

On the primary model to explain the relation between a rigidity-dependent spectral hardening of proton and helium spectra and a sharp knee of the all-particle spectrum

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Cosmic Ray (CR) spectrum in some energy region can not be well expressed by a simple power-law as it has a structure. Recently, AMS-02, PAMELA, ATIC-2 and CREAM-2 presented a rigidity dependent CR spectra with a remarkable hardening in excess of around 300 GV. On the other hand, the all-particle energy spectrum of primary CRs observed in a wide range from 10^{14} eV to 10^{17} eV with the Tibet-III AS array clearly shows a sharp knee at around 4 PeV. Based on these results, in this paper, we propose a phenomenological model capable of well explaining both the data by AMS-02 and others in the lower energy region and that by the Tibet AS γ experiment simultaneously. We discuss some details of our model in which extra nearby CR sources are responsible for creating a sharp knee. This model also predicts a dominance of CR nuclei heavier than helium at the knee energies.

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1. Introduction

The spectrum of CRs has offered a few clues to its origin so far. The only features observed are at very high and ultrahigh energies, the so-called knee at a few times 10^{15} eV, the second "knee" at $\sim 10^{18}$ eV, the "ankle" at higher energies, and a spectral steepening above 10^{18} eV [1][2][3][4]. But, recently CR measurements by the new-generation experiments such as the AMS-02, ATIC-2, CREAM-2 and PAMELA seem to indicate that the CR spectrum deviates from a single power law in the rigidity region lower than 1 TeV[5][6][7][8]. These measurements have caused a wide range of interest. Many models were proposed to account for the hardening of the flux based on different sources, acceleration mechanisms, diffusive propagation effects, and their superposition [9]. On the other hand, The Tibet AS γ Collaboration has measured the all-particle spectrum in a wide energy range between 10^{14} eV and 10^{17} eV by air-shower array of the area of 37000 m^2 located at 4300m above sea level (Tibet, Yangbajing, China, atmospheric depth of 606 g/cm^2) and provided the most detailed spectrum around the knee which clearly shows a sharp knee at around 4 PeV[10]. The origin of the knee in the energy spectrum of cosmic rays is an outstanding problem in astroparticle physics. Proposals for its origin range from astrophysical scenarios such as a change of acceleration mechanisms at the sources of cosmic rays (supernova remnants, pulsars, etc.), to a single-source assumption or effects due to propagation inside the galaxy (diffusion, drift, escape from the Galaxy), to particle physics models such as the interaction with relic neutrinos during transport or new processes in the atmosphere during air-shower development[11][12][13][14][15]. In order to resolve the origin of the knee, the detailed study of the chemical composition provides a key information since the change of the energy spectrum must be closely related to its chemical composition.

Based on these results, we here propose a phenomenological model capable of well explaining both the data by AMS-02 and others in the lower energy region and that by the Tibet AS γ experiment simultaneously. We discuss some details of our model show that extra nearby CR sources can well explain a sharp knee. This model also predicts a dominance of cosmic-ray nuclei heavier than helium at the knee.

2. Direct measurements and hardening spectrum

Measurements of the energy spectra of different nuclei show a consistent concave shape over the common range 30-1000 GeV/nucleon (virtually the same rigidity) from Helium to Iron. Protons show a similar pattern up to a higher energy per nucleon. Appropriate approximation for the energy spectra of primary nuclei taking into account the rate of change of spectral slope was reported in [16] where E is the energy of a primary nucleus A with charge Z , $R = E_k / Z$ is the rigidity dependent energy at which the asymptotic energy spectral power index γ_1 for $E \ll E_k$ is changed to the asymptotic value γ_2 ($\gamma_2 = \gamma_1 + \Delta\gamma_1$) for $E \gg E_k$ at sharpness parameter ξ correlating with the rate of change of the spectral slope. So the broken power-law energy spectrum can be expressed as

$$h(E) = \Phi_z * E^{-\gamma_1} * (1 + (E/E_k)^{\xi})^{(-\Delta\gamma_1/\xi)} \quad (2.1)$$

Above about 10 GeV, the solar modulation is negligible. The reason for γ_1 being different for proton and other nuclei (heavier than proton) is beyond the scope of the present work, and

