

# Observation and measurements of the production of prompt and non-prompt $J/\psi$ mesons in association with a $Z$ boson in $pp$ collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector

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A key observable for understanding the quarkonium production mechanism is the associated production of a vector boson with heavy quarkonia. ATLAS Collaboration at LHC observed the associated production of prompt and non-prompt  $J/\psi$  mesons with  $Z$  boson with  $5\sigma$  and  $9\sigma$  significance, respectively. In this poster the measurement of the production rate of  $Z + J/\psi$  over inclusive  $Z$  is discussed. Additionally, contributions from single and double parton interactions are evaluated and the results are compared to latest theoretical calculations in the colour singlet and colour octet formalisms. Finally, a lower limit of the double parton scattering effective cross section is extracted.

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## 1. Introduction

A process that ignited the interest of both the theoretical [1–3] and experimental communities is the associated production of vector boson ( $\mathcal{V}$ ) with heavy quarkonia ( $\mathcal{Q}$ ). Such processes are particularly interesting, enabling the study of the contributions from colour-singlet (CS) and colour-octet (CO) models [4], the heavy flavour production (non-prompt  $\mathcal{Q}$  originating from a  $b$ -hadron) in association with a  $\mathcal{V}$  and finally, the multi-parton interactions in hadronic collisions (since the two final state particles can be produced either from a single parton interaction, SPS, or from the interaction of two different pairs of partons, DPS).

The ATLAS Collaboration contributed in the study of the  $\mathcal{V} + \mathcal{Q}$  processes, observing the production of prompt  $J/\psi(\rightarrow \mu\mu)$  mesons in association with  $W(\rightarrow \mu\nu)$  bosons [5], using  $4.5 \text{ fb}^{-1}$  of  $\sqrt{s} = 7 \text{ TeV}$   $pp$  collisions at the LHC and the associated production of prompt and non-prompt  $J/\psi(\rightarrow \mu\mu)$  mesons with  $Z(\rightarrow \ell\ell)$  bosons ( $\ell = \mu, e$ ) [6]. The latter analysis used  $20.3 \text{ fb}^{-1}$  of  $\sqrt{s} = 8 \text{ TeV}$   $pp$  collisions and is the process discussed here.

## 2. Event selection and reconstruction

The ATLAS detector [7] is a general purpose detector with cylindrical geometry<sup>1</sup> and forward-backward symmetric coverage in pseudorapidity ( $\eta$ ). The detector consists of inner tracking detectors (ID), calorimeters, the muon spectrometer (MS) and has a three-level trigger system. The ID directly surrounds the beam pipe and is immersed in a 2 T axial magnetic field generated by a superconducting solenoid.

Both  $Z$  and  $J/\psi$  particles are reconstructed in their di-lepton decay mode ( $Z \rightarrow \ell\ell$ ,  $\ell = \mu, e$  and  $J/\psi \rightarrow \mu\mu$ ). Events with two opposite-charged lepton pairs are selected, with at least one lepton having  $p_T > 24 \text{ GeV}$ . Each lepton pair is then fitted in a common vertex, and only events where the invariant mass of the first pair is between  $2.6 - 3.6 \text{ GeV}$  ( $J/\psi$ ) and the second between  $81.2 - 101.2 \text{ GeV}$  ( $Z$ ) are selected for the analysis. Additional requirements of the  $J/\psi$  candidate are  $p_T^{J/\psi} > 8.5 \text{ GeV}$  and rapidity<sup>2</sup>  $|y_{J/\psi}| < 2.1$ . In order to reduce contamination from pileup ( $Z$  and  $J/\psi$  produced from two independent inelastic collisions), the  $Z$  and  $J/\psi$  vertices are required to be closer than 10 mm in the  $z$ -direction.

