

Models of Walking Technicolor on the Lattice

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We study QCD with 2 colour-sextet quarks as a walking-Technicolor candidate. As such it provides a description of the Higgs sector of the standard model, in which the Higgs field is replaced by the Goldstone ‘pions’ of this QCD-like theory, and the Higgs itself is the σ . Such a theory will need to be extended if it is to also give masses to the quarks and leptons. What we are attempting to determine is whether it is indeed QCD-like and hence walking, or if it has an infrared fixed point making it a conformal field theory. We do this by simulating its lattice version at finite temperature and observing the running of the bare (lattice) coupling at the chiral transition, as the lattice spacing is varied, and comparing this running with that predicted by 2-loop perturbation theory. Our results on lattices with temporal extents (N_t) up to 12 indicate that the coupling runs, but not as fast as asymptotic freedom predicts. We discuss our program for studying the zero-temperature phenomenology of this theory.

QCD with 3 colour-sextet quarks, which is believed to be conformal is studied for comparison. Simulations of this theory at finite temperature, on lattices with N_t as large as 8 indicate that the coupling still runs, and shows no sign of approaching a finite limit for large N_t , in contrast to what is expected for a conformal theory. We are now extending these runs to $N_t = 12$. While it is too early to draw any conclusions from our $N_t = 12$ runs, there is a hint that the running of the coupling might have slowed.

Finally we looked at a new candidate theory, $SU(2)$ gauge theory with 3 Majorana adjoint fermions. However, this did not appear to allow embedding of the electroweak group ($SU(2) \times U(1)$) to give the correct masses to the W s and Z .

The 32nd International Symposium on Lattice Field Theory
23-28 June, 2014
Columbia University New York, NY

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[†]This research was supported in part by US Department of Energy contract DE-AC02-06CH11357

1. Introduction

We study models for the Higgs sector in which the Higgs is composite. In particular, we study Technicolor models – QCD-like theories with massless fermions, where the Goldstone pion-like excitations play the role of the Higgs field, giving mass to the W^\pm and Z [1, 2]. Of particular interest are walking-Technicolor models, where there is a range of mass scales over which the running coupling evolves very slowly [3, 4, 5, 6] Such models can avoid the phenomenological problems with naive Technicolor.

QCD with 2 colour-sextet quarks is a candidate walking-Technicolor model. (For a summary of the properties of $SU(N)$ gauge theories which lead one to such a conclusion see [7, 8].) We need to distinguish whether this theory walks or is conformal. It is attractive because it has just the right number of Goldstone bosons (3) to give mass to the W^\pm and Z , with none left over. Other groups are studying this model, in particular, DeGrand, Shamir & Svetitsky [9, 10, 11, 12, 13, 14] and the Lattice Higgs Collaboration [15, 16, 17, 18, 19, 20] We study this theory at finite temperature to see if the coupling at the chiral transition evolves as predicted by asymptotic freedom for a finite-temperature transition [21, 22].

QCD with 3 colour-sextet quarks, which is believed to be conformal, is studied for comparison [23]. For a conformal theory, the chiral transition is a first-order bulk transition at fixed coupling separating the chirally-symmetric conformal theory from the chirally broken phase, which is a lattice artifact.

We simulate these theories, latticized with unimproved staggered fermions, using the RHMC method.

If QCD with 2 sextet quarks walks, we need to answer the following questions. Does QCD with 2 colour-sextet quarks have a light Higgs with standard-model properties? What other light particles are in its spectrum? Can any of its particles be dark-matter candidates? Do its S (T and U) parameter(s) pass precision electroweak tests.

We have also considered $SU(2)$ Yang-Mills with 3 Majorana/Weyl fermions. However, we have been unable to embed the electroweak gauge group ($SU(2) \times U(1)$) into this theory to give physical masses to the W s and Z .