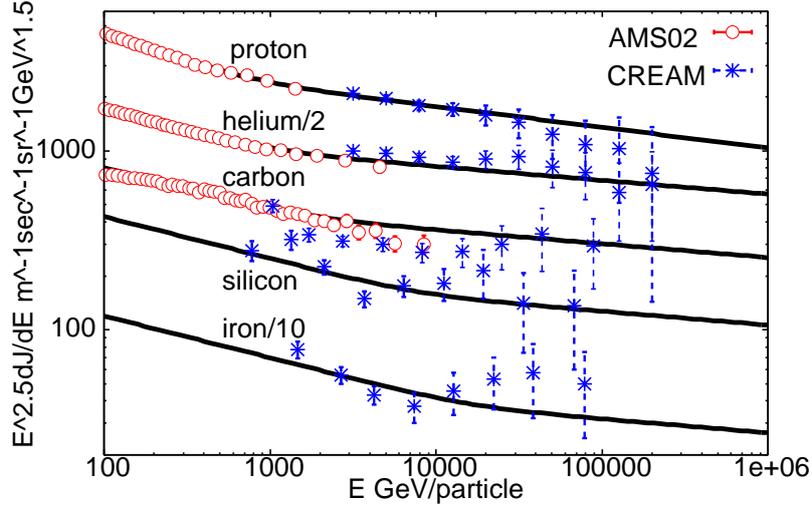


Figure 1: Elemental spectrum expressed by a broken power-law function at lower energy region (<100 TeV). The cited data are AMS-02[5] and CREAM-2 [7].

Table 1: Parameters of the Chemical Composition Used in the Present Work. The adopted values in function(2.1) for Spectra index γ_1 , hardening energy E_k and spectral index difference $\Delta\gamma_1$ are used to describe hardening spectrum at low energy , another two parameters: cutoff energy E_c and spectral index difference $\Delta\gamma_2$ are obtained by fitting the function (3.1) to characterise the knee structure at high energy.

nuclei	γ_1	E_k	$\Delta\gamma_1$	E_c	$\Delta\gamma_2$
p	2.8	700GeV	0.2	200TeV	0.3
other	2.7	$5*Z$ TeV	0.16	$200*Z$ TeV	0.3

we divide proton and other nuclei as two groups. The values of E_k , γ_1 and $\Delta\gamma_1$ are taken from experimental data in the energy range from several 10^{10} eV to several 10^{14} eV obtained by AMS-02 and CREAM-2. The adopted values for E_k , γ_1 and $\Delta\gamma_1$ are listed in Table 1, these parameters can be used to describe hardening spectrum at lower energy region (<100 TeV). Our model use another two parameters: cutoff energy E_c and spectral index difference $\Delta\gamma_2$ to characterise the knee structure at higher energy region (>100 TeV) which can be seen below. Some experimental results of the direct observation for individual elements (proton, helium, carbon, oxygen, neon, magnesium, silicon and iron) are summarized in Fig. 1, in which the solid line represents a broken power-law function.

3. Indirect experiments

Individual spectra of primary CRs as directly measured at the top of the atmosphere for energies below 100 TeV, are extrapolated to high energies and fitted to results from air-shower measurements. We require this model should not only can well explain the data by AMS-02 and others in the lower energy region, but also simultaneously coincide to the all-particle spectrum measured by the Tibet AS γ experiment which is the highest statistical and the best systematics-controlled measurement covering the widest energy range around the knee energy region . The energy spectra of

