

ATLAS IBL operational experience

Yosuke Takubo on behalf of the ATLAS collaboration*

High Energy Accelerator Research Organization (KEK)

1-1 Oho Tsukuba Ibaraki, 305-0801, Japan

E-mail: yosuke.takubo@kek.jp

The Insertable B-Layer (IBL) is the inner most pixel layer in the ATLAS experiment, which was installed at a 3.3 cm radius from the beam axis in 2014 to improve the tracking performance. To cope with the high radiation and hit occupancy due to proximity to the interaction point, a new read-out chip and two different silicon sensor technologies (planar and 3D) have been developed for the IBL. After the long shut-down period over 2013 and 2014, the ATLAS experiment started data-taking in May 2015 for Run-2 of the Large Hadron Collider (LHC). The IBL has been operated successfully since the beginning of Run-2 and shows excellent performance with a low dead module fraction, high data-taking efficiency and improved tracking capability. The experience and challenges in operation of the IBL are described and results on the performance are presented.

The 25th International workshop on vertex detectors

September 26-30, 2016

La Biodola, Isola d' Elba, ITALY

*Speaker.

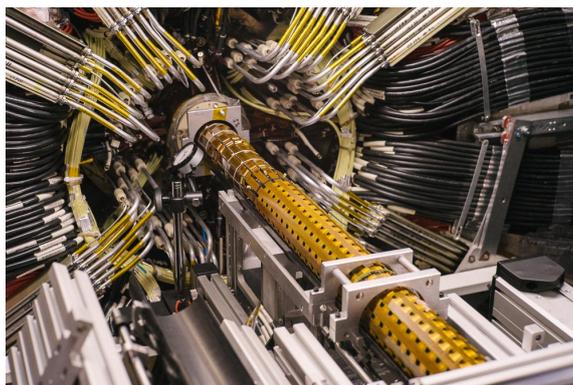


Figure 1: Picture of IBL insertion into the ATLAS detector.

1. Introduction

The Pixel Detector used in Run-1 is designed to work at the instantaneous luminosity of $1 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$. The instantaneous luminosity of the Large Hadron Collider (LHC) is expected to increase up to $3 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$ after long shutdown 2 (LS2) which is planned in 2019. As the luminosity increases the hit data rate becomes a serious issue for the front-end electronics. To cope with the increase during Run-2, a new pixel sensor layer, IBL [1], was installed as the inner most layer in 2014. This was the first upgrade project in the ATLAS detector [2].

The installation of the IBL was strongly motivated by physics analysis as well. For example, b -jet tagging is used in about 65% of the analysis results published by the ATLAS collaboration at the ICHEP conference in 2016 [3]. Providing additional hit information at the closest position to the beam collision point, the IBL significantly improves the performance of b -jet tagging.

After the long shut-down period over 2013 and 2014 (LS1), the ATLAS experiment started data-taking in May 2015 for Run-2 of the LHC. The IBL has been successfully operated and shows excellent performance. In the following sections, the experience and challenges in the operation of the IBL are described as well as its performance.

2. Insertable B-Layer (IBL)

In the IBL, two different sensor technologies are adopted: n^+ -in- n planar sensors [4] and 3D sensors [5]. The pixel size of the sensors is $50 \times 250 \mu\text{m}^2$ which is 60% of the pixel size used for the Pixel Detector ($50 \times 400 \mu\text{m}^2$). The IBL has about 12 million pixels in total. The thickness of the planar and 3D sensors is $200 \mu\text{m}$ and $230 \mu\text{m}$, respectively. In the planar sensor, the slim-edge of $200 \mu\text{m}$ with long pixels under the guard-ring is adopted to minimize inactive region. The sensors were designed to work at a fluence of up to $5 \times 10^{15} \text{ 1 MeV n}_{\text{eq}}$, which is the expected fluence by the end of LHC Phase-1 operation around 2022.

A new front-end chip (FE-I4B) was developed for the IBL. The FE-I4B is produced with IBM 130 nm technology, and 80×336 pixels are embedded per chip. It is designed to be radiation hard up to a 250 Mrad ionizing dose. The signals from the sensor are amplified and digitized in analog and digital circuits respectively, and finally the charge information is output with 4-bit value of the time over threshold (TOT).

