

Odd and Even Partial Waves of $\eta\pi^-$ and $\eta'\pi^-$ in 191 GeV/c $\pi^- p$

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In the year 2008 COMPASS recorded diffractive events of the signature $\pi^-(191\text{ GeV})p \rightarrow X_{\text{fast}}p$. We present results of the analysis of the subsystems $X = \eta^{(\prime)}\pi^-$. Besides the known resonances $a_2(1320)$, $a_4(2040)$, we study the properties of the spin-exotic P_+ wave, and all other natural-parity exchange partial waves up to spin $J = 6$. We find a striking difference between the two final states: whereas the even partial waves 2, 4, 6 in the two systems are related by phase-space factors, the odd partial waves are relatively suppressed in the $\eta\pi^-$ system. The relative phases between the even waves appear identical whereas the phase between the D and P waves behave quite differently, suggesting different resonant and non-resonant contributions in the two odd-angular-momentum systems. Branching ratios and parameters of the well-known resonances a_2 and a_4 are measured. We find

$$m(a_2) = 1315 \pm 12 \text{ MeV}, \quad \Gamma(a_2) = 119 \pm 14 \text{ MeV},$$

and

$$m(a_4) = 1900_{-20}^{+80} \text{ MeV}, \quad \Gamma(a_4) = 300_{-100}^{+80} \text{ MeV}.$$

(consistent with COMPASS's 3π analyses.) The relative branchings we measure are

$$\frac{BR(a_2 \rightarrow \eta'\pi)}{BR(a_2 \rightarrow \eta\pi)} = (5 \pm 2)\%, \quad \frac{BR(a_4 \rightarrow \eta'\pi)}{BR(a_4 \rightarrow \eta\pi)} = (23 \pm 7)\%.$$

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1. Introduction

The systems $\eta\pi$ and $\eta'\pi$ are attractive laboratories for strong-interaction physics because of their simplicity and clear experimental signature. Besides the well-known resonances $a_2(1320)$ and $a_4(2040)$, resonance-like behavior was observed in the P -wave, whose neutral isospin member carries the exotic quantum numbers $J^{PC} = 1^{-+}$ (see e.g. Ref. [1]). In this contribution, we discuss an analysis of the $\eta\pi^-$ and $\eta'\pi^-$ systems, diffractively produced off a proton target during the 2008 run of the COMPASS experiment. Previous work on this analysis was discussed in Refs. [2, 3]. A journal publication is in progress.

The COMPASS experiment is a fixed-target experiment installed at the CERN SPS. Its two-stage spectrometer allows for high-resolution particle detection and reconstruction over a wide range in angles and momenta, both for charged and neutral particles [4]. The data recorded for the analysis under discussion was produced by having a 191 GeV π^- beam impinging on a LH2 target. The target was surrounded by a recoil proton detector which together with a veto detector surrounding the spectrometer entry formed a trigger ensuring a clean sample of diffractive excitation reactions with momentum transfer $|t| \gtrsim 0.08 \text{ GeV}^2$ [5]. Samples of approximately 35×10^3 exclusive $\pi^- \eta'$ and 110×10^3 exclusive $\pi^- \eta$ events with invariant masses from threshold up to several GeV is obtained in the reaction $\pi^- p \rightarrow \pi^- \eta^{(\prime)} \rightarrow \pi^- \pi^- \pi^+ \gamma \gamma p$, where the two photons result from the decay of an intermediate η (π^0) in the η' (η) decay. Acceptances for the two reactions as a function of mass and polar angle (Gottfried-Jackson system) are shown in Fig. 1.

