

Di-boson production and limits on triple gauge boson couplings at the Tevatron

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The CDF and D0 experiments have observed all varieties of diboson events at the Fermilab Tevatron. These observations are all in statistical agreement with Standard Model expectations, and limits have been placed on non-Standard-Model anomalous triple gauge couplings. New limits from CDF on $Z\gamma$ anomalous couplings are at least twice as stringent as previously reported couplings.

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1. Introduction

The structure of self couplings between gauge bosons in the Standard Model (SM) is a direct consequence of electroweak symmetry breaking, so studying these couplings not only confirms the structure of the electroweak sector of the SM, but also presents an opportunity to search for beyond-the-Standard-Model interactions. While direct couplings exist in the SM between W bosons and photons, and between W and Z bosons, there are no Standard Model couplings between the Z boson and the photon, or between the Z boson and itself. Another way to say this is that in the Standard Model, a W boson can radiate a photon or a Z boson, but neither a photon nor a Z boson can radiate either itself or the other. This means that the SM only includes t -channel production of $Z\gamma$ and ZZ diboson final states, while $W\gamma$, WW , and WZ diboson final states can be produced by both s - and t -channel processes at tree level. In this paper we examine diboson production in 1.96 TeV $p\bar{p}$ collisions using the CDF and D0 detectors, compare this production to Standard Model predictions, and set limits on the strength of some anomalous couplings.

2. $W\gamma$ Production

The Standard Model does have a coupling between the W boson and the photon, but the SM form of this coupling leads to an unusual feature. Destructive interference between diagrams leads to a “radiation amplitude zero” (RAD): no radiation is emitted in a particular direction. This RAD was first predicted in 1979 [2], and it was observed by the D0 collaboration in 2008 [3]. In this publication the D0 collaboration also put limits on non-SM $W\gamma$ couplings, parameterized by κ_γ and λ_γ , which are 1 and 0 respectively in the SM. These 95% C.L. limits are $0.49 < \kappa_\gamma < 1.51$ and $-0.12 < \lambda_\gamma < 0.13$.

3. $Z\gamma$ Production

Since there is no direct $Z\gamma$ coupling in the Standard Model, the only tree-level diagram for a $Z\gamma$ final state is from initial-state radiation (ISR), where a photon is radiated off an incoming quark. If the Z boson decays into two charged leptons (electrons or muons), then the dilepton mass will generally be near the Z -boson mass for ISR. In addition, when the Z decays to charged leptons, it is possible to get radiation off one of these final-state leptons (FSR for final-state radiation). In the case of FSR, the three-body mass $M_{ll\gamma}$ will generally be near the Z -boson mass. Were there an anomalous coupling between the photon and the Z boson, it would be possible for a Z to radiate a photon or for a photon to radiate a Z , and in this case the two leptons from the Z decay would have a mass near the Z mass, similar to ISR. In particular, anomalous couplings would lead to an increase in the number of $Z\gamma$ events with large E_T^γ .

The CDF and D0 collaborations have used two methods to look for $Z\gamma$ events. If the Z boson decays into electrons or muons, it is possible to find $Z\gamma$ events with low backgrounds by first identifying the Z decay and then looking for the additional photon. On the other hand, if the Z boson decays into a pair of neutrinos, these will exit the detector without being seen, so we only see the photon and an imbalance of transverse energy. This signature has significant backgrounds, and these backgrounds are overwhelming at low E_T^γ , so the searches for $Z\gamma \rightarrow \nu\nu\gamma$ requires a larger photon energy than the searches where the Z boson decays to a pair of charged leptons.

