

## Half-life measurements of neutron-deficient isotopes using laser Compton scattering Gamma-rays at NewSUBARU

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We measure half-lives of the ground state of  $^{184}\text{Re}$  and an isomer in  $^{164}\text{Ho}$ , which are populated by  $(\gamma, n)$  reactions with laser Compton scattering (LCS)  $\gamma$ -ray source at NewSUBARU. These neutron-deficient isotopes are located on nucleosynthesis flows of the supernova  $\gamma$  process. The measured half-life of  $35.4 \pm 0.3$  d for  $^{184}\text{Re}$  is shorter than the previous half-life of  $T_{1/2} = 38.0 \pm 0.5$  by about 7%. The half-life of the  $^{164}\text{Ho}$  isomer is  $36.4 \pm 0.3$  min. This is about 3% shorter than a recommended value  $T_{1/2} = 37.5^{+1.5}_{-0.5}$  min. These results indicate that half-lives of all unstable nuclei near the  $\beta$  stability line may be not robust and that the LCS  $\gamma$ -rays are useful for precise determination of the half-lives.

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## 1. Introduction

The ( $\gamma$ ,n) reactions on neutron-deficient isotopes are important for understanding nucleosynthesis by photodisintegration reactions in supernova explosions ( $\gamma$  process or  $p$  process) [1–7]. For this reason, the ( $\gamma$ ,n) reaction cross-sections were measured in the rare-earth region [8–10]. An activation method has been used for measurements of nuclear reaction cross sections in these studies [8,9,11]. The half-life  $T_{1/2}$  of the populated nucleus is crucial for the activation method since the evaluated cross-section is proportional to  $T_{1/2}$  of the populated nucleus.

The progress of the relativistic engineering (for example see Ref. [12]) provides a new  $\gamma$ -ray source with a MeV energy range. These  $\gamma$ -rays are generated by Compton scattering of relativistic electrons by laser photons. These laser Compton scattering  $\gamma$ -rays (LCS  $\gamma$ -rays) have advantages that the maximum energy is sharply determined and that the  $\gamma$ -ray flux with high energy is relatively high. The Duke Free Electron Laser Laboratory at Duke University [14] and the National Institute of Advanced Industrial Science and Technology [13] have provided the LCS  $\gamma$ -rays in the MeV energy range and they have been widely used for applications with photon-induced reactions. Recently, a new LCS  $\gamma$ -ray source was installed at an electron storage ring NewSUBARU in SPring-8[15,16]. We measure half-lives of  $^{184}\text{Re}$ [17] and  $^{164}\text{Ho}$  isomer[18] populated by photodisintegration reactions with the LCS  $\gamma$ -rays at NewSUBARU.

## 2. Laser Compton scattering $\gamma$ -ray source at NewSUBARU

A Q-switch Nd:YVO<sub>4</sub> laser system and a nuclear experiment room with a heavy shield locate at the outside of the electron storage ring NewSUBARU as shown in Fig. 1. The collision of the relativistic electrons and the laser photons creates a high energy  $\gamma$ -ray, whose energy depends on an angle between the direction of the incident electrons and the generated  $\gamma$ -rays. The energy distribution of the LCS  $\gamma$ -rays is determined in the basic QED process. The diameter of the LCS  $\gamma$ -rays is about 20 mm without a collimator at the target position, which is located at about 20 m from the collision point. The electron storage ring NewSUBARU can store electrons with an energy of 978 MeV up to 230 mA in a top-up mode. The 198 electron bunches circulate in the storage ring with a frequency of 2.5 MHz. An interval time of the electron bunches is about 2 ns. The Nd:YVO<sub>4</sub> laser system provide laser photons with a wavelength of 1064 nm at 100 kHz. A single laser pulse with a pulse length of 10 ns has a chance to collide with four or five electron bunches in the collision region. The laser power is typically 4 W and the estimated  $\gamma$ -ray flux is  $0.5\text{--}1.5 \times 10^6$  photons/s with an energy range from 3.3 MeV to 16.7 MeV. This maximum energy of the LCS  $\gamma$ -rays is higher than the peak energy of the giant dipole resonance (GDR), and thus neutron-deficient isotope is effectively populated by the GDR excitation.