2. QCD with 2 colour-sextet quarks at finite temperature

We simulate QCD with 2 color-sextet quarks at finite temperature by simulating on an $N_s^3 \times N_t$ lattice with $N_s \gg N_t$. Since $T = 1/N_t a$, increasing N_t with T fixed decreases a . Assuming the chiral phase transition is a finite-temperature transition, this yields a convenient T , T_χ . Measuring g or $\beta = 6/g^2$ at T_χ gives a running coupling at a sequence of a s which approach zero as $N_t \rightarrow \infty$.

2.1 $N_t = 12$

Much of the past year has been devoted to increasing the statistics for our simulations on $24^3 \times 12$ lattices at quark masses $m = 0.0025$ and $m = 0.005$, close to the chiral transition. For our largest mass $m = 0.01$ we have extended our simulations at low β s to determine the position of the deconfinement transition.

Because the β dependence of the chiral condensate is so smooth for the masses we use, we determine the position β_χ of this transition from the peaks in the (disconnected) chiral susceptibility:

$$\chi_{\bar{\psi}\psi} = V [\langle (\bar{\psi}\psi)^2 \rangle - \langle \bar{\psi}\psi \rangle^2] \quad (2.1)$$

extrapolated to $m = 0$. V is the space-time volume.

For $m = 0.01$ in the range $5.7 \leq \beta \leq 5.9$, near the deconfinement transition, we run for 50,000 trajectories for each β with β s spaced by 0.02. In the range $6.6 \leq \beta \leq 6.9$, near the chiral transition we run for 25,000 trajectories per β with β s spaced by 0.02. Elsewhere in the range $5.7 \leq \beta \leq 7.2$ we run for 10,000 trajectories for β s spaced by 0.1.

For $m = 0.005$ in the range $6.6 < \beta \leq 6.9$, we run for 50,000 trajectories per β at β s spaced by 0.02. At $\beta = 6.6$ we run for 100,000 trajectories. Elsewhere in the range $6.4 \leq \beta \leq 7.2$, we run 10,000 trajectories per β for β s spaced by 0.1.

For $m = 0.0025$ in the range $6.7 \leq \beta \leq 6.9$, we run for 100,000 trajectories per β with β s spaced by 0.02. In the range $6.6 \leq \beta < 6.7$, we run for 50,000 trajectories per β . We are currently extending our run at $\beta = 6.68$ to 100,000 trajectories. Elsewhere in the range $6.5 \leq \beta \leq 7.2$ we run 10,000 trajectories per β at β s spaced by 0.1.

While the chiral condensates measured in these simulations suggest that this condensate will vanish in the chiral limit for large enough β values, they do not allow a precise determination of β_χ where this phase transition occurs. For this we turn to the chiral susceptibilities. Figure 1 shows the chiral susceptibilities from these runs. The peak of the $m = 0.0025$ susceptibility yields an estimate of β_χ , namely $\beta_\chi = 6.77(1)$. Combining this with our $N_t = 8$ results yields:

$$\beta_\chi(N_t = 12) - \beta_\chi(N_t = 8) = 0.08(2) , \quad (2.2)$$

significantly smaller than the 2-loop perturbative prediction:

$$\beta_\chi(N_t = 12) - \beta_\chi(N_t = 8) \approx 0.12 . \quad (2.3)$$

Figure 2 shows histograms of the magnitudes of Wilson Lines for $m = 0.01$ near to the deconfinement transition. From this we deduce that $\beta_d = 5.81(1)$ for $m = 0.01$. Previous experience indicates that this should be close to the value for $m = 0$.

Table 1 shows our estimates for the positions of the deconfinement (β_d) and chiral (β_χ) transitions from these and previous simulations.

N_t	β_d	β_χ
4	5.40(1)	6.3(1)
6	5.54(1)	6.60(2)
8	5.65(1)	6.69(1)
12	5.81(1)	6.77(1)

Table 1: $N_f = 2$ deconfinement and chiral transitions for $N_t = 4, 6, 8, 12$.