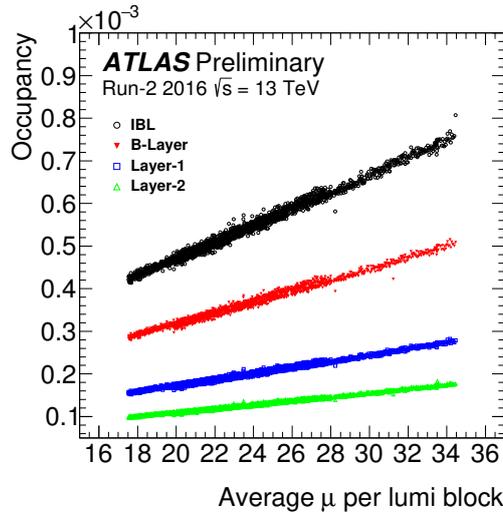


Figure 2: Hit occupancy at each pixel layer as a function of the number of pile-up events (μ). IBL, B-Layer, Layer-1 and Layer-2 are placed at the radius of 33, 50.5, 88.5 and 122.5 mm, respectively, centered around the beam axis [6].

29 An IBL module consists of a sensor bump-bonded to an FE-I4B chip with flex cables which
 30 provide electric lines. The modules are mounted on the staves where planar sensors are used for
 31 75% of the middle region in η and 3D sensors used in the remaining 25% (high η). The full
 32 IBL consists of 14 staves integrated into a support tube. The IBL surrounds a new beam pipe,
 33 and is contained in a package of 7 m long in the limited radial envelope of 10 mm. The IBL was
 34 transported to the ATLAS experimental hall 90 m underground on May 5, 2014 and installed into
 35 the ATLAS detector on May 7 as shown in Fig. 1. To install the IBL in the limited space inside
 36 B-Layer, an insertion clearance of less than 0.1 mm was available.

37 3. Performance of IBL in ATLAS Run-2 operation

38 The ATLAS experiment started data-taking in May 2015 for Run-2 of the LHC. The IBL
 39 has been operating successfully since the beginning of Run-2 and shows excellent performance
 40 with low dead module fraction, high data-taking efficiency and improved tracking capability as
 41 summarized in this section.

42 3.1 Detector stability in data-taking

43 One of the big achievements in the IBL operation is that 100% of the sensors and FE-I4B chips
 44 are kept operational in the tracking coverage region of $|\eta| < 2.5$ during Run-2. It allows the IBL
 45 to show full tracking performance.

46 Figure 2 [6] shows the hit occupancy at each pixel layer as a function of the number of pile-up
 47 events (μ). The hit occupancy at the IBL increases linearly with μ , and has been successful even
 48 at high μ . Actually, the rate of desynchronization in the readout is of the order of 0.1% at the IBL
 49 and this value is about one order of magnitude less than those of the other Pixel layers.

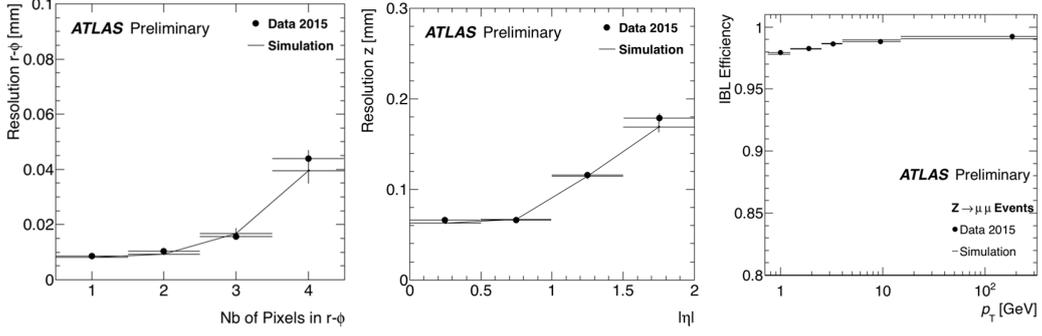


Figure 3: The hit position resolution on the IBL in the $r-\phi$ (left) and z (middle) projections as a function of cluster width along ϕ and the track $|\eta|$, respectively, and efficiency of associating an IBL hit to a reconstructed particle track as a function of the particle p_T (right) [7]. The selected $Z \rightarrow \mu\mu$ events in collision data (circles with error bars) and simulation (lines) are used for these plots.

