

## Fragmentation contributions to hadroproduction of prompt $J/\psi$ , $\chi_{cJ}$ , and $\psi(2S)$ states

---

**Hee Sok Chung\***

*Theory Department, CERN, 1211 Geneva 23, Switzerland*

*E-mail: hee.sok.chung@cern.ch*

We compute leading-power fragmentation corrections to hadroproduction of charmonium states  $J/\psi$ ,  $\chi_{cJ}$ , and  $\psi(2S)$  in the nonrelativistic QCD factorization formalism. We include fragmentation functions through order  $\alpha_s^2$  and parton production cross sections through order  $\alpha_s^3$ . We also resum leading logarithms of the transverse momentum divided by the charm-quark mass to all orders in  $\alpha_s$ . We find that the fragmentation corrections have a significant impact on the hadroproduction cross section of charmonia. We obtain good fits to the hadroproduction cross sections measured at the Tevatron and the LHC. Using the long-distance matrix elements obtained from the fits, we make predictions for prompt  $J/\psi$  polarization that are in good agreement with the LHC data.

*38th International Conference on High Energy Physics  
3-10 August 2016  
Chicago, USA*

---

\*Speaker.

## 1. Introduction

In the nonrelativistic QCD (NRQCD) factorization conjecture, the inclusive production cross section of a charmonium  $H$  is written as [1]

$$d\sigma_{H+X} = \sum_n d\sigma_{c\bar{c}(n)+X} \langle \mathcal{O}^H(n) \rangle, \quad (1.1)$$

where  $c$  and  $\bar{c}$  are the charm quark and charm antiquark, respectively. The  $d\sigma_{c\bar{c}(n)+X}$  are perturbative short-distance cross sections (SDCSs), which are the cross sections to produce a  $c\bar{c}$  pair in a color and angular-momentum state  $n$ .  $\langle \mathcal{O}^H(n) \rangle$  are long-distance matrix elements (LDMEs) that describe the evolution of the state  $c\bar{c}(n)$  into  $H$ . The LDMEs scale with powers of  $v$ , where  $v$  is a typical velocity of a charm quark or a charm antiquark inside the  $H$  [1]. For charmonia,  $v^2 \approx 0.3$ .

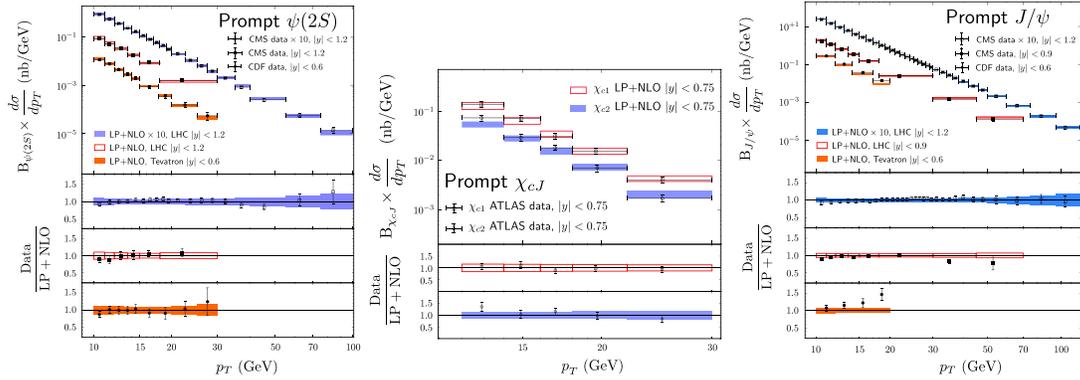
In current-level quarkonium phenomenology, the sum in Eq. (1.1) is truncated at relative order  $v^4$ . For  $H = J/\psi$  or  $\psi(2S)$ , Eq. (1.1) involves four LDMEs to relative order  $v^4$ :  $\langle \mathcal{O}^H(^3S_1^{[1]}) \rangle$ ,  $\langle \mathcal{O}^H(^3S_1^{[8]}) \rangle$ ,  $\langle \mathcal{O}^H(^1S_0^{[8]}) \rangle$ , and  $\langle \mathcal{O}^H(^3P_J^{[8]}) \rangle$ . Here, the color and angular-momentum states of the  $c\bar{c}$  pair are given in spectroscopic notation. For  $H = \chi_{cJ}$ ,  $\langle \mathcal{O}^{\chi_c}(^3P_J^{[1]}) \rangle$  and  $\langle \mathcal{O}^{\chi_c}(^3S_1^{[8]}) \rangle$  contribute in Eq. (1.1). The color-singlet LDMEs can be measured in lattice QCD, computed in potential models, or extracted from decay widths. The color-octet LDMEs are usually obtained by comparing Eq. (1.1) with data.

