

## Vertexing and Tracking Software at Belle II

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Belle II is a  $B$  factory experiment aiming to start physics data taking in 2017. It is currently being set up at the SuperKEKB accelerator at the KEK facility in Tsukuba (Japan), an asymmetric  $e^+e^-$  collider which aims to achieve an unprecedented instantaneous luminosity of  $8 \cdot 10^{35} \text{Hz/cm}^2$ . This forty-fold increase over predecessor experiments is achieved by employing a novel nano-beam scheme. Originally developed for the now-defunct SuperB experiment, this scheme allows a significant increase in luminosity at only modest increases of beam currents. Challenges for the vertex detector result from increased data and background rates. At full luminosity, physics data will be recorded at a rate of 30 kHz. The radiation-hard DEPFET-sensors of the innermost layer of the vertex detector will be read out employing a novel data-reduction scheme using selective detector read out based on online reconstruction of event data. Belle II uses a software framework in which data handling is unified between various data processing modules. In this way, tasks can be divided flexibly and the same software framework can be used for a diversity of tasks ranging from the high-level trigger to final processing of plots for publications. Tracking and vertexing modules are currently under development and we will discuss their features.

PACS: 29.40.Gx, 29.85.Ca, 29.85.-c

*The 23rd International Workshop on Vertex Detectors*

*15-19 September 2014*

*Macha Lake, The Czech Republic*

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\*This research was supported by the DFG cluster of excellence 'Origin and Structure of the Universe' and under BMBF Contract 05H12WM8.

## 1. Introduction

Our experimental knowledge of the flavor structure of the Standard Model is in large parts due to the  $B$ -factories Belle and BaBar. Situated at the asymmetric  $e^+e^-$  colliders KEKB and PEP-II, respectively, they collected data throughout the first decade of the 2000s, confirming in great detail the CKM theory, and the related violation of the fundamental discrete symmetries  $C$ ,  $P$  and  $T$  [1]. Belle II is a successor to these experiments [2]. It employs the same concept of asymmetric  $e^+e^-$  collisions at the  $\Upsilon(10580)$  resonance which decays into two mesons  $B\bar{B}$ , but compared to the predecessor experiments Belle II uses a slightly reduced beam momentum asymmetry. Typical travel distances of the  $B$  mesons before their decays are  $200\mu\text{m}$ . Belle II is an exploratory experiment aiming to test the predictions of the flavor sector of the Standard Model at unprecedented precision, thus probing new physics contributions well into the TeV mass range, given sufficient misalignment of the new physics flavor structure with that of the Standard Model [3]. The Belle II experiment is currently assembled at the SuperKEKB accelerator at the KEK facility in Tsukuba, Japan. The experiment is organized as an international collaboration of some 600 researchers and will start taking physics data in 2017, quickly integrating more luminosity than its predecessor experiments, then aiming for a total integrated luminosity  $50\text{ab}^{-1}$  accumulated over a period of approximately eight years.

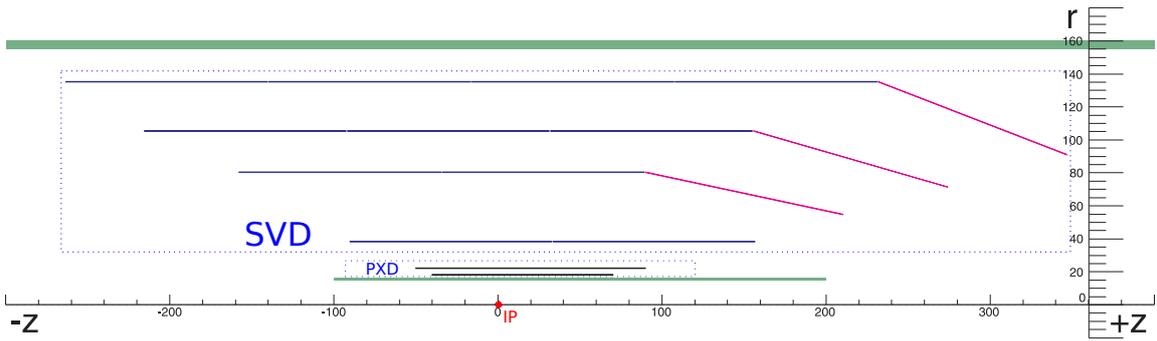
The target luminosity of  $8 \cdot 10^{35}\text{Hz}/\text{cm}^2$  is approximately forty times larger than that of the predecessor facility KEKB. This is achieved by a novel nano-beam scheme for the interaction region supplemented by a two-fold increase in beam currents [4]. The increased luminosity leads to a corresponding increase in background levels, mainly QED backgrounds (which scale proportionally to luminosity), and in-beam scattering (Touschek effect, due to stronger focussing). Other machine backgrounds only scale as the beam current.

