

Production of 91 Nb for a measurement of the 91 Nb(p, γ) reaction at FRANZ

B. Thomas*, K. Sonnabend, C. Arda, A. Endres, P. Erbacher, S. Fiebiger, O. Hinrichs, M. Reich, R. Reifarth, and S. Schmidt[†]

Institut für Angewandte Physik, Goethe Universität Frankfurt a. M., Germany E-mail: thomas@iap.uni-frankfurt.de

J. Glorius

GSI Helmholtzzentrum für Schwerionenforschung GmbH, Darmstadt, Germany, and Institut für Angewandte Physik, Goethe Universität Frankfurt a. M., Germany

U. Giesen

Physikalisch Technische Bundesanstalt (PTB), Braunschweig, Germany

R. Dressler, D. Schumann

Paul Scherrer Institut, Villingen, Switzerland

H. Kim

Institut für Kernphysik, Goethe Universität Frankfurt a. M., Germany

In order to answer open questions about the production mechanisms of the p nucleus 92 Mo, the cross section of the 91 Nb(p, γ) reaction must be known precisely. The challenge in measuring this reaction is the instability of 91 Nb. However, it will be possible to perform measurements in standard kinematics in combination with the summing crystal technique using the high proton current FRANZ (Frankfurter Neutronen Quelle am Stern-Gerlach-Zentrum) will deliver. To perform this measurement, it is necessary to produce 91 Nb. One possible way is to activate 92 Mo with protons at $E_p \approx 20$ MeV. A first step towards the production run has been done by measuring the cross sections of the reactions nat Mo(p,x) around this energy.

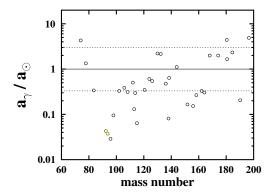
XIII Nuclei in the Cosmos 7-11 July, 2014 Debrecen, Hungary

^{*}Speaker.

 $^{^{\}dagger}$ We thank the accelerator staff and A. Heiske of the PTB, Braunschweig, Germany, and the Institut für Kernphysik at Goethe Universität Frankfurt a. M., Germany. This work is supported by Deutsche Forschungsgemeinschaft (SO907/2-1) and by HIC for FAIR within the framework of LOEWE launched by the State of Hesse, Germany. J.G., S.F. and R.R. acknowledge support by the European Research Council under the European Union's Seventh Framework Programme (FP/2007-2013) / ERC Grant Agreement no. 615126.

1. Motivation

The origin of the elements from hydrogen to iron can be explained by the primordial nucleosynthesis and by fusion reactions during the different burning stages in stars. The synthesis of the elements heavier than iron can be explained by neutron capture processes, namely the s and r processes [1]. However, there are around 35 proton-rich isotopes, which cannot be produced by neutron capture processes. These isotopes are called the p nuclei [2]. Nowadays, it is assumed that these p nuclei are produced during core-collapse supernovae or type Ia supernovae by a series of photo-disintegration reactions – the so called γ process [3]. Network simulations of these scenarios showed that the γ process could not reproduce the expected abundance ratios [4]. The p isotopes of Mo and Rb were in particular underproduced by several orders of magnitude (compare Fig. 1). However, simulations showed, that the experimentally observed abundance ratios may be reproduced by including a series of proton capture reactions [5,6]. To confirm this assumption, the cross sections of the crucial reactions (see Fig. 1) have to be known with precision.



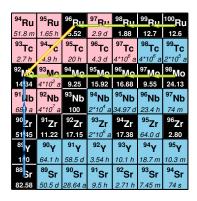


Figure 1: The figure on the left shows the normalized abundances a_{γ} of the p nuclei from a Type II supernova network simulation relative to their normalized solar abundances a_{\odot} . Data taken from [7]. Most of the ratios are within a factor of 3 (dashed lines). The highly underproduced 92,94 Mo isotopes are marked green. The figure on the right shows a part of the reaction network of the γ process with (γ, n) and (γ, α) reactions leading to 92 Mo. Also shown is the possible proton capture path with (p, γ) reactions leading to 92 Mo.

For the p nucleus with the highest relative isotopic abundance, 92 Mo, these reactions are the radiative proton captures of 90 Zr and 91 Nb, due to the fact that these are the last two steps in this production path and 91 Nb is unstable. The reaction 90 Zr(p, γ) was recently remeasured [8] and is under investigation. Details on the studies of this cross section can be found in the contribution by K. Sonnabend in this volume. In contrast, there is no experimental data for the cross section of the 91 Nb(p, γ) reaction. Due to the instability of 91 Nb, the common way would be an experiment in inverse kinematics, e.g., with a proton gas-jet target as shown in the contribution by J. Glorius in this volume. Due to the half-life of 91 Nb of 680 yr [9], it seems feasible to produce a 91 Nb target to perform a measurement in standard kinematics. Radioactive targets suffer three challenges: Firstly, a sufficient amount of target material must be produced with as few impurities as possible. Secondly, the small amount of target material results in a low reaction yield. And at last, the decay radiation from the target must be handled.

