## PROCEEDINGS OF SCIENCE

# PoS

## $|V_{us}|$ and Charged Lepton Universality from au Decays

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 $\tau$  decays provide a unique opportunity to make precision measurements of the weak interaction between the first and second generation of quarks, the Cabibbo-Kobayashi-Maskawa matrix (CKM) element  $|V_{us}|$ , and tests charged lepton universality, the assumption that all leptons have the same coupling strength to the gauge bosons. We present the recent measurements of  $\tau$  decays from BELLE and BABAR and the improvements in determining  $|V_{us}|$  and testing lepton universality.

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#### **1.** $\tau$ Decay Measurements from BELLE and BABAR

The general strategy used at B-Factories to select clean sample of  $\tau^+\tau^-$  events is to split the event into two hemisphere using the thrust in the CM frame. One hemispheres is required to have an isolated electron or muon from  $\tau \to ev_{\tau}\overline{v}_{e}$  or  $\tau \to \mu v_{\tau}\overline{v}_{\mu}$  to tag the event, while the other hemisphere, the signal hemisphere, is required to decay hadronically. Particle identification is applied to the charged tracks and selection of  $\gamma$ s,  $\pi^{0}$ s and  $\eta$ s are applied to identify the hadronic decay mode.

The BABAR Collaboration has published a measurement of  $BR(\tau^- \to K^- \pi^0 v_\tau)$  [1]. This analysis used a lepton tag with one well reconstructed  $\pi^0$  and one kaon identified in the signal hemisphere. The branching fraction was measured to be  $BR(\tau^- \to K^- \pi^0 v_\tau)=(4.16 \pm 0.03 \pm 0.18) \times 10^{-3}$ .

BABAR and BELLE have both measured  $\tau^- \to K_S^0 \pi^- \nu_{\tau}$ , where the BABAR result is preliminary [2] and the BELLE result is published[3]. In this analysis, the signal hemisphere was required to contain a  $K_S^0$  reconstructed from a  $\pi^+\pi^-$  pair and a bachelor pion. In the  $\tau^- \to K_S^0\pi^-\pi^0\nu_{\tau}$  channel, the signal hemisphere was also required to contain a  $\pi^0$  reconstructed from two photons[2]. BELLE measured  $BR(\tau^- \to K_S^0\pi^-\nu_{\tau})=(8.08 \pm 0.04 \pm 0.26) \times 10^{-3}$ , while BABAR measured  $BR(\tau^- \to K_S^0\pi^-\nu_{\tau})=(8.40 \pm 0.23) \times 10^{-3}$  and  $\tau^- \to K_S^0\pi^-\pi^0\nu_{\tau}=(3.42 \pm 0.06 \pm 0.15) \times 10^{-3}$ .

Both the BELLE and BABAR Collaborations have measured  $\tau^- \rightarrow h^- h^- h^+ v_\tau$  where *h* has been identified as a pion or kaon[4, 5]. After selecting the event with a leptonic tag and three charged tracks in the signal hemisphere, particle identification was applied and a matrix technique was used to extract the branching fractions shown in Table 1. The BABAR paper also measured  $BR(\tau^- \rightarrow \pi^- \phi v_\tau) = (3.42 \pm 0.55 \pm 0.25) \times 10^{-5}$ ,  $BR(\tau^- \rightarrow K^- \phi v_\tau) = (3.39 \pm 0.20 \pm 0.28) \times 10^{-5}$ and  $BR(\tau^- \rightarrow K^- K^- K^+ v_\tau [ex. \phi]) < 2.5 \times 10^{-6}$  at 90% CL. whereas BELLE measured  $BR(\tau^- \rightarrow K^- \phi v_\tau) = (4.06 \pm 0.25 \pm 0.26) \times 10^{-5}$  in [6]

**Table 1:** The measured  $\tau^- \rightarrow h^- h^- h^+ v_{\tau}$  branching fractions.

Decay Mode	BELLE	BABAR
$BR(\tau^- \to \pi^- \pi^- \pi^+ \nu_\tau \ [ex. \ K_s^0])$	$(8.42 \pm 0.01 \substack{+0.26 \\ -0.25}) \times 10^{-2}$	$(8.83\pm 0.01\pm 0.13)\times 10^{-2}$
$BR(\tau^- \to K^- \pi^- \pi^+ v_\tau \ [ex. \ K_s^0])$	$(3.30 \pm 0.01 \stackrel{+0.16}{-0.17}) \times 10^{-3}$	$(2.73 \pm 0.02 \pm 0.09) \times 10^{-3}$
$BR(\tau^- \to K^- \pi^- K^+ v_{\tau})$	$(1.55 \pm 0.010 \stackrel{+0.06}{-0.05}) \times 10^{-3}$	$(1.346 \pm 0.010 \pm 0.036) \times 10^{-3}$
$BR(\tau^- \to K^- K^- K^+ \nu_{\tau})$	$(3.29 \pm 0.17^{+0.19}_{-0.20}) \times 10^{-5}$	$(1.58\pm0.13\pm0.12)\times10^{-5}$

