New physics opportunities in the charm/tau region: the BESIII - experiment in Beijing

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A new facility for precision physics in the Charm/Tau energy regime has become operational at the Institute for High Energy Physics (IHEP) in Beijing, China. The dual ring $e^+e^−$ collider BEPCII is operating at luminosities up to $10^{33} \, cm^{-2}s^{-1}$ and at a centre-of-mass energy between 2 GeV and 4.6 GeV. BESIII, a new state-of-the-art detector has been build and successfully commissioned. This contribution describes the facility, presents a brief overview of the physics program and shows preliminary results from a first run in spring 2009 which collected a sample of $156 \, pb^{-1}$ at the $\psi(2S)$ resonance.
1. Introduction

Electron-positron colliders have been at the forefront of research in hadron and particle physics for many years. Depending on the CM energy of the collider, different regions of interest can be accessed. The Phi factory at Frascati, the BEPCII collider at IHEP Beijing and the KEKb collider at Tsukuba are examples of contemporary colliders focusing on Phi meson and Kaon production, Charm/Tau physics and B physics, respectively. Due to the versatility of such machines, there is a very long history of $e^+e^-$ colliders.

The new BEPCII $e^+e^-$ collider at IHEP Beijing became recently operational, operating already at record luminosities. A new state-of-the-art experiment, the BESIII detector has been built, successfully commissioned and is ready to address a wide range of physics topics in hadron physics and particle physics.

2. The BEPCII $e^+e^-$ collider

BEPCII is a double-ring $e^+e^-$ collider designed for a peak luminosity of $10^{33} \cdot cm^{-2}s^{-1}$ at a beam current of 0.93 A. The luminosity of BEPCII is a factor of 100 larger compared to its predecessor, the BEPC collider. BEPCII consists of two storage rings with 237.5 m circumference. Each of the two rings stores 93 bunches with a bunch length of 1.5 cm. At the interaction point the beams are colliding with a horizontal crossing angle of 11 mrad and a bunch spacing of 8 ns.

The energy per beam ranges between 1.0 GeV and 2.3 GeV, with the optimum beam energy at 1.89 GeV. The energy spread amounts to $5 \cdot 10^{-4}$ GeV. In Summer 2008, the BEPC2 accelerator and the BESIII experiment at IHEP Beijing were successfully commissioned. A first run took place in November 2008 collecting about 14 Million $\psi(2S)$ events. In March and April 2009, a sample of $156 pb^{-1}$ corresponding to 106M $\psi(2S)$ events was collected with BEPCII operating at about 30 % of the peak luminosity.

3. The BESIII detector

The BESIII detector [1] features a large geometrical acceptance of 93% of $4\pi$. Charged particle tracking is performed with a 43-layer mini drift chamber (MDC) system with 6796 signal and 21884 field wires. The gas mixture of He (40%) and $C_3H_8$ (60%) has been optimized for position resolution and large radiation length. The average single wire resolution is 135 $\mu$m. The chambers operate in a solenoidal magnetic field of 1 Tesla which is produced by a superconducting magnet. The momentum resolution for charged particles amounts to 0.5% at 1 GeV/c.

The high performance electromagnetic calorimeter (EMC) is divided into a barrel section and two end cap sections with a total of 6240 CsI(Tl) crystals. At 1 GeV, the energy resolutions of 2.5% and 5% in the barrel and the end caps, respectively, are achieved. A position resolution of 6 mm and 9 mm for the barrel and the end caps, respectively, is obtained.

A time-of-flight system (TOF) constructed of 5 cm thick plastic scintillation detectors serves for particle identification. The barrel portion employs 176 detectors with a length of 2.4 m arranged in two layers. The endcaps are covered by 96 fan-shaped scintillators. The time resolution amounts
Figure 1: The BESIII detector [1], featuring tracking with Mini Drift Chambers (MDC), an electromagnetic calorimeter (EMC) consisting of 6240 CsI(Tl), a time-of-flight system for particle identification and an RPC based muon detection system.

Figure 2: The first event recorded with BESIII on July 19, 2008.

to 80 ps for the barrel and 110 ps for the endcaps, respectively. This provides a $2\sigma$ $K/\pi$ separation up to about 1 GeV/c.

The muon system (MUC) consists of 1000 $m^2$ of Resistive Plate Chambers (RPCs) with a position resolution of about 2 cm, comprising nine barrel and eight endcap layers.

BESIII is a state-of-the-art detector which is in many respects similar to the CLEOc detector. A schematic view of the detector with its main components is shown in Fig. 1. Fig. 2 shows the first event which was recorded with BESIII on July 19, 2008.
4. Physics program

BESIII features a comprehensive physics program benefiting from the unprecedented luminosity of BEPCII and the high performance BESIII detector which is able to measure charged particles as well as photons over the relevant momentum range with excellent resolution and particle identification capabilities.