2. Production of 91Nb

A possible way to produce ⁹¹Nb is to irradiate ⁹²Mo with protons above an incident energy of 15 MeV. Theoretical calculations with TALYS 1.6 [10] show that, above this energy, the reaction channels ⁹²Mo(p,pn) and ⁹²Mo(p,2n) are open leading to ⁹¹Mo and ⁹¹Nb. The subsequent decay of ⁹¹Mo will also produce ⁹¹Nb. Figure 2 shows these production paths and the calculated cross sections. Other niobium isotopes in the target would lead to a high background and could not be separated chemically. But all other open reaction channels will not efficiently produce any other stable or long-lived niobium isotope directly. Any production through subsequent decays is shielded by stable or long-lived molybdenum isotopes. This allows the production of a high-purity ⁹¹Nb sample after chemical separation of the ⁹¹Nb from the molybdenum. Using the theoretical cross sections, 1.2×10^{16} 91Nb nuclei can be produced by irradiation of 300 mg sample material enriched to 95 % in ⁹²Mo with a total charge of about 7 C. This can be realized for example with the Cyclotron at Physikalisch Technische Bundesanstalt, Braunschweig, Germany, in four days or at the Oslo Cyclotron Laboratory, Norway in one day. After the chemical separation of 91Nb from the ⁹²Mo matrix at Paul Scherrer Institut, Villingen, Switzerland, it is estimated that 50 % of the produced 91 Nb can be recovered. With this 91 Nb material, one could perform a 91 Nb(p, γ) experiment around 2 MeV using the high proton current of 2 mA soon available at FRANZ, Frankfurt a. M., Germany. With the theoretical cross sections of 91 Nb(p, γ), a count rate of 60 Hz is expected in a 4π BaF₂ calorimeter [11]. The radiation from the active target will be on the order of 1 MBq. Under these conditions, an experiment using the summing crystal technique should be possible.

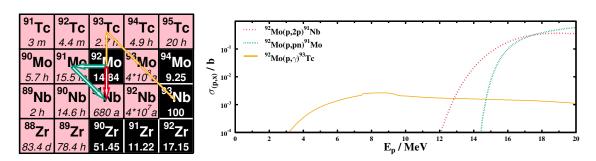


Figure 2: Relevant reaction channels for 92 Mo + p around 15 MeV and their subsequent decays. The cross sections were calculated using TALYS [10]. The production of 93 Nb is suppressed by a small cross section and shielded by the long-lived 93 Mo (yellow arrows).

3. Determination of the production yields of 92 Mo + p

The first step in the sample production was the experimental check of the predicted cross sections. In cooperation with the PTB, Braunschweig, Germany, the PTB cyclotron was used to irridate molybdenum foils of natural isotopic composition. The thickness of the foils was approximately $10\,\mu m$, confirmed by RBS measurements using the Van-de-Graaff accelerator at Goethe Universität Frankfurt a. M., Germany, as well as at PTB, Braunschweig, Germany, at proton energies of $2.4\,MeV$ and $3.0\,MeV$, respectively. Five different energies $E_p = 15-19\,MeV$ and currents

of approximately 500 nA were used. At each energy, there was an activation of 15 minutes to determine the cross section of the reaction 92 Mo(p, pn) leading to 91 Mo with a half-life of 15.49 minutes. Between the end of the activations and the beginning of γ -ray spectroscopy with a large volume high purity germanium detector there was a delay of approximately 6 minutes. A layer of 2.5 cm lead between the target and the detector was reducing the rate in the detector by the low energy γ and X-rays. Figure 3 shows a measurement after an irridation at $E_p = 19$ MeV. Due to the long half-lives of 91 Nb (680 yr) and its isomer (60.86 d) an eight-hour activation was performed at each energy as well to gain higher yields. The spectroscopic analysis of these activation runs is currently carried out using a LEPS (Low Energy Photon Spectrometer) at Goethe Universität Frankfurt a. M., Germany. A first spectrum of the target irridated with protons of $E_p = 15$ MeV for 8 h is depicted in Fig. 4. The first preliminary spectra show the expected counts in the peaks and it should, therefore, be possible to produce the needed amount of 91 Nb.