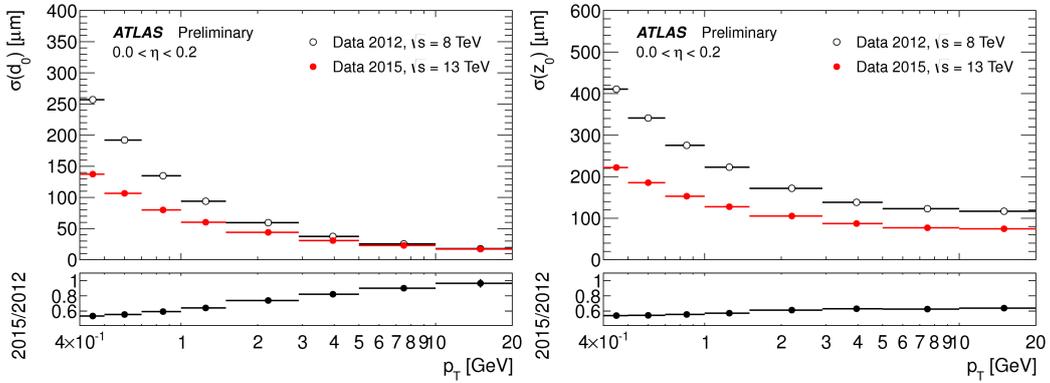


Figure 4: d_0 (left) and z_0 resolution (right) as a function of p_T in ATLAS Run-1 (without IBL) and Run-2 (with IBL) [8].

50 3.2 Position resolution and hit efficiency

51 Figure 3 [7] shows the hit position resolution on the IBL in the $r-\phi$ (left) and z (middle)
 52 projections as a function of cluster width along ϕ and the track $|\eta|$, respectively. In Fig. 3 (left),
 53 the clusters that are broader than two pixels along the $r-\phi$ projection typically include secondary
 54 ionization and thus exhibit a degraded hit spatial resolution along that coordinate, mostly due to
 55 the effect of δ -rays, as predicted by simulation. The evolution of the measured resolution along the
 56 longitudinal coordinate in Fig. 3 (middle) is interpreted as a convolution of the change in resolution
 57 with the increasing number of pixels used for charge interpolation and the increase of the multiple
 58 scattering effects.

59 Figure 3 (right) [7] shows efficiency of associating an IBL hit to a reconstructed particle track
 60 as a function of the particle p_T . The efficiency is computed as the fraction of selected tracks with
 61 $p_T > 0.7$ GeV and $|\eta| < 2.5$ that have an associated IBL cluster. The efficiency is better than 98%
 62 and is measured as expected by the MC simulation.

63 The IBL provides the hit information at the closest position to the beam collision point, and
 64 it improves the impact parameter resolution. Figure 4 [8] shows d_0 and z_0 resolution in Run-1

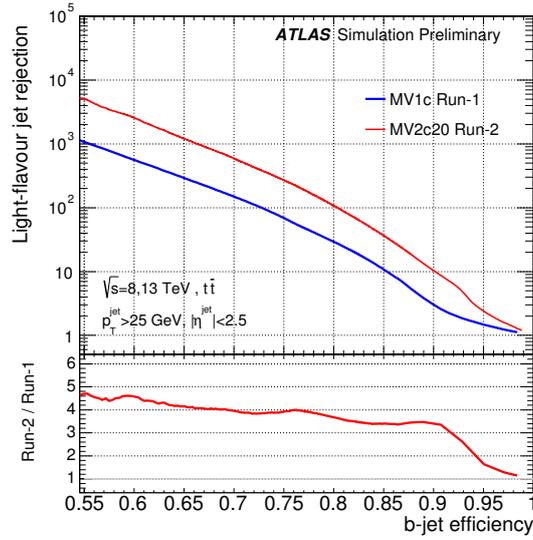


Figure 5: The light jet rejection versus b -jet tagging efficiency for the MV1c b -tagging algorithm using the Run-1 detector and reconstruction software (blue) compared to the MV2c20 b -tagging algorithm using the Run-2 setup (red) [9].

65 and Run-2. The impact parameter resolutions are better in Run-2 with the IBL. At low transverse
 66 momentum (p_T) region, and it is improved by about 40% at p_T below 1 GeV for both d_0 and z_0 .

67 3.3 b -jet tagging performance

68 The better impact parameter resolution with the IBL improves performance of b -jet tagging.
 69 Figure 5 shows simulation results for two software algorithms: MV1c and MV2c20. The light jet
 70 rejection as a function of b -jet tagging efficiency for MV1c b -tagging algorithm using the Run-1
 71 detector and reconstruction software (blue) compared to the MV2c20 b -tagging algorithm using
 72 the Run-2 setup [9]. The light jet rejection shows better performance with the IBL together with
 73 improvements in the algorithm of b -jet tagging.

