

# D-meson measurements in Pb-Pb collisions with the ALICE detector at the LHC

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> Heavy quarks (charm and beauty) are effective probes to investigate the properties of the hot and dense strongly-interacting medium created in heavy-ion collisions as they are produced in partonic scattering processes occurring in the early stages of the collision. Due to their long life time, they probe all the stages of the medium evolution and they interact with its constituents, losing energy via gluon radiation and elastic collisions. The measurement of the D-meson nuclear modification factor provides a key test of parton energy-loss models. These models predict that beauty quarks lose less energy than charm quarks and the latter experience less in-medium energy loss than light quarks and gluons.

> D-meson production was measured with ALICE in Pb–Pb collisions at  $\sqrt{s_{\rm NN}} = 2.76$  TeV. D mesons were reconstructed via their hadronic decays at central rapidity. We will discuss the latest results of the measurement of the D-meson nuclear modification factor as a function of transverse momentum ( $p_{\rm T}$ ) and collision centrality. In addition, the measurement of the azimuthal anisotropy ( $v_2$ ) of charmed-meson production as a function of  $p_{\rm T}$  in different centrality intervals will be presented.

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# **3** 1. Introduction

Heavy quarks are unique probes to study the Quark-Gluon Plasma produced in heavy-ion col-4 lisions at the LHC. Due to their large masses, they are produced predominantly in hard scatterings, 5 during the initial phase of the collision. Therefore, they experience the entire evolution of the 6 medium created in the collision and can act as probes of its properties. One of the key methods 7 used to characterise the medium (i.e its medium density, temperature and heavy-quark transport 8 coefficient) is the study of energy loss of the partons traversing it. In a QCD picture, radiative 9 in-medium energy loss is one of the main mechanisms expected to contribute, with dependence 10 on the mass and the color charge of the particle. The radiation is suppressed at small angles for 11 massive partons because of the dead-cone effect [1] and is larger for gluons, which have stronger 12 color charge with respect to quarks (Casimir coupling factor). Therefore, a hierarchy in the  $R_{AA}$ 13 is expected to be observed when comparing the mostly gluon-originated light-flavor hadrons (e.g. 14 pions) to D and to B mesons. The measurement and comparison of these different probes of the 15 medium provides a unique test of the colour charge and mass dependence of parton energy loss [2]. 16 Further insight into the medium properties is provided by the measurement of the anisotropy in 17 the azimuthal distribution of particle momenta. In non-central heavy-ion collisions, the anisotropy 18 in the spatial distribution of the nucleons participating in the collision is converted into a momen-19 tum anisotropy, if sufficient rescatterings with the medium constituents occur. Hence the azimuthal 20 distribution of the particles in the final state reflects the initial anisotropy and the medium char-21 acteristics. The azimuthal anisotropy of produced particles is characterized by the second Fourier 22 coefficient  $v_2 = \langle \cos[2(\varphi - \Psi_2)] \rangle$ , where  $\varphi$  is the azimuthal angle of the particle momentum, and 23  $\Psi_2$  is the azimuthal angle of the initial-state symmetry plane for the second harmonic [3]. At low 24  $p_{\rm T}$ , the  $v_2$  of heavy-flavour hadrons is sensitive to the degree of thermalization of charm and beauty 25 quarks in the deconfined medium. At higher  $p_{\rm T}$ , the measurement of  $v_2$  carries information on the 26 path length dependence of in-medium parton energy loss. The measurement of heavy-flavour  $v_2$ 27 offers a unique opportunity to test whether also quarks with large mass participate in the collective 28 expansion dynamics and possibly thermalize in the QGP. 29

## 30 2. D-meson measurement with ALICE

ALICE (A Large Ion Collider Experiment) is the LHC experiment dedicated to heavy-ion 31 studies. ALICE [4] consists of: a barrel at central rapidity, a muon spectrometer at forward rapidity 32 and a set of detectors for global properties determination. For the present analysis, we use the 33 information from a subset of the central barrel detectors, namely the Inner Tracking System (ITS), 34 the Time Projection Chamber (TPC) and the Time Of Flight detector (TOF) for charged particle 35 tracking and identification, the T0 for collision time measurement and the VZERO scintillator for 36 triggering and centrality measurement. The two tracking detectors, ITS and TPC, allow the recon-37 struction of charged-particle tracks in the pseudorapidity range  $-0.9 < \eta < 0.9$  with a momentum 38 resolution better than 2% for  $p_{\rm T} < 20$  GeV/c. The TPC provides the particle identification via a 39 dE/dx measurement. The ITS, in particular, is a key detector for open heavy-flavour studies be-40 cause it allows us to measure the track impact parameter (i.e. the distance of closest approach of the 41 track to the primary vertex) with a resolution better than 65  $\mu$ m for  $p_{\rm T} > 1$  GeV/c, thus providing 42

the capability to reconstruct secondary vertices originating from heavy-flavour hadron decays. The

<sup>44</sup> TOF detector provides particle identification by time of flight measurement.

