

Initial state anisotropies in ultrarelativistic heavy-ion collisions from the Monte Carlo Glauber model

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In hydrodynamical modeling of heavy-ion collisions the initial state spatial anisotropies translate into momentum anisotropies of the final state particle distributions. Thus, understanding the origin of the initial anisotropies and quantifying their uncertainties is important for the extraction of specific QCD matter properties, such as viscosity, from the experimental data. In this work we study the wounded nucleon approach in the Monte Carlo Glauber model framework, focusing especially on the uncertainties which arise from the modeling of the nucleon-nucleon interactions between the colliding nucleon pairs and nucleon-nucleon correlations inside the colliding nuclei. We compare the black disk model and a probabilistic profile function approach for the inelastic nucleon-nucleon interactions, and study the effects of initial state correlations using state-of-the-art modeling of these.

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1. Introduction

In ultrarelativistic heavy-ion collisions performed at the Relativistic Heavy Ion Collider (RHIC) and Large Hadron Collider (LHC) significant azimuthal momentum distribution anisotropies have been measured [1, 2]. These anisotropies can be explained with relativistic hydrodynamics: the initially produced QCD-matter contains spatial anisotropies and during the hydrodynamical evolution these anisotropies are transferred to the momentum distributions of final state particles.

Simulations with viscous hydrodynamics have shown that the shear viscosity of the QCD-matter produced in ultrarelativistic heavy-ion collisions can be estimated from the final state momentum anisotropies. Since the origin of these anisotropies is in the initial state, uncertainties related to the initial state must be charted before reliable estimates for the viscosity can be made. In this work [3] we consider two sources of uncertainties related to the Monte Carlo Glauber (MCG) model which is often used to initialise the hydrodynamical simulations (see *e.g.* [4], [5], [6]).

The Glauber model [7] is usually a key element in computing the initial states for hydrodynamical modeling of ultrarelativistic heavy-ion collisions. Some years back, most hydrodynamical calculations assumed smooth initial states where the (energy or entropy) densities were assumed to scale with the density of binary collisions or wounded nucleons computed from the optical Glauber model; see, *e.g.* [8]. Now that the importance of the initial density fluctuations has been realized, Monte Carlo Glauber (MCG) modeling has become more frequently used. So far the black disk (hard-sphere) modeling of the nucleon-nucleon (NN) interactions has been the standard choice [4, 9, 10, 11] in these studies, although also more involved probabilistic ways to model the NN interactions have been known for a long time [7, 12, 13, 14]. In Ref. [3] we studied *black disk* and *profile function* models.

In the MCG modeling one needs to know the positions of the initial state nucleon configurations. In most cases the nucleon positions inside the colliding nuclei are just sampled using the Woods-Saxon potential, neglecting nucleon-nucleon correlations [15]. However, there exist calculations which show that high-momentum components of the nuclear wave function are generated by the two-body NN correlations [16, 17].

In Ref. [3] we studied for Au-Au collision at RHIC center of mass energy $\sqrt{s_{NN}}=200$ GeV, two different uncertainties in computing the initial state asymmetries from the MCG model: one related to the modeling of the inelastic NN collisions between nucleons from different nuclei, and one related to the NN correlations in the nucleon configurations in each of the colliding nuclei. In this contribution we extend the results from Ref. [3] about the anisotropy moments ϵ_n with $n = 1, 2, 3, 4, 5$: dipole asymmetry, eccentricity, triangularity [3], quadangular and pentagonal asymmetries. Results will be shown for Au+Au collisions at RHIC with an inelastic nucleon-nucleon cross section $\sigma_{NN} = 42$ mb.

2. MCG framework: nucleon configurations

The initial state of a nucleus in the MCG calculations is usually taken as a collection of particles distributed according to a probability distribution given by the corresponding (Woods-Saxon) number density distribution measured in electron scattering experiments. Given the complexity of the nuclear many-body problem, the effects of spatial, spin and isospin dependent correlations

among the nucleons are usually overlooked and the nucleons are positioned randomly for each of the simulated events. Recently, in Ref. [18] it was shown how such an approach can be modified by including initial states, which are prepared in advance, in the commonly used computer codes. Also, a method to produce such configurations was introduced. The method is based on the notion of a nuclear wave function ψ , which contains the nucleonic degrees of freedom and which is used to iteratively modify the positions of randomly distributed nucleons using the Metropolis method so that the final positions correspond to the probability density given by $|\psi|^2$. The method is constructed to reproduce the same nucleon number-density distribution as the usual one and, in addition, to reproduce the basic features of the two-nucleon density in the presence of the NN correlations. The model wave function is taken in the form