74 3.4 dE/dx measurement

75 dE/dx information is important especially for physics analysis to find new heavy particles
 76 predicted in physics beyond the Standard Model [10]. The IBL gives another measurement of
 77 charge deposition in the pixel sensors which improves the dE/dx measurement. Figure 6 [11]
 78 shows a distribution of dE/dx obtained with (black dots) and without the IBL (red triangles) (left)
 79 and bi-dimensional distribution of dE/dx and momentum (right). It can be seen that the distribution
 80 of dE/dx in Fig. 6 (left) becomes narrower with the IBL. In addition, the particle types are clearly
 81 separated in Fig. 6 (right), and this information is used for particle identification in physics analyses.

82 4. Challenges in IBL operation

83 The IBL has faced several challenges like wire-bond oscillation, distortion of the staves with
 84 changing temperature and low-voltage (LV) current increase due to total ionizing dose (TID) effect.

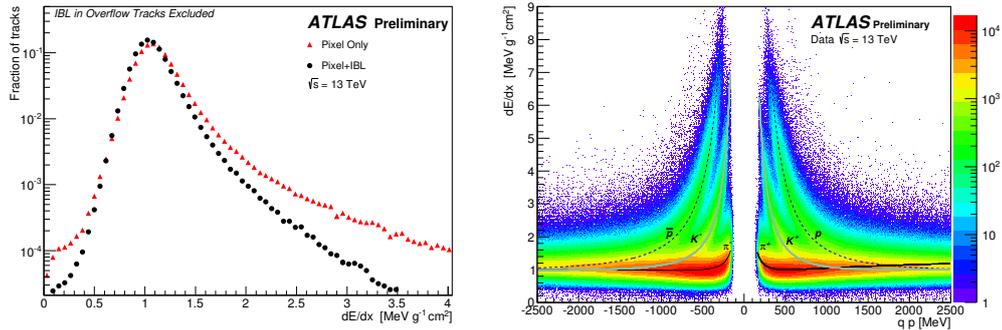


Figure 6: Distribution of dE/dx obtained with (black dots) and without the IBL (red triangles) (left) and bi-dimensional distribution of dE/dx and momentum (right) [11]. In the right plot, the distributions of the most probable value for the fitted probability density functions of pions, kaons and protons are superimposed.

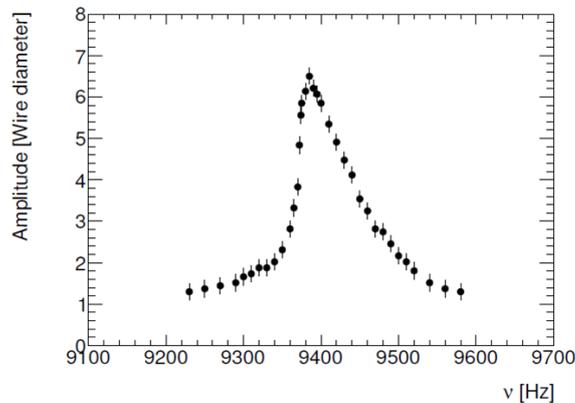


Figure 7: Oscillation amplitude expressed in wire diameter ($25 \mu\text{m}$) as a function of the frequency for a wire of 2.8 mm length obtained under IBL-like conditions [12].

85 Coping with these challenges, the detector has been successfully operated in stable condition during
 86 Run-2 in 2015 and 2016. In this section, the challenges faced in the operation and the treatment for
 87 them are described.

88 4.1 Wire bond oscillation

89 Wire bonding is used to connect the FE-I4B chip to the flex cables on the IBL module. The
 90 digital circuit of the FE-I4B chip consumes current which is susceptible to consecutive triggers or
 91 calibration scans [12]. Accordingly, the wire bonds may oscillate at relatively high frequency in
 92 the magnetic field of 2 T where the IBL is located. Figure 7 shows the oscillation amplitude as a
 93 function of the frequency obtained experimentally for a 2.8 mm wire under IBL-like conditions.
 94 Depending on the oscillation amplitude and number of cycles, micro-cracks can develop at the wire
 95 bond heal and, even more, the wire can cross the material elastic limit. Both aspects may lead to
 96 possible failures on the bond.