## 2. The Belle II detector

The Belle II detector is an upgrade of the Belle detector: it is housed in the same structure, reusing the spectrometer magnet. The design aimed for performance comparable to the performance of the Belle detector in spite of the increased background levels expected at the higher-luminosity accelerator. Where most parts of the detector were refined in various ways, two major new components were introduced in the design: a pixel vertex detector based on the DEPFET technology very close to the interaction point which is capable of withstanding the increased radiation levels while providing excellent resolution, and a complete overhaul of the particle identification detectors which are of Cherenkov imaging type. Additionally, the beam pipe is replaced with a smaller, thinner device, which allows installing the innermost layer of the vertex detector at a distance of only 1.4cm from the interaction point.

The components of the Belle II vertex detector are discussed in detail in Peter Kodys's [5], Michael Schnell's [6] and Lorenzo Vitale's [7] contributions to this conference, so we will restrict ourselves to an executive summary.

Belle II is designed for detecting and reconstructing particle trajectories for transverse momenta exceeding  $50\text{MeV}/c$  while providing excellent momentum resolution all the way up to  $p_T$  at the physical limit which is half the center of mass energy  $m(\Upsilon(4S)) = 10.58\text{GeV}/c^2$ . The vertex

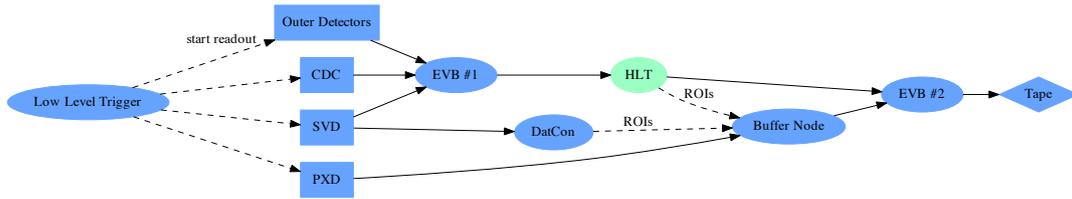


**Figure 1:** Layout of the Belle II vertex detector (VXD). Distances are indicated in mm relative to the interaction point (IP). The two layers of the pixel detector (PXD) and the four layers of the silicon vertex detector (SVD) are indicated. The forward parts of the outer three layers of the SVD are slanted to reduce the amount of silicon. The VXD volume is bounded by the beam pipe on the inside and by the inner wall of the drift chamber on the outside (both indicated in green).

detector consists of six layers of silicon detectors, situated at radial distances from the interaction point between 1.4 cm and 14 cm. It is surrounded by a wire-chamber that extends up to a radius of 1.1 m.

The vertex detector itself consists of two subdetectors, cf. Fig. 1. The two innermost layers are built from DEPFET-type active sensors with pixel readout [8], able to withstand the high radiation levels. These semiconductor detectors are thinned to  $75\ \mu\text{m}$  in the active area corresponding to a radiation length of  $X/X_0 = 0.21\%$ . The surrounding four layers are built using double-sided silicon strip detectors of thickness  $300\ \mu\text{m}$ , read out using APV25 chips [9]. Readout has to deal with the long integration times of the DEPFET sensors ( $\approx 20\ \mu\text{s}$ ) which is comparable to the physics trigger rate (30 kHz). Given the expected occupancy in the range of 1%, full readout at every trigger would lead to a data rate of 20 GB/s for the pixel layers alone, exceeding the amount of data collected by the other parts of Belle II by an order of magnitude. In order to reduce this to a level supportable, an online data reduction scheme has to be implemented. This is implemented by a two-level event building/high-level triggering scheme, illustrated in Fig. 2. Upon occurrence of a low-level trigger signal, all subdetectors besides the pixel detector are read out, events are built and fed to the high-level trigger. Simultaneously, the pixel-detector data is read out and stored on a buffer node [10]. The high-level trigger performs a full reconstruction of the event excluding pixel data. Reconstructed tracks are then extrapolated back to the pixel layers, where the vicinities of the estimated crossing points define regions-of-interest which are read from the buffer node and added to the event in a second event-building step [11]. As an alternative, an FPGA-based implementation of this region-of-interest definition is implemented on dedicated readout hardware of the silicon strip detector. This approach was discussed in another contribution to this conference [6]. In order to recover tracks with very low transverse momenta, additional pixel data corresponding to very large ionisation is recorded irrespective of whether it can be associated with a region-of-interest. In this way a ten-fold reduction of the amount of data recorded can be achieved.

