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Particle production at HERA

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> Results on particle production in deep-inelastic scattering in *ep* collision at HERA, obtained with the H1 and ZEUS detectors, are presented. The underlying parton dynamics are investigated by looking at the hardness of the transverse momentum distribution of charged particles and comparing the measurements with Monte Carlo generators based on different approaches for simulating the parton cascade. Studies of the scaled momentum distribution for for K_S^0 and Λ particles in the current fragmentation region of the Breit frame in the context of hadronisation are presented. The data are compared to models and to next-to-leading order QCD calculations. In addition, the production of photon at very small angles with respect to the proton beam direction is presented. The cross sections are compared to the predictions of models of deep-inelastic scattering and models of the hadronic interactions of high energy cosmic rays.

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1. Parton dynamics

1.1 Charged particle transverse momentum spectra

It was suggested in [1] that the transverse momentum spectrum of charged particles is sensitive to whether partons are emitted in a p_T -ordered cascade (DGLAP) [2] or in an unordered way (beyond DGLAP). The H1 Collaboration has presented measurements [3] of the transverse momenta and pseudorapidities of charged particles and compared them to predictions of various Monte Carlo (MC) generators using different approaches to simulate the parton cascade: the RAPGAP generator [4] based on leading-log DGLAP parton showers; the DJANGOH [5] MC which uses the Colour Dipole Model (CDM) [6] as implemented in ARIADNE [7], which provides a BFKL-like [8, 9] approach; the CASCADE [10] MC generator based on the CCFM model [11]. All generators use the Lund string model [12] for hadronisation as implemented in PYTHIA [13]. The fragmentation parameter set are tuned by the ALEPH collaboration to fit LEP data [14].

The analysed data set corresponds to an integrated luminosity of $L = 88.64 \text{ pb}^{-1}$. The analysis is performed in the kinematic phase space region defined by $5 < Q^2 < 100 \text{ GeV}^2$ and 0.05 < y < 0.6. The data are corrected for detector effects using MC event samples, as well as for the charged decay products of K_S^0 , Λ and for other weakly decaying particles. The transverse momentum and pseudorapidity distributions are presented in the hadronic centre-of-mass system (HCM), i.e. in the rest frame of the virtual photon and proton, and are labeled as p_T^* and η^{*1} , respectively. All distributions shown are normalised to the total number of DIS events in the analysed phase space.

The normalised pseudorapidity distributions for $p_T^* < 1$ GeV and for $p_T^* > 1$ GeV, separately, are shown in Fig. 1. The data are compared to the RAPGAP, DJANGOH and CASCADE generators. As argued in [1], hadronisation effects should be relevant at small p_T^* , while hard parton radiation should manifest itself in the tail of the p_T^* distribution. To check the sensitivity to hadronisation effects, RAPGAP predictions with default PYTHIA fragmentation parameters and with parameters tuned by ALEPH are shown in Fig. 1. Significant differences between these two models are seen in the soft p_T^* region, while for particles with $p_T^* > 1$ GeV this difference is much smaller. Predictions from generators with different approaches for QCD radiation are also shown in Fig. 1. The best description of the data is achieved by DJANGOH. At large p_T^* RAPGAP predicts a spectrum which



Figure 1: Measured η^* spectra in the HCM system for the charged particles with $p_T^* < 1$ GeV (left) and with $p_T^* > 1$ GeV (right) together with RAPGAP, DJANGOH and CASCADE Monte Carlo predictions.

¹The definition of pseudorapidity in the HCM frame is $\eta^* = -\ln(\tan(\theta^*/2))$, where θ^* is the angle with respect to the virtual photon direction, i.e. in the positive z^* direction.

is too soft towards small η^* . CASCADE predicts too high multiplicities away from the photon direction.

The normalised transverse momentum spectra of charged particles in the pseudorapidity interval of $0 < \eta^* < 1.5$ and in eight different x and Q² bins are presented in Fig. 2. In the range



Figure 2: Measured p_T^* spectra of charged particles in the hadronic centre-of-mass system (HCM) in $0 < \eta^* < 1.5$ range for eight intervals of Q² and x_{Bj} together with RAPGAP, DJANGOH and CASCADE Monte Carlo predictions.

 $0 < \eta^* < 1.5$ the sensitivity to the parton shower may be studied, while for higher η^* values the sensitivity to the hard scatter is mainly expected. The CDM model provides a good description of the data over the full kinematic range. The RAPGAP strongly undershoots the data in the lowest *x* and Q² bin. The description becomes better at larger values of *x* and Q². CASCADE describes the data in the lowest *x* and Q² bin and high p_T^* but has a much harder spectrum at high *x* and Q².

