

Observation and measurements of the production of prompt and non-prompt J/ψ mesons in association with a Z boson in pp collisions at $\sqrt{s}=8\,\mathrm{TeV}$ with the ATLAS detector

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The associated production of a vector boson with heavy quarkonia is a key observable for understanding the quarkonium production mechanisms. In this poster the observation of the production of the Z boson in association with a prompt or with a non-prompt J/ψ meson with the ATLAS detector at LHC is presented and its production rate is measured in comparison of the inclusive Z production. Relative contributions to the signal from single and double parton scattering are estimated. Single parton scattering cross-sections are compared to cutting-edge theoretical calculations in the colour singlet and colour octet formalisms. Finally, a lower limit in the double parton scattering effective cross section is extracted.

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1. Introduction

The associated production of vector boson (\mathcal{V}) with heavy quarkonia (\mathcal{Q}) is a subject of interest from both the experimental [1,2] and the theoretical point of view [3-5]. These final states do not only provide constraints on the contributions from colour-singlet and colour-octet production processes, and their properties [6], but can also be used to study heavy flavour production (non-prompt \mathcal{Q} originating from a b-hadron) in association with a \mathcal{V} . Finally, since the $\mathcal{V}+\mathcal{Q}$ can be produced either from a single parton interaction (SPS) or from the interaction of two different pairs of partons (DPS), it can be used for modelling multi-parton interactions in pp collisions.

The ATLAS Collaboration observed the production of prompt $J/\psi(\to \mu\mu)$ mesons in association with $W(\to \mu\nu)$ bosons [2], using $4.5\,\mathrm{fb^{-1}}$ of $\sqrt{s}=7\,\mathrm{TeV}$ pp collisions at the LHC. Here the observation of the associated production of prompt and non-prompt $J/\psi(\to \mu\mu)$ mesons with $Z(\to \ell\ell)$ bosons ($\ell=\mu,e$) is presented, using $20.3\,\mathrm{fb^{-1}}$ of $\sqrt{s}=8\,\mathrm{TeV}$ pp collisions [7].

2. The ATLAS detector and event selection

The ATLAS detector [8] is a general purpose detector with cylindrical geometry¹ and forward-backward symmetric coverage in pseudorapidity (η). The detector consists of inner tracking detectors (ID), calorimeters, the muon spectrometer (MS) and has a three-level trigger system. The ID directly surrounds the beam pipe and is immersed in a 2T axial magnetic field generated by a superconducting solenoid.

For the search of the associated production of $Z+J/\psi$, events with two opposite-charged lepton pairs are selected, with at least one lepton having $p_T>24\,\mathrm{GeV}$. Each pair is fitted in a common vertex, and only events where the invariant mass of the first pair is between $2.6-3.6\,\mathrm{GeV}$ (J/ψ) and the second between $81.2-101.2\,\mathrm{GeV}$ (Z) are considered. Additionally, the J/ψ candidate must have $p_T^{J/\psi}>8.5\,\mathrm{GeV}$ and $|y_{J/\psi}|<2.1$. In order to reduce contamination from pileup (Z and J/ψ produced from two independent inelastic collisions), the Z and J/ψ vertices are required to be closer than 10 mm in the z-direction.

Muons from Z boson decay are required to have $p_T > 15\,\text{GeV}$ and $|\eta| < 2.5$ and electrons to have $p_T > 15\,\text{GeV}$, $|\eta| < 2.47$ and satisfy isolation requirements based on tracking information (scalar sum of p_T inside an $\eta - \phi$ cone of size $\Delta R = 0.2$ around the lepton to be less than 15% of the lepton p_T). One of the Z boson leptons must be matched with the lepton that fired the trigger and that lepton must have $p_T > 25\,\text{GeV}$ and $|\eta| < 2.4$. For the J/ψ muons, at least one of them must have $p_T > 4\,\text{GeV}$ and an additional requirement of $p_T > 2.5(3.5)\,\text{GeV}$ is applied for those with $|\eta| > 1.3(< 1.3)$. After all selections are applied a total of 290 events are found (139 with $Z \to \mu\mu$ and 151 with $Z \to ee$).

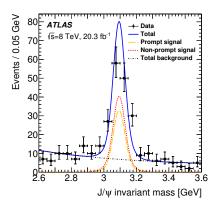
The same requirements described above, applied to the Z boson, are also applied for the inclusive Z candidate selection. Estimates of the background in the inclusive Z sample are obtained

¹ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the *z*-axis along the beam pipe. The *x*-axis points from the IP to the centre of the LHC ring, and the *y*-axis points upward. Cylindrical coordinates (r,ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the beam pipe. The pseudorapidity η is defined in terms of the polar angle θ as $\eta = -\ln\tan(\theta/2)$ and the transverse momentum p_T is defined as $p_T = p\sin\theta$. The rapidity is defined as $y = 0.5\ln((E+p_z)/(E-p_z))$, where *E* and p_z refer to energy and longitudinal momentum, respectively. The $\eta-\phi$ distance between two particles is defined as $\Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2}$.

using both MC simulation and data-driven techniques. The total number of inclusive Z candidates, after background subtraction (which was found to be $0.4 \pm 0.4\%$) is 16.15 million (8.20 million with $Z \to \mu\mu$ and 7.95 million $Z \to ee$).