$$\psi(x_1, \dots, x_A) = \prod_{i < j}^A \hat{f}_{ij} \phi(x_1, \dots, x_A), \quad (2.1)$$

where ϕ is the uncorrelated wave function and \hat{f}_{ij} are correlation operators; here, x_i denotes the position, spin and isospin projection of the i -th nucleon. The correlation operator contains a detailed spin-isospin dependence. In the most general case, this dependence includes a number of channels that are the same as the one appearing in modern nucleon-nucleon potentials used to successfully describe a variety of properties of light and medium-heavy nuclei within different *ab initio* approaches. The major merit of this approach is that the two-body densities resulting from configurations obtained using these correlations are clearly more realistic than the completely uncorrelated ones [19]. In this paper we have used configurations generated by two-body correlations, including the tensor operator, and three-body clusters surrounding each of the nucleons, induced by full correlations. The effects of genuine three-body correlations have been discussed in Ref. [3].

2.1 MCG framework: modeling the inelastic interactions

In each simulated event, given the impact parameter of the $A + A$ collision, the nucleon-nucleon interactions must be modeled. We work in the Glauber model framework [7], neglecting the effects of inelastic diffraction that lead to fluctuations of the strength of the NN interactions [20, 21]. To generate the inelastic NN collisions of interest here, we use the following two different approximations for establishing whether a collision between the nucleons i and j from different nuclei takes place:

- *Black disk approximation*, used recently, e.g., in Ref. [4], where one assumes the two nucleons to interact inelastically with a probability one if their transverse separation b_{ij} is within a radius defined by the inelastic NN cross section σ_{NN}^{in} ,

$$b_{ij}^2 \leq \frac{\sigma_{NN}^{in}}{\pi}; \quad (2.2)$$

- *Profile function approach*, where the probability of an inelastic interaction between the nucleons i and j is given by

$$P(b_{ij}) = 1 - |1 - \Gamma(b_{ij})|^2, \quad (2.3)$$

and where the profile function Γ is expressed in terms of the total and elastic NN cross sections as follows:

$$\Gamma(b_{ij}) = \frac{\sigma_{NN}^{tot}}{4\pi B} e^{-b_{ij}^2/(2B)}, \quad (2.4)$$

with $B = (\sigma_{NN}^{tot})^2 / (16\pi\sigma_{NN}^{el})$.

The probability distribution $P(b_{ij})$ can be derived in the Born approximation of the potential-scattering formalism [7, 20, 22]; details are given in Ref. [3]. The nucleon-nucleon elastic, total and inelastic cross sections are given in this formalism by:

$$\sigma_{NN}^{el} = \frac{(\sigma_{NN}^{tot})^2}{16\pi B}, \quad (2.5)$$

$$\sigma_{NN}^{tot} = \sigma_{NN}^{el} + \sigma_{NN}^{in}, \quad (2.6)$$

$$\sigma_{NN}^{in} = \int d^2b_{ij} \left(1 - |1 - \Gamma(b_{ij})|^2\right). \quad (2.7)$$

After the nucleon-nucleon interactions have been determined, we can calculate the initial asymmetries from the positions of the nucleons which had experienced at least one collision. These are called wounded (or participant) nucleons.

2.2 Spatial asymmetries and their fluctuations

We compute the spatial asymmetries from the wounded nucleon positions which are obtained from the MCG model as follows:

$$\varepsilon_n = -\frac{\langle w(r) \cos(n(\phi - \psi_n)) \rangle}{\langle w(r) \rangle}, \quad (2.8)$$

where $w(r)$ is a weight and we choose $w(r) = r^2$ for $n = 1, 2, 3, 4, 5$, following Ref. [23] (at variance with Ref. [3] where we used $w(r) = r^3$ for $n = 1$ [26] and $w(r) = r^n$ for $n = 2, 3$).

