

Phase diagram and critical point evolution in NLO and NNLO strong coupling lattice QCD

Akira Ohnishi*

Yukawa Institute for Theoretical Physics, Kyoto University, Kyoto 606-8502, Japan
E-mail: ohnishi@yukawa.kyoto-u.ac.jp

Kohtaroh Miura

Yukawa Institute for Theoretical Physics, Kyoto University, Kyoto 606-8502, Japan

Takashi Z. Nakano

Department of Physics, Faculty of Science, Kyoto University, Kyoto 606-8502, Japan

Noboru Kawamoto

Department of Physics, Faculty of Science, Hokkaido University, Sapporo 060-0810, Japan

We investigate the chiral phase transition in the strong coupling lattice QCD (SC-LQCD) at finite temperature and density with finite coupling effects. We adopt one species of staggered fermion, and develop an analytic formulation based on strong coupling and cluster expansions. We derive the effective potential as a function of two order parameters, the chiral condensate σ and the vector potential ω_τ , in a self-consistent treatment of the next-to-leading order (NLO) and the next-to-next-to-leading order (NNLO) effective action terms. Finite coupling effects lead to modifications of quark mass, chemical potential and the quark wave function renormalization factor. Finite coupling effects suppress the critical temperature at $\mu = 0$ ($T_{c,\mu=0}$), while critical temperature at $T = 0$ ($\mu_{c,T=0}$) is not affected much. NNLO corrections does not significantly affect $T_{c,\mu=0}$ and $\mu_{c,T=0}$, but the phase diagram shape including the position of the critical point is sensitive to the NNLO effects. Partially chiral restored matter is found to exist in NLO and NNLO SC-LQCD.

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*Speaker.

1. Introduction

Understanding the quantum chromodynamics (QCD) phase diagram is one of the most interesting problems in quark and hadron physics. Since the lattice QCD Monte-Carlo (MC) simulation has a sign problem at finite chemical potential μ [1], it is necessary to invoke some approximations in QCD or to apply effective models in order to elucidate the whole phase boundary, including the position of the critical point [2]. The strong-coupling lattice QCD (SC-LQCD) is one of the most instructive approximations especially to investigate the phase structure at finite temperature T and chemical potential μ . Based on the successes of SC-LQCD in pure Yang-Mills theory [3], chiral transition at finite T and μ has been investigated at strong coupling [4, 5, 6, 7]. For example, the phase diagram structure has been predicted in the strong coupling limit (SCL) [6, 7], and it is recently confirmed qualitatively in MC simulations [8] based on the monomer-dimer-polymer (MDP) formalism [9]. In order to make a step forward towards the true phase diagram, it is necessary to develop the formalism to include the plaquette effects both in MC simulations and SC-LQCD.

In this proceedings, we evaluate the effective potential in SC-LQCD with one species of unrooted staggered fermion at finite T and μ including the next-to-leading (NLO, $\mathcal{O}(1/g^2)$) [10, 11] and the next-to-next-to-leading (NNLO, $\mathcal{O}(1/g^4)$) [12] effects, and investigate the coupling dependence of the phase diagram and the critical point.

2. Effective Potential in NNLO SC-LQCD

The effective potential \mathcal{F}_{eff} is obtained in two steps in a finite T treatment of SC-LQCD. We first integrate out the spatial links U_j and obtain the effective action S_{eff} . In the next step we integrate out the quark field χ and the temporal links U_0 , and obtain the effective potential \mathcal{F}_{eff} . We define the effective action and potential (S_{eff} and \mathcal{F}_{eff}) on a lattice with spatial (temporal) size $L(N_\tau)$ at chemical potential μ as,

$$\exp[-S_{\text{eff}}(\chi, \bar{\chi}, U_0)] = \int \mathcal{D}U_j e^{-S_{\text{LQCD}}} = \int \mathcal{D}U_j e^{-S_F - S_G} = e^{-S_{\text{SCL}}} \langle e^{-S_G} \rangle, \quad (2.1)$$