3. Signal extraction

The $Z+J/\psi$ candidates might originate from pileup, be fake particles mimicking the Z boson or the J/ψ meson and the J/ψ may originate either from prompt QCD interactions or by a b-hadron decay. First step in selecting true $Z+J/\psi$ events is to distinguish prompt and non-prompt J/ψ from fake J/ψ mesons. Prompt J/ψ mesons are separated from non-prompt using a two-dimensional unbinned maximum likelihood fit in the J/ψ invariant mass and pseudo-proper time (see figure 1). Due to the small statistics of the $Z+J/\psi$ associated production sample and to improve the stability of the fit process, the pseudo-proper time and invariant mass of the associated production J/ψ candidates are fitted simultaneously with a high statistics inclusive J/ψ sample, selected with the same criteria as the $Z+J/\psi$ sample.



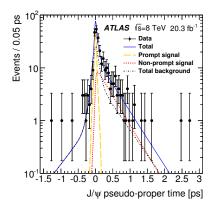


Figure 1: Projection of the unbinned mass (left) and pseudo-proper time (right) maximum-likelihood fit [7].

Performing the fit to the $Z+J/\psi$ sample, 56 ± 10 prompt and 95 ± 12 non-prompt J/ψ mesons are found to be produced in association with a Z boson candidate. After the fit is performed, the sPlot tool [9] is used in order to extract per-event weights for each of the four yield components of the fit (prompt J/ψ , non-prompt J/ψ , prompt background and non-prompt background). Applying the weights coming from the prompt J/ψ and non-prompt J/ψ to the invariant mass distribution of Z boson candidates (see figure 2), the contamination of background to the Z candidates associated with prompt and non-prompt J/ψ mesons can be evaluated. For this, signal and multi-jet templates, extracted from the Powheg MC generator [10] and data respectively, are compared with the sPlot weighted distributions. The number of background events was found to be $0\pm4(1\pm4)$ and $1\pm5(0\pm5)$ for the $Z\to ee(\mu\mu)$ candidates, produced in association with prompt and non-prompt J/ψ mesons.

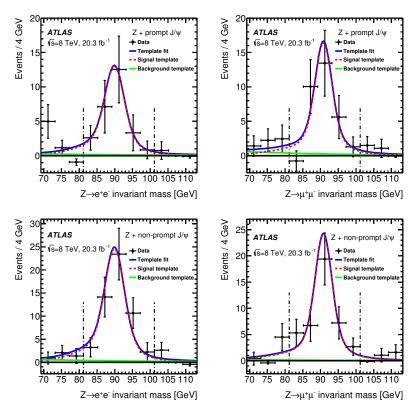


Figure 2: $Z \to ee$ (left) and $Z \to \mu\mu$ (right) mass distributions for Z bosons produced in association with prompt (top) and non-prompt (bottom) J/ψ mesons [7].

As discussed in section 2, the Z and J/ψ vertices are required to be closer than 10 mm, aiming to reduce pileup contamination. The number of pileup candidates are estimated according to the formula $N_{\text{pileup}} = N_{\text{extra}} N_Z P_{J/\psi}$, where N_{extra} is the number of additional vertices which lie within 10 mm of a vertex which produced a Z boson, N_Z is the number of inclusive Z candidates in the fiducial region and $P_{J/\psi}$ is the probability for a J/ψ to be produced at a given pileup vertex. The total number of pileup events were found to be $5.2^{+1.8}_{-1.3}$ and $2.7^{+0.9}_{-0.6}$ for the prompt and non-prompt cases respectively.

The DPS contribution to the $Z+J/\psi$ events, which is treated as part of the signal, is estimated using the assumption that the DPS effective cross-section ($\sigma_{\rm eff}$) is process-independent and that the two hard interactions are uncorrelated. Based on that, for a collision where a Z boson is produced, the probability that a J/ψ meson will be produced in association with the Z is $P_{J/\psi|Z}=\sigma_{J/\psi}/\sigma_{\rm eff}$. $\sigma_{J/\psi}$ is the cross-section of the J/ψ production and the value of $\sigma_{\rm eff}$ is taken to be $\sigma_{\rm eff}=15\pm 3({\rm stat.})^{+5}_{-3}({\rm sys.})$ mb, based on the ATLAS measurement [11]. The estimated number of DPS events were found to be $11.1^{+5.7}_{-5.0}$ and $5.8^{+2.8}_{-2.6}$ for prompt and non-prompt J/ψ mesons produced in association with a Z boson.