Recently, the BELLE collaboration has published a paper on precision measurements of  $\tau$  decays containing an  $\eta$ [7]. In this paper, the signal hemisphere and  $\pi^0$  and  $\eta$  are reconstructed from two photons, where the  $\tau^- \to K^- \eta v_{\tau}$  used the additional  $\eta$  to  $\pi^- \pi^+ \pi^0$  decay mode. Using a fit of the  $\eta$  invariant mass, BELLE measured  $BR(\tau^- \to K^- \eta v_{\tau}) = (1.48 \pm 0.05 \pm 0.09) \times 10^{-4}$ ,  $BR(\tau^- \to K^- \pi^0 \eta v_{\tau}) = (4.6 \pm 1.1 \pm 0.4) \times 10^{-5}$ , and  $BR(\tau^- \to K_s^0 \pi^- \eta v_{\tau}) = (4.4 \pm 0.7 \pm 0.2) \times 10^{-5}$ . For the  $\tau^- \to K^* (892)^- \eta v_{\tau}$  decays the  $\eta$ s are selected within a mass window, and a fit is applied to the  $K_s^0 \pi^-$  and  $K^- \pi^0$  invariant mass to extract  $BR(\tau^- \to K^* (892)^- \eta v_{\tau}) = (1.34 \pm 0.12 \pm 0.09) \times 10^{-4}$ .

The BABAR Collaboration has recently published measurements of the one prong decay channels:  $\tau \to e v_{\tau} \overline{v}_e$ ;  $\tau \to \mu v_{\tau} \overline{v}_{\mu}$ ;  $\tau \to \pi v_{\tau}$  and  $\tau \to K v_{\tau}$  [8]. This was a blind analysis that used a

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 $\tau^- \to \pi^- \pi^- \pi^+ v_{\tau}$  decay to tag, instead of the typical leptonic tag, and required an isolated signal track in the signal hemisphere that was topology consistent with coming from a one prong  $\tau$  decay. A high purity selection is achieved by applying a tight particle identification criteria to suppress cross-feed between channels and other  $\tau$  backgrounds. The backgrounds, cross-feed and particle identification are carefully estimated using a variety of Monte Carlo and data control samples. To reduce systematics that are common between channels, *BABAR* measured the  $BR(\tau \to \mu v_\tau \overline{v_\mu})/BR(\tau \to e v_\tau \overline{v_e}) = (0.9796 \pm 0.0016 \pm 0.0035)$ ,  $BR(\tau \to \pi v_\tau)/BR(\tau \to e v_\tau \overline{v_e}) = (0.5945 \pm 0.0014 \pm 0.0061)$ ,  $BR(\tau \to K v_\tau)/BR(\tau \to e v_\tau \overline{v_e}) = (0.06531 \pm 0.00056 \pm 0.00093)$ . Using the world average  $BR(\tau \to e v_\tau \overline{v_e}) = (17.82 \pm 0.05) \times 10^{-2}$ , the branching fractions  $BR(\tau \to \mu v_\tau \overline{v_\mu}) = (17.46 \pm 0.09) \times 10^{-2}$ ,  $BR(\tau \to \pi v_\tau) = (10.59 \pm 0.11) \times 10^{-2}$  and  $BR(\tau \to K v_\tau) = (6.92 \pm 0.12) \times 10^{-3}$  were determined.

#### 2. Cabibbo-Kobayashi-Maskawa matrix element $|V_{us}|$ from $\tau$

The CKM matrix element  $|V_{us}|$  can be determined from  $\tau$  decays using multiple techniques. One technique to extract  $|V_{us}|$  from  $\tau$  is to use flavor breaking difference with Finite Energy Sum Rules to extract  $|V_{us}|$ :

$$\frac{R_{\tau,strange}}{|V_{us}|^2} - \frac{R_{\tau,nonstrange}}{|V_{ud}|^2} = \delta R_{\tau,SU3} \text{ breaking}$$

