

## Studies of Higgs boson production in the four-lepton final state at $\sqrt{s} = 13$ TeV on CMS detector

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Studies of Higgs boson production are presented using the  $H \rightarrow ZZ \rightarrow 4\ell$  ( $\ell = e, \mu$ ) decay channel. These studies are performed using a data sample corresponding to an integrated luminosity of  $2.8 \text{ fb}^{-1}$  of pp collisions at a center-of-mass energy of 13 TeV collected by the CMS experiment at the LHC during 2015. The observed significance for the standard model Higgs boson with  $m_H = 125.09 \text{ GeV}$  is  $2.5\sigma$ , where the expected significance is  $3.4\sigma$ . The model independent fiducial cross section is measured to be  $\sigma_{\text{fid.}} = 2.48_{-1.14}^{+1.48}(\text{stat.} \oplus \text{sys.})_{-0.04}^{+0.01}(\text{model dep.}) \text{ fb}$ . In addition, a search for an additional Higgs boson is performed for a range of masses up to 1 TeV and with various widths, and no significant excess is observed. The results of this search are interpreted in the context of the two Higgs doublet model.

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## 1. Introduction

The start of the LHC Run 2 in 2015, at an increased center-of-mass energy of  $\sqrt{s} = 13$  TeV, opens the way for an era of new precision measurements of the Higgs boson. This note presents the first CMS studies of Higgs boson production in the  $H \rightarrow 4\ell$  decay channel at  $\sqrt{s} = 13$  TeV, using  $2.8 \text{ fb}^{-1}$  of proton-proton collision data collected with the CMS experiment at the LHC in 2015. Detail of the studies can be found in Ref. [1] with references therein.

## 2. Background estimation

The irreducible background to the Higgs signal in the  $4\ell$  channel comes from the production of  $ZZ$  via  $q\bar{q}$  annihilation or gluon fusion. The contribution from these backgrounds are estimated using simulation. Higher order QCD and electroweak corrections are applied by K factors as a function of four-lepton mass.

Reducible backgrounds arise from  $Z + \text{jets}$  and  $t\bar{t}$  production where heavy-flavor jets produce secondary leptons, and from  $Z + \text{jets}$ ,  $Z\gamma + \text{jets}$ ,  $WW + \text{jets}$  and  $WZ + \text{jets}$  events where decays of heavy-flavor hadrons, in-flight decays of light mesons within jets, or (for electrons) the decay of charged hadrons overlapping with  $\pi^0$  decays are misidentified as leptons. Two independent control regions in data are formed by requiring the presence of two leptons which satisfy the full identification criteria, plus two additional opposite-sign or same-sign leptons satisfying relaxed identification requirements. The event yields in the control regions are weighted by the lepton misidentification probability to obtain the expected yield in the signal region.

