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Inert Model and evolution of the Universe

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> We consider evolution of the Universe after EWSB leading to the present inert phase, containing a SM-like Higgs boson and scalar dark particles, among them a Dark Matter candidate. We address the question, if there is a possibility to have a sequence of the phase transitions instead of a single one leading directly from EW symmetric phase to the inert one.

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We consider the 2HDM Lagrangian for two SU(2) scalar doublets φ_S, φ_D with Yukawa interaction set to Model I and the potential V:

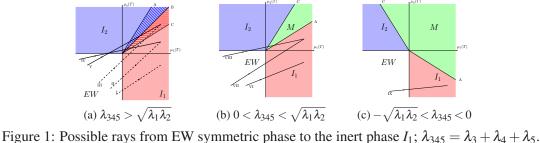
$$V = -\frac{m_{11}^2}{2}|\varphi_S|^2 - \frac{m_{22}^2}{2}|\varphi_D|^2 + \frac{\lambda_1}{2}|\varphi_S|^4 + \frac{\lambda_2}{2}|\varphi_D|^4 + \lambda_3|\varphi_S|^2|\varphi_D|^2 + \lambda_4|\varphi_S^{\dagger}\varphi_D|^2 + \frac{1}{2}\left(\lambda_5(\varphi_S^{\dagger}\varphi_D)^2 + h.c\right).$$

V is invariant under a Z_2 transformation: $\varphi_S \to \varphi_S$, $\varphi_D \to -\varphi_D$. In the Inert Model also vacuum state is Z_2 -symmetric and the Dark Matter candidate, from the Z_2 -odd doublet φ_D , appears.

The most general EWSB solution $\langle \varphi_S \rangle^T = \frac{1}{\sqrt{2}}(0, v_S), \langle \varphi_D \rangle^T = \frac{1}{\sqrt{2}}(u, v_D)$ gives three neutral extrema (u = 0): *inert* $(I_1; v_D = 0, v_S^2 = v^2)$ with SM-like Higgs *h* from φ_S and DM candidate from φ_D (eg. *H*); *inert-like* $(I_2; v_S = 0, v_D^2 = v^2)$ with massless fermions (Model I: only φ_S couples to fermions) and no candidate for DM; *mixed* $(M; v_D, v_S \neq 0, v^2 = v_S^2 + v_D^2)$ – a standard 2HDM extremum. The lowest energy extremum, fulfilling positivity constraints, is the vacuum.

The one-loop thermal corrections to *V* are $m_{ii}^2(T) = m_{ii}^2 - c_i T^2$ (i = 1, 2) where $c_i = c_i (\lambda_{1-4}; g, g'; g_t^2 + g_b^2$ for i = 1), with fixed λ_i (g, g' - EW gauge couplings, $g_t, g_b - \text{SM}$ Yukawa couplings).

The possible sequences of phase transitions between different vacua are shown as rays in plots (1a-1c), for (μ_1, μ_2) plane $(\mu_i(T) = m_{ii}^2(T)/\sqrt{\lambda_i}, i = 1, 2)$. For $EW \to I_1$, I_1 is the only vacuum that existed after EWSB. For **rays I**, **VI**, **IX** I_2 is not an extremum, for **rays II**, **VII** I_2 is an extremum, but not a local minimum; for **ray III** I_2 is a local minimum. The sequence: $EW \to I_2 \to I_1$, with transition between I_2 and I_1 vacuum, is possible for **ray IV**, where I_2 is not a local minimum and for **ray V**, where I_2 is a local minimum, that coexists with the global minimum I_1 . For **ray VIII** there is a possibility of going through the mixed vacuum $EW \to I_2 \to M \to I_1$. In this case there is only one minimum at any temperature.



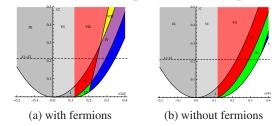


Figure 2: Example of rays for $M_h = 120 \text{ GeV}, M_H = 60 \text{ GeV}, M_A = 68 \text{ GeV}, M_{H^{\pm}} = 110 \text{ GeV}.$

Plots 2a,2b show the different vacua sequences (rays I-IX) leading to I_1 in $(\lambda_2, \lambda_{345})$ plane. For $\lambda_2 \leq \lambda_1$ different types of vacua in the past are possible only if the fermionic part of c_1 is included.

References

 Evolution of Universe to the present inert phase, I.F. Ginzburg, K.A. Kanishev, M. Krawczyk, D. Sokołowska [arXiv:1009.4593]