The wire chamber construction was recently completed. The detector is now being tested and first calibrations using cosmic rays are being performed. It has a total of 14k sense wires and 42k field wires, distributed over the  $8000\ \ell$  volume. The signal wires are organized in 56 radial layers



**Figure 2:** The two-step event-building procedure of the Belle II experiment. The Low Level Trigger causes readout of the various subdetectors. Data from the pixel detector (PXD) is stored on a buffer node, the data from the wire-chamber (CDC), the silicon strip detector (SVD) and the other detectors is combined on the first event builder (EVB #1). The High Level Trigger (HLT) processes the data, and transmits regions-of-interest to the buffer node, which then passes the relevant data from the pixel detector to the second event Builder (EVB #2), which combines the pixel data in the regions-of-interest with the data from the other detectors thus creating the final data samples that are stored on disk. An alternative approach to the definition of regions-of-interest is implemented in the DatCon hardware which employs Hough-transform-based reconstruction of the SVD data implemented on FPGA hardware.

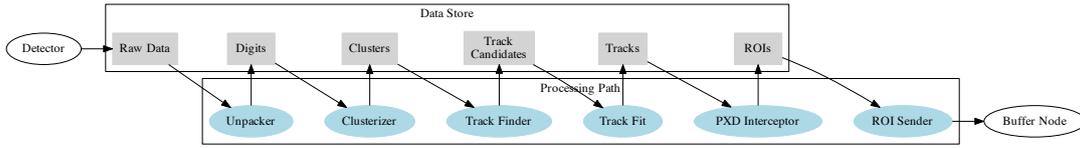
which are subdivided into nine super-layers, of which five have wires along the direction of the magnetic field (“axial wires”) and four are inclined in order to determine the track coordinates in the direction of the field (“stereo wires”) [12].

### 3. The Belle II software framework

Belle II unifies all data processing tasks in a single software package, `basf2` [13–16] which runs on standard Linux systems. This package consists of a flexible framework of individual modules which process the data. The modules are organized in a user-defined chain. Each module can read data provided by the preceding modules via the so-called data store, and it can also add additional data to the data store. Special modules can read data from disk or from detectors, or send and receive data over the network. As an illustration of this, consider a track-finding module: it will read detector hits from the data store and write track candidates to the data store. If it processes MC data, it may also read a relation between the hits and MC truth information, and use this to establish a relation between track candidates and the particles created during the simulation. These track candidates can then be processed by a track fitting module.

Data store contents can be serialized or read at any point in the chain. This functionality is used in several ways. Besides storage of the processed data, this allows distributing the data over multiple cores or networked hosts transparently, and it allows an arbitrary splitting of the tasks between different computer jobs. E.g. an alternative implementation of a module can be tested without repeating any of the work done before: one processes the chain up to, but not including, the module under test, and stores the result to disk. Now any change to the module can be tested immediately, by reading the data from disk and processing the chain starting from the module under inspection.

The design is driven in particular by the needs of the high-level trigger of the Belle II experiment, which uses a networked computing farm running the full reconstruction software, and which



**Figure 3:** Example block-diagram visualization of the data flow in the High-Level Trigger processing chain. Each module is indicated as a blue oval. Grey rectangles indicate the data store arrays. The flow of data is indicated by the arrows.

plays an essential role in the readout of the pixel vertex detector, as explained above. As a complement to Fig. 2, we show as an example of such a processing path the sequence of modules used in a recent beam test where pixel and silicon strip layers as well as the full readout architecture including the high-level trigger were employed [11]. Fig. 3 shows how the incoming raw detector data is sequentially processed to yield the regions-of-interest, sent to the buffer node by the “ROI Sender” module in the end. In the test-beam experiment, this chain ran on the high-level trigger node, indicated by HLT in Fig. 2.