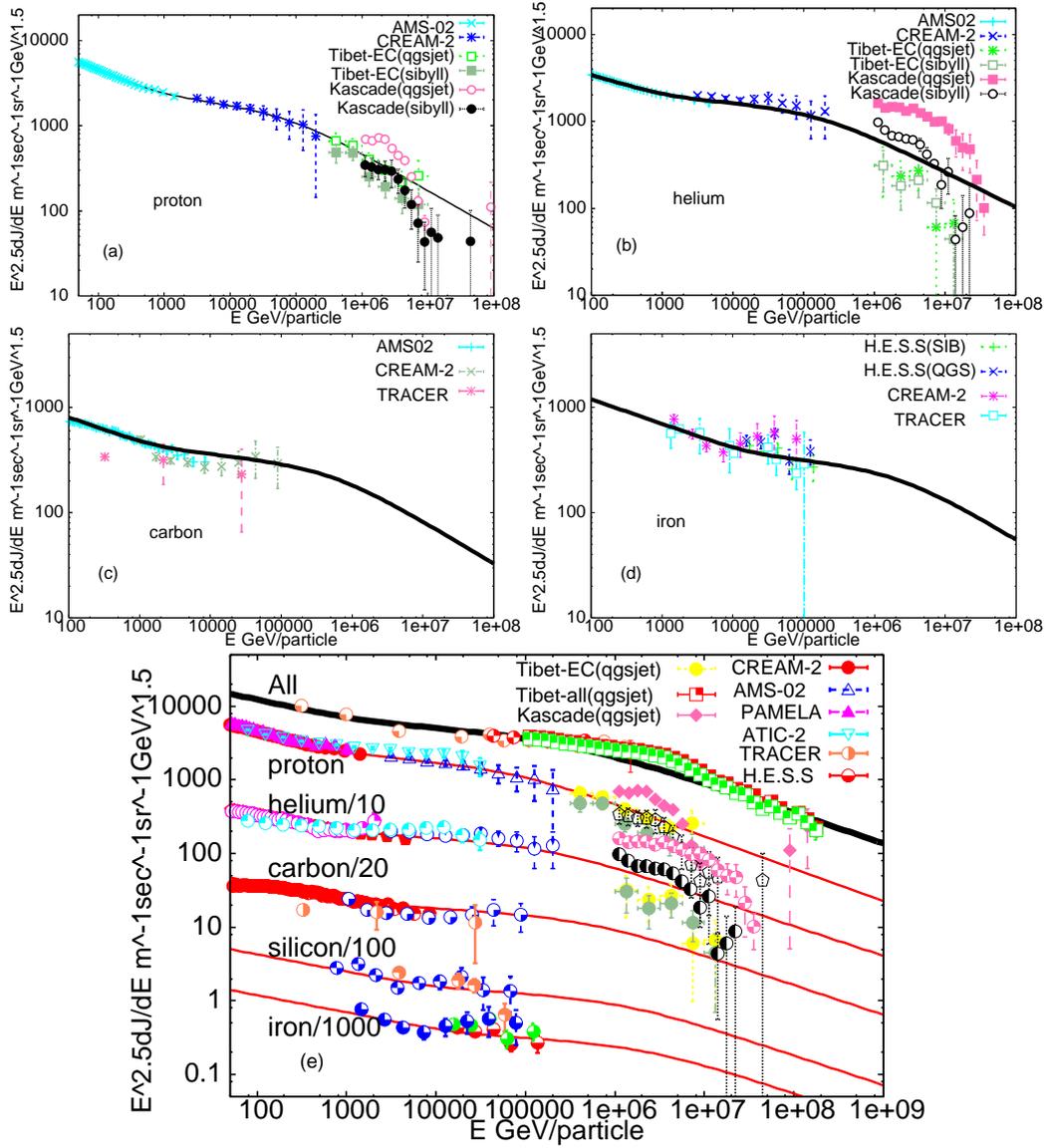


Figure 2: Energy spectra of proton(a), helium(b), carbon(c),iron(d) and the all-particle(e).References of the data are proton: AMS-02[5], CREAM-2[7], Tibet-EC[17], Cascade[18]; Helium: AMS-02[5], CREAM-2[7], Tibet-EC[17], Cascade[18]; Carbon: AMS-02[5], CREAM-2[7], TRACER[19]; iron: CREAM-2[7], HESS[20], TRACER[19]; all-particle: Tibet-III[10]. The solid line in each panel represents a broken power-law type behavior of the energy spectra.

individual nuclei derived from extensive air shower experiments can also be described by a broken power-law function with a cut-off at high energies. We use this kind of function to parameterize the observed differential energy spectra for individual elements of CRs, which can be written as

$$f(E) = h(E) * (1 + (E/E_c)^s)^{-\Delta\gamma/2/s} \quad (3.1)$$

According to the Diffusive Shock Acceleration model, which is considered to be the most promising model for the acceleration of the galactic cosmic rays at SNRs, the acceleration limit