Selection requirements to muons from  $Z$  boson decay are  $p_T > 15 \text{ GeV}$  and  $|\eta| < 2.5$  and electrons are required to have  $p_T > 15 \text{ GeV}$ ,  $|\eta| < 2.47$  and must satisfy isolation requirements based on tracking information (scalar sum of  $p_T$  inside an  $\eta - \phi$  cone of size  $\Delta R = 0.2$  around the lepton to be less than 15% of the lepton  $p_T$ ). One of the leptons originating from the  $Z$  boson must be matched with the lepton that fired the trigger and must have  $p_T > 25 \text{ GeV}$  and  $|\eta| < 2.4$ . For the  $J/\psi$  muons, at least one of them must have  $p_T > 4 \text{ GeV}$  and an additional requirement of  $p_T > 2.5(3.5) \text{ GeV}$  is applied for those with  $|\eta| > 1.3(< 1.3)$ . A total of 290 events are found (139 with  $Z \rightarrow \mu\mu$  and 151 with  $Z \rightarrow ee$ ) after all selections are applied.

Since this analysis measures the  $Z + J/\psi$  production ratio over inclusive  $Z$ , the same requirements described above, applied to the  $Z$  boson, are also applied for the inclusive  $Z$  candidate selection. Both MC simulation and data-driven techniques are employed to estimate background contributions in the inclusive  $Z$  sample. After background subtraction, which was found to be  $0.4 \pm 0.4\%$ , the total number of inclusive  $Z$  candidates is 16.15 million (8.20 million with  $Z \rightarrow \mu\mu$  and 7.95 million  $Z \rightarrow ee$ ).

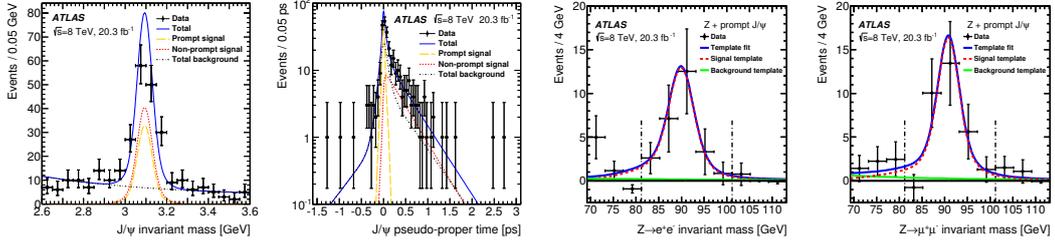
## 3. Signal extraction

Possible sources of  $Z + J/\psi$  candidates are pileup, fake particles that mimic the  $Z$  boson or the  $J/\psi$  meson and that the  $J/\psi$  may originate either from prompt QCD interactions or by a  $b$ -hadron decay. This analysis first selects true  $Z + J/\psi$  events by distinguishing prompt and non-prompt  $J/\psi$  from fake

<sup>1</sup>ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the  $z$ -axis along the beam pipe. The  $x$ -axis points from the IP to the centre of the LHC ring, and the  $y$ -axis points upward. Cylindrical coordinates  $(r, \phi)$  are used in the transverse plane,  $\phi$  being the azimuthal angle around the beam pipe. The pseudorapidity  $\eta$  is defined in terms of the polar angle  $\theta$  as  $\eta = -\ln \tan(\theta/2)$  and the transverse momentum  $p_T$  is defined as  $p_T = p \sin \theta$ . The  $\eta$ - $\phi$  distance between two particles is defined as  $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$ .

<sup>2</sup>The rapidity is defined as  $y = 0.5 \ln((E + p_z)/(E - p_z))$ , where  $E$  and  $p_z$  refer to energy and longitudinal momentum, respectively.

$J/\psi$  mesons. This is realised with a two-dimensional unbinned maximum likelihood fit in the  $J/\psi$  invariant mass and pseudo-proper time (see figure 1), since non-prompt  $J/\psi$  mesons are expected to have longer pseudo-proper times. The shape related parameters of the fit to the  $Z + J/\psi$  sample are driven from a high statistics inclusive  $J/\psi$  sample, which is fitted simultaneously with the  $Z + J/\psi$  sample. This prevents instabilities that may occur from the low statistics of the  $Z + J/\psi$  sample.