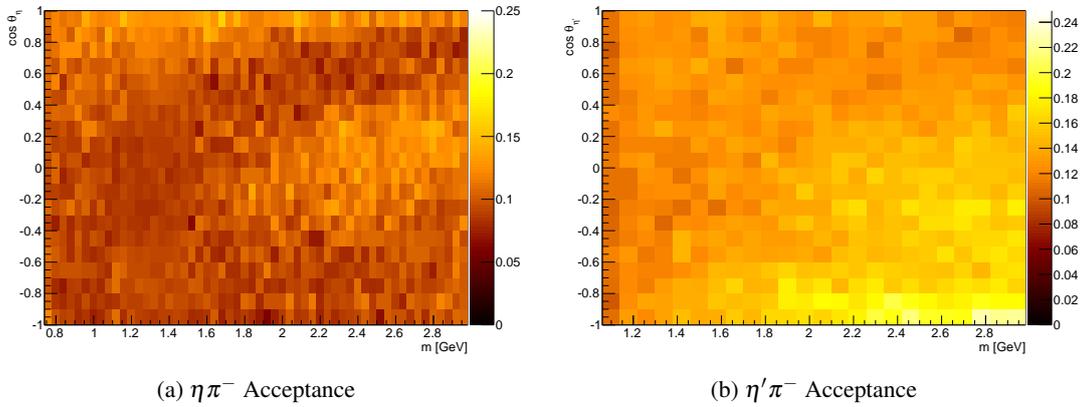


Figure 1: Acceptance evaluated from Monte Carlo simulation. An azimuthal distribution $\propto \sin^2 \phi$ (i.e. natural-parity exchange, $M = 1$) and the experimental t' distribution were used.

In the flavor basis, η - η' mixing is described by an angle $\phi \approx 39^\circ$ [6]. One expects in particular for branching ratios of the a_2 and a_4 resonance decays to pseudoscalars

$$\text{BR}(a_J \rightarrow \pi\eta')/\text{BR}(a_J \rightarrow \pi\eta) = F(J, q'(m), q(m)) \tan^2 \phi,$$

where J is angular momentum, and $q^{(\prime)}(m)$ are the breakup momenta at invariant mass m . For $q^{(\prime)} \rightarrow 0$, the behavior of each cross-section has to follow $(q^{(\prime)})^{2J+1}$ from analyticity of the partial-wave series. Therefore, the simplest form of dynamical term, which we use in the following, is