The energy regime of BESIII is of particular interest in that it allows precision studies of numerous narrow resonances in the charmonium region. Going beyond improved measurements of known branching fractions, near threshold the possibility of tagging permits almost background-free studies of rare channels. Moreover, many of the recently discovered X,Y,Z - states are accessible at BESIII and could be studied with improved precision.

Measurements of the total cross section for $e^+ e^-$ annihilation into hadrons are indispensable input for the determination of the non-perturbative hadronic contribution to the running QED fine structure constant and an essential input parameter in precision electroweak measurements.

Furthermore, the charm region presents a challenge for lattice QCD (LQCD) calculations. Measured properties such as D meson decay constants or transition form factors can be compared to results of LQCD, thereby probing the precision of such calculations. Such studies are crucial to cut down errors on hadronic observables for precision CKM physics.

Since charmonium decays are a rich source of gluons, there is a significant discovery potential for QCD exotica such as glueballs, multi-quark states and hybrids. Furthermore, meson and baryon spectroscopy will benefit from the high luminosity and high quality detection including neutral particles.

Improved precision for the $\tau$ lepton mass can be achieved by threshold scans employing a laser backscattering system for precision energy measurement. A comprehensive description of the physics program can be found in the BESIII physics book [2].

5. First results

The data which will be presented in the following have been collected in a run where 106M $\psi(2S)$ events have been recorded. It should be noted that all results given below are preliminary.

5.1 $\chi_c$ states

Fig. 3 shows part of the charmonium level scheme below the open charm threshold. The arrows mark radiative decays between the $\psi(2S)$ state and various $\chi_c$ states as well as radiative decays of $\chi_c$ states to $J/\psi$. Fig. 4 shows the inclusive gamma ray spectrum observed at the $\psi(2S)$ resonance. E1-transitions corresponding to the decays $\psi(2S) \rightarrow \chi_{c2}(P_2), \chi_{c1}(1^3P_1), \chi_{c0}(1^3P_0)$ are observed. Furthermore, the decays $\chi_{c2}(1^3P_2), \chi_{c1}(1^3P_1) \rightarrow \gamma J/\psi$ are seen.

5.2 Observation of the $h_c$

The spectrum of charmonium states can be quite well described with potential models assumes a one gluon exchange Coulomb-like term and a linear confining term. In a more refined models, spin-dependent terms have to be considered. Information about the spin-dependent interaction can be obtained from a precision measurement of the 1P hyperfine mass splitting. Despite many
extensive studies of the charmonium system, the properties of the $c\bar{c}$ singlet state $h_c(11P_1)$ state are not well known. CLEOc has measured the mass of the $h_c$ [3] both in the exclusive process $\psi(2S) \rightarrow \pi^0 h_c \rightarrow \pi^0\gamma\eta_c$ as well as in the inclusive process $\psi' \rightarrow \pi^0 h_c$ where the $h_c$ shows up as a peak in the $\pi^0$ recoil mass distribution. E835 has observed [4] the decay $h_c \rightarrow \gamma\eta_c$.

With BESIII, we have studied $h_c$ production and decay at the $\psi(2S)$ resonance, both in the recoil mass distribution of inclusive $\pi^0$ production and in $\gamma$-tagged events from the process $\psi' \rightarrow \pi^0 h_c \rightarrow \pi^0\gamma\eta_c$. Fig. 5 shows the $\pi^0$ recoil mass distribution of $E1 - \gamma$-tagged events of the type $\psi(2S) \rightarrow \pi^0 h_c \rightarrow \pi^0\gamma\eta_c$. A clear $h_c$ signal is observed with the mass of $(3525.16 \pm 0.16 \pm 0.10)$ MeV and the width of $(0.89 \pm 0.57 \pm 0.23)$ MeV can be extracted. The latter has been measured for the first time.

Combining the inclusive measurement shown in Fig. 6 with the $\gamma$-tagged measurement, the absolute branching ratios of $\psi' \rightarrow \pi^0 h_c$ and $h_c \rightarrow \gamma\eta_c$ have been determined for the first time as $(8.42 \pm 1.29(stat)) \times 10^{-4}$ and $(55.7 \pm 6.3(stat))%$. As noted above, these results are preliminary and systematic errors are currently under study.

5.3 $\chi_c$ decays into $\pi^0\pi^0$ and $\eta\eta$ final states

Since the electromagnetic interaction conserves C-parity, $\chi_c$ states cannot be directly formed in
Figure 5: $\pi^0$ recoil mass distribution of $E1-\gamma$-tagged events of the type $\psi(2S) \rightarrow \pi^0h_c \rightarrow \pi^0\eta_c$ (left) and with subtracted background (right). A clear $h_c$ signal is observed.

