

Exploring the Masses of Exotic Heavy Pentaquarks

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The masses of crypto exotic heavy pentaquarks have been probed considering a di-hadronic state consisting of a meson and a baryon . The interaction between the hadrons is assumed to be Van der Waals' type of molecular interaction and we have estimated the binding energies of the states. A spin interaction has also been considered. Masses of the crypto heavy pentaquarks such as $P_s^*(1), P_s^*(2), P_b^*(1), P_b^*(2)$, have been predicted . Moreover we have also calculated the masses of exotic pentaquarks for the charm and the bottom families constituting a heavy quark.

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1. Introduction

The existence of pentaquark charmonium states with the decay of $\Lambda_b^0(\Lambda_b^0 \rightarrow J/\Psi K^- p)$ has been reported by LHCb [1] recently. The intermediate states have been identified as $P_c^*(4380)$ and $P_c^*(4450)$. New impetus has been put forward to the study of the properties and dynamics of the multi-quark states by Jaffe et al [2]. The works of Lipkin [3], Gignoux et al [4], Monemzadeh et al [5], Diakonov [6] in the context of the exotic doublets ($cuud\bar{s}$) ($cudd\bar{s}$), Bicudo [7] for possible crypto heptaquark hadrons with exotic flavors, Kopeliovich et al [8] for cryptoexotic pentaquarks with hidden beauty are noteworthy.

2. Methodology

In the present work exotic heavy pentaquarks have been described as di-hadronic molecules consisting of a meson and a baryon held together by Van der Waals' type of molecular interaction. The mass formula for the low-lying di-hadronic molecule is taken as

$$M_{Tot} = M_1 + M_2 + E_{BE} + E_{SD} \quad (2.1)$$

where M_1 and M_2 represent the masses of the constituent hadrons respectively, E_{BE} represents the binding energy of the di-hadronic system and E_{SD} represents the spin dependent term. The binding energy can be expressed as

$$E_{BE} = \langle \Psi(r_{12}) | V(r_{12}) | \Psi(r_{12}) \rangle \quad (2.2)$$

where r_{12} is the radius parameter of the di-hadronic molecule and $V(r)$ is the interaction potential taken as Van der Waals' type [9] which is expressed as

$$V(r_{12}) = -\frac{k_{mol}}{r_{12}} e^{-C^2 r_{12}^2 / 2} \quad (2.3)$$

where k_{mol} is the residual strength of the strong interaction molecular coupling and C is the effective color screening of the confined gluons. $\Psi(r_{12})$ is the wave function of the dihadronic state and is assumed to be in the context of the statistical model as [10],

$$|\Psi(r_{12})|^2 = \frac{315}{64\pi r_0^{9/2}} (r_0 - r_{12})^{3/2} \theta(r_0 - r_{12}) \quad (2.4)$$

corresponding to the linear type of background potential. r_0 is the radius of the di-hadronic molecule, $\theta(r_0 - r_{12})$ is the usual step function. Employing the additive rule for the radii of constituent hadrons i.e. $r_0 = r_1 + r_2$, r_1 and r_2 representing the individual radii of the hadrons constituting the molecule, we have executed the binding energy using equations (2.2), (2.3) and (2.4), which yields

$$E_{BE} = \frac{2.25 k_{mol}}{r_0} [{}_2F_2[(1.5, 1), (2.75, 2.25), -\beta]] \quad (2.5)$$

if $Re\beta > 0$, ${}_2F_2$ is the relevant hypergeometric function and $\beta = C^2 r_{12}^2 / 2$. Input values are $C = 50 MeV$ [11] $k_{mol} = 0.65$ [10]. The radii considered are $r(p) = 6 GeV^{-1}$, $r(n) = 4.7 GeV^{-1}$, $r(K) = 4.77 GeV^{-1}$, $r(\Phi) = 5.0 GeV^{-1}$, $r(Y) = 1.63 GeV^{-1}$, $r(D^-) = 4.97 GeV^{-1}$, $r(B_s) = 3.67 GeV^{-1}$, $r(D_s) = 4.8 GeV^{-1}$, $r(B^o) = 3.82 GeV^{-1}$, $r(\Sigma) = 3.9 GeV^{-1}$, $r(\Delta) = 5.977 GeV^{-1}$.