where  $R_{\tau,strange} = \Gamma(\tau^- \to X_{strange}v_{\tau})/\Gamma(\tau \to ev_{\tau}\overline{v}_e)$  is the strange hadronic width,  $R_{\tau,nonstrange} = \Gamma(\tau^- \to X_{nonstrange}v_{\tau})/\Gamma(\tau \to ev_{\tau}\overline{v}_e)$  is the nonstrange hadronic width and  $\delta R_{\tau,SU3}$  breaking is the theoretical SU(3)flavor breaking correction determined using Operator Product Expansion (OPE). Using the HFAG values shown in Table 2, which combine the recent measurements from BELLE and *BABAR* with previous  $\tau$  results, one obtains  $|V_{us}| = 0.2188 \pm 0.0023$ . The FESR approach for extracting  $|V_{us}|$  has theoretical errors which range from 0.23%-0.47% depending on the FESR weight. This is the smallest theoretical error of the  $\tau$  decay methods and smaller than the theoretical error from  $K_{l3}$  decays (0.58%)[9] and  $K_{l2}$  decays (0.5%)[10]. If the strange branching fractions and spectral functions are updated with the data currently available at BELLE and *BABAR* this method has the potential for making the most precise measurement of  $|V_{us}|$  [11].

Another method which can be used to determine  $|V_{us}|$  from  $\tau$  decays is from the ratio:

$$\frac{BR(\tau \to K v_{\tau})}{BR(\tau \to \pi v_{\tau})} = \frac{f_K^2 |V_{us}|^2}{f_{\pi}^2 |V_{ud}|^2} \frac{\left(1 - \frac{m_K^2}{m_{\tau}^2}\right)^2}{\left(1 - \frac{m_{\pi}^2}{m_{\tau}^2}\right)^2} (1 + \delta_{LD}),$$

where  $f_K/f_{\pi} = 1.189 \pm 0.007$  [12] is determined from lattice QCD,  $|V_{ud}|$  [13], and the longdistance correction  $\delta_{LD} = (0.03 \pm 0.44)\%$  is estimated [14] using corrections to  $\tau \rightarrow hv_{\tau}$  and  $h \rightarrow \mu v_{\mu}$  [15, 16]. This method is orthogonal to the inclusive sum of strange  $\tau$  decays approach and has a theoretical uncertainty of 0.5%. Using this method, BABAR [8] determined  $|V_{us}| =$  $0.2255 \pm 0.0024$ . This value is consistent with CKM unitarity [13] and 2.5 $\sigma$  higher than  $|V_{us}|$ from the inclusive sum of strange  $\tau$  decays.

A third method for determining  $|V_{us}|$  from  $\tau$  decays is from the  $BR(\tau^- \to K^- \nu_{\tau})$  directly.

$$BR(\tau^{-} \to K^{-} \nu_{\tau}) = \frac{G_F^2 f_K^2 |V_{us}|^2 m_{\tau}^3 \tau_{\tau}}{16\pi\hbar} \left(1 - \frac{m_K^2}{m_{\tau}^2}\right)^2 S_{EW},$$

Decay Mode	Branching Fraction (%)	BELLE	BABAR	
$\frac{1}{BR(\tau \to K \nu_{\tau})}$	0.697±0.010		[8]	
$BR(\tau^- \to K^- \pi^0 v_{\tau})$	$0.431 \pm 0.015$		[1]	
$BR(\tau^- \to K^- \pi^0 \pi^0 \nu_\tau \ (ex. \ K^0))$	$0.060 \pm 0.022$			
$BR(\tau^- \to K^- \pi^0 \pi^0 \pi^0 \nu_\tau \ (ex. \ K^0, \eta))$	$0.039 \pm 0.022$			
$BR(\tau^- \to K^0 \pi^- v_{\tau})$	$0.831 \pm 0.018$	[3]	[2]	
$BR(\tau^- \to K^0 \pi^- \pi^0 \nu_{\tau})$	$0.350 \pm 0.015$		[17]	
$BR(\tau^- \to K^0 \pi^- \pi^0 \pi^0 \nu_{\tau})$	$0.031 \pm 0.023$			
$BR(\tau^- \to K^0 h^- h^- h^+ \nu_{\tau})$	$0.029 \pm 0.020$			
$BR(\tau^- \to K^- \pi^- \pi^+ \nu_\tau \ (ex. \ K^0))$	$0.294 \pm 0.007$	[4]	[5]	
$BR(\tau^- \to K^- \pi^- \pi^+ \pi^0 v_\tau \ (ex. \ K^0, \eta))$	$0.078 \pm 0.012$			
$BR(\tau^- \to K^- \eta \nu_{\tau})$	$0.016 \pm 0.001$	[7]		
$BR(\tau^- \to K^- \eta \pi^0 v_{\tau})$	$0.005\pm0.001$	[7]		
$BR(\tau^- \to K^0 \eta \pi^- v_{\tau})$	$0.009 \pm 0.001$	[7]		
$BR(\tau^- \to K^- K^- K^+ \nu_{\tau})$	$0.0024 \pm 0.0008$	[4]	[5]	
$BR(\tau^- \to K^- K^0 K^0 \nu_{\tau}) \text{ from } \tau^- \to K^- K^- K^+ \nu_{\tau} \times \frac{\phi \to K^0 K^0}{\phi \to K^- K^+}$	$0.0015 \pm 0.0001$			
Total	$2.9\overline{1(55)\pm0.05(10)}$			
Branching Fractions from HFAG constrained fit [18] $\chi^2$ /d.o.f.=135.2/115 CL=9.6%				