2. Hadronisation

2.0.1 $K_{\rm S}^0$ and Λ scaled momentum distributions

An investigation of the production of K_S^0 mesons and Λ baryons in the current fragmentation region of the Breit frame is presented. Multiplicity distributions are measured as a function of Q^2 per unit of the scaled momentum, $x_p = 2p^{\text{Breit}}/Q$. Here, p^{Breit} denotes the momentum of a hadron in the Breit frame. The aim is to check the universality of the quark fragmentation function and the factorisation theorem approach used to predict hadron production in different processes.

The data presented in this section are corrected for detector acceptance, efficiency and resolution effects. All distributions shown are normalised to the total number of DIS events, N, in the analysed phase space. Scaled momentum distributions for K_S^0 mesons and Λ baryons were measured by ZEUS [15] in the kinematic range $10 < Q^2 < 40000 \text{ GeV}^2$ and 0.001 < x < 0.75. The data correspond to an integrated luminosity of $L = 290 \text{ pb}^{-1}$.

Figure 3 shows the scaled momentum spectra, as a function of Q^2 in different regions of x_p , for K_S^0 (left) and Λ (right) production. Large scaling violations are seen: with increasing Q^2 the phase



Figure 3: Scaled momentum distributions as a function of Q^2 in different regions of x_p for K_S^0 (left) and Λ (right) and together with predictions of NLO QCD and LO Monte Carlo programs.

space for soft gluon radiation gets enlarged, which leads to a rise in the number of soft particles at small x_p and to a decrease in the number of those with high x_p .

The data are compared to two sets of the NLO QCD predictions based on different parametrisation of the fragmentation functions. The first set was obtained from fits to e^+e^- data and based on the program CYCLOPS [16], called "AKK+CYCLOPS" [17]. The second set was obtained from a global fit to e^+e^- , pp and ep data, called "DSS" [18]. Both NLO QCD predictions describe the data only in certain regions of the phase space. For the AKK+CYCLOPS calculation this is limited to $0.6 < x_p < 1$, whereas the DSS calculation describes the K_S^0 data adequately, except for regions of low x_p and Q^2 . The scaling violation in both NLO calculations is less pronounced than in the data. Together with the NLO QCD calculations the predictions of MC generators ARIADNE based on CDM and the MEPS model of LEPTO [19], based on DGLAP parton shower, are shown in Fig. 3. Both predictions give a reasonable description of the data.

2.0.2 Production of very forward photon in DIS

Measurements of particle production at very small angles with respect to the proton beam direction (forward direction) in positron-proton collisions are important for the understanding the fragmentation of the proton remnant. These measurements also provide important constraints for the modeling of the high energy air showers.

In this analysis [20] photons are detected at very small angles, $\eta > 7.9$. It relies on the upgraded H1 Forward Neutron Calorimeter (FNC) [21]. The FNC consist of the Main Calorimeter and the Preshower Calorimeter. Events are restricted to the kinematic range $6 < Q^2 < 100 \text{ GeV}^2$ and 0.05 < y < 0.6. Forward photons are required to have a signal only in the Preshower calorimeter, and to have $\eta > 7.9$ and an energy above 92 GeV, which corresponds to the longitudinal momentum fraction $x_L = E_{\gamma}/E_p > 0.1$, where E_p and E_{γ} are the proton beam and forward photon energy, respectively. The data are corrected for the detector effects and QED radiation.

The normalised differential cross sections for the most energetic forward photons (leading photon) are presented in Fig. 4 as a function of x_L^{lead} . The data are compared to the predictions of the MC generators ARIADNE based on CDM and LEPTO as well as the predictions of several hadronic interaction models which are commonly used for the simulation of cosmic ray air shower cascades: EPOS [22], QGSJET 01 [23], QGSJET II-03 [24] and SIBYLL [25]. In all of these models, the main source of forward photons is the decay of π^0 mesons produced from the hadronisation of the proton remnant. The ratios of MC model predictions to the measurements are shown separately. All models overestimate the total rate of forward photons. The best description is achieved by



Figure 4: Normalised cross sections of forward photon production as a function of x_L^{lead} . The data are compared to two predictions of the DJANGOH Monte Carlo simulation, using LEPTO and CDM to simulate higher orders. The lower row shows the ratios of the Monte Carlo predictions to the data.

the QGSJETs models, which are consistent with the data at high x_L^{lead} values, but in general have steeper behavior than data. The shapes of all measured distributions are well described by LEPTO.

The measurement of forward photons allows a test of the limiting fragmentation hypothesis, according to which the production of forward photons in DIS is insensitive to Q^2 and x_{Bj} . The ratio of the forward photon production cross section to the inclusive DIS cross section is measured as a function of Q^2 and x_{Bj} , as shown in Fig. 5. The fraction of DIS events with forward photons is independent from Q^2 and x_{Bj} in agreement with the limiting fragmentation hypothesis.

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Figure 5: Fraction of DIS events with forward photons as a function of Q^2 and x_{Bj} . The expectation from the LEPTO and CDM models are also shown.

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