$$F(J, q', q) = (q'/q)^{2J+1}. \quad (1.1)$$

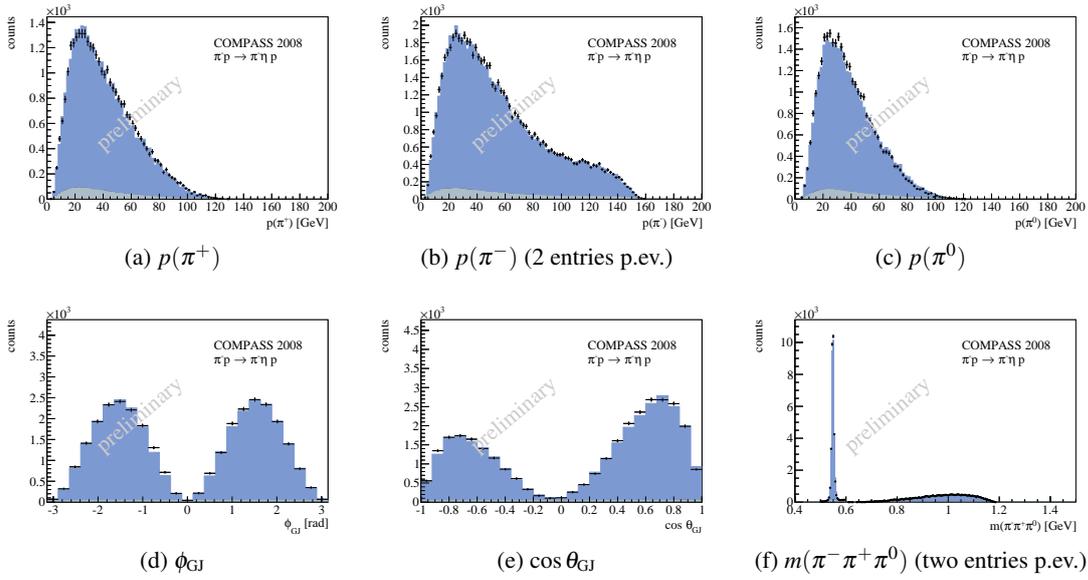


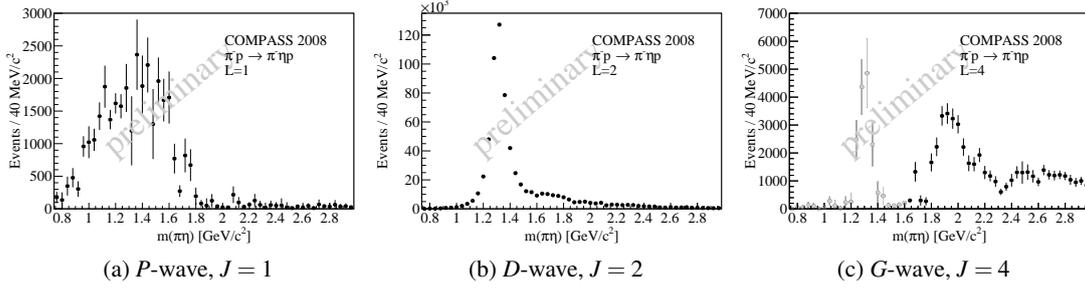
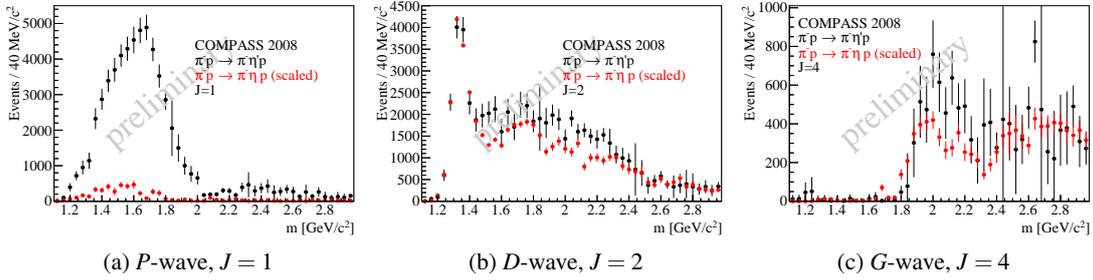
Figure 2: Comparison of fit results to the data. First row: lab momenta of pions, second row: two-body angular variables and invariant three-body mass with the $\eta(548)$ peak. Black dots: data. The fit result is decomposed into the incoherent contributions as follows: Light blue: natural-parity exchange waves. Gray: non- η background. Dark blue: unnatural-parity exchange waves (negligible).

2. Partial-wave Analysis Procedure

The data were subjected to partial-wave analysis. Here, an acceptance-corrected partial-wave model was fit to the data in 40 MeV wide bins of $m(\eta^{(\prime)}\pi^-)$ from threshold up to 3 GeV, separately for the two final states. The formalism used was an extended log-likelihood fit where partial waves were parametrized in the reflectivity basis. Natural-parity exchange partial waves with spin projection $M = 1$ for angular momenta up to $J = 6$ were included. For $J = 2$ in $\eta\pi$, an $M = 2$ partial wave was also included in the analysis. It was found to contribute 3% of the a_2 intensity. The complete four-body information was used to distinguish the three-body $\eta^{(\prime)}$ peak from background reactions which were modeled by a partial wave isotropic in four-body phase space. The partial wave model consisted of three incoherent contributions: natural-parity exchange waves, unnatural-parity exchange waves and the four-body phase-space background. Consistent with the expectation of a dominant Pomeron contribution, unnatural parity exchange is found to be suppressed. In Fig. 2 we illustrate the procedure by overlaying the data and the fit results for the $\eta\pi^-$ data in the vicinity of the $a_2(1320)$ resonance. For details see Ref. [3]. Physical hypotheses are then tested by fitting mass-dependent (resonance) models to the partial-wave intensities and phases extracted in the mass-binned fit. The Breit-Wigner fit results for the known resonances are given in the abstract.

3. Partial-wave Results

We show the intensities of the main waves of the $\eta\pi^-$ data in Fig. 3. The $J = 1$ P -wave shows a broad bump and vanishing intensity above 1.8 GeV. The $J = 2$ wave is dominated by the well-


 Figure 3: Main waves of the $\eta(\rightarrow \pi^- \pi^+ \gamma\gamma)\pi^-$ data.