<sup>45</sup> D<sup>0</sup>, D<sup>+</sup>, D<sup>\*+</sup> and D<sup>+</sup><sub>s</sub> mesons are reconstructed in the central rapidity region from their <sup>46</sup> hadronic decay channels D<sup>0</sup>  $\rightarrow$  K<sup>-</sup> $\pi^+$ , D<sup>+</sup>  $\rightarrow$  K<sup>-</sup> $\pi^+\pi^+$ , D<sup>\*+</sup>  $\rightarrow$  D<sup>0</sup> $\pi^+ \rightarrow$  K<sup>-</sup> $\pi^+\pi^+$  and D<sup>+</sup><sub>s</sub>  $\rightarrow$ <sup>47</sup>  $\phi\pi^+ \rightarrow$  K<sup>+</sup>K<sup>-</sup> $\pi^+$  (and charge conjugates). The D-meson yield is measured with an invariant mass <sup>48</sup> analysis of fully reconstructed decay topologies displaced from the interaction vertex selected by <sup>49</sup> requiring a large decay length and a good alignment between the reconstructed D-meson momen-<sup>50</sup> tum vector and its flight line. The identification of charged kaons in the TPC and TOF detectors <sup>51</sup> helps to further reduce the background at low  $p_{\rm T}$ .

<sup>52</sup> In this proceeding, we will discuss the ALICE measurements of D-meson production in Pb–Pb applicipate of  $\sqrt{2}$  = 2.76 TeV

collisions at  $\sqrt{s_{\rm NN}}$  = 2.76 TeV.

# 54 **2.1 Nuclear Modification Factor** ( $R_{AA}$ )

The nuclear modification factor is defined as the ratio of the transverse momentum spectrum measured in nucleus-nucleus (AA) collisions to the one measured in pp collisions at the same centre of mass energy, rescaled by the average number of binary nucleon-nucleon collisions ( $N_{coll}$ ) expected in heavy-ion collisions. The ratio can be expressed also in terms of average nuclear overlap function ( $T_{AA}$ ) estimated within the Glauber model [5].

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$$R_{\mathrm{AA}}(p_{\mathrm{T}}) = rac{1}{\langle T_{\mathrm{AA}} 
angle} rac{dN_{\mathrm{AA}}/dp_{\mathrm{T}}}{d\sigma_{\mathrm{pp}}/dp_{\mathrm{T}}}$$

The pp reference is scaled from 7 to 2.76 TeV using pQCD calculations (FONLL) [6]. The scaling was validated comparing the scaled results with the available measurement in pp collisions at  $\sqrt{s} = 2.76$  TeV [7] (this sample was not used in the R<sub>AA</sub> measurement due to the limited statistics).



**Figure 1:** Left:  $R_{AA}$  of  $D^0$ ,  $D^+$ ,  $D^{*+}$  and  $D_s^+$  mesons as a function of  $p_T$  in central collision [9, 10]. Right: Average  $R_{pPb}$  of D mesons as a function of  $p_T$  compared to D-meson  $R_{AA}$  in 20% most central and 40-80% Pb-Pb collisions.

Figure 1 (left) represents the nuclear modification factor of  $D^0$ ,  $D^+$  and  $D^{*+}$  mesons as a function of  $p_T$  in most central collisions [9]. The  $R_{AA}$  of the three D-meson species are compatible

- within uncertainties. A suppression up to a factor five is seen at  $p_{\rm T} \sim 10 \,{\rm GeV}/c$ . The first measure-67 ment of  $D_s^+ R_{AA}$  [10] in heavy-ion collisions is also shown. In the highest measured  $p_T$  bin (8-12) 68 GeV/c), the  $R_{AA}$  of  $D_s^+$  mesons is compatible with that of non-strange charmed mesons. At lower 69  $p_{\rm T}$ , the  $R_{\rm AA}$  of  $D_{\rm s}^+$  seems to increase, but with the current statistical and systematic uncertainties 70 no conclusion can be drawn on the expected enhancement of  $D_s^+$ -meson production with respect 71 to that of non-strange D mesons at low  $p_{\rm T}$ , due to c-quark coalescence with the abundant strange 72 quarks [8]. 73 Figure 1 (right) shows the average D-meson nuclear modification factor measured in minimum-74 bias p-Pb collisions at  $\sqrt{s_{\rm NN}}$  = 5.02 TeV [11]. Since no significant modification of the D-meson 75
- <sup>76</sup> yield is observed in p-Pb collisions for  $p_T > 2$  GeV/c, it can be concluded that the strong suppres-
- <sup>77</sup> sion observed in central Pb-Pb collisions is due to the interaction of heavy quarks with the hot and
- 78 dense medium.