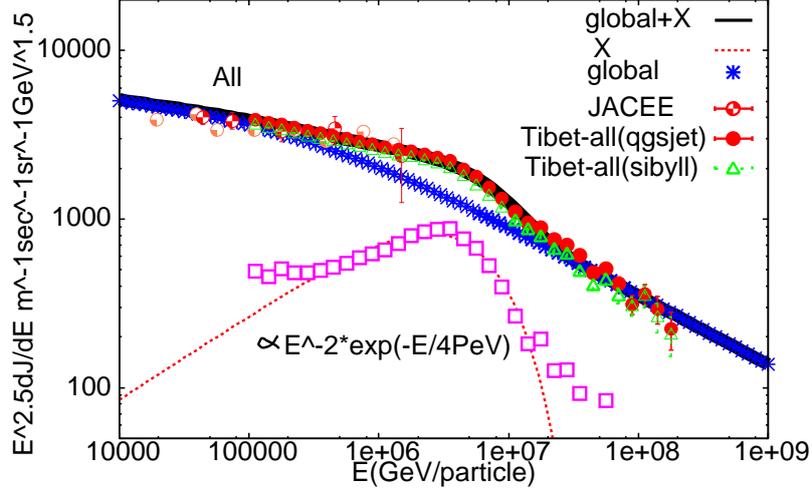


Figure 3: All-particle spectrum around the knee. Black line is the sum of the global component and an extra component whose functional form is expressed as $\propto E^{-2} \exp(-E/4\text{PeV})$.

is expected at Z times 100 TeV, where Z is the atomic number. Inspired by these theories, the following ansatz is adopted to describe the energy dependence of the flux for particles with charge Z . Here, we assume that the cut-off energy of proton E_c is about 200 TeV. The cut-off energy for each component depends on its charge Z , namely $E_c(z) = Z * E_c(p)$ is assumed. We then sum up individual spectrum to obtain the all-particle spectrum called a global component. The numerical values of fit parameters for each chemical component are listed in Table 1. The calculated energy spectra of proton, helium, carbon, iron and the total spectrum, together with the observational data are shown in Fig. 2. The results show a good agreement with the data of each elemental species, but the sharpness of the knee is more prominent in the experimental data.

4. Excess Component at the knee

The all-particle spectrum above 10^{15} eV has been obtained by many air-shower experiments, all of which have shown a sudden change in the spectral power index around 4×10^{15} eV. Although our model described here has been designed to account for this change, the sharpness of the knee is more prominent in the experimental data, as shown in Fig. 2(e). Subtracting the global component from all-particle spectrum measured by the Tibet AS γ experiment, we can see a hard spectrum of power index close to -2, which is just the expected value of source spectrum before the modulation by the propagation, and it also shows a cutoff feature as shown in Fig. 3 indicating that the excess component is due to the contribution of nearby source(s). The enhancement around the knee can be reproduced well whose functional form is expressed as $\propto E^{-2} \exp(-E/4\text{PeV})$, as shown in Fig. 3. The sharp peak of the excess component indicates that the chemical composition is not mixed as much as global component but its main component consists of elements with limited range of atomic numbers if we adopt the rigidity dependent acceleration limit. Individual spectra of elements is calculated assuming that the excess component with certain nuclei called X component consisting entirely of proton, helium, carbon, or iron at the knee, as shown in Fig. 4. Because the experimental data points are still poor in the high energy range, it seems difficult to judge which is