Particles	Approach	Mass (GeV)	Expt. Mass (GeV)	Others Mass (GeV)
$P_s^*(1/2)^+$	$p + \Phi$	2.110	—	2.303 [8]
$P_s^*(3/2)^+$ ($uuds\bar{s}$)	$\Delta + \Phi$	2.381	—	2.373 [8]
$P_b^*(1/2)^+$	$p + \Upsilon$	10.639	—	10.743 [8]
$P_b^*(3/2)^+$ ($uudb\bar{b}$)	$\Delta + \Upsilon$	10.881	—	10.813 [8]
Θ_c^0 ($uudd\bar{c}$)	$p + D^-$	2.656	$3.099 \pm 0.003 \pm 0.005$	2.650 [6]
	$n + \bar{D}^0$	2.670	[13]	2.710 [3]
N_c^0 ($uuds\bar{c}$)	$p + D^0$	2.752	—	2.870[3]
Ξ_c^0 ($uuss\bar{c}$)	$\Sigma^+ + D^0$	2.905	—	3.135[3]
Θ_{cs}^{++} ($uudc\bar{s}$)	$p + \bar{D}^0$	2.751	—	2.427[5]
Θ_b^+ ($uudd\bar{b}$)	$p + B^0$	6.043	—	6.050 [3]
	$n + B^+$	6.067	—	
N_b^+ ($uuds\bar{b}$)	$p + B_s$	6.127	—	6.210 [3]
Ξ_b^+ ($uuss\bar{b}$)	$\Sigma^+ + B_s$	6.366	—	6.351 [3]
Θ_{bs}^+ ($uudb\bar{s}$)	$p + \bar{B}_s$	6.128	—	5.752 [5]

Table 1: Masses of cryptoexotic heavy pentaquarks

The spin hyperfine interaction can be expressed as,

$$E_{SD} = \frac{8\alpha_s}{9M_1M_2} \mathbf{S}_1 \cdot \mathbf{S}_2 |\Psi(0)|^2 \quad (2.6)$$

α_s is the strong interaction constant, \mathbf{S}_1 and \mathbf{S}_2 are the spins of the hadrons involved. $|\Psi(0)|^2$ is the di-hadronic wave function at the origin. With $\alpha_s = 0.59$ for light hadrons and 0.2 for heavier sector [12] E_{SD} has been estimated subsequently. $P_s^*(1/2)^+$ and $P_s^*(3/2)^+$ have been assumed to have configuration as proton- Φ state and $\Delta - \Phi$ state and $P_b^*(1/2)^+$ and $P_b^*(3/2)^+$ as proton- Υ state and $\Delta - \Upsilon$ state respectively. Based on this formulation several cryptoexotic heavy pentaquark masses such as Θ_c^0 , Θ_b^+ , Θ_{cs}^{++} , Θ_{bs}^+ , N_c^0 , N_b^+ , Ξ_c^0 , Ξ_b^+ have also been calculated and have been displayed in the Table I along with the comparison with others.

3. Conclusion

We have investigated several cryptoexotic pentaquark systems as hadronic composites of a meson and a baryon. Our results are in good agreement with the other theoretical works as well as experiments wherever available. Experimental mass of Θ_c^0 [13] is available but that is not confirmed

later [14]. We hope that this ambiguity could be understood in near future experiments. In this context it may be pertinent to state that in our previous work we have already estimated the masses of recently reported $P_c^*(4380)$ and $P_c^*(4450)$ as 4168 MeV and 4491 MeV describing these as $p + J/\Psi$ and $\Delta + J/\Psi$ di-hadronic states respectively [15]. We want to explore the properties of other exotic pentaquark families. Although the production rates are probably very low, these can be looked for at LHC, Fermilab, B-factories, RHIC etc. The predictions made in this work may inspire the future experiments.

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