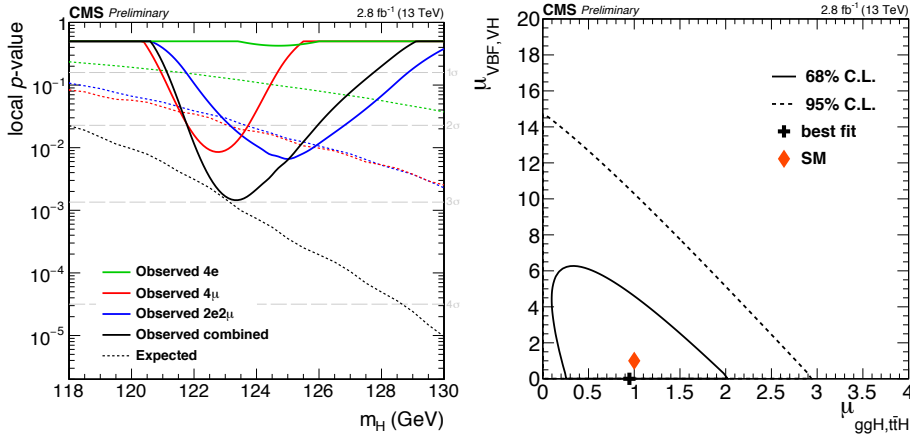
## 3. Signal modeling

The signal lineshape of a narrow resonance around  $m_H \sim 125$  GeV and mass resolution for higher-mass resonances are parametrized using a double-sided Crystal Ball function. In addition, a Landau function is also added in the total probability density function for the non-resonant part of the signal for the case of  $WH$ ,  $ZH$  and  $t\bar{t}H$  production modes. In the case of a search for an additional high-mass resonance, a general parameterization of the signal and background  $gg \rightarrow 4\ell$  process is used which includes all contributing amplitudes for  $gg \rightarrow \text{bkg} + H(125)^* + X(m_X) \rightarrow 4\ell$ , and their interference.

## 4. Results

### 4.1 Significance and signal strength

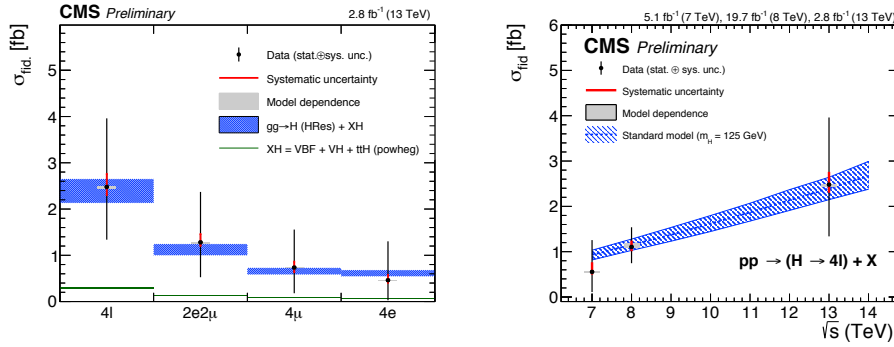
Figure 1 (left) shows the significance of the local fluctuation with respect to the SM background-only expectation as a function of  $m_H$ . The minimum of the local  $p$ -value is reached at 123.4 GeV and corresponds to a local significance of  $3.0\sigma$ , while  $3.1\sigma$  are expected for the SM Higgs boson. Two signal-strength modifiers  $\mu_{ggH, t\bar{t}H}$  and  $\mu_{VBF, VH}$  are defined as scale factors for the fermion and vector-boson induced contributions to the expected SM cross section. A two-dimensional fit is performed assuming a mass of  $m_H = 125.09$  GeV and profiling the likelihood for all nuisance parameters, leading to the measurements of  $\mu_{ggH, t\bar{t}H} = 0.95_{-0.49}^{+0.64}$  and  $\mu_{VBF, VH} = 0.0_{-0.0}^{+2.5}$ . The 68% and 95% CL contours in the  $(\mu_{ggH, t\bar{t}H}, \mu_{VBF, VH})$  plane are shown in Fig.1 (right).



**Figure 1:** (left) Significance of the local fluctuation with respect to the background-only expectation as a function of the Higgs boson mass. (right) Result of the 2D likelihood scan for the  $\mu_{\text{ggH, ttH}}$  and  $\mu_{\text{VBF, VH}}$  signal-strength modifiers. The solid and dashed contours show the 68% and 95% CL regions, respectively.

#### 4.2 Fiducial cross section

The fiducial volume for the cross section measurement is defined to match closely the reconstruction level selection and is very similar to the definition used in Ref. [2]. The only difference in the fiducial volume definition with respect to the Run 1 analysis is that the leptons are dressed by adding the four-momenta of photons within  $\Delta R < 0.4$  to the bare leptons. The integrated fiducial cross section is measured to be  $\sigma_{\text{fid.}} = 2.48^{+1.46}_{-1.13}(\text{stat.})^{+0.28}_{-0.18}(\text{sys.})^{+0.01}_{-0.04}(\text{model dep.})$  fb. This can be compared to the SM expectation  $\sigma_{\text{fid.}}^{\text{SM}} = 2.39 \pm 0.25$  fb. The integrated fiducial cross section as a function of  $\sqrt{s}$  is also shown in Fig. 2 (right).