**Figure 1:** Projection of the unbinned mass (left) and pseudo-proper time (middle-left) maximum-likelihood fit.  $Z \rightarrow ee$  and  $Z \rightarrow \mu\mu$  mass distributions for  $Z$  bosons produced in association with prompt  $J/\psi$  mesons (middle-right and right) [6].

A total of  $56 \pm 10$  prompt and  $95 \pm 12$  non-prompt  $J/\psi$  mesons are found to be produced in association with a  $Z$  boson candidate from the fit procedure. After the fit is performed, the sPlot tool is used, so per-event weights for each of the four yield components of the fit (prompt  $J/\psi$ , non-prompt  $J/\psi$ , prompt background and non-prompt background) can be extracted. Weights from prompt  $J/\psi$  and non-prompt  $J/\psi$  yields are further applied to the invariant mass distribution of  $Z$  boson candidates (see figure 1). From these distributions, showing the  $Z$  boson candidates produced in association with either a prompt or a non-prompt  $J/\psi$  meson, the contamination from background sources is evaluated. Signal and multi-jet templates, extracted from the Powheg MC generator and data respectively, are applied and compared with the sPlot weighted distributions. Backgrounds were estimated to be  $0 \pm 4(1 \pm 4)$  and  $1 \pm 5(0 \pm 5)$  for the  $Z \rightarrow ee(\mu\mu)$  candidates, produced in association with prompt and non-prompt  $J/\psi$  mesons.

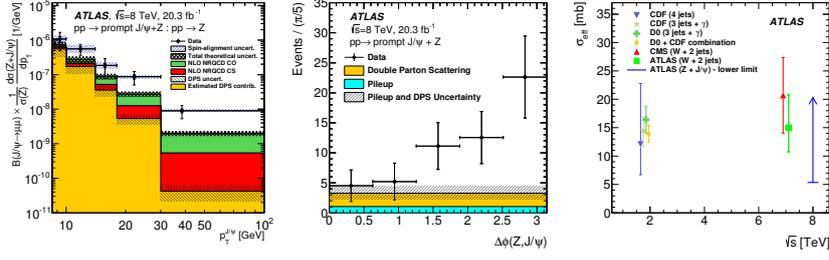
In order to reduce pileup contamination, the  $Z$  and  $J/\psi$  vertices are required to be closer than 10 mm in the  $z$  direction. The number of pileup candidates are estimated according to the formula  $N_{\text{pileup}} = N_{\text{extra}} N_Z P_{J/\psi}$ , where  $N_{\text{extra}}$  is the number of additional vertices which lie within 10 mm of a vertex which produced a  $Z$  boson,  $N_Z$  is the number of inclusive  $Z$  candidates in the fiducial region and  $P_{J/\psi}$  is the probability for a  $J/\psi$  to be produced at a given pileup vertex. Using the formula above, the total number of pileup events were found to be  $5.2^{+1.8}_{-1.3}$  and  $2.7^{+0.9}_{-0.6}$  for the prompt and non-prompt cases respectively.

The DPS is not considered as a background source and, as it is treated as part to the signal, its contribution is estimated. This estimation is using the assumption that the DPS effective cross-section ( $\sigma_{\text{eff}}$ ) is process-independent and that the two hard interactions are uncorrelated. Based on that, for a collision where a  $Z$  boson is produced, the probability that a  $J/\psi$  meson will be produced in association with the  $Z$  is  $P_{J/\psi|Z} = \sigma_{J/\psi} / \sigma_{\text{eff}}$ .  $\sigma_{J/\psi}$  is the cross-section of the  $J/\psi$  production and the value of  $\sigma_{\text{eff}}$  is taken to be  $\sigma_{\text{eff}} = 15 \pm 3(\text{stat.})^{+5}_{-3}(\text{sys.})$  mb, based on the ATLAS measurement [8]. The estimated number of DPS events were found to be  $11.1^{+5.7}_{-5.0}$  and  $5.8^{+2.8}_{-2.6}$  for prompt and non-prompt  $J/\psi$  mesons produced in association with a  $Z$  boson.