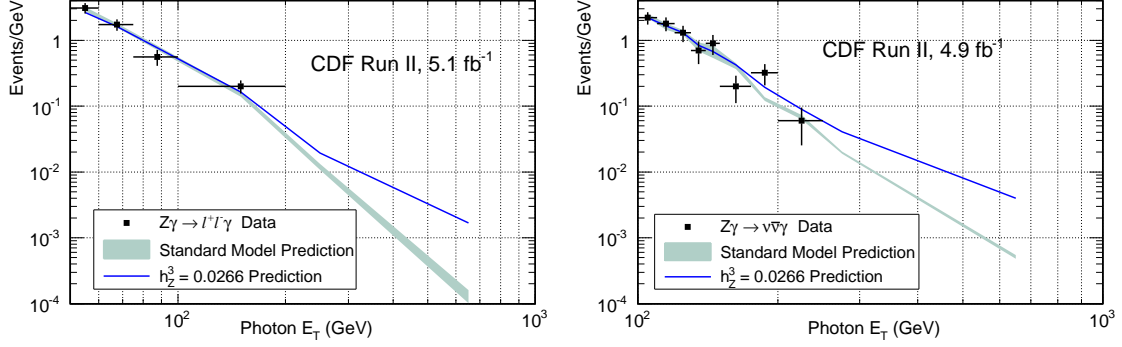


Figure 1: Photon E_T distribution in the $Z \rightarrow \mu\mu + Z \rightarrow ee$ datasets for photons with $E_T > 50$ GeV (left). Photon E_T distribution in the $Z \rightarrow \nu\nu$ dataset for photons with $E_T > 100$ GeV (right).

Parameter	CDF ($\Lambda = 1.2$ TeV)	CDF ($\Lambda = 1.5$ TeV)	D0 ($\Lambda = 1.5$ TeV)
h_3^Z	-0.018, 0.020	-0.017, 0.016	< 0.033
h_4^Z	-0.0009, 0.0009	-0.0006, 0.0005	< 0.0017
h_3^γ	-0.022, 0.021	-0.017, 0.016	< 0.033
h_4^γ	-0.0009, 0.0010	-0.0006, 0.0006	< 0.0017

Table 1: Upper limits (95% C.L.) on anomalous $Z\gamma$ couplings using notation from reference [4] for $\Lambda = 1.2$ and 1.5 TeV. The D0 limits are from reference [5]; the CDF limits are new.

The most recent Tevatron $Z\gamma$ analysis, using 5 fb^{-1} from CDF, requires $E_T^\gamma > 50$ GeV for $Z \rightarrow e^+e^-$ and $Z \rightarrow \mu^+\mu^-$, and $E_T^\gamma > 100$ GeV for $Z \rightarrow \nu\nu$. The $E_T^\gamma > 50$ GeV requirement eliminates most of the FSR events. The distributions of E_T^γ for these two cases are shown in Figure 1 along with the SM expectations. There is no evidence for an excess at high E_T^γ , so limits were set on the anomalous coupling parameters h_3^Z , h_4^Z , h_3^γ , and h_4^γ (see reference [4] for an explanation of this notation). These limits are shown in Table 1 along with the published limits from the D0 collaboration. An important note is that a cutoff (Λ) is required when making the SM predictions in order to avoid unitarity violations, and the limits are a function of this cutoff. Previous results have been published with cutoffs of $\Lambda = 1.2$ TeV and $\Lambda = 1.5$ TeV, so the new results are reported using both of these values.

There are significant correlations between the h_3 and h_4 parameters. Figure 2 shows the two-dimensional limits for the new CDF analysis when both parameters are varied at the same time.

4. WZ Production

The Standard Model includes a WWZ coupling, but it is possible to have anomalous WWZ couplings with a different form, and both collaborations have searched for these anomalous triple gauge couplings. Candidate events are found by searching for three charged leptons and missing transverse energy, where two oppositely charged leptons of the same flavor have a mass consistent with coming from a Z boson, and the third lepton has a transverse mass when combined with the

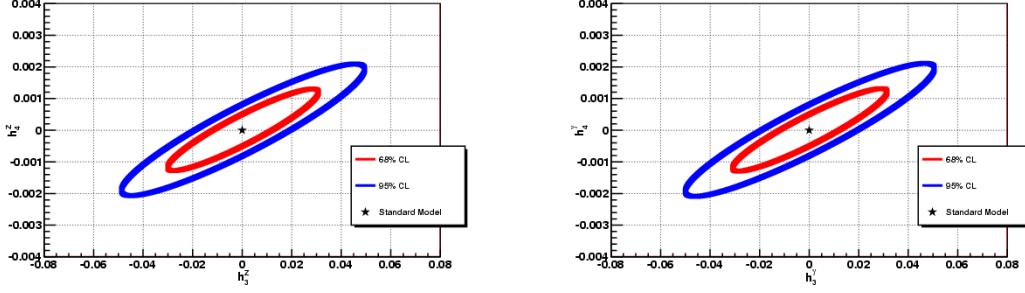


Figure 2: Limits on h_3^Z vs. h_4^Z showing significant correlation between the two (left). Limits on h_3^γ vs. h_4^γ also showing significant correlation (right). Both plots are for the cutoff parameter $\Lambda = 1.5$ TeV.