## 3. Measurements of half-lives of $^{164}\text{Ho}$ isomer and $^{184}\text{Re}$

In an experiment of Re, we used three natural Re metallic foils. The individual Re foil had a thickness of 0.2 mm and a size of 25 mm  $\times$  25mm. The three stacked natural Re metallic foils were

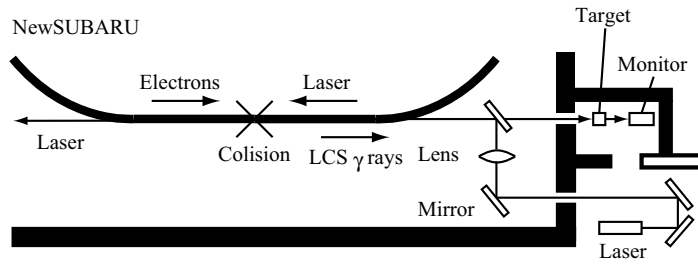


Figure 1: The laser Compton scattering  $\gamma$ -ray source at NewSUBARU

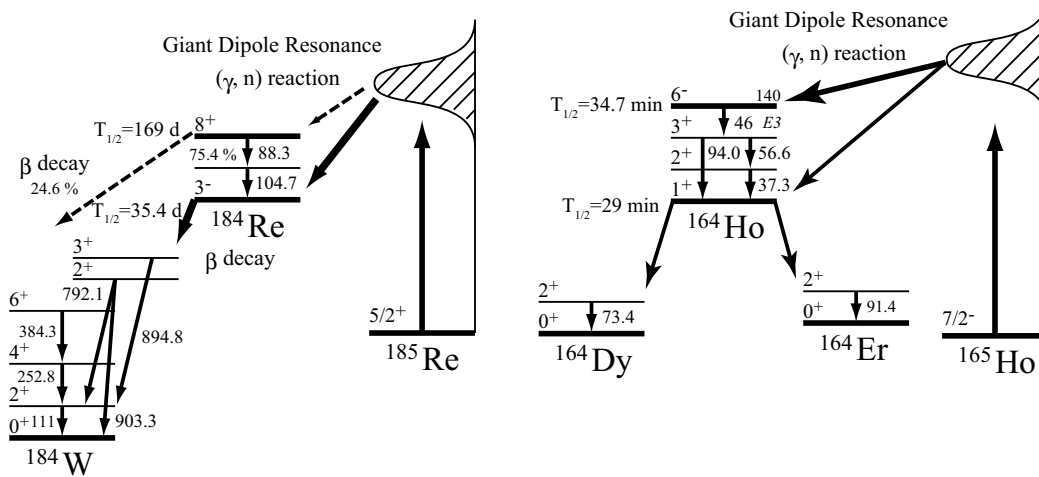
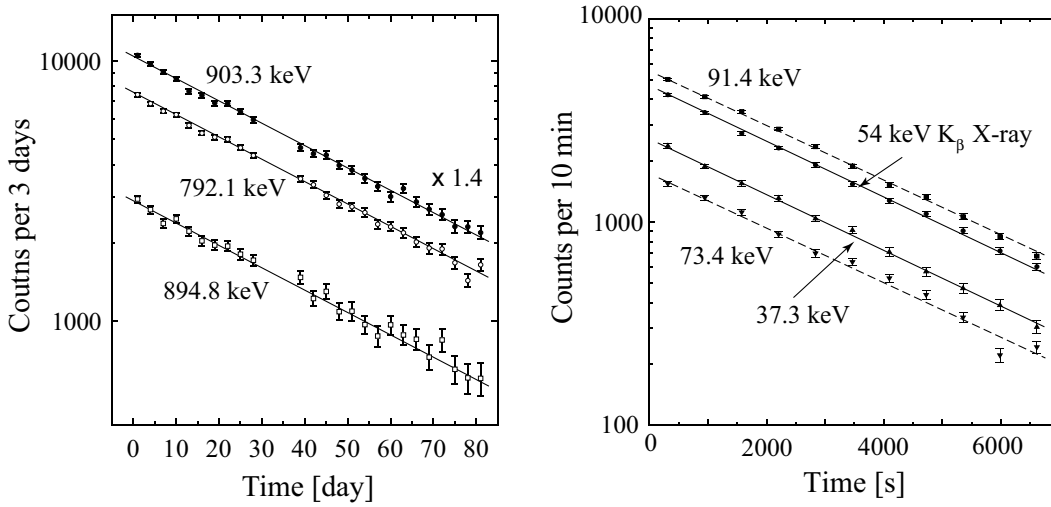


Figure 2: Nuclear reactions and decay scheme of  $^{184}\text{Re}$  (left) and  $^{164}\text{Ho}$  isomer (right).

irradiated by the LCS  $\gamma$ -rays for about 9 hours. The irradiated targets were cooled for a period of 23 days to reduce the background from short-lived radioactivities such as  $^{186}\text{Re}$  ( $T_{1/2} = 90.64$  h) and to obtain the stability of the electronics system. To evaluate the half-life of  $^{184}\text{Re}$ , time dependence of  $\gamma$ -ray intensities from the activities was measured for a period of 83 days. The  $\gamma$ -rays emitted after the  $\beta$  decay were measured by a HPGe detector with lead shields. The three Re foils were located on a plain in the front of the HPGe detector. The efficiency of the HPGe detector was larger than 70% relative to a  $3'' \times 3''$  NaI detector and its energy resolution was 2.1 keV at 1.3 MeV. The measurement system was almost stable and the peaks of the  $\gamma$ -rays shifted by only one or two channel/s relative to about 3000 channels during the measurement of 83 days.

