

Dark Matter Signature from the Sky and at Colliders

Shou-hua Zhu* †

ITP, Peking University E-mail: shzhu@pku.edu.cn

> In this talk, we briefly review our recent investigations on the properties of dark matter (DM) particle. (1) In 2008 PAMELA released their first results on the positron and antiproton ratios. Stimulated by the new data, we studied systematically [1] the cosmic ray propagation models and calculated the secondary positron and antiproton spectra. Above ~ 10 GeV PAMELA data shows a clear excess on the positron ratio. However, the secondary antiproton is roughly consistent with data. The positron excess may be a direct evidence of dark matter. We compared the positron and anti-proton spectra with data by assuming dark matter annihilates or decays into different final states. The PAMELA data actually excludes quark pairs being the main final states, disfavors gauge boson final states. Only in the case of leptonic final states the positron and anti-proton spectra can be explained simultaneously. Finally we considerd a decaying neutralino dark matter model in the frame of supersymmetry with R-parity violation. The PAMELA data is well fitted with neutralino mass $600\sim 2000~\text{GeV}$ and life time $\sim 10^{26}$ seconds. We also investigated the corresponding high energy neutrino [2] and the synchrotron and inverse Compton radiation [3] signatures and found those signature may distinguish different scenarios which can account for the PAMELA observations. (2) The DM sector can be complicated. Motivated by recent observations and theoretical progress, the possible whole picture of DM sector emerges, which includes O(TeV) DM and O(GeV) or less light particle. Via the so-called Sommerfeld enhancement effects, the new light particle can fill the gap between the DM thermal relic density and the observed SM particle flux induced by nowadays DM annihilating. The O(TeV) may be detected by LHC. The related O(GeV) or less new particles can be studied by low energy linear colliders (e.g. BESIII) [4, 5] or via neutrino telescope [6].

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*Speaker.

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1. Introduction

Synergy of non-accelerator- and accelerator-based experiments has shaped our knowledge of fundamental particles and interactions. The famous examples were the discoveries of positron and the 'V-particle' in the cosmic rays (CRs) in the early stage of particle physics. Since 2008, the PAMELA, ATIC, Fermi, HESS, CDMSII etc. have released their new observations. Some new phenomena were reported. The investigations on their origin and possible connection became extremely necessary. Furthermore whether the resulting underlying picture can be checked by colliders and astrophysical experiments is also worthy of study. This proceeding paper is not the thorough review, instead it contains merely our recent several works [1, 2, 3, 4, 5, 6]. In section II, we reported our investigation on PAMELA data interpretation. In section III, we describe how to find the new light particle at low-energy collider (BESIII) and via neutrino telescope. Such new light particle at O(GeV) or less is the necessary component to get boost factor via the so-called Sommerfeld enhancement effects.

2. PAMELA data and leptonically decaying dark matter

PAMELA is a satellite borne experiment designed to measure CRs in a wide energy range with unprecedented accuracy. Recently, the PAMELA collaboration released the first data about antiprotons and positrons. Usually it is thought that antiprotons and positrons are produced when CRs propagate in the Milky Way and collide with the interstellar medium (ISM). The abundance of these secondaries can be calculated with relatively high precision. However, the PAMELA results show an obvious excess in the faction of $e^+/(e^+ + e^-)$ at energies above ~ 10 GeV. Interestingly the excess keeps to rise up to energy ~ 100 GeV. On the other hand, the spectrum of antiprotons fits the prediction quite well. These results confirm the previous results by HEAT and AMS within the error bars.

As a systematic study we investigated firstly the cosmic ray propagation model carefully so that we can predict the positron and antiproton spectra to compare with the PAMELA data. Our studies [1] showed that it is barely to account for PAMELA observations on positron and anti-proton by adjusting only the propagating parameters. Secondly, we adopted a relatively model-independent approach to constrain the DM annihilation or decay products from the PAMELA data. We found the PAMELA data actually excludes the annihilation or decay products being quark pairs, strongly disfavors the gauge bosons and favors dominant leptonic final states. Details analysis can be found in Ref. [1]. In some case, the heavier DM can also fit Pamela data and predict obviously larger flux for not-reached higher energy anti-proton, which will be checked by future experiment.

Thirdly, we turn to discuss the specific decaying DM model in supersymmetry scenario. We find, if the positron excess at PAMELA is indeed of DM origin, the decaying DM is superior to annihilating DM [1]. In SUSY model, the discrete symmetry of R-parity can be invoked to avoid dangerous baryon-number violation terms which drive unexpected proton decay. Defined as $R = (-1)^{2S+3B+L}$, the R-parity of a SM particle is even while its superpartner is odd. Then the lightest supersymmetric particle (LSP) is stable. As the LSP, the neutralino can easily induce correct relic density and act as the suitable DM candidate.