4. Results

Using $20.3 \,\mathrm{fb^{-1}}$ of $\sqrt{s} = 8 \,\mathrm{TeV} \,pp$ data, ATLAS Collaboration observed the prompt and non-prompt J/ψ meson production in association with a Z boson with a 5 σ and 9 σ significance respec-

tively. After correcting for the J/ψ muon reconstruction efficiency, the fiducial cross-section ratio, defined as $R_{Z+J/\psi}^{\rm fid} = \sum_{P_{\rm T} \ \rm bins} [N^{\rm ec}(Z+J/\psi)-N^{\rm ec}_{\rm pileup}]/N(Z)$ is measured to be $(36.8\pm6.7({\rm stat.})\pm2.5({\rm syst.}))\times10^{-7}$ for the prompt and $(65.8\pm9.2({\rm stat.})\pm4.2({\rm syst.}))\times10^{-7}$ for the non-prompt $J/\psi+Z$ production. Assuming unpolarised J/ψ decays and correcting for geometric acceptance losses due to J/ψ muon $p_{\rm T}$ and η requirements, the inclusive cross-section ratio is measured to be $(63\pm13({\rm stat.})\pm5({\rm syst.})\pm10({\rm pol.}))\times10^{-7}$ and $(102\pm15({\rm stat.})\pm5({\rm syst.})\pm3({\rm pol.}))\times10^{-7}$ for prompt and non-prompt J/ψ production respectively.

Subtracting DPS contributions, in order to compare the SPS measurement with theoretical models, the Z+ prompt J/ψ production cross-section ratio was found to be $(45\pm13(\text{stat.})\pm6(\text{syst.})\pm10(\text{pol.}))\times10^{-7}$ and Z+ non-prompt J/ψ $(94\pm15(\text{stat.})\pm5(\text{syst.})\pm3(\text{pol.}))\times10^{-7}$. Leading Order (LO) colour-singlet mechanism calculations vary between $(11.6\pm3.2)\times10^{-8}$ [5] and $46.2^{+6.0}_{-6.5}\times10^{-8}$ [4] and Next to LO (NLO) NRQCD model predicts a cross-section ratio of $(45.7^{+10.5}_{-9.6})\times10^{-8}$ [4], underestimating ATLAS measurement, as illustrated on figure 3.

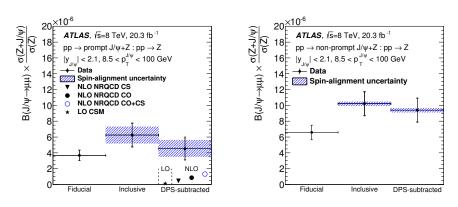
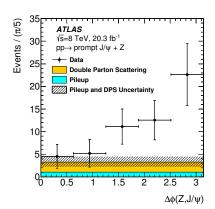


Figure 3: Production cross-section ratios of prompt (left) and non-prompt(right) J/ψ mesons produced in association with a Z boson relative to inclusive Z production [7].

5. Double Parton Scattering

Figure 4 (left) shows the azimuthal angle between the Z boson and the prompt J/ψ momentum vectors ($\Delta\phi$) after the application of sPlot weights corresponding to the prompt J/ψ signal component. The estimated DPS and pileup contribution are overlaid with yellow and cyan colours respectively. DPS events are expected to be distributed uniformly along the $\Delta\phi$ variable because the Z and J/ψ particles are produced from two independent scatters. SPS events are expected to show a back-to-back correlation $\Delta\phi=\pi$, since they originate from a single parton interaction, with a smearing due to detector effects, presence of gluons in the final state and radiation from the leptons and by the intrinsic properties of the protons.

Based on above, the low $\Delta\phi$ region which is sensitive to DPS, can be used to limit the maximum allowed DPS contribution to the observed signal, which corresponds to a lower limit on $\sigma_{\rm eff}$. Using the assumption that the first bin $\Delta\phi(Z,J/\psi)<\pi/5$ is DPS dominated, the data uncertainties and uncertainties inherent in the DPS estimate allow a lower limit $\sigma_{\rm eff}>5.3\,{\rm mb}(3.7\,{\rm mb})$ at 68%(95%) confidence level to be extracted (see figure 4 right).



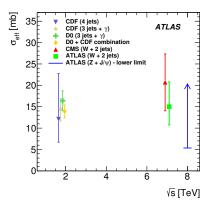


Figure 4: Left: Azimuthal angle between the Z boson and the J/ψ meson after the application of J/ψ signal prompt sPlot weights. Right: Measurements and limits on $\sigma_{\rm eff}$ as a function of \sqrt{s} [7] (JHEP 03 032 (2014), New J. Phys. 15 (2013) 033038, Phys. Rev. D47 4857-4871 (1993), Phys. Rev. D56 3811-3832 (1997), Phys. Rev. D81 052012 (2010), JHEP 113 1310 (2013)).

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