97 To mitigate this effect, the trigger cycle sent to the FE-I4B chip was manipulated to avoid
 98 the resonance frequency of the wire bonds. In operation of the IBL, a fixed frequency trigger

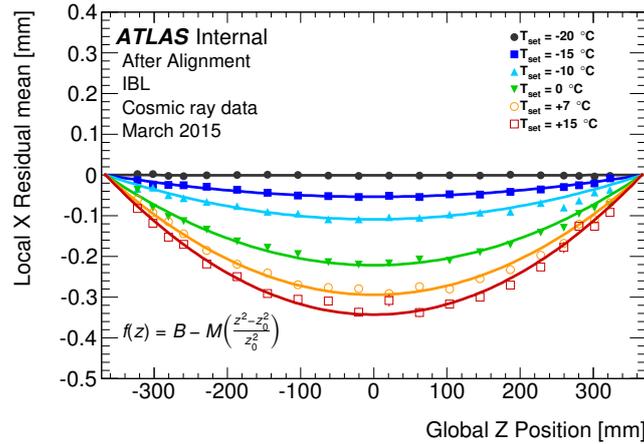


Figure 8: The track-to-hit residual mean in the local x direction [13]. The residual mean is averaged over all hits of modules at the same global z position. The alignment corrections derived at -20 °C are applied to the local positions in the module frames.

99 veto (FFTV) was newly implemented in data-taking system to limit the number of triggers in the
100 resonance region.

101 4.2 Mechanical distortion of IBL staves

102 During the commissioning of the IBL, distortion of the staves was found [13]. The magnitude
103 of the distortion depends on the operating temperature as shown in Fig. 8, and it is caused by
104 a mismatch between the coefficients of thermal expansion (CTE) of a bare stave made with the
105 carbon foam and the flex cable attached on the bare stave. The maximum displacement due to the
106 distortion is more than $300 \mu\text{m}$ at the planned operating temperature of -20 °C with respect to the
107 nominal position at the room temperature.

108 In the operation of the IBL, the impact of this effect is mitigated by temperature control at the
109 level of 0.2 K and the regular alignment correction in the offline reconstruction.

110 4.3 Effect of total ionizing dose on FE-I4B

111 The LV current consumption in FE-I4B chips increased in 2015. This was caused by the
112 effect of TID on NMOS transistors in the chip. Significant leakage current was induced by positive
113 charges trapped in the bulk of the shallow trench isolation (STI) in NMOS transistor [14].

114 The increase of the leakage current causes several issues in the detector operation: In the worst
115 case, and if only one power wire bond is left, the IBL module approaches the maximum LV current
116 of 0.55 A. In 2015 the current started to approach this limit before mitigating actions were taken. In
117 addition, it was observed that this phenomenon causes detuning of ToT and threshold values [15].
118 The power consumption in a FE-I4B increases with more LV current, and it changes temperature
119 on the module. As described in Sec. 4.2, the different temperature on the modules causes distortion
120 of the staves accordingly.

121 To study this effect as a function of TID the measurement of LV current in a FE-I4B chip was
122 performed after irradiating it with an x-ray source and proton beam [16, 17]. The results showed

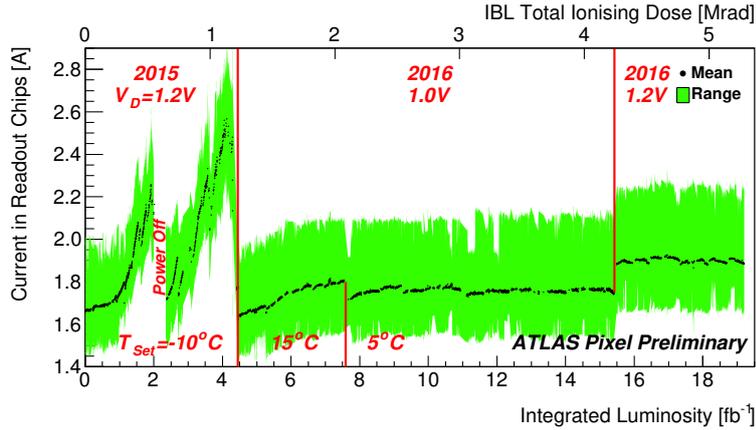


Figure 9: Mean LV current of four FE-I4B chips in one LV channel group against integrated luminosity and TID during Run-2 [18].

123 that LV current increases until about 1 Mrad where it begins to decrease again[16]. This can be
 124 explained by the interface traps filling with electrons in front of STI after the positive charge has
 125 been trapped there. This layer of negative charge compensates for the effect of trapped holes in the
 126 STI. In addition, it was found that the LV current can be decreased by operating the chip at higher
 127 temperature due to the annealing effect [16].