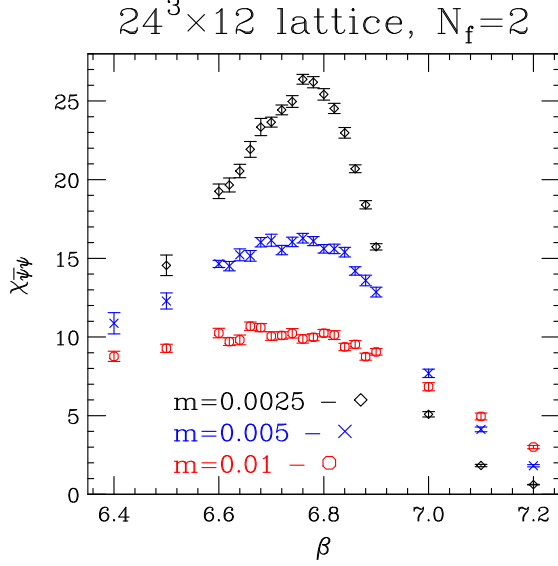


Figure 1: Chiral susceptibilities on a $24^3 \times 12$ lattice.

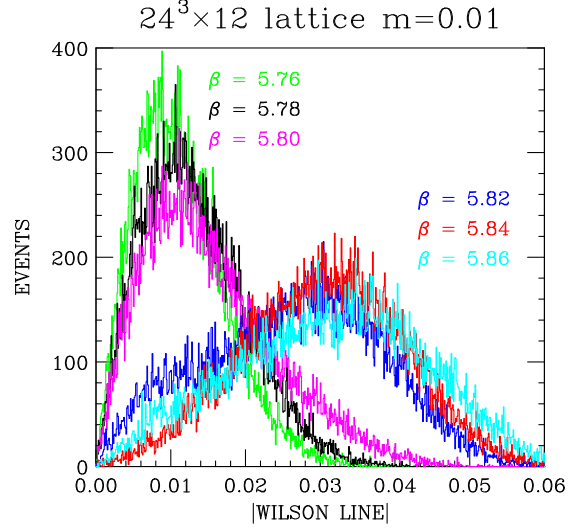


Figure 2: Histograms of magnitudes of Wilson Lines for β s close to the deconfinement transition for $m = 0.01$.

3. Planned simulations and measurements

We plan simulations at zero temperature to understand the phenomenology of this theory, starting with simulations on a $36^3 \times 72$ lattice at β values below the deconfinement transition for $N_t = 36$. Unfortunately, such β s are still too small to access the continuum limit. However, we hope that we will be able to get results which are qualitatively correct. We then plan to move to a $48^3 \times 96$ lattice which will allow us to move to larger β s.

We will measure f_π and the meson spectrum, including disconnected contributions. Here we will look for a light Higgs, and measure ‘taste’-breaking in the ‘pion’ spectrum. The ‘glueball’ spectrum will be measured since we suspect that the low-lying glueballs could be light, and that mixing with the mesons could produce a light Higgs-like particle. We will then need to check if any Higgs candidate has the correct couplings to the W^\pm and Z . We plan to measure S -parameter contributions, and to determine the scaling behaviour of the chiral condensate to extract γ_m .

4. QCD with 3 colour-sextet quarks at finite temperature

We simulate lattice QCD with 3 colour-sextet quarks at finite temperature for comparison with the 2-flavour case. This theory is believed to be conformal with an infrared fixed point. The chiral transition should be a bulk transition fixed at a finite constant β_χ for N_t sufficiently large. We have simulated this theory at $N_t = 4, 6$ and 8 , and are now performing $N_t = 12$ simulations.

For $N_t = 6$ we simulate on a $12^3 \times 6$ lattice at $m = 0.02, m = 0.01$ and $m = 0.005$. Close to the chiral transition ($6.2 \leq \beta \leq 6.4$), at the lowest quark mass ($m = 0.005$), we simulate at β s separated by 0.02 , with $100,000$ trajectories per β . We estimate the position of the chiral transition as the peak in the chiral susceptibility for $m = 0.005$.