Coupling relation	CDF ($\Lambda = 2$ TeV)	D0 ($\Lambda = 2$ TeV)
$\Delta g_1^Z = \Delta \kappa_Z = 0$	$-0.13 < \lambda_Z < 0.14$	$-0.075 < \lambda_Z < 0.093$
$\lambda_Z = \Delta \kappa_Z = 0$	$-0.13 < \Delta g_1^Z < 0.23$	$-0.053 < \Delta g_1^Z < 0.156$
$\lambda_Z = \Delta g_1^Z = 0$	$-0.68 < \Delta \kappa_Z < 1.00$	$-0.376 < \Delta \kappa_Z < 0.686$

Table 2: Limits (95% C.L.) on anomalous WZ couplings using notation from reference [6] for form factor $\Lambda = 2$ TeV. The CDF limits use 1.9 fb^{-1} of data; the D0 limits use 4.1 fb^{-1} [7].

missing transverse energy consistent with coming from a W boson. Limits on anomalous triple gauge couplings are listed in Table 2.

The cross section for WZ production has been measured by both collaborations. The D0 collaboration has recently reported observing 34 WZ candidate events with an estimated background of 6.0 ± 0.4 events in 4.1 fb^{-1} of data for a measured cross section $\sigma(p\bar{p} \rightarrow WZ) = (3.90_{-0.90}^{+1.06})$ pb [7]. The CDF collaboration used 6 fb^{-1} of data to find 50 $WZ \rightarrow 3l + \nu$ candidate events with 11.2 ± 1.6 expected background and to measure the WZ production cross section to be $\sigma(p\bar{p} \rightarrow WZ) = (4.1 \pm 0.7)$ pb (preliminary). In a separate measurement CDF used 5.9 fb^{-1} of data to measure $\sigma(p\bar{p} \rightarrow WZ) = (3.7 \pm 0.8)$ pb (preliminary).

5. ZZ Production

The ZZ cross section is small, and when this is combined with the small branching fraction for the Z to decay to charged leptons, very few events are expected for $ZZ \rightarrow eeee$, $ZZ \rightarrow ee\mu\mu$, or $ZZ \rightarrow \mu\mu\mu\mu$. Anomalous couplings can produce additional ZZ events, so both collaborations have searched for ZZ events, and when no excess over SM expectations were found, both collaborations set limits on anomalous couplings. These limits, based upon 1 fb^{-1} of data for D0 and 1.9 fb^{-1} of data for CDF are shown in Table 3.

More recent analyses by both collaborations use more data. The D0 collaboration used 1.7 fb^{-1} of data to report [10] the first ZZ observation with a significance exceeding 5 standard deviations. The CDF collaboration has used 6 fb^{-1} of data to measure the production cross section

Parameter	CDF ($\Lambda = 1.2$ TeV)	D0 ($\Lambda = 1.2$ TeV)
f_4^γ	-0.1, 0.1	-0.26, 0.26
f_4^Z	-0.12, 0.12	-0.28, 0.28
f_5^γ	-0.11, 0.11	-0.30, 0.28
f_5^Z	-0.13, 0.12	-0.31, 0.29

Table 3: Limits (95% C.L.) on anomalous ZZ couplings using notation from reference [8] for $\Lambda = 1.2$. The D0 limits, using 1 fb^{-1} , are from reference [9]; the CDF limits use 1.9 fb^{-1} .

$\sigma(p\bar{p} \rightarrow ZZ) = (1.7_{-0.7}^{+1.2}(\text{stat}) \pm 0.2(\text{syst})) \text{ pb}$ (preliminary). For this analysis, 4 $ZZ \rightarrow 4l$ candidates were observed with ZZ mass below 300 GeV. The analysis remains blind to candidate events with larger mass with the intent of searching for heavy resonances that decay into a pair of Z bosons. Neither collaboration used these larger datasets to set limits on anomalous couplings.

6. Conclusions

All varieties of diboson events have been observed at the Tevatron. Production rates and kinematic distributions match Standard Model expectations, and limits have been set on anomalous triple-gauge-coupling parameters.

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