obtained on HESS J1825-137 observed in the same field of view.



Figure 4: Phasogram (Integral run-by-run γ -ray flux above 1 TeV as function of orbital phase) of LS 5039 from H.E.S.S. data from 2004 to 2005, using the orbital ephemeris[8]. Each run is ~ 28 minutes. Two full phase periods are shown for clarity. The vertical blue arrows depict the respective phases of minimum (*periastron*) and maximum (*apastron*) binary separation. The vertical dashed red lines show the respective phases of inferior and superior conjunction, when the star and the compact object are aligned along the observer's line of sight. From [3].

Trying to go to smaller phase interval, Fig 6 shows the results (photon index and differential flux at 1 TeV) of a pure power-law fit of the high energy spectra in 0.1 orbital phase bins (restricted to energies below 5 TeV to avoid systematic effect introduced by the high state cutoff). The flux normalisation and photon index are strongly correlated, the flux being higher when the spectrum is harder and vice-versa. Interestingly, a similar effect, however in a smaller variation range and a different phasogram, was found in X rays[7].

3. Interpretation and Conclusion

The basic paradigm of VHE γ -ray production requires the presence of particles accelerated to multi-TeV energies and a target comprising photons (for γ -ray production through Inverse Compton effect) and/or matter of sufficient density (for γ -ray production through pion decay in hadronic processes). Several model classes are available to explain VHE emission from microquasars, differencing one from the other by the nature of accelerated particles and/or the location of the acceleration region. In jet-based models, particle acceleration could take place directly inside and along the jet, e.g. [6, and references therein], and also in the jet termination shock regions[12]. Non-jet scenarios are also available, e.g. [13, 11], where the emission arises from the interaction of a pulsar wind with the stellar companion equatorial wind.

New observations by HESS have established orbital modulation of the VHE γ -ray flux and energy spectrum from the XRB LS 5039. The observed VHE modulation indicates that the emission most probably takes place close (within $\sim 1 \text{ AU}$) to the massive stellar companion, where modu-



Figure 5: Very high energy spectral energy distribution of LS 5039 for the two broad orbital phase intervals defines in the text, **INFC** (red circles) and **SUPC** (blue triangles). The shades regions represent the 1σ confidence bands on the fitted functions. A clear spectral hardening is occurring in the 200 GeV to a few TeV range during the **INFC** phase interval. From [3].



Figure 6: Top: Fitted pure power-law photon index vs. phase interval of width $\Delta \phi = 0.1$. Bottom: Differential flux at 1 TeV for the same phase interval. From [3].

lated γ -ray absorption via pair production (e^+e^-) on the intense stellar photon field is unavoidable (e.g. [11]) The observed spectral modulation is however incompatible with a pure absorption scenario, which in particular predicts a maximum variability around 300 GeV and a VHE spectral hardening in the low flux state, inconsistent with observations.

Modulation could also arise from a modulation of the acceleration and cooling timescales along the orbit due to varying magnetic field and photon field densities (e.g. [3, and references therein]) which could modify the maximum electron energy and therefore induce a phase-dependent energy break in the γ -ray spectrum. Modulation of the accretion rate due to interaction of the stellar wind with the compact object in the microquasar scenario (e.g. [18]) could be another ingredient of the observed modulation.

A detailed study is now required to fully explain these new observations and understand the complex relationship between γ -ray absorption and production processes within these binary systems.

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