For $N_t = 8$ we simulate on a $16^3 \times 8$ lattice at $m = 0.01$ and $m = 0.005$. Close to the chiral transition ($6.28 \leq \beta \leq 6.5$) at the lowest quark mass ($m = 0.005$) we simulate at β s separated by 0.02, with 100,000 trajectories per β . We estimate the position of the chiral transition as the peak in the chiral susceptibility for $m = 0.005$. The results for the positions of the chiral and deconfinement transitions for $N_t = 4, 6$ and 8 are given in table 2.

N_t	β_d	β_χ
4	5.275(10)	6.0(1)
6	5.375(10)	6.278(2)
8	5.45(10)	6.37(1)

Table 2: $N_f = 3$ deconfinement and chiral transitions for $N_t = 4, 6, 8$. In each case we have attempted an extrapolation to the chiral limit.

Since

$$\beta_\chi(N_t = 8) - \beta_\chi(N_t = 6) = 0.09(1) \quad (4.1)$$

we have yet to see evidence of a bulk transition. We are therefore starting $N_t = 12$ simulations on a $24^3 \times 12$ lattice. Figure 3 shows the $m = 0.005$ chiral susceptibilities for $N_t = 6$, $N_t = 8$ and preliminary results for $N_t = 12$.

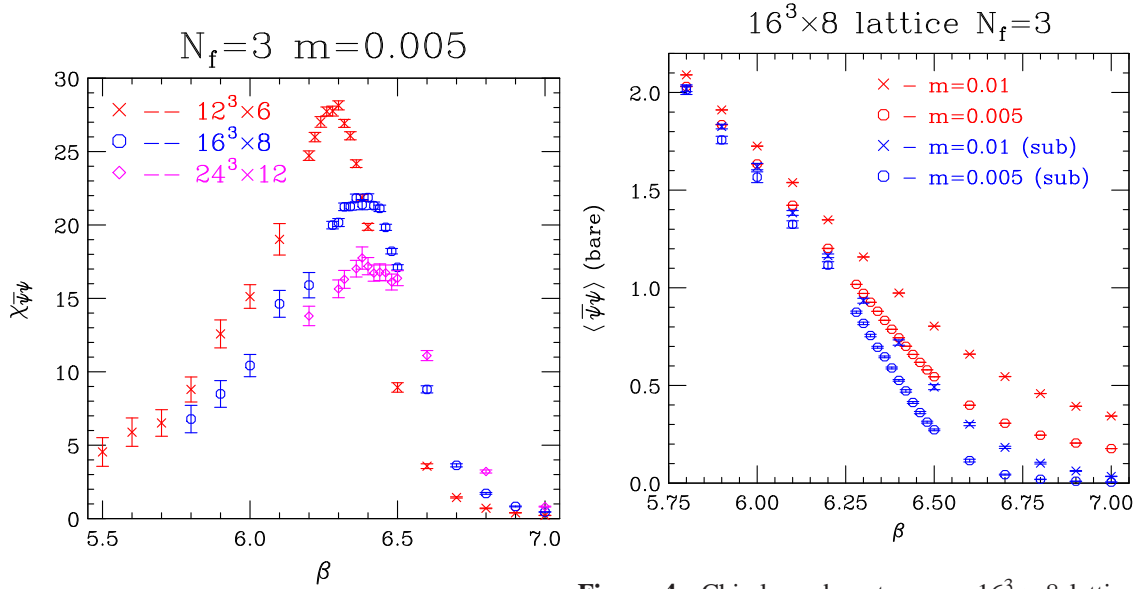


Figure 3: Chiral susceptibilities for $N_f = 3$, $m = 0.005$ on $12^3 \times 6$, $16^3 \times 8$ and $24^3 \times 12$ lattices.

Figure 4: Chiral condensates on a $16^3 \times 8$ lattice for $m = 0.005$ and $m = 0.01$. The red graphs are unsubtracted, lattice regulated condensates. The blue graphs have been subtracted.