 Figure 4: Main waves of the $\eta'(\rightarrow \pi^- \pi^+ \gamma\gamma)\pi^-$ data. In red: the $\eta\pi^-$ data multiplied by the mass-dependent phase-space factor from Eq. 1.1, taking into account final-state branching fractions.

known $a_2(1320)$ resonance with a shoulder at high mass. Besides leakage from the dominant $J = 2$ wave, the $J = 4$ wave exhibits a clear $a_4(2040)$ signal, followed by a broad structure at high mass.

The same partial waves are depicted for the $\eta'\pi^-$ data in Fig. 4. Again, we see a broad structure in the $J = 1$ P -wave, this time vanishing near 2 GeV with some intensity reappearing at higher masses. In the $J = 2$ wave, the relative height of the high-mass shoulder compared to the peak is enhanced compared to the $\eta\pi^-$. Similarly, the peak in the $J = 4$ wave stands out less in the $\eta'\pi^-$ data. Overlaid on the $\eta'\pi^-$ data are the $\eta\pi^-$ data from Fig. 3, where the content of each bin has been multiplied with the factor from Eq. 1.1 and a factor taking into account the final-state decays $\eta^{(\prime)} \rightarrow \pi^- \pi^+ \gamma\gamma$ [6]. We find a surprising difference between different partial waves: whereas the even waves with $J = 2, 4$ show very similar behavior, the odd $J = 1$ wave is relatively enhanced in the $\eta'\pi^-$ data. These properties extend also to the waves with spins $J = 3, 5, 6$ (not shown): odd waves are relatively enhanced in $\eta'\pi^-$, even waves largely agree after phase-space multiplication.

For the phases, a similar behavior is observed, shown in Fig. 5: the phases between the even-spin waves $J = 2$ and $J = 4$ agree between the two channels. The phases between the $J = 1$ and $J = 2$ partial waves disagree in the region of the $J = 1$ intensity peaks. A particularly intriguing feature is the agreement of this ($J = 1$) – ($J = 2$) phase near the $\eta'\pi^-$ threshold.

The difference in even/odd behavior is detailed in Tab. 1, where we show the relative intensities of the various partial waves and the ratios of their integrals after phase-space scaling.

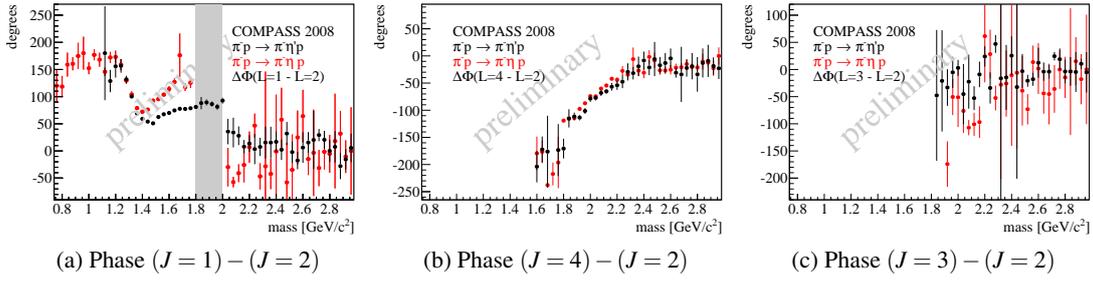
Figure 5: Relative phases in $\eta'\pi^-$ and $\eta\pi^-$ (red) for selected partial waves.

Table 1: Relative intensities of the $J = 1$ to 6 partial waves resulting from the PWA fits integrated over the mass range up to 3 GeV. Experimental acceptance is taken into account. The total $\eta'\pi^-$ to $\eta\pi^-$ intensity ratio in this mass range is 0.19 ± 0.02 . The ratio of the integrals of the red and black histograms in Fig. 4, and similar for the other waves, is given in the last row.

J	1	2	3	4	5	6
$\frac{I_J(\eta\pi^-)}{I_{\text{total}}(\eta\pi^-)}$ [%]	4.4	81.9	0.3	6.9	0.1	0.7
$\frac{I_J(\eta'\pi^-)}{I_{\text{total}}(\eta'\pi^-)}$ [%]	41.7	42.3	3.7	8.4	0.9	1.2
R_{corr}	0.17 ± 0.01	0.94 ± 0.02	0.16 ± 0.05	0.83 ± 0.07	0.15 ± 0.12	0.68 ± 0.15

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