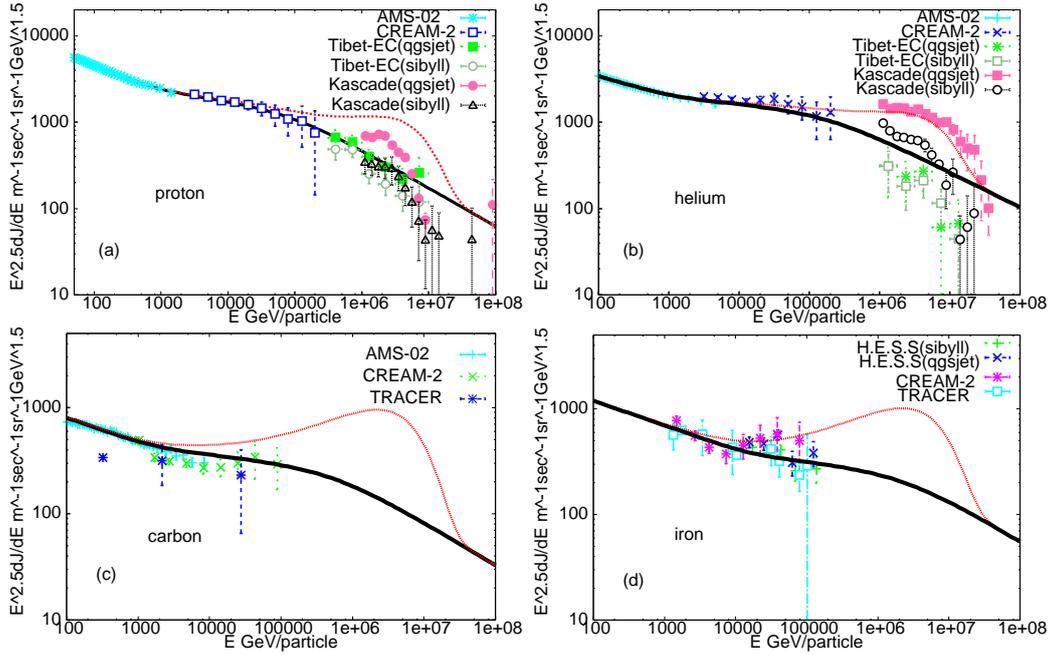


Figure 4: Individual spectra of elements is calculated assuming that the excess component with a certain nuclei called X component comprised only of proton (a), helium (b), carbon (c) and iron (d) each at the knee. References of the data are proton: AMS-02[5], CREAM-2[7], Tibet-EC[17], Cascade[18]; Helium: AMS-02[5], CREAM-2[7], Tibet-EC[17], Cascade[18]; Carbon: AMS-02[5], CREAM-2[7], TRACER[19]; iron: CREAM-2[7], HESS[20], TRACER[19]; all-particle: Tibet-III[10].

the X component except protons due to its much higher flux compared with present experimental data.

5. Chemical composition around the knee

In such a simple scenario, average mass $\langle \ln A \rangle$ is calculated assuming that the excess X component at the knee which is described in Fig.4 (a), (b), (c) and (d) as shown in Fig. 5. While the experimental data are still divergent because of the difficulty of the composition measurement at very high energy, some results fairly agree with the lines calculated by the present work. The result of Tibet-EC data for proton and helium shows that the excess component is heavier than helium. Fig. 5 may suggest that the most possible nuclei of the excess component is heavier than helium, which is not contradict with the result of Tibet-EC data[17].

6. Conclusion and discussion

We propose a phenomenological model capable of well explaining both the data by AMS-02 and others in the lower energy region and that by the Tibet AS γ experiment simultaneously. In this model, an excess component is assumed to overlap the global component and the shape of the energy spectrum of the excess component suggests that it may be attributed to CRs from nearby source(s) since the spectral shape is very close to that expected by DSA model of SNR. This model also predicts a dominance of cosmic-ray nuclei heavier than helium at the knee.