#### 4. Tracking and vertexing software

Physics analysis in Belle II will focus on activity near the interaction point, the objects of interest in a typical analysis are the decay products of the  $B$  and  $\bar{B}$  mesons produced in the initial reactions. The  $B$  mesons decay in close vicinity of a few hundred  $\mu\text{m}$  to the interaction point which in the employed nanobeam scheme is itself very well-defined. In this vein, the reconstructed objects used for analysis are parameterized in the vicinity of the interaction point. In the mDST data format used for analysis, track parameters are stored in a perigee helix parametrization together with information from particle identification.  $V^0$  objects which decay at appreciable distances from the interaction point are reconstructed from their decay tracks before being stored in the mDST files. Photons and  $K_L^0$  are stored using the information from the respective detectors. This way, the original  $B$  decay vertices can be reconstructed efficiently on the level of physics analysis and file sizes are minimized. In the following we will discuss the reconstruction steps that lead to the mDST data.

##### 4.1 Track finding

Track candidates are identified independently in the vertex detector and the wire chamber. In the vertex detector, a cellular-automaton algorithm is used to cope with the potentially large combinatorics, especially in the presence of background [11, 17]. Adjacent pairs of hits are combined into what is called a cell. If two different cells share a hit, they become neighbors in the space of the cellular automaton. The cellular automaton is then evolved where the cell state develops according to the number of geometrically allowed neighbors and the state of the neighbors. From this procedure results a set of possible combinations of hits, where the same hits may appear in several combinations. A Hopfield network, trained with Monte Carlo simulations, is then used to extract the best non-overlapping set of track candidates, which are forwarded to the track fit.

In the wire chamber, besides an adapted version of the track finding algorithm used in the predecessor experiment Belle, two newly developed track finders are used: a global track finder based on Legendre transformations which is limited to tracks emitted in the vicinity of the interaction point, and a local track finder which uses a cellular automaton to follow tracks through the wire chamber volume.

After track fitting, compatible tracks from both subdetectors are merged, and tracks for which no compatible track in the other subdetector could be found are extrapolated, looking for unassociated hits which can be added. The stored tracks are updated to reflect these operations.

## 4.2 Track fitting

The track candidates are then fitted using a largely overhauled version of the GENFIT track fitting package [18, 19]. On the high-level trigger, the default tracking algorithm is the usual Kalman fitter [20]. On the other hand, in physics analysis the adaptive deterministic annealing algorithm [21] is used. It is an extension of the Kalman fitter that uses an annealing procedure to reject background-induced hits that were mistakenly added by the track finding algorithms, and it can also resolve left/right ambiguities in the wire chamber measurements. For the purpose of alignment, tracks are reconstructed with the General Broken Lines algorithm [22], which is a global formulation of the Kalman fit with a discretized model of material. The alignment procedure is described in another contribution to this conference [23].

The fitted tracks are extrapolated outwards to the various particle identification detectors (Cherenkov detectors, electromagnetic calorimeter, and muon/ $K_L^0$  id chambers). The ionization deposits in the silicon detector layers and the wire chamber are also used for  $dE/dx$ -based particle identification. The Cherenkov detector places stringent limits on the precision of the time-of-flight determination which is evaluated during the track fit based on the individual segments between the individual hits. The tracks are extrapolated towards the interaction point, and the helix information in the perigee parametrization is stored in the mDST.

## 4.3 Vertex finding and fitting

In order to reconstruct  $V_0$  type decays, all negative tracks are paired up with all positive tracks, and a vertex fit is performed using the standard Kalman fit implemented in the RAVE package [24]. If a vertex that is separated from the interaction region could be reconstructed by this procedure, a  $V_0$  object is added to the mDST. Physics analysis then uses the reconstructed tracks, the  $V_0$ s and additional calorimeter and  $K_L^0$  objects stored in the mDSTs to recover the  $B$  mesons. The  $B$  vertex is reconstructed using either the RAVE package or the vertexing tool adapted from the Belle experiment, KFIT. As a promising new development, the increased data sample of the Belle II experiment allows to perform with unprecedented statistics analyses where all decay particles of at least one of the  $B$ ,  $\bar{B}$  mesons are reconstructed, and thus the complete event (up to invisible particles from one of the decays) can be reconstructed. The small size of the beam crossing region allows additional strong constraints as the transverse momentum of the  $B$  mesons is strongly correlated with the observed displacement from the interaction point.

## 5. Outlook

The Belle II experiment will begin its physics data taking in 2017. Owing to its significantly increased luminosity, it will allow great improvements in measurements that were thus far limited by statistics. New developments such as the full-event reconstruction will allow to put stringent limits on invisible  $B$  decays or decays with only one charged particle. Not discussed in this proceeding, but worth mentioning are also the large samples of  $D\bar{D}$  and  $\tau^+\tau^-$  events whose analysis will also benefit from Belle II's advanced vertexing capabilities.

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