

A proposal to solve some puzzles in semileptonic B decays

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About 25% of all B mesons decay semileptonically via the tree-level $b \rightarrow c$ quark transition and allow the precision determination of the Cabibbo-Kobayashi Maskawa matrix element $|V_{cb}|$. Although these decays were the major focus of both B -Factory experiments over the last decade, about 10% of the semileptonic remain poorly known what limits our understanding of rare decays, such as the semileptonic $b \rightarrow u$ transitions. We investigate the possibility of a large contribution from radial excited states to the semileptonic branching fraction and discuss its consequences.

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1. Introduction

About 25% of all B mesons decay semileptonically via the tree-level $b \rightarrow c$ quark transition. The study of such decays allow for the precision determination of the Cabibbo-Kobayashi Maskawa matrix element $|V_{cb}|$ and provide important background estimates or input for the study of rare decays. For the past ten years several puzzling features of this decay mode were observed, which individually are not several sigma problems, but remained unresolved.

In the quark model picture, if one considers $1S$, $1P$, and $2S$ excitations, in total eight mesons should exist: We refer to the ground two charm meson states as $D^{(*)}$, the four orbitally excited as D^{**} , and the radial states as $D'^{(*)}$. Several tensions arise in decay modes involving the two broad members of the $1P$ excitations, either between different measurements or between measurements and theory predictions. The relevant points for our discussion are:

	[%]	HFAG
$B^+ \rightarrow D^0 \ell^+ \nu_\ell$		2.30 ± 0.10
$B^+ \rightarrow \bar{D}^{*0} \ell^+ \nu_\ell$		5.34 ± 0.12
$B^+ \rightarrow \bar{D}_0^0 \ell^+ \nu_\ell$		
$\hookrightarrow D\pi$	0.44 ± 0.08	
$B^+ \rightarrow \bar{D}_1^0 \ell^+ \nu_\ell$		
$\hookrightarrow D^* \pi$	0.20 ± 0.06	
$B^+ \rightarrow \bar{D}_1^0 \ell^+ \nu_\ell$		
$\hookrightarrow D^* \pi$	0.43 ± 0.03	
$B^+ \rightarrow \bar{D}_2^0 \ell^+ \nu_\ell$		
$\hookrightarrow D^{(*)} \pi$	0.28 ± 0.05	
$B^+ \rightarrow \bar{D}_1^0 \ell^+ \nu_\ell$		
$\hookrightarrow D\pi\pi$	0.29 ± 0.08	
$B^+ \rightarrow D^{**}(1P) \ell^+ \nu_\ell$		
$\hookrightarrow D^{(*)} \pi(\pi)$	1.64 ± 0.18	
$\sum D^{(*)} + \sum D^{**}(1P)$		9.28 ± 0.24
$B^+ \rightarrow \bar{D} \pi \ell^+ \nu_\ell$	0.66 ± 0.08	
$B^+ \rightarrow \bar{D}^* \pi \ell^+ \nu_\ell$	0.87 ± 0.09	
$\sum D^{(*)} + \sum D^{**} \pi$		9.17 ± 0.20
Incl. $B^+ \rightarrow X_c \ell \nu$		10.91 ± 0.14
'Gap' Incl. v excl.	1.36 ± 0.30	
'Gap' Incl. v semi-incl./excl.	1.74 ± 0.24	

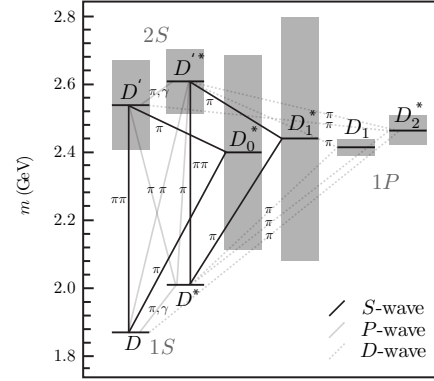


Figure 1: [Left] Branching fractions from HFAG [1] and our private averages. For $B \rightarrow D^{(*)} \pi \ell \bar{\nu}_\ell$: weighted average, assuming a 100% correlation, of both $B^{0,+}$ isospin modes. [Right] Strong decays of the $D'^{(*)}$ involving pion emissions, omitting possible near off-shell transitions with a ρ and η . Grey bands correspond to the measured widths of the excited states.

1. The sum of the measured exclusive rates is less than the inclusive one, where the measurements are listed in Fig. 1. Combining the measurements we have

- With the semi-inclusive $[\sum D^{(*)} + \sum D^* \pi]$ branching fractions the gap is

$$(1.74 \pm 0.24)\%$$

- The measured $1P$ decay $[\sum D^{(*)} + \sum D^{**} \rightarrow D^{(*)} \pi]$ amounts to a gap of

$$(1.36 \pm 0.30)\%$$

- The quoted numbers slightly differ from the ones in our publication [2]: a fraction of $B \rightarrow D_1 \ell \bar{\nu}_\ell$ with $D_1 \rightarrow D\pi\pi$ estimated by the ratio of the non-leptonic $B \rightarrow D_1 \pi$ with $D_1 \rightarrow D\pi\pi$ and $D_1 \rightarrow D^* \pi$ was added and the Belle lower limit on $B \rightarrow D_1' \ell \bar{\nu}_\ell$ was fully included.