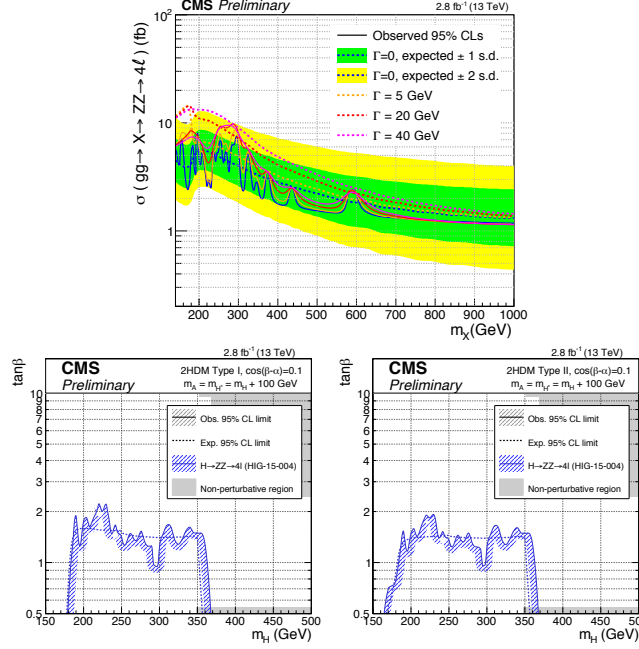


**Figure 2:** Result of the measured fiducial cross section in each final state (left) and the measured fiducial cross section as a function of  $\sqrt{s}$  (right).

#### 4.3 Search for additional resonances

Figure 3 shows observed and expected 95% CL limits on the  $\text{gg} \rightarrow \text{X} \rightarrow \text{ZZ} \rightarrow 4\ell$  cross section including the four-lepton branching fraction, as a function of  $m_X$ , for several values of  $\Gamma_X$ .

In addition to the generic cross section limit, the results are also interpreted in the context of a two Higgs doublet model (2HDM) benchmark scenario. The parameters  $m_H$ , the mass of the heaviest CP-even Higgs boson, and  $\tan\beta$  are chosen for a two-dimensional scan. The results are shown in Fig. 3.



**Figure 3:** Top: Upper limits at the 95% CL on the  $gg \rightarrow X \rightarrow ZZ \rightarrow 4\ell$  cross section  $\sigma$ , including four-lepton branching fraction, as a function of  $m_X$ , shown for several values of  $\Gamma_X$ . Bottom: Exclusion limits at the 95% CL in the  $m_H - \tan\beta$  plane for the benchmark 2HDM scenario with Type I (left) and Type II (right) couplings.

## 5. Summary

Several studies of Higgs boson production in the four-lepton final state at  $\sqrt{s} = 13$  TeV have been presented, using data samples corresponding to an integrated luminosity of  $2.8 \text{ fb}^{-1}$ . The observed significance for the SM Higgs boson at a mass of  $m_H = 125.09$  GeV is  $2.5\sigma$ . The measured signal strength modifiers and the model-independent fiducial cross section for this boson are consistent, within their uncertainties, with the expectations for the SM Higgs boson. In addition, upper limits at a 95% CL are set on the production of an additional Higgs boson for masses up to 1 TeV and for various widths, and these results are interpreted in two 2HDM scenarios.

## References

- [1] CMS Collaboration, CMS Physics Analysis Summary Report No. CMS-PAS-HIG-15-004, 2015. [<http://cdsweb.cern.ch/record/2139978>].
- [2] CMS Collaboration, "Measurement of differential cross sections for Higgs boson production in the diphoton decay channel in pp collisions at  $\sqrt{s}=8$  TeV", Eur. Phys. J. C 76 (2015) 13