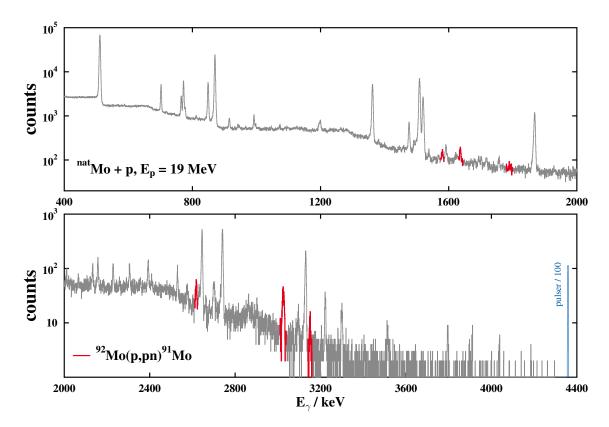


Figure 3: Spectrum of a naturally composed molybdenum target shortly after the 15 minutes proton irradiation with $E_p = 19 \, \text{MeV}$. Red marked peaks correspond to the decay of $^{91} \text{Mo}$. Most of the other peaks result from the decay of $^{94} \text{Tc}$, produced, e.g., by the $^{94} \text{Mo}(p,n)$ reaction.

4. Outlook

After the determination of the production cross sections, the next step will be the production of ⁹¹Nb. In parallel, the investigation of a way to separate the niobium from the molybdenum matrix takes place at PSI, Villingen, Switzerland. At the same time the FRANZ facility at Frankfurt

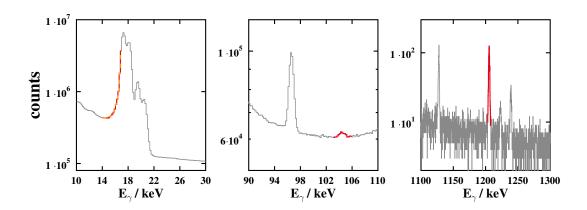


Figure 4: Details of the spectrum of a naturally composed molybdenum target two months after the activation with $E_p = 15 \,\text{MeV}$. Red marked peaks correspond to the decay of the isomer, while the yellow range indicates the energy range which shows the decay of the ground state of ^{91}Nb . The other peaks mainly result from the decays of the long-lived technetium isotopes. This spectrum was measured for 72 hours.

a. M., Germany, is finalized. Aspects of focusing the high-current beam, cooling the target and beam diagnostics are currently studied to optimize the experimental area. For more details, see the contribution by K. Sonnabend in this volume.

References

- [1] D. L. Lambert, Astron. Astroph. Rev. 3 (1992) 201
- [2] E. Burbidge, G. Burbidge, W. Fowler, and F. Hoyle, Rev. Mod. Phys. 29 (1957) 547
- [3] S. E. Woosley, W. M. Howard, Astrophys. J. Suppl. 36 (1978) 285
- [4] M. Arnould, S. Goriely, Phys. Rep. 384 (2003) 1
- [5] M. Kusakabe, N. Iwamoto, K. Nomoto, Astrophys. J. 726 (2011) 25
- [6] C. Travaglio, F. Röpke, R. Gallino, W. Hillebrandt, Astrophys. J. 739 (2011) 93
- [7] W. Rapp, J. Görres, M. Wiescher, H. Schatz, and F. Käppeler, Astrophys. J. 653 (2006) 474
- [8] A. Spyrou, S. J. Quinn, A. Simon, T. Rauscher, A. Battaglia, A. Best, B. Bucher, M. Couder, P. A. DeYoung, A. C. Dombos, X. Fang, J. Görres, A. Kontos, Q. Li, L. Y. Lin, A. Long, S. Lyons, B. S. Meyer, A. Roberts, D. Robertson, K. Smith, M. K. Smith, E. Stech, B. Stefanek, W. P. Tan, X. D. Tang, M. Wiescher, *Phys. Rev. C* 88 (2013) 045802
- [9] C. Baglin, *Nucl. Data Sheets* **114** (2013) 1293, online version (http://www.nndc.bnl.gov, 10th September 2014)
- [10] A. J. Koning, S. Hilaire, M. C. Duijvestijn, in O. Bersillon et al., Eds., Proceedings of the International Conference on Nuclear Data for Science and Technology - ND2004, Sep. 26 - Oct. 1, 2004, Santa Fe, USA, AIP Conf. Proc. 769 (2005) 1154
- [11] K. Wisshak, K. Guber, F. Käppeler, J. Krisch, H. Muller, G. Kupp, and F. Voss, *Nucl. Instr. Methods A* **292** (1990) 3