 Searches for hadronically decaying Dark Matter mediator particles at ATLAS

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Searches for hadronic resonances of the Dark Matter (DM) particles in the sub-TeV mass region remain as a viable target at ATLAS. However, due to the bandwidth limitation, the events that available for performing an analysis were statistically limited. Reducing the event size by recording a fraction of the full event information overcomes this limitation. An analysis that is performed on those events is called Trigger-Level Analysis (TLA). This poster highlights the TLA strategy used to search for low-mass dijet resonances. No significant excesses are found in a region between 450 and 950 GeV. As an addition, limits are set on a simplified leptophobic $Z'$ model of DM mediator with axial coupling to quarks and DM particles as well as on Gaussian resonances.

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1. Introduction

A search for DM mediator particles via dijet resonance in the sub-TeV region benefits greatly from the high luminosity run in ATLAS Run-2[1]. However, a trigger prescaling needs to be applied to that region due to a bandwidth limitation. As a consequence, numerous potential events are discarded, limiting events availability for this kind of search.

With fixed bandwidth allocation, recording only a subset of full jet information will improve event rate considerably. Limited information in these events is partially-reconstructed and calibrated to be used later on in the Trigger-Level Analysis (TLA).

2. Trigger-level analysis

Trigger-Level Analysis (TLA) is the usage of the High Level Trigger (HLT) objects which collected via Data Scouting (DS) technique to perform a physics analysis. Figure 1 shows the DS stream has higher rates compared to the sum of single trigger rates.

3. High level trigger calibration and performance

Jets, as a collection of particles, which reconstructed at the trigger level are calibrated to be as close as possible to jets in the particle level. Since DS stream does not save tracking information, calibration procedures such as origin correction and Global Sequential Correction (GSC)[2] which require tracking are not applied. Figure 2a shows dijet invariant mass of leading and sub-leading HLT jets compared to its offline counterpart in data. Black line indicates the average $m_{jj}$ response between HLT and offline jets. It is shown that the difference is within 1% of the unity value. Additional $\eta$ inter calibration is also performed to correct for detector dependence effect as shown in Figure 2b.
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4. Search phase and limit setting

The search is performed by selecting events with at least two jets event with $|\eta| < 2.8$. The minimum leading jet $p_T$ is 185 GeV, with sub-leading jet should have 85 GeV of transverse momentum as a minimum value. Two signal regions then defined as $y^* < 0.6$ and $y^* < 0.3$. Figure 3a shows the benefit of unprescaled trigger in the TLA jets. The number of events below 800 GeV are significantly increased. One can also see the agreement between the shape of $m_{jj}$ distribution between HLT and offline jets.

Most discrepant regions that are identified by the Bumphunter algorithm [3] lies between 574 and 685 GeV as shown in Figure 3b. It has a p-value of 0.44 for $y^* < 0.6$ and 0.19 for $y^* < 0.3$. Hence, there is no significant excesses found.
The observed upper limit on $m_{jj}$ on the cross-section, $\sigma$ times acceptance, $A$, times branching ratio, $BR$ of a leptophobic $Z\bar{\nu}A'Z$ simplified model are shown in Figure 4b. Gaussian contributions of $Z'$ signal ranging from 3 pb at 450 GeV, to 9 pb at 600 GeV, to 0.7 pb at 850 GeV are excluded with 95% credibility level. Figure 4a shows the upper limits for the $Z'$ model obtained from the $m_{jj}$ distribution plotted in quark couplings, $g_q$, as a function of $Z'$ mass, $m_{Z'}$.

References