$$\exp[-L^d N_\tau \mathcal{F}_{\text{eff}}] = \int \mathcal{D}[\chi, \bar{\chi}, U_0] \exp[-S_{\text{eff}}], \quad (2.2)$$

$$S_F = \frac{1}{2} \sum_x [V_x^+(\mu) - V_x^-(\mu)] + m_0 \sum_x M_x + \frac{1}{2} \sum_{j,x} \eta_{j,x} [\bar{\chi}_x U_{j,x} \chi_{x+\hat{j}} - \bar{\chi}_{x+\hat{j}} U_{j,x}^\dagger \chi_x], \quad (2.3)$$

$$S_G = \frac{2N_c}{g^2} \sum_P \left[1 - \frac{1}{2N_c} [U_P + U_P^\dagger] \right], \quad (2.4)$$

where $d = 3$, m_0 and $\eta_{j,x} = (-1)^{x_0 + \dots + x_{j-1}}$ are the spatial dimension, the bare quark mass and the staggered phase factor, respectively, $S_{\text{LQCD}} = S_F + S_G$ is the lattice QCD action, and U_P denotes the trace of a plaquette P . Mesonic composites are defined as $M_x = \bar{\chi}_x \chi_x$, $V_x^+ = \bar{\chi}_x e^{\mu} U_{0,x} \chi_{x+\hat{0}}$ and $V_x^- = \bar{\chi}_{x+\hat{0}} e^{-\mu} U_{0,x}^\dagger \chi_x$.

In order to perform the strong coupling expansion at finite T systematically, the cumulant (or coupled cluster) expansion is indispensable. By using the cumulant expansion, the expectation value of $\exp[-S_G]$ is found to be,

$$\langle e^{-S_G} \rangle \equiv \frac{1}{e^{-S_{\text{SCL}}}} \int \mathcal{D}U_j e^{-S_F} e^{-S_G} = \exp \left[\sum_{n=1}^{\infty} \frac{(-1)^n}{n!} \langle S_G^n \rangle_c \right], \quad e^{-S_{\text{SCL}}} = \int \mathcal{D}U_j e^{-S_F}. \quad (2.5)$$

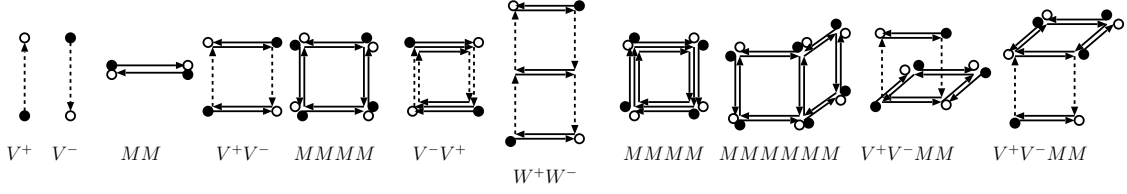


Figure 1: Diagrams contributing to the SCL, NLO and NNLO effective action terms. Quarks (anti-quarks) are shown in open (filled) circles, spatial (temporal) link variables are represented by solid (dotted) arrows, and plaquettes are represented by squares. Note that for NLO and NNLO diagrams, we also consider the hermite conjugate of these diagrams.

The bracket $\langle \dots \rangle_c$ is called a cumulant and shows the connected diagram contributions, e.g. $\langle S_G^2 \rangle_c = \langle S_G^2 \rangle - \langle S_G \rangle^2$. We find that the sum in the exponent corresponds to the strong coupling expansion of the effective action,

$$S_{\text{eff}} = S_{\text{SCL}} - \sum_{n=1}^{\infty} \frac{(-1)^n}{n!} \langle S_G^n \rangle_c = S_{\text{SCL}} + \Delta S_{\text{NLO}} + \Delta S_{\text{NNLO}} + \mathcal{O}(1/g^6, 1/\sqrt{d}). \quad (2.6)$$

The n -th term in the sum is proportional to $1/g^{2n}$, and we can identify $n = 1$ and $n = 2$ terms as NLO and NNLO corrections (ΔS_{NLO} and ΔS_{NNLO}), and S_{SCL} shows the SCL effective action.