In an experiment of Ho, twenty stacked metallic Ho foils were irradiated by the LCS  $\gamma$ -rays for 41 min and subsequently the targets were moved to a measurement position. The size of each Ho foil was  $10 \text{ mm} \times 10 \text{ mm} \times 1 \text{ mm}$ . After 8 min from the LCS  $\gamma$ -ray irradiation, decay  $\gamma$ -rays from the foils were measured by a HPGe detector for 110 min. The twenty Ho targets were located on a plain in the front of the HPGe detector with lead shields. The efficiency of the HPGe detector was about 45%.

Three  $\gamma$ -rays of 792.1, 894.8 and 903.3 keV of  $^{184}\text{Re}$  are clearly observed. The half-life was



**Figure 3:** Measured decay curves of  $\gamma$ -rays of  $^{184}\text{Re}$  (left) and  $^{164}\text{Ho}$  isomer (right).

evaluated from the decay curves of these three  $\gamma$ -rays. Since the  $\gamma$ -rays following the decay of the isomer on  $^{184}\text{Re}$  were not observed the ground state of  $^{184}\text{Re}$  is dominantly populated in this reaction. The decay curves are well fitted by a straight line as shown in Fig. 3. The individual spectrum is recorded for a period of three days. We obtain the half-life of the individual  $\gamma$ -ray by using  $\chi$ -square fitting and the results are  $35.1 \pm 0.5$ ,  $36.0 \pm 0.9$  and  $35.6 \pm 0.5$  d for 792.1, 894.8 and 903.3 keV  $\gamma$ -rays, respectively. These three half-lives are identical within the uncertainty. Finally we obtain the half-life of  $35.4 \pm 0.3$  d as the average value of these three  $\gamma$ -rays.

Historically, the measurement of the half-life,  $38 \pm 1$  d, was reported in 1960 [19]. The most precise half-life,  $38.0 \pm 0.5$  d, was reported in 1962 [20] and this was widely taken as the recommended value. In these two studies, Re activities were prepared by using deuteron-induced reactions. After these studies, the isomer with a half-life of 169 d was found by a measurement of decay of activities populated by the neutron-induced reactions in a nuclear reactor. Therefore radioactive samples in the two historical studies [19,20] may include the isomer, but the effect of the isomer was not taken into account.

We obtain the decay curves of the 37.3-keV  $\gamma$ -ray and 54-keV  $K_\beta$  X rays of Ho as shown in Fig. 3. The half-lives obtained by  $\chi$  square fitting are  $36.1 \pm 0.4$  min and  $36.6 \pm 0.4$  min for 37 keV  $\gamma$  ray and 54 keV X ray, respectively. We take the average value of  $36.4 \pm 0.3$  min as the half-life of the  $^{164}\text{Ho}$  isomer in the present experiment. This half-life is about 3% shorter than the previous value of  $37.5^{+1.5}_{-0.5}$  min, which was measured by a NaI(Tl) detector in 1966 [22]. Note that the decay curves of 91.4 keV and 73.4 keV from the decay of  $^{164}\text{Ho}$  must be composed of feeding from the  $^{164}\text{Ho}$  isomer and  $\beta$ -decay.

#### 4. Conclusion

We report half-lives of the  $^{184}\text{Re}$  ground state and  $^{164}\text{Ho}$  isomer, which are populated via photodisintegration reactions with laser Compton scattering (LCS)  $\gamma$ -rays at electron storage ring

NewSUBARU. The ground state of  $^{184}\text{Re}$  is dominantly populated in this reaction. The measured half-life is  $35.4 \pm 0.3$  d. This is about 7% shorter than a recommended value  $T_{1/2} = 38.0 \pm 0.5$ , which was reported in 1962 before a discovery of an isomer with  $J^\pi = 8^+$  in  $^{184}\text{Re}$ . Our result provides essential information for applications using an activation method because the cross-section should be smaller by about 7% than that based on the previous value. The measured half-life of the  $^{164}\text{Ho}$  isomer is  $36.4 \pm 0.3$  min. This is about 3% shorter than a recommended value  $T_{1/2} = 37.5^{+1.5}_{-0.5}$  min, which was measured by a NaI(Tl) detector in 1966. These experiments indicate that measured half-lives of all unstable nuclei near the  $\beta$  stability line may be not robust and that the LCS  $\gamma$  rays are useful for a precise measurement even if a high spin isomer exists in a nucleus of interest.

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