2. The exclusive vs. inclusive determination of $|V_{cb}|$ differs as [3]

$$|V_{cb}| = (41.9 \pm 0.7) \times 10^{-3} \quad (\text{inclusive})$$

$$|V_{cb}| = (39.6 \pm 0.9) \times 10^{-3} \quad (\text{exclusive})$$

3. "1/2 vs 3/2 puzzle" [4]: Theory prediction in conflict with data, see Fig. 1

$$\mathcal{B}(B^+ \rightarrow D_{1/2=\text{broad}}^{**} \ell^+ \nu) / \mathcal{B}(B^+ \rightarrow D_{3/2=\text{narrow}}^{**} \ell^+ \nu) \sim 0.1 - 0.2 \quad [\text{Theory}]$$

$$\mathcal{B}(B^+ \rightarrow D_{1/2=\text{broad}}^{**} \ell^+ \nu) / \mathcal{B}(B^+ \rightarrow D_{3/2=\text{narrow}}^{**} \ell^+ \nu) \sim 1 \quad [\text{Data}].$$

Although the size of the gap depends on the actual interpretation of the measurement and the used data, it remains significant in a statistical sense. In experimental analysis this gap is often filled up with a mix of known states, or the analysis is restricted in regions of phase space where decays making up this 'Gap' no longer contributes significantly due to kinematic restrictions. But this is unsatisfactorily for many reasons and a thorough understanding of the matter is highly desirable. Here we investigate the viability of a proposal, which could solve or at least ease some of these tensions.

2. Proposal and its Viability

The allowed decay modes of the considered excitations are displayed in Fig. 1. The most important feature is, that the radially excited $2S$ mode can decay via a s-wave to the orbitally excited broad $1P$ modes, but only via d-waves into the narrow $1P$ states. The pion emitted in this strong decay has a soft momentum of $p_\pi \sim 0.01 - 0.5$ GeV, allowing for the possibility of missing detection. We investigate the possibility of a substantial branching fraction into radial states of the order

$$\mathcal{B}(B \rightarrow D^{(*)} \ell \bar{\nu}) \sim \mathcal{O}(1\%). \quad (2.1)$$

If true, this could ease many of the above mentioned tensions in the following way:

1. It would be sufficient to saturate inclusive rate closing the gap.
2. Decays involving the production of experimentally challenging soft pions could enhance the observed decay rate to the broad states $s_l^{\pi_l} = \frac{1}{2}^+$ states, enhancing their population and thus ease the "1/2 vs 3/2 puzzle"
3. The mass gap of the $1S$ and $2S$ is relatively small and thus the charged lepton energy spectrum stays hard, which is in agreement with observations.
4. There is no direct conflict between the hypothesis and the $\mathcal{B}(B \rightarrow D^{(*)} \pi \ell \bar{\nu})$ measurement: The $D^{(*)}$ decay would yield two or more pions most of the time.

In order to investigate the viability, we want to estimate the possible branching fractions. The decay distributions are the same as for the ground-states $D^{(*)}$ up to different form factors:

$$\begin{aligned} \frac{d\Gamma_{D^*}}{dw} &= \frac{G_F^2 |V_{cb}|^2 m_B^5}{48\pi^3} r^3 (1-r)^2 \sqrt{w^2-1} (w+1)^2 \left[1 + \frac{4w}{w+1} \frac{1-2rw+r^2}{(1-r)^2} \right] [F(w)]^2 \\ \frac{d\Gamma_{D'}}{dw} &= \frac{G_F^2 |V_{cb}|^2 m_B^5}{48\pi^3} r^3 (1+r)^2 (w^2-1)^{3/2} [G(w)]^2. \end{aligned} \quad (2.2)$$

The momentum transfer is small $1 \leq w = v \cdot v' \lesssim 1.3$ and thus we investigate a linear and a quadratic interpolation of the Isgur Wise function in the heavy quark limit. At zero recoil $\xi_2(1) = 0$ and the rate at $w = 1$ comes entirely from Λ_{QCD}/m_b corrections. We expect the slope to be positive, because the only change from $1S$ to $2S$ is, that the expectation value of the distance from the heavy quark of a spherically symmetric wave function is increased. For the estimate we use the quark model estimate [5], hoped to be valid at $w = 1$. We further modify an existing light-cone sum rules calculation [6] to project out the ground-state, with the hope to be reasonably valid at w_{max} . We