We consider zero, one, and two-connected plaquette configurations, and integrate out U_j after putting fermionic elements ($\bar{\chi}U\chi$ and $\bar{\chi}U^\dagger\chi$). In the leading order of the $1/d$ expansion [5], we find that the diagrams shown in Fig. 1 contribute to the effective action,

$$S_{\text{eff}} = \frac{1}{2} \sum_x (V_x^+ - V_x^-) - \frac{b_\sigma}{2d} \sum_{x,j>0} [MM]_{j,x} \quad (2.7)$$

$$+ \frac{1}{2} \frac{\beta_\tau}{2d} \sum_{x,j>0} [V^+V^- + V^-V^+]_{j,x} - \frac{1}{2} \frac{\beta_s}{d(d-1)} \sum_{x,j>0,k>0,k \neq j} [MMMM]_{jk,x} \quad (2.8)$$

$$- \frac{\beta_{\tau\tau}}{2d} \sum_{x,j>0} [W^+W^- + W^-W^+]_{j,x} - \frac{\beta_{ss}}{4d(d-1)(d-2)} \sum_{\substack{x,j>0,|k|>0,|l|>0 \\ |k| \neq j, |l| \neq j, |l| \neq |k|}} [MMMM]_{jk,x} [MM]_{j,x+\hat{l}}$$

$$+ \frac{\beta_{\tau s}}{8d(d-1)} \sum_{x,j>0,|k| \neq j} [V^+V^- + V^-V^+]_{j,x} \left([MM]_{j,x+\hat{k}} + [MM]_{j,x+\hat{k}+\hat{0}} \right), \quad (2.9)$$

We have introduced a short-hand notation, $[AB]_{j,x} = A_x B_{x+\hat{j}}$, $[ABCD]_{jk,x} = A_x B_{x+\hat{j}} C_{x+\hat{j}+\hat{k}} D_{x+\hat{k}}$. Additional mesonic composites appears as $W_x^+ = e^{2\mu} \bar{\chi}_x U_{0,x} U_{0,x+\hat{0}} \chi_{x+2\hat{0}}$ and $W_x^- = e^{-2\mu} \bar{\chi}_{x+2\hat{0}} U_{0,x+\hat{0}}^\dagger U_{0,x}^\dagger \chi_x$. The coefficients are defined as, $b_\sigma = d/2N_c$, $\beta_\tau = d(1 + 1/2g^2)/N_c^2 g^2$, $\beta_s = d(d-1)(1 + 1/2g^2)/8N_c^4 g^2$, $\beta_{\tau\tau} = d/2N_c^3 g^4$, $\beta_{ss} = d(d-1)(d-2)/16N_c^7 g^4$, $\beta_{\tau s} = d(d-1)/2N_c^5 g^4$.

We can convert the above effective action in a spatially local and bi-linear form of quark fields by using the extended Hubbard-Stratonovich (EHS) transformation [10, 11]. The fermion part of the transformed effective action is found to be in the form of,

$$S_{\text{eff}}^{(F)} = \frac{1}{2} \sum_x (C V_x^+ - \bar{C} V_x^-) + \sum_x m M_x - \beta_{\tau\tau} \sum_x (\psi_{\tau\tau} W_x^+ + \bar{\psi}_{\tau\tau} W_x^-), \quad (2.10)$$

where $\psi_{\tau\tau}$ and $\bar{\psi}_{\tau\tau}$ are the auxiliary fields and C, \bar{C}, m contain auxiliary field contributions. We further reduce this effective action by introducing a gluonic dressed fermion,

$$\chi'_x = \chi_x - 2\beta_{\tau\tau}\psi_{\tau\tau}e^\mu U_{0,x}\chi_{x+\hat{0}}/C, \quad \bar{\chi}'_x = \bar{\chi}_x + 2\beta_{\tau\tau}\bar{\psi}_{\tau\tau}\bar{\chi}_{x+\hat{0}}e^{-\mu}U_{0,x}^\dagger/\bar{C}. \quad (2.11)$$

With χ' and $\bar{\chi}'$, we can absorb the next-to-nearest neighbor (NNN) interaction terms via W^\pm up to $\mathcal{O}(1/g^6)$ in the modification of mass and the coefficients of V^\pm, m, C and \bar{C} .