The SDCSs for the hadroproduction of  $J/\psi$  have been calculated to next-to-leading order (NLO) in  $\alpha_s$  [2, 3, 4, 5, 6, 7]. At large  $p_T$ , the SDCSs can be written as a convolution of parton production cross sections (PPCSs) and fragmentation functions (FFs) [8, 9]. In this work, we compute fragmentation contributions to hadroproduction of  $J/\psi$ ,  $\chi_{cJ}$  and  $\psi(2S)$  at leading power (LP) in  $m_c^2/p_T^2$ , where  $m_c$  is the charm quark mass and  $p_T$  is the transverse momentum of the charmonium. Our approach is based on Refs. [10, 11]. We refer the reader to Ref. [11] for the details of the calculation. We use the PPCSs through NLO in  $\alpha_s$  (order  $\alpha_s^3$ ), FFs through order  $\alpha_s^2$ , and resum the leading logarithms of  $p_T/m_c$  to all orders in  $\alpha_s$ . This goes beyond the fixed-order NLO calculations of the SDCSs. We find that the LP fragmentation contributions, combined with the NLO calculations, have a significant effect on the shapes of the SDCSs.

The remainder of this paper is organized as follows. We present our results in Section 2, followed by summary in Section 3.

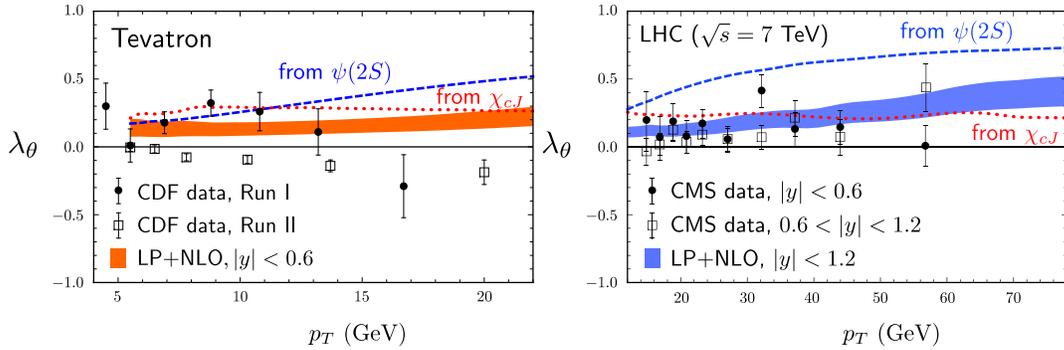
## 2. Results

By using the LP fragmentation contributions, combined with the NLO SDCSs, we obtain good fits to cross section data [12, 13, 14, 15]. The color-singlet LDMEs for the  $J/\psi$  and  $\psi(2S)$  were obtained from leptonic decay rates, and the color-singlet LDME for the  $\chi_{cJ}$  and the color-octet LDMEs were extracted from fits to data. The  $\chi^2/\text{d.o.f.}$  is 1.71/29, 1.19/8, and 8.20/40 for  $\psi(2S)$ ,  $\chi_{cJ}$ , and  $J/\psi$ , respectively. In the case of the  $J/\psi$ , the feeddown contribution from decays of  $\psi(2S)$  and  $\chi_{cJ}$  have been taken into account. In order to suppress non-factorizing contributions that may arise at small  $p_T$ , we use only the data with  $p_T$  larger than 3 times the charmonium mass in the fit. We use a correlation matrix method to compute uncertainties in the LDMEs. We refer the reader to Ref. [11] for details of the fit. In Fig. 1, we show the LP+NLO prediction for the cross section against the experimental data.



**Figure 1:** Prompt  $\psi(2S)$  and  $J/\psi$  cross section at the Tevatron and the LHC, and prompt  $\chi_{cJ}$  cross section at the LHC.

Using the LDMEs obtained from the fit, we make predictions for  $J/\psi$  polarization at the Tevatron and at the LHC. Our results are shown in Fig. 2 against CDF [16, 17] and CMS [18] data. Our results agree well with the CMS data [18], whereas the agreement is only fair with the CDF data [16, 17].



**Figure 2:** Polarization of prompt  $J/\psi$  at the Tevatron (left) and at the LHC (right). The polarization of  $J/\psi$ 's from decays of  $\psi(2S)$  and  $\chi_{cJ}$  states are shown with dashed and dotted lines, respectively.