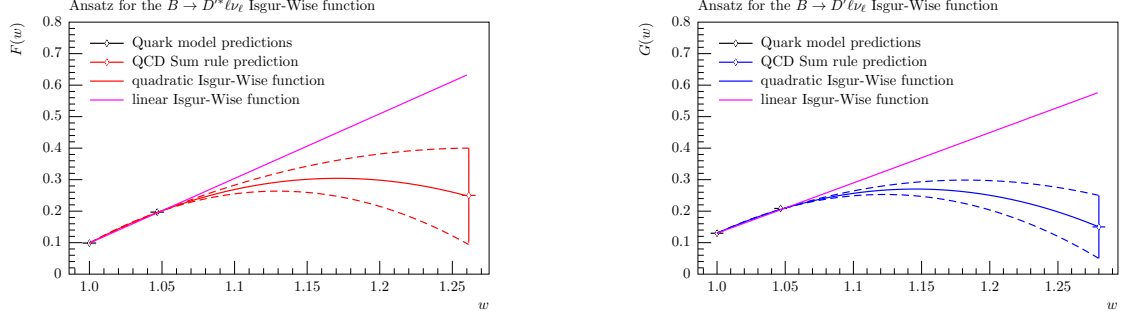


Figure 2: Isgur-Wise function of the $D^{(*)}$ for $F(w)$ [left] and $G(w)$ [right]

obtain from the calculation and the quark model estimate the numbers

$$\begin{aligned} F(1.0) &= 0.10 & F(1.05) &= 0.20 & G(1.0) &= 0.13 & G(1.05) &= 0.21 \\ F(w_{\text{max}}) &= 0.25 \pm 0.15 & G(w_{\text{max}}) &= 0.15 \pm 0.1. \end{aligned} \quad (2.3)$$

With a linear [quadratic] parameterization for the form factors as

$$G(w)\{F(w)\} = \beta_0^{\{*\}} + (w-1)\beta_1^{\{*\}} + [(w-1)^2\beta_2^{\{*\}}] \quad (2.4)$$

we obtain the parameters the values Eq. (2.3), illustrated in Fig. 2,

$$\begin{aligned} \beta_0^* &= 0.10, & \beta_1^* &= 2.1 \\ \beta_0 &= 0.13, & \beta_1 &= 1.6 \\ \beta_0^* &= 0.10, & \beta_1^* &= 2.3 - 2.5, & \beta_2^* &= -(4.2 - 9.8) \\ \beta_0 &= 0.13, & \beta_1 &= 1.9 - 2.0, & \beta_2 &= -(5.1 - 8.2). \end{aligned} \quad (2.5)$$

Using these interpolations, we obtain the branching fractions in the proposed order of magnitude, which would help to ease or solve the puzzles

$$\begin{aligned} \mathcal{B}(B \rightarrow (D' + D'^*)\ell\nu_\ell) &\sim 1.4\% & \text{Linear Interpolation} \\ \mathcal{B}(B \rightarrow (D' + D'^*)\ell\nu_\ell) &\sim (0.3 - 0.7)\% & \text{Quadratic Interpolation.} \end{aligned} \quad (2.6)$$

3. Discussion

If future measurements find a substantial $B \rightarrow D^{(*)}\ell\bar{\nu}$ decay rate, the precise determination of the branching fraction, the shape of the $F(w)$ and $G(w)$ functions in Eq. (2.2), and data on the corresponding nonleptonic two-body decays with a pion would be able to test this picture. Especially

LHCb could contribute by measuring the non-leptonic decay $B \rightarrow D^{(*)}\pi$, which is related to the form factors by factorization

$$\Gamma(B \rightarrow D^{(*)}\pi) = \frac{3\pi^2 C^2 |V_{ud}|^2 f_\pi^2}{m_B m_{D^{(*)}}} \left. \frac{d\Gamma(B \rightarrow D^{(*)}\ell\bar{\nu})}{dw} \right|_{w_{\max}} \quad (3.1)$$

The strong channel $D^{(*)} \rightarrow D^{(*)}\eta$ could be studied by the present day B-Factories. A considerable radial contribution may also impact other measurements and the theory of semileptonic decays, e.g., it may yield

- a better understanding of the $b \rightarrow c$ background in fully inclusive $b \rightarrow u$ measurements, i.e., lead to a more precise determination of $|V_{ub}|$;
- a better understanding of the semileptonic $b \rightarrow c$ background in the exclusive $|V_{cb}|$ measurements using $B \rightarrow D^{(*)}\ell\bar{\nu}$;
- a better understanding of the missing exclusive contributions to the inclusive $B \rightarrow X_c\ell\bar{\nu}$ rate, and the lepton energy and hadronic mass spectrum;
- a better understanding of the measured $B \rightarrow D^{(*)}\tau\bar{\nu}$ branching fraction and its tension with respect to the Standard Model expectation [7];
- a more precise determination of the semileptonic branching fractions of the $s_l^{\pi_l} = \frac{1}{2}^+$ and $\frac{3}{2}^+$ states, thus maybe help resolve the “1/2 vs. 3/2 puzzle”;
- a stronger sum rule bound [[8, 9, 10], [11]] on the $B \rightarrow D^*\ell\bar{\nu}$ form factor, $\mathcal{F}(1)$, relevant for the determination of $|V_{cb}|$ from exclusive decay.

There are a number of measurements that should be possible using the $B\bar{A}B\bar{A}R$, Belle, LHCb, and future $e^+e^- B$ factory data samples, which could shed light on whether this possibility is realized in nature.

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