Figure 6: $\pi^0$ recoil mass distribution of inclusive events of the type $\psi(2S) \rightarrow \pi^0X$ (left) and with subtracted background (right). A clear $h_c$ signal is observed, albeit with larger background when compared to the $E1-\gamma$-tagged data.

e^+e^-$ annihilation. However, they can easily be studied as products of radiative $\psi(2S)$ decays (see Fig. 4) in a nearly background-free environment. Among the various hadronic decays of $\chi_c$ states, channels with two neutral pseudoscalar mesons are the most accessible, both from the experimental point of view [5] as well as in model descriptions [6].

Fig. 7 shows the radiative photon spectrum for events with two reconstructed $\pi^0$ mesons (left) and two reconstructed $\eta$ mesons (right). The two observed peaks correspond to the processes $\psi(2S) \rightarrow \gamma\chi_{c2}$ and $\psi(2S) \rightarrow \gamma\chi_{c0}$.

The extracted branching fractions are shown in Tab. 1 and compared to previous results from CLEOc [5] and PDG values [7]. Again, systematic errors are under study.

<table>
<thead>
<tr>
<th>Process</th>
<th>BR ($10^{-3}$)</th>
<th>$\chi_{c0}$</th>
<th>$\chi_{c2}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\pi^0\pi^0$ BESIII</td>
<td>3.25±0.03(stat)</td>
<td>0.86±0.02(stat)</td>
<td></td>
</tr>
<tr>
<td>PDG</td>
<td>2.43±0.20</td>
<td>0.71±0.08</td>
<td></td>
</tr>
<tr>
<td>CLEO-c</td>
<td>2.94±0.07±0.35</td>
<td>0.68±0.03±0.08</td>
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</tr>
<tr>
<td>$\eta\eta$ BESIII</td>
<td>3.11±0.1(stat)</td>
<td>0.59±0.05(stat)</td>
<td></td>
</tr>
<tr>
<td>PDG</td>
<td>2.4±0.4</td>
<td>&lt;0.5</td>
<td></td>
</tr>
<tr>
<td>CLEO-c</td>
<td>3.18±0.13±0.35</td>
<td>0.51±0.05±0.06</td>
<td></td>
</tr>
</tbody>
</table>

Table 1: Extracted branching fractions (preliminary) and comparison to previous results from CLEOc [5] and PDG values [7]. Systematic errors are currently under study.
Figure 7: Radiative photon spectrum for events with two reconstructed $\pi^0$ mesons (left) and two reconstructed $\eta$ mesons (right). The two observed peaks correspond to radiative decays of $\psi(2S)$ into $\chi_{c2}$ (left peak) and $\chi_{c0}$ (right peak). The background has been studied with Monte Carlo simulations including potential background channels taken from the PDG [7].

5.4 Confirmation of the threshold enhancement in $J/\psi \to \gamma p\bar{p}$ observed by BESII

BESII has reported evidence for a near-threshold enhancement in $J/\psi \to \gamma p\bar{p}$ [8]. A possible explanation involves the existence of a narrow resonances $X(1860)$ below threshold corresponding to a $p\bar{p}$ bound state. There are no well established mesons that could be associated with such a state. With BESIII, we have confirmed the signal seen by BES2. Fig. 8 presents the invariant mass distribution of $p\bar{p}$ pairs, shifted by twice the proton mass. The figure’s left part shows the BES2 data. On the right hand side, Fig. 8 presents new data for the decay $\psi(2S) \to \pi\pi J/\psi$, $J/\psi \to \gamma p\bar{p}$, measured with BESIII. The previously observed near-threshold enhancement is clearly confirmed.

It is interesting to study other radiative decays such as $\psi(2S) \to \gamma p\bar{p}$. Fig. 9 presents the previously published BESII data [8] (left) and new data taken with BESIII, showing that both experiments consistently observe no near-threshold enhancement, in striking contrast to the result.
Figure 9: Absence of any near-threshold enhancement in the $p\bar{p}$ system observed in radiative decays of $\psi(2S) \to \gamma p\bar{p}$. The left figure shows data taken by the BESII collaboration, confirming the BESIII data.

obtained for the decay $J/\psi \to \gamma p\bar{p}$.

It should be noted that no enhancement is seen in radiative upsilon decays \cite{ref9} measured by CLEO. In order to clarify the origin of the enhancement observed in $\psi(2S) \to \gamma p\bar{p}$, further studies are required.

6. Conclusions

With the BEPC2 collider and the BESIII experiment, a new state-of-the art facility for physics in the charm/tau region has become operational. Record luminosities and high quality detection systems for neutral and charged particles will provide new high statistics data improving the experimental knowledge in a multitude of fields such as light hadron spectroscopy, charmonium spectroscopy, open charm physics and the search for rare and forbidden decays.

7. Acknowledgements

The author gratefully acknowledges the contributions of his colleagues in the BESIII collaboration. He thanks the staff of BEPC2 and the IHEP computing center for their enthusiastic support.

References