We now find that the NLO and NNLO corrections lead to the coefficient modification of V^+, V^- and M in the fermionic effective action. These modifications are interpreted as the modification of the wave function renormalization factor Z_χ , quark mass m_q and chemical potential $\tilde{\mu}$. The effective potential \mathcal{F}_{eff} is obtained in a similar way to that in SCL, and is found to be,

$$\mathcal{F}_{\text{eff}} = \mathcal{F}_{\text{eff}}^{(X)} + \mathcal{V}_q(m_q; \tilde{\mu}, T) - N_c \log \sqrt{Z_+ Z_-}, \quad (2.12)$$

$$\mathcal{F}_{\text{eff}}^{(X)} = \frac{1}{2}b'_\sigma \sigma^2 + \frac{1}{2}\beta'_\tau \bar{\psi}_\tau \psi_\tau + \frac{1}{2}\beta'_s \varphi_s^2 + \beta_{\tau\tau} \bar{\psi}_{\tau\tau} \psi_{\tau\tau} + \beta_{ss} \bar{\psi}_{ss} \psi_{ss} + \frac{1}{2}\beta_{\tau s} \bar{\psi}_{\tau s} \psi_{\tau s}, \quad (2.13)$$

$$m_q = \frac{m'}{\sqrt{Z_+ Z_-}}, \quad m' = b'_\sigma \sigma + m_0 - \beta_{\tau\tau}(\bar{\psi}_{\tau\tau} + \psi_{\tau\tau}), \quad \tilde{\mu} = \mu - \log \sqrt{Z_+/Z_-}, \quad (2.14)$$

$$\mathcal{V}_q = -T \log \left[\frac{\sinh[(N_c + 1)E_q/T]}{\sinh[E_q/T]} + 2 \cosh(N_c \tilde{\mu}/T) \right], \quad E_q(m_q) = \text{arcsinh}(m_q), \quad (2.15)$$

$$\beta'_\tau = \beta_\tau + \beta_{\tau s} \psi_{\tau s}, \quad \beta'_s = \beta_s + 2\beta_{ss} \bar{\psi}_{ss}, \quad b'_\sigma = b_\sigma + 2\beta'_s \varphi_s + 2\beta_{ss} \psi_{ss} + 2\beta_{\tau s} \bar{\psi}_{\tau s}, \quad (2.16)$$

$$Z_+ = 1 + \beta'_\tau \bar{\psi}_\tau + 4\beta_{\tau\tau} m' \bar{\psi}_{\tau\tau}, \quad Z_- = 1 + \beta'_\tau \psi_\tau + 4\beta_{\tau\tau} m' \psi_{\tau\tau}. \quad (2.17)$$

The auxiliary fields introduced during the bosonization procedure have to satisfy the stationary condition, $\partial \mathcal{F}_{\text{eff}}/\partial \Phi = 0$, where Φ represents one of the auxiliary fields. The stationary condition of \mathcal{F}_{eff} tells us that auxiliary fields other than σ and ω_τ are expressed as functions of (σ, ω_τ) . Thus the effective potential is a function of T, μ , and two order parameters; the chiral condensate σ and ω_τ . We can regard ω_τ is a vector potential field for quarks; the chemical potential shift is mainly determined by ω_τ , and ω_τ contributes repulsively to the effective potential in equilibrium. This two order parameter feature may be a natural consequence from the potential term from quarks, $\mathcal{V}_q(m_q; \tilde{\mu}, T)$. There are two independent derivatives, $\partial \mathcal{V}_q/\partial m_q$ and $\partial \mathcal{V}_q/\partial \tilde{\mu}$, which appear in the equilibrium condition, then we have two degrees of freedom.