Figure 4 shows the chiral condensates, both unsubtracted and subtracted for our $16^3 \times 8$ simulations. The subtracted condensates use the definition of the Lattice Higgs Collaboration [16]:

$$\langle \bar{\psi}\psi \rangle_{sub} = \langle \bar{\psi}\psi \rangle - \left(m_V \frac{\partial}{\partial m_V} \langle \bar{\psi}\psi \rangle \right)_{m_V=m}, \quad (4.2)$$

where m_V is the valence-quark mass. Note, although it is clearer that the subtracted condensate will vanish in the continuum limit for β sufficiently large than is the case for the unsubtracted condensate, it still does not yield an accurate estimate of β_χ .

5. $SU(2)$ gauge theory with 3 Majorana/Weyl colour-adjoint fermions

The symmetries of this theory are easiest to see in terms of 2-component (Weyl) fermions.

$$\mathcal{L} = -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} + \frac{1}{2}\psi^\dagger i\sigma^\mu \overleftrightarrow{D}_\mu \psi + \frac{m}{2}[\psi^T i\sigma_2 \psi - \psi^\dagger i\sigma_2 \psi^*] \quad (5.1)$$

where ψ is a 3-vector in colour₂ space and in flavour space.

If $m = 0$, the chiral flavour symmetry is $SU(3)$. The Majorana mass term reduces this flavour symmetry to the real elements of $SU(3)$, i.e. to $SO(3)$. Thus when $m = 0$ and the chiral symmetry breaks spontaneously, the chiral condensate is $\langle \psi^T i\sigma_2 \psi - \psi^\dagger i\sigma_2 \psi^* \rangle$. and the spontaneous symmetry breaking pattern is $SU(3) \rightarrow SO(3)$. The unbroken generators of $SU(3)$ are the 3 imaginary generators. These form a spin-1 representation under the unbroken $SO(3)$. The 5 broken generators are the 5 real generators. They, as well as the 5 corresponding Goldstone bosons, form a spin-2 representation of $SO(3)$.

The problem occurs when one tries to embed the weak $SU(2) \times U(1)$ group in such a way as to give masses to W^\pm and Z . This is easiest to see if we consider the case where the Weinberg angle is zero. Then we need to embed $SU(2)$ in such a way that all 3 components are broken spontaneously. Thus we would need to make a set of $SU(2)$ generators from the 5 real $SU(3)$ generators. However, the $SU(2)$ algebra requires at least one of its generators to be complex, so this is impossible. The only Weinberg angle which would work is $\pi/2$ where the photon is pure $SU(2)$ and the Z is pure $U(1)$.

6. Discussion and Conclusions

We simulate lattice QCD with 2 colour-sextet quarks at finite temperature to distinguish whether it is QCD-like and hence walks, or if it is a conformal field theory. We run on lattices with $N_t = 4, 6, 8, 12$. β_χ increases by 0.08(2) between $N_t = 8$ and $N_t = 12$. While this increase favours the walking scenario, it is significantly smaller than the 2-loop prediction of ≈ 0.12 . Is this because 2-loop perturbation theory is inadequate for this lattice action and β ? Are there sizable finite volume corrections? Will the theory finally prove to be conformal?

If walking, this theory is a promising walking-Technicolor theory. It has just the right number of Goldstone bosons to give masses to the W^\pm and Z , with no extras to be explained away. We have outlined a program for checking its zero-temperature properties. Does it have a light Higgs? Does it satisfy the precision electroweak constraints? Does it have a Dark Matter candidate? What about its particle spectrum?

We simulate QCD with 3 colour-sextet quarks which should be conformal. The increase in β_χ between $N_t = 6$ and $N_t = 8$ is still appreciable (0.09(1)), so we don't yet have evidence for β_χ approaching a finite constant as $N_t \rightarrow \infty$. We are now simulating at $N_t = 12$. This shows some promise.

QCD₂ with 3 Majorana/Weyl quarks does not appear to be a Technicolor candidate.

Acknowledgements

These simulations were performed on Hopper, Edison and Carver at NERSC, and Kraken at NICS and Stampede at TACC under XSEDE project TG-MCA99S015, and Fusion and Blues at LCRC, Argonne. NERSC is supported by DOE contract DE-AC02-05CH11231. XSEDE is supported by NSF grant ACI-1053575.

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