Although R-parity symmetry is well motivated for SUSY phenomenology, there is no reason for this symmetry to be exact. One can introduce some R-parity violation terms in the Lagrangian which make the LSP decay into SM particles. The general gauge invariant superpotential of the minimal supersymmetric standard model can be written as

$$W = W_{MSSM} + \lambda_{ijk} L_i L_j \bar{E}_k + \lambda'_{ijk} L_i Q_j \bar{D}_k + \lambda''_{ijk} \bar{U}_i \bar{D}_j \bar{D}_k + \mu'_i L_i H_u$$
(2.1)

where i, j, k are generation indices. In general, we should consider all R-parity violation terms in Eq. 2.1. However, they might not appear in theory at the same time. In this work, we only consider lepton-number violation term $LL\bar{E}$ in order to account for the PAMELA positron excess. The chosen parameters are in Tab. 1.

DC	$ au(10^{26}s)$	$\lambda'(10^{-25})$	DR	$\tau(10^{26}s)$	$\lambda'(10^{-25})$
А	4.5	3.2	Α	3.9	3.4
В	2.6	14.8	В	2.3	15.7
С	1.7	16.1	С	1.5	17.1
D	1.2	59.5	D	1.1	62.1
Е	1.0	253.2	E	0.9	266.9
F	0.6	159.7	F	0.6	159.7
G	0.5	156.7	G	0.5	156.7

Table 1: Life time τ in unit of $10^{26}s$ and R-parity violation parameter λ' in unit 10^{-25} for different benchmark points. "DC" and "DR" refer to the two propagation models we adopted.

The induced positron and antiproton fractions as function of energy for different benchmark points are shown in Fig. 1. From the curves we can see clearly that PAMELA observations can be accommodated in this R-parity-violated model.

3. New light particle at low-energy collider and neutrino telescope

In the annihilating DM scenario, there is a long-standing gap between the DM thermal relic density and the observed SM particle flux induced by nowadays DM annihilating. In the earlier studies the weak bosons can be treated as the light particle provided that DM is O(10) TeV or higher. However recent observations implied that the DM is not so heavy, instead at O(TeV). Therefore the *possible* whole picture of dark matter sector contains O(TeV) DM and O(GeV) or less new light particle. At the powerful LHC, O(TeV) DM may be detected. Here we only focus on the detection of light particle at low energy collider [4, 5] and neutrino telescope [6].

In our studies, the observability of such kind of light particle at BESIII detector is investigated through the processes $e^+e^- \rightarrow U\gamma$, followed by $U \rightarrow e^+e^-$, $U \rightarrow \mu^+\mu^-$ and $U \rightarrow v\overline{\nu}$. In the invisible channel where light boson decays into neutrinos, BESIII can measure the coupling of the extra U(1) down to $O(10^{-4}) \sim O(10^{-5})$ because of the low Standard Model backgrounds. In the visible channel where light particle decays into charged lepton pair, BESIII can only measure the coupling down to $O(10^{-3}) \sim O(10^{-4})$ due to the large irreducible QED backgrounds.





Figure 1: Positron fraction as a function of energy for different benchmark points. The neutralino mass of each model is given, while the life time of neutralino in each model is given in Tab. 1. The neutralino mass in each line is increasing from left to right.

We also investigate the possibility of detecting light long-lived particle (LLP) produced by high energy cosmic ray colliding with atmosphere. The LLP may penetrate the atmosphere and decay into a pair of muons near/in the neutrino telescope ['V-particle' again?]. Such muons can be treated as the detectable signal for neutrino telescope. The dark sector may be complicated, and there may exist more than one light particles, for example the dark gauge boson A' and associated dark Higgs boson h'. In our study, we discuss the scenario with A' heavier than h' and h' is treated as LLP. Based on our numerical estimation, we find that the large volume neutrino telescope IceCube has the capacity to observe several tens of di-muon events for favorable parameters if the decay length of LLP can be comparable with the depth of atmosphere. The challenge here is how to suppress the muon backgrounds induced by cosmic rays and atmospheric neutrinos.

Conclusions and discussions: Single experiment can hardly give us the complete picture of the nature. Synergy of non-accelerator- and accelerator-based experiments as well as the theoretical investigations have shaped and will keep on shaping our knowledge of fundamental particles and interactions.

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