The orientation angle ψ_n is determined as

$$\psi_n = \frac{1}{n} \arctan \frac{\langle w(r) \sin(n\phi) \rangle}{\langle w(r) \cos(n\phi) \rangle} + \frac{\pi}{n}, \quad (2.9)$$

where the arctan must be placed in the correct quadrant. Since the experimental methods typically measure the root mean square (rms) of the flow coefficients, we also present the rms values of initial state anisotropies $\sqrt{\langle \varepsilon_n^2 \rangle}$ [23].

3. Results and discussion

In Ref. [3] we charted some of the uncertainties in the computation of the initial state anisotropies from the Monte Carlo Glauber model. A summary of the results is plotted in Fig. 1, now using rms values. Also harmonics $n = 4, 5$ [24, 25], which were not shown in the original article, are shown here. We used two different ways of modeling the inelastic interactions between the colliding nucleons. The difference between these two cases gives us an estimate about the uncertainties related to this part of the model: in central collisions the details of the interaction model play a minor role, but in the peripheral collisions such details can cause uncertainties up to 10% in the first three harmonics ε_1 , ε_2 and ε_3 . The situation is similar for the rms of ε_n , as shown in Fig. 1, where $n = 1, 2, 3, 4, 5$ are considered with the choice of the weight function, $w(r) = r^2$; the effect is at about 10% over the N_{part} range and it is mostly due to the probabilistic profile function approach. We also checked in Ref. [3] that with these two interaction models the difference in

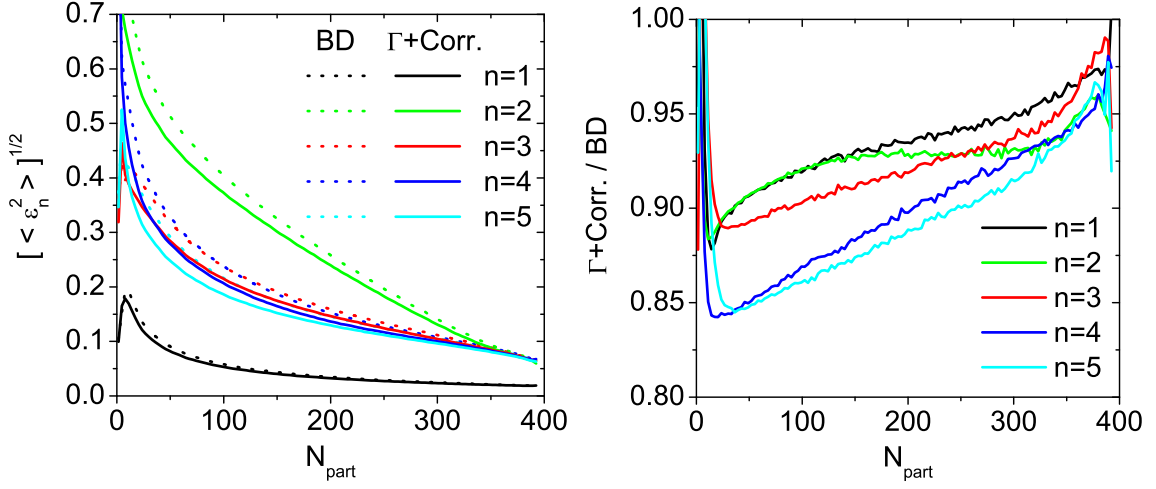


Figure 1: Left panel: the root mean square of the initial anisotropies, $\sqrt{\langle \varepsilon_n^2 \rangle}$, defined in Eq. (2.8), as a function of the number of participants. Results are relative to the black disk, uncorrelated approximation (dotted lines) and probabilistic interaction plus correlated configurations (solid lines). Right panel: the ratio of the Γ +correlations over black disk, uncorrelated results shown in the left panel.

the number of wounded nucleons and binary collisions remains small in central collisions, but at impact parameters 10-15 fm the difference can be around 10%.

We also presented a study of the effects of NN correlations with an update of correlated configurations and extended discussion as compared with the previous published papers on this subject. We confirmed that the inclusion of centrally correlated nucleon configurations produce the effects to eccentricity and its relative variance as was claimed by Ref. [28]. As a new result, we observed that the inclusion of realistically correlated configurations (two-body full correlations, three-body chains) seems to essentially cancel this effect and bring the results back close to the no correlations case. The effect is similar for dipole asymmetry and triangularity as for eccentricity. However, we also showed that there are still uncertainties caused by the truncation done in the nucleon configuration calculation with full correlations and we expect three-body correlations to play a role.