**Figure 2:** D-meson nuclear modification factor,  $R_{AA}$ , in Pb-Pb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV. Left:  $R_{AA}$  as a function of  $p_T$  compared to charged hadrons and pions. Right:  $R_{AA}$  as a function of  $N_{part}$  [13] compared to non-prompt J/ $\psi$  measured by the CMS collaboration [14].

<sup>79</sup> D-meson nuclear modification factor is similar to that of charged pions and charged particles <sup>80</sup> in the measured  $p_{\rm T}$  interval within uncertainties as shown in Fig. 2 (left). It should be noted that <sup>81</sup> the  $R_{\rm AA}$  of D mesons and pions is also sensitive to the shape of the parton momentum distribution <sup>82</sup> and their fragmentation functions. Model calculations including those effects and a colour-charge <sup>83</sup> hierarchy in parton energy loss are able to describe the measurements, see [12] as example.

In Fig. 2 (right),  $R_{AA}^{D}$  as a function of collision centrality (quantified by the average number 84 of participant nucleons) [13] is shown. This measurement is compared with results from the CMS 85 collaboration of non-prompt  $J/\psi$  [14] and theoretical predictions from [15, 16, 21]. For D mesons, 86 a smaller suppression in peripheral than in central collisions is observed. A larger suppression 87 in central collisions is seen for D mesons than for non-prompt  $J/\psi$ , indicating a different energy 88 loss for charm and beauty quarks. This observation is supported by predictions from energy-loss 89 models, where the difference between the  $R_{AA}$  of D and B mesons arises from the different masses 90 of c and b quarks. 91

#### 92 **2.2 Elliptic Flow** (*v*<sub>2</sub>)

The  $v_2$  of prompt D<sup>0</sup>, D<sup>+</sup> and D<sup>\*+</sup> mesons at mid-rapidity was measured in the three centrality classes 0-10%, 10-30% and 30-50% [17, 18]. The D<sup>0</sup> meson  $v_2$  as a function of  $p_T$  in the three centrality classes is shown in the Fig.3, compared with the measured  $v_2$  of charged particles.



**Figure 3:**  $v_2$  of  $D^0$  as a function of  $p_T$  in three centrality ranges, compared with the  $v_2$  of charged particles [17, 18].

The average of the  $v_2$  of D<sup>0</sup>, D<sup>+</sup> and D<sup>\*+</sup> in the centrality class 30-50% is larger than zero with 96  $5\sigma$  significance in the range  $2 < p_T < 6$  GeV/c. A positive  $v_2$  is also observed for  $p_T > 6$  GeV/c, 97 which most likely originates from the path length dependence of the in-medium partonic energy 98 loss, although the present statistics does not allow to give a firm conclusion on this. The measured 99 D-meson  $v_2$  is comparable in magnitude with that of the charged particles, which are mostly light-100 flavour hadrons. This result indicates that low- $p_{\rm T}$  charm quarks take part in the collective motion 10 of the system. The  $v_2$  of D mesons decreases from peripheral to central collisions as expected due 102 to the decreasing initial geometrical anisotropy. 103

 $R_{AA}$  and  $v_2$  are two complementary measurements to gain insight into the heavy-quark transport coefficient of the medium. Several theoretical model calculations are available for the  $R_{AA}$ and  $v_2$  of heavy flavour hadrons. Fig. 4 shows the D-meson  $R_{AA}$  (left) and  $v_2$  (right) compared to predictions from various models [19, 20, 21, 22, 23, 24, 25, 26]. A simultaneous description of the  $R_{AA}$  and  $v_2$  starts to provide constraints to the models themselves.

#### 109 3. Conclusions

The ALICE detector provides excellent tracking, vertexing and particle identification to allow measurements of charmed mesons via their hadronic decays, over a wide range of transverse momentum. The nuclear modification factor and the elliptic azimuthal anisotropy of prompt D mesons were measured in Pb-Pb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV. The first measurement of the D<sup>+</sup><sub>s</sub>-meson production has also been presented. The D-meson nuclear modification factor is similar to that of charged particles and pions in central collisions. An ordering is observed for nuclear modification factor of prompt D and non-prompt J/ $\psi$  in central collisions, in agreement with the expectation of a



**Figure 4:** D-meson  $R_{AA}$  and  $v_2$  compared to model predictions. Only models with predictions for both  $R_{AA}$  and  $v_2$  are shown. Left: D-meson  $R_{AA}$  as a function of  $p_T$ . Right: D-meson  $v_2$  as a function of  $p_T$ .

larger in-medium energy loss of charm compared to beauty quarks. A positive  $v_2$  is measured, suggesting that charm quarks take part in the collective expansion of the medium. The models based on parton energy loss in the medium describe the data within the uncertainties, but the simultaneous description of  $R_{AA}$  and  $v_2$  remains challenging. The improvement of statistical and systematic

<sup>121</sup> uncertainties expected in Run 2 will help to provide further constraints to models.

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