#### 4. Results

ATLAS Collaboration, using  $20.3 \text{ fb}^{-1}$  of  $\sqrt{s} = 8 \text{ TeV}$   $pp$  data, observed the prompt and non-prompt  $J/\psi$  meson production in association with a  $Z$  boson with a  $5\sigma$  and  $9\sigma$  significance respectively. The fiducial cross-section ratio, defined as  $R_{Z+J/\psi}^{\text{fid}} = \sum_{p_T^{J/\psi} \text{ bins}} [N^{\text{ec}}(Z + J/\psi) - N_{\text{pileup}}^{\text{ec}}] / N(Z)$ , where  $N^{\text{ec}}(Z + J/\psi)$  is the yield of  $Z + J/\psi$  events after correcting for  $J/\psi$  muon reconstruction efficiency,  $N(Z)$  is the background-subtracted yield of inclusive  $Z$  events and  $N_{\text{pileup}}^{\text{ec}}$  is the efficiency-corrected expected pileup background contribution in the fiducial  $J/\psi$  acceptance, is measured to be  $(36.8 \pm 6.7(\text{stat.}) \pm 2.5(\text{syst.})) \times 10^{-7}$  for the prompt and  $(65.8 \pm 9.2(\text{stat.}) \pm 4.2(\text{syst.})) \times 10^{-7}$  for the non-prompt  $J/\psi + Z$  production. Assuming unpolarised  $J/\psi$  decays and correcting for geometric ac-

ceptance losses due to  $J/\psi$  muon  $p_T$  and  $\eta$  requirements, the inclusive cross-section ratio is measured to be  $(63 \pm 13(\text{stat.}) \pm 5(\text{syst.}) \pm 10(\text{pol.})) \times 10^{-7}$  and  $(102 \pm 15(\text{stat.}) \pm 5(\text{syst.}) \pm 3(\text{pol.})) \times 10^{-7}$  for prompt and non-prompt  $J/\psi$  production respectively.



**Figure 2:** (Left) Production cross-section of  $J/\psi$  in association with a  $Z$  boson as a function of the prompt  $J/\psi$   $p_T$ , normalised to the inclusive  $Z$  cross-section. (Middle) Azimuthal angle between the  $Z$  boson and the  $J/\psi$  meson after the application of  $J/\psi$  signal prompt sPlot weights. (Right) Measurements and limits on  $\sigma_{\text{eff}}$  as a function of  $\sqrt{s}$  (right) (JHEP 03 032 (2014), New J. Phys. 15 (2013) 033038, Phys. Rev. D47 4857-4871 (1993), Phys. Rev. D56 3811-3832 (1997), Phys. Rev. D81 052012 (2010), JHEP 113 1310 (2013)) [6].

The inclusive  $Z$ + prompt or non-prompt  $J/\psi$  production over inclusive  $Z$  cross-section ratio is also measured as a function of the  $p_T^{J/\psi}$ , as illustrated in figure 2. Together with the inclusive differential cross-section ratio, the differential DPS contribution is shown. The measurement is compared with NLO CS and CO calculations. The combination of the DPS and NLO CO+CS contributions underestimates the rate observed in data.

## 5. Double Parton Scattering

A sensitive variable to study DPS is the azimuthal angle between the  $Z$  boson and the prompt  $J/\psi$  momentum vectors ( $\Delta\phi$ ). The  $\Delta\phi$  distribution, after the application of sPlot weights corresponding to the prompt  $J/\psi$  signal component, is shown in figure 2 (middle). SPS events are expected to show a back-to-back correlation  $\Delta\phi = \pi$ , while DPS events are expected to be distributed uniformly along the  $\Delta\phi$  variable, because the  $Z$  and  $J/\psi$  particles are produced from two independent scatters.

Based on the above, the low  $\Delta\phi$  region ( $\Delta\phi(Z, J/\psi) < \pi/5$ ), which is dominated by DPS, is used to place a lower limit to the maximum allowed DPS contribution to the observed signal. The lower limit was found to be  $\sigma_{\text{eff}} > 5.3 \text{ mb}(3.7 \text{ mb})$  at 68%(95%) confidence level (see figure 2 right).

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