In this Proceedings contributions, we added the study of the root mean square of initial state anisotropies for $n = 1, 2, 3, 4, 5$ and presented results in Fig. 1, with similar sensitivity to the two sources of uncertainties studied, which was found to be maximally of the order of 10%. Now that, thanks to the recent developments in event-by-event hydrodynamics and high-precision data, more precise comparisons of flow coefficients between the data and the theory are becoming possible, it is important to quantify all the relevant uncertainties to this precision, so that the QCD matter properties could eventually be determined from the measured particle spectra and their azimuthal asymmetries.

References

- [1] Y. Schutz, U. Wiedemann, CERN Cour. September 2011.

- [2] K. Aamodt *et al.* (ALICE Collaboration), Phys. Rev. Lett. **107**, 032301 (2011).
- [3] M. Alvioli, H. Holopainen, K. J. Eskola and M. Strikman, Phys. Rev. **C85**, 034902 (2012).
- [4] H. Holopainen, H. Niemi, K. J. Eskola, Phys. Rev. **C83**, 034901 (2011).
- [5] G.-Y. Qin, H. Petersen, S. A. Bass, B. Muller, Phys. Rev. **C82**, 064903 (2010).
- [6] M. Rybczynski and W. Broniowski, Phys. Rev. C **84**, 064913 (2011).
- [7] R. J. Glauber, in: Lectures in theoretical physics, ed. W. E. Brittin et al. (Interscience Publishers, New York, 1959) vol. I, p. 315.
- [8] P. F. Kolb, U. W. Heinz, P. Huovinen, K. J. Eskola, K. Tuominen, Nucl. Phys. **A696**, 197-215 (2001).
- [9] B. Alver, M. Baker, C. Loizides, P. Steinberg, arXiv:0805.4411 [nucl-ex].
- [10] B. Alver, B. B. Back, M. D. Baker, M. Ballintijn, D. S. Barton, R. R. Betts, R. Bindel, W. Busza *et al.*, Phys. Rev. **C77**, 014906 (2008).
- [11] T. Hirano, Y. Nara, Phys. Rev. **C79**, 064904 (2009).
- [12] H. Pi, Comput. Phys. Commun. **71**, 173-192 (1992).
- [13] X.-N. Wang, Phys. Rev. **D43**, 104-112 (1991).
- [14] R. J. Glauber, G. Matthiae, Nucl. Phys. **B21**, 135-157 (1970).
- [15] R. Subedi *et al.*, Science **320**, 1476-1478 (2008).
- [16] M. Alvioli, C. Ciofi degli Atti, H. Morita, Phys. Rev. **C72**, 054310 (2005).
- [17] M. Alvioli, C. Ciofi degli Atti and H. Morita, Phys. Rev. Lett. **100**, 162503 (2008).
- [18] M. Alvioli, H.-J. Drescher, M. Strikman, Phys. Lett. **B680**, 225-230 (2009).
- [19] M. Alvioli, M. Strikman, Phys. Rev. **C83**, 044905 (2011).
- [20] B. Blaettel, G. Baym, L. L. Frankfurt, H. Heiselberg, M. Strikman, Phys. Rev. **D47**, 2761-2772 (1993).
- [21] G. Baym, B. Blattel, L. L. Frankfurt, H. Heiselberg, M. Strikman, Phys. Rev. **C52**, 1604-1617 (1995).
- [22] M. Alvioli, C. Ciofi degli Atti, I. Marchino, V. Palli, H. Morita, Phys. Rev. **C78**, 031601 (2008).
- [23] R. S. Bhalerao, M. Luzum and J.-Y. Ollitrault, Phys. Rev. C **84**, 054901 (2011).
- [24] N. Borghini, J.-Y. Ollitrault, Phys. Lett. **B642**, 227-231 (2006).
- [25] J.-Y. Ollitrault, A. M. Poskanzer, S. A. Voloshin, Phys. Rev. **C80**, 014904 (2009).
- [26] D. Teaney, L. Yan, Phys. Rev. **C83**, 064904 (2011).
- [27] B. H. Alver, C. Gombeaud, M. Luzum and J.-Y. Ollitrault, Phys. Rev. C **82**, 034913 (2010).
- [28] W. Broniowski, M. Rybczynski, Phys. Rev. **C81**, 064909 (2010).