### 3. Summary

We computed leading-power (LP) fragmentation corrections to hadroproduction cross sections of the charmonium states  $J/\psi$ ,  $\chi_{cJ}$  and  $\psi(2S)$ . We employed parton production cross sections at next-to-leading order (order  $\alpha_s^3$ ) and fragmentation functions at order  $\alpha_s^2$ , and resummed leading logarithms of  $p_T/m_c$  to all orders. The LP fragmentation contributions we computed goes beyond the fixed-order NLO calculations of the short-distance cross sections (SDCSs), and give part of the non-logarithmic contributions at next-to-next-to-leading order in  $\alpha_s$ . The LP fragmentation corrections we found have a significant effect on the shapes of the SDCSs.

We obtain good fits to measured charmonium cross sections at large  $p_T$ . By using the LDMEs obtained from the fit, we make predictions for the polarization of  $J/\psi$  at the Tevatron and at the LHC. Our predictions are in good agreement with the CMS data [18], and in fair agreement with the CDF data [16, 17].

Although we have obtained good agreement with the data at large  $p_T$ , a complete calculation of order  $\alpha_s^5$  may be necessary in order to reduce the theoretical uncertainties. Measurements of the cross sections and polarizations of charmonium states at much larger values of  $p_T$  at the LHC can provide good tests of the theoretical predictions.

## References

- [1] G. T. Bodwin, E. Braaten, and G. P. Lepage, Phys. Rev. D **51**, 1125 (1995); **55**, 5853(E) (1997) [hep-ph/9407339].
- [2] Y. Q. Ma, K. Wang, and K. T. Chao, Phys. Rev. D **84**, 114001 (2011) [arXiv:1012.1030 [hep-ph]].
- [3] M. Butenschoen and B. A. Kniehl, Phys. Rev. Lett. **106**, 022003 (2011) [arXiv:1009.5662 [hep-ph]].
- [4] Y. Q. Ma, K. Wang, and K. T. Chao, Phys. Rev. Lett. **106**, 042002 (2011) [arXiv:1009.3655 [hep-ph]].
- [5] K. T. Chao, Y. Q. Ma, H. S. Shao, K. Wang, and Y. J. Zhang, Phys. Rev. Lett. **108**, 242004 (2012) [arXiv:1201.2675 [hep-ph]].
- [6] M. Butenschoen and B. A. Kniehl, Phys. Rev. Lett. **108**, 172002 (2012) [arXiv:1201.1872 [hep-ph]].
- [7] B. Gong, L. P. Wan, J. X. Wang, and H. F. Zhang, Phys. Rev. Lett. **110**, 042002 (2013) [arXiv:1205.6682 [hep-ph]].
- [8] J. C. Collins and D. E. Soper, Nucl. Phys. B **194**, 445 (1982).
- [9] G. C. Nayak, J.-W. Qiu, and G. F. Sterman, Phys. Rev. D **72**, 114012 (2005). [hep-ph/0509021].
- [10] G. T. Bodwin, H. S. Chung, U-R. Kim, and J. Lee, Phys. Rev. Lett. **113**, 022001 (2014) [arXiv:1403.3612 [hep-ph]].
- [11] G. T. Bodwin, K. T. Chao, H. S. Chung, U. R. Kim, J. Lee and Y. Q. Ma, Phys. Rev. D **93**, no. 3, 034041 (2016) doi:10.1103/PhysRevD.93.034041 [arXiv:1509.07904 [hep-ph]].
- [12] D. Acosta *et al.* [CDF Collaboration], Phys. Rev. D **71**, 032001 (2005) [hep-ex/0412071].
- [13] S. Chatrchyan *et al.* [CMS Collaboration], JHEP **1202**, 011 (2012) [arXiv:1111.1557 [hep-ex]].
- [14] V. Khachatryan *et al.* [CMS Collaboration], Phys. Rev. Lett. **114**, 191802 (2015) [arXiv:1502.04155 [hep-ex]].
- [15] G. Aad *et al.* [ATLAS Collaboration], JHEP **1407**, 154 (2014) [arXiv:1404.7035 [hep-ex]].
- [16] T. Affolder *et al.* [CDF Collaboration], Phys. Rev. Lett. **85**, 2886 (2000) [hep-ex/0004027].
- [17] A. Abulencia *et al.* [CDF Collaboration], Phys. Rev. Lett. **99**, 132001 (2007) [arXiv:0704.0638 [hep-ex]].
- [18] S. Chatrchyan *et al.* [CMS Collaboration], Phys. Lett. B **727**, 381 (2013) [arXiv:1307.6070 [hep-ex]].