3. Chiral Phase Transition in NNLO SC-LQCD

The effective potential determines the vacuum and the phase structure of QCD matter. The vacuum is determined by solving the stationary condition of \mathcal{F}_{eff} , which is equivalent to searching for the saddle point of \mathcal{F}_{eff} in the (σ, ω_τ) plane [11]. In Fig. 2, we show the the critical temperature at $\mu = 0$ ($T_{c,\mu=0}$) and the critical chemical potential at $T = 0$ ($\mu_{c,T=0}$). In both NLO and NNLO, the phase transition on the T -axis is found to be the second order, and $T_{c,\mu=0}$ decreases as β increases. NLO and NNLO give almost the same values of $T_{c,\mu=0}$. We also show the MC results on T_c in SCL and the critical coupling (β_c) for given values of $N_\tau = 1/T$ at $\mu = 0$ (filled triangles) [8, 13]. The suppression of $T_{c,\mu=0}$ in NLO and NNLO is not enough to explain the MC results. The phase transition on μ -axis is numerically found to be the first-order in NLO in the coupling range studied here ($\beta \leq 6$), and for $\beta \lesssim 5.5$ in NNLO. The first order phase transition is determined by the

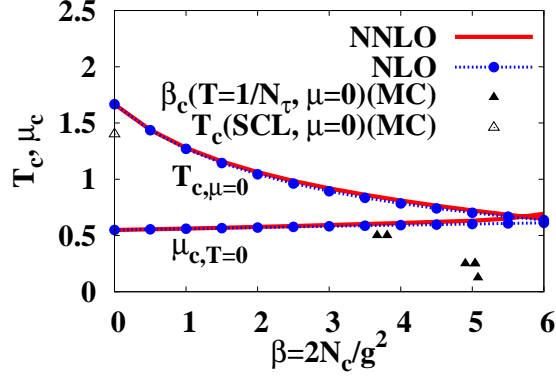


Figure 2: The β dependences of a critical temperature at $\mu = 0$ and a critical chemical potential at $T = 0$. We also show the MC results of T_c in the strong coupling limit and the critical coupling β_c for a given N_τ . MC results are taken from [8, 13].

magnitude relation between E_q and $\tilde{\mu}$. As β increases, both of E_q and $\tilde{\mu}$ are suppressed, and the suppression effects on $\tilde{\mu}$ is slightly larger than that in E_q . As a result, $\mu_{c, T=0}^{(1st)}$ is a slightly increasing function of β .

In Fig. 3, we show the β dependence of the phase diagram in NLO and NNLO. While $T_{c, \mu=0}$ and $\mu_{c, T=0}$ are very similar in NLO and NNLO, the shape of the phase boundary and the critical point are sensitive to the NNLO corrections. In NLO, the critical point (CP) temperature (T_{CP}) decreases and the CP chemical potential (μ_{CP}) stays almost the same as β increases. The tricritical point at strong coupling starts to deviate from the second order phase boundary at $\beta \simeq 4.5$, and it becomes the critical end point (CEP) even in the chiral limit. In NNLO, both T_{CP} and μ_{CP} decreases as β increases, and the CP stays on the second order phase transition boundary.

The CP behavior in NNLO is considered to be an improvement over NLO. Since one species of staggered fermion corresponds to $N_f = 4$, the phase transition at $\mu = 0$ is expected to be the first order in the continuum region due to anomaly contributions [14]. In the strong coupling region, the chiral $SU_f(4)$ symmetry is not manifest with staggered fermions. The lattice QCD action with staggered fermion is invariant under the chiral transformation, $\chi_x \rightarrow \exp(i\theta \epsilon_x) \chi_x$, where $\epsilon_x = (-1)^{x_0+x_1+\dots+x_d}$ is a γ_5 related phase [4]. This transformation corresponds to the $\gamma_5 \otimes \gamma_5$ rotation in the spinor-flavor space in the continuum region, which is only a small subgroup of $SU_A(4)$. As a result, effective number of flavors is small at strong coupling, and the second order phase transition at $\mu = 0$ emerges. If we can take account of the finite coupling effects correctly, the order of the phase transition at $\mu = 0$ should change from the second to the first order at a certain coupling strength. The behavior of CP in NNLO suggests that the first order transition boundary comes closer to the T -axis, and agrees with the above expectation while the shift is not enough.