**Table 2:** The current status of the branching fraction for the strange  $\tau$  decays.

Using the kaon decay constant  $f_K = 157 \pm 2$  MeV [12], and  $S_{EW} = 1.0201 \pm 0.0003$  [19], BABAR determined  $|V_{us}| = 0.2193 \pm 0.0032$ . This method has a theoretical uncertainty of 1.27%.

Figure 1 shows a summary of  $|V_{us}|$  measurements, including the HFAG values of  $|V_{us}|$  for the three methods mentioned above.



**Figure 1:** A summary of  $|V_{us}|$  measurements. The FlaviaNet measurements come from [10], the hyperon decays come from [20], the BABAR  $\tau$  measurements come from [8] and the HFAG results come from [18].

### 3. Charged Lepton Universality

Charged lepton universality  $\mu - e$  can be tested using

$$\left(\frac{g_{\mu}}{g_{e}}\right)_{\tau}^{2} = \frac{BR(\tau^{-} \to \mu^{-} \overline{\nu}_{\mu} \nu_{\tau})}{BR(\tau^{-} \to e^{-} \overline{\nu}_{e} \nu_{\tau})} \frac{f(m_{e}^{2}/m_{\tau}^{2})}{f(m_{\mu}^{2}/m_{\tau}^{2})},$$

where  $f(x) = 1 - 8x + 8x^3 - x^4 - 12x^2 \log x$ , assuming that the neutrino masses are negligible [21]. The *BABAR* paper [8] gives  $\left(\frac{g_{\mu}}{g_e}\right)_{\tau} = 1.0036 \pm 0.0020$ , yielding a new world average from  $\tau$  decays of  $1.0018 \pm 0.0014$  [18]. Figure 2 shows the current status of  $\mu - e$  lepton universality.

 $\tau$ - $\mu$  universality is tested with

$$\left(\frac{g_{\tau}}{g_{\mu}}\right)_{h}^{2} = \frac{BR(\tau \to h\nu_{\tau})}{BR(h \to \mu\nu_{\mu})} \frac{2m_{h}m_{\mu}^{2}\tau_{h}}{(1+\delta_{h})m_{\tau}^{3}\tau_{\tau}} \left(\frac{1-m_{\mu}^{2}/m_{h}^{2}}{1-m_{h}^{2}/m_{\tau}^{2}}\right)^{2}$$

where the radiative corrections are  $\delta_{\pi} = (0.16 \pm 0.14)\%$  and  $\delta_{K} = (0.90 \pm 0.22)\%$  [15]. BABAR [8] determined  $\left(\frac{g_{\tau}}{g_{\mu}}\right)_{\pi(K)} = 0.9856 \pm 0.0057 \ (0.9827 \pm 0.0086)$  and  $\left(\frac{g_{\tau}}{g_{\mu}}\right)_{h} = 0.9850 \pm 0.0054$ , where the world averaged mass, lifetime values and meson decay rates were taken from [22]. A summary of  $\tau$ - $\mu$  universality tests are shown in Figure 2. It is interesting to note, that the measurements of  $\left(\frac{g_{\tau}}{g_{\mu}}\right)_{h}$  for both kaon and pions are low compared to the HFAG values from the constrained fit and Standard Model expectation.



**Figure 2:** (left) The current status of  $g_{\mu}/g_e$  lepton universality measurements. The HFAG average is the weighted average of previous  $\tau$  results with the recent *BABAR*  $g_{\mu}/g_e$  measurement. (right) The current status of  $g_{\tau}/g_{\mu}$  lepton universality measurements. The HFAG values are presented for:  $g_{\tau}/g_{\mu}$  determined from  $\Gamma(\tau^- \to e^- v_{\tau} \bar{v}_e) \times \tau_{\mu}/\tau_{\tau}$  using the HFAG fit for  $\tau^- \to e^- v_{\tau} \bar{v}_e$ ;  $(g_{\tau}/g_{\mu})_{\pi}$  using the HFAG fit value for  $\tau^- \to \pi^- v_{\tau}$ ;  $(g_{\tau}/g_{\mu})_K$  using the HFAG fit value for  $\tau^- \to K^- v_{\tau}$ ; the average of the HFAG  $(g_{\tau}/g_{\mu})_h$  values; and the HFAG average for  $\tau$  decays.

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