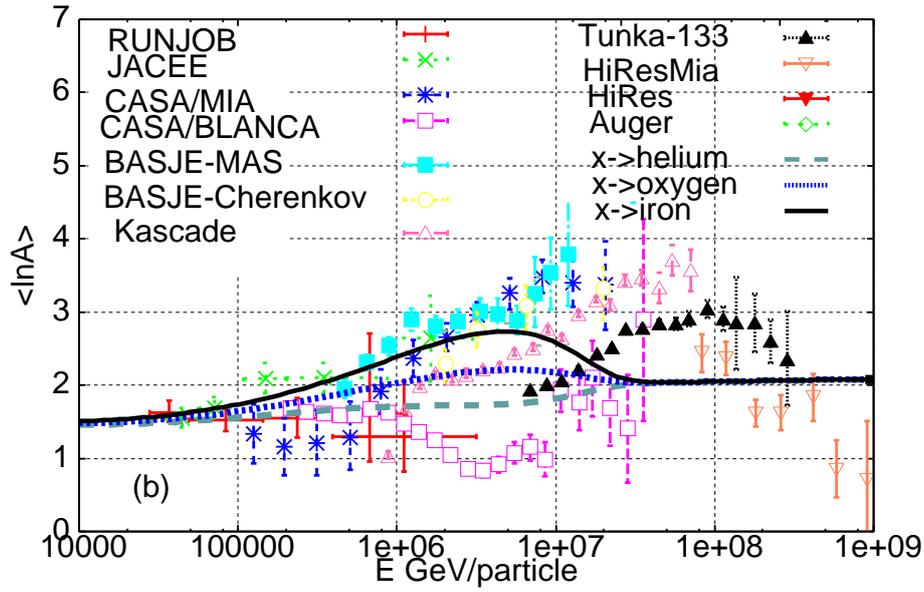


Figure 5: Average mass calculated in this model. The three lines show a case of the excess component comprised only of helium (green broken line), oxygen (blue dash line) and iron (black solid line) each.

To identify such source(s), which might be diffused, it is interesting to investigate primary gamma rays beyond 10^{13} eV, which can be produced by nuclear interactions of accelerated particles in the source region. This can be examined by observing air showers with a capability of p/γ separation. Further measurements of the chemical composition of cosmic rays will improve the model parameters used in the present analysis and enable more detailed discussion of the structure of the knee and the contribution of extra galactic cosmic rays.

7. Acknowledgements

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References

- [1] Kulikov, G. V., & Khristiansen, G. B. 1958, Soviet Physics JETP, 35, 635
- [2] Haungs, A., Rebel, H., & Roth, M. 2003, Reports on Progress in Physics, 66, 1145
- [3] Abbasi, R., et al. 2005, Physics Letters B, 619, 271
- [4] Abbasi, R. U., et al. 2009, Astroparticle Physics, 32, 53
- [5] AMS-02 Collaboration, "AMS Days at CERN γ and Latest Results", 15, April, 2015 Abraham, J., et al. 2010, Physics Letters B, 685, 239

- [6] For the ATIC experiment see A. D. Panovet et al., *Bull. Russ. Acad. Sci. Phys.* 73, 564 (2009)
- [7] For the CREAM experiment see Y. S. Yoon et al., *Astrophys. J.* 728, 122 (2011).
- [8] For the PAMELA experiment see O. Adriani et al., *Astrophys. J.* 765, 91 (2013); *Science* 332, 69 (2011).
- [9] See, for example, G. Bernard, T. Delahaye, P. Salati, and R. Taillet, *Astron. Astrophys.* 555, A48 (2013); V.S. Ptuskin, V. Zirakashvili, and E. S. Seo, *Astrophys. J.* 763, 47 (2013); N. Tomassetti, *Astrophys. J. Lett.* 752, L13 (2012); P. Blasi, E. Amato, and P. D. Serpico, *Phys. Rev. Lett.* 109, 061101 (2012); A. E. Vladimirov, G. Jóhannesson, I. V. Moskalenko, and T. A. Porter, *Astrophys. J.* 752, 68 (2012).
- [10] Amenomori, M et al, 2008, *Astrophys. J.*, 678, 1165-1179
- [11] Berezhko, E. G., & Ksenofontov, L. G. 1999, *J. Exp. Theor. Phys.*, 89, 391
- [12] Erlykin, A. D., & Wolfendale, A. W. 2005, *Astropart. Phys.*, 23, 1
- [13] Ptuskin, V. S., et al. 1993, *A&A*, 268, 726
- [14] Wigmans, R. 2003, *Astropart. Phys.*, 19, 379
- [15] Nikolsky, S. I., & Romachin, V. A. 2000, *Phys. At. Nuclei*, 63, 1799
- [16] Horandel, J. R. 2003, *Astropart. Phys.*, 19, 193
- [17] Amenomori, M., et al. 2006, *Phys. Lett. B*, 632, 58
- [18] T. Antoni et al., *Astropart. Phys.* 24, 1 (2005).
- [19] Ave, M., et al. 2008, *ApJ*, 678, 262
- [20] Aharonian, F. A., et al. 2007, *Phys. Rev. D*, 75, 042004