Another interesting point is the existence of the region where $\mu_c^{(1st)}(T) < \mu_c^{(2nd)}(T)$ is satisfied. Between these critical chemical potentials, $\mu_c^{(1st)}(T) < \mu < \mu_c^{(2nd)}(T)$, chiral condensate is suppressed compared to its vacuum value but the chiral symmetry is not fully restored. This partially chiral restored (PCR) matter suggested in NLO SC-LQCD [10] persists to exist also with

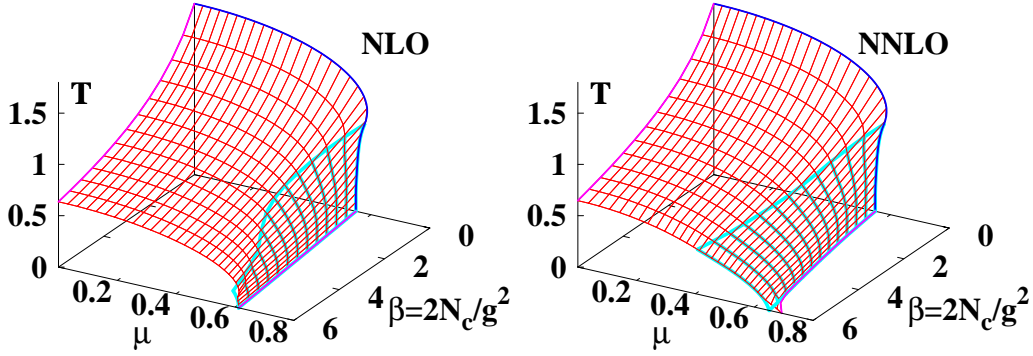


Figure 3: Phase diagram evolution in NLO and NNLO SC-LQCD. Phase boundary are shown by thin line surface, and thick surfaces show the first order phase transition boundary.

NNLO corrections, and would correspond to the quarkyonic matter suggested at large N_c [15].

4. Summary

In this proceedings, we have derived an analytic expression of the effective potential in the strong coupling lattice QCD (SC-LQCD) with one species of unrooted staggered fermion for color $SU(N_c)$ including the next-to-leading (NLO) and the next-to-next-to-leading (NNLO) effects at finite temperature and density, and investigated finite coupling effects on the chiral phase transition and the phase diagram. Effective action terms have been systematically evaluated based on the strong coupling expansion in the leading order of the $1/d$ expansion. NLO and NNLO effective action terms have been found to come from one plaquette and two connected plaquette configurations. We have applied the extended Hubbard-Stratonovich transformation to bosonize fermion interaction terms. The next-to-nearest neighbor interaction appears in NNLO corrections, and is evaluated by introducing a gluonic dressed fermion. We have obtained the effective potential as a function of temperature (T), chemical potential (μ), and the two order parameters: the chiral condensate (σ) and the vector potential (ω_τ). The vacuum is determined from the stationary condition of the effective potential with respect to each auxiliary field. NLO and NNLO effects result in modification of the wave function renormalization factor, quark mass and chemical potential.

The critical temperature at $\mu = 0$ ($T_{c,\mu=0}$) and chemical potential at $T = 0$ ($\mu_{c,T=0}$) are found to be similar in NLO and NNLO. The critical temperature at $\mu = 0$ ($T_{c,\mu=0}$) is a decreasing function of β in both NLO and NNLO, while the suppression is not enough to explain the Monte-Carlo results at $N_\tau = 1/T = 2, 4$ and 8 . The critical chemical potential at $T = 0$ ($\mu_{c,T=0}$) is found to stay almost constant in both NLO and NNLO, due to the cancellation of the effects from mass and effective chemical potential reduction. The shape of the phase boundary and the position of the critical point (CP) are found to be sensitive to the NNLO effects. As β increases, while CP moves towards the μ -axis in NLO, CP moves towards the T -axis in NNLO.

Inclusion of the Polyakov loop and higher order terms in the $1/d$ expansion would be interesting directions of study.

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