128 Taking into account the results of these measurements, the operating temperature of the IBL
 129 was increased to 15 °C at the beginning of 2016 from -10 °C in 2015 and then eventually decreased
 130 again to 5 °C. In addition, the digital supply-voltage was decreased to 1.0 V from 1.2 V until the
 131 TID was more than 4 Mrad. Figure 9 [18] shows the mean LV current of four FE-I4B chips in one
 132 LV channel group against integrated luminosity and TID. Although there were significant increases
 133 in LV current during 2015, it became stabilized and contained in 2016.

134 5. Summary and conclusions

135 The IBL is the inner most pixel layer in the ATLAS experiment, which was installed at a 3.3
 136 cm radius from the beam axis in 2014 to improve the tracking performance. After LS1, the ATLAS
 137 experiment started data-taking in May 2015 for Run-2 of the LHC. The IBL has been operated
 138 successfully since the beginning of Run-2 and shows excellent performance with the low dead
 139 module fraction, high data-taking efficiency and improved tracking capability.

140 The additional hit information provided by the IBL at the closest position to the beam collision
 141 point provides invaluable information. It significantly improves the tracking performance, b -jet
 142 tagging and dE/dx measurements for example.

143 Operating the IBL has faced several challenges like wire-bond oscillation, distortion of the
 144 staves with changing temperature and LV current increase due to TID effect. Coping with these
 145 challenges, the detector has been successfully operated in stable conditions during Run-2 in 2015
 146 and 2016.

147 **References**

- 148 [1] The ATLAS Collaboration, *ATLAS Insertable B-Layer Technical Design Report*,
149 CERN-HLCC-2010-013, ATLAS TDR **19**, 15 September 2010.
- 150 [2] ATLAS Collaboration, *The ATLAS Experiment at the CERN Large Hadron Collider*, JINST **3**
151 S08003, 2008.
- 152 [3] 38th INTERNATIONAL CONFERENCE ON HIGH ENERGY PHYSICS (ICHEP 2016),
153 <https://www.ichep2016.org>.
- 154 [4] C. Goessling et al., *Planar n^+ -in-n silicon pixel sensors for the ATLAS IBL upgrade*, Nucl. Instrum.
155 Meth. A **650** (2011), 198-201.
- 156 [5] Cinzia Da Via et al., *3D silicon sensors: Design, large area production and quality assurance for the*
157 *ATLAS IBL pixel detector upgrade*, Nucl. Instrum. Meth. A **694** (2012), 321-330.
- 158 [6] ATLAS Collaboration, Public Pixel Tracker Plots for Collision Data, PIX-2016-007,
159 <http://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PLOTS/PIX-2016-007>.
- 160 [7] The ATLAS Collaboration, *Efficiency and Hit Spatial Resolution of ATLAS IBL Sensors in LHC Run*
161 *2 Collision Events*, ATL-INDET-PUB-2016-001, <http://cds.cern.ch/record/2203893>.
- 162 [8] The ATLAS Collaboration, Approved Plots of the Tracking Combined Performance Group,
163 INTR-2015-007, <http://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PLOTS/IDTR-2015-007>.
- 164 [9] The ATLAS Collaboration, *Expected performance of the ATLAS b-tagging algorithms in Run-2*,
165 ATL-PHYS-PUB-2015-022, <http://cds.cern.ch/record/2037697>.
- 166 [10] The ATLAS Collaboration, *Search for metastable heavy charged particles with large ionization*
167 *energy loss in pp collisions at $\sqrt{s} = 13$ TeV using the ATLAS experiment*, Phys. Rev. D **93**, 112015
168 (2016).
- 169 [11] The ATLAS Collaboration, Public Pixel Tracker Plots for Collision Data, PIX-2015-002,
170 <http://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PLOTS/PIX-2015-002>.
- 171 [12] D. A. Feito, A. Honma, B. Mandelli, *Studies of IBL wire bonds operation in a ATLAS-like magnetic*
172 *field*, PH-EP-Tech-Note-2015-002, <http://cds.cern.ch/record/2010249>.
- 173 [13] The ATLAS Collaboration, *Study of the mechanical stability of the ATLAS Insertable B-Layer*,
174 ATL-INDET-PUB-2015-001, <http://cds.cern.ch/record/2022587>.
- 175 [14] L. Gonella et al., *Total Ionizing Dose effects in 130-nm commercial CMOS technologies for HEP*
176 *experiments*, Nucl. Instrum. Meth. A **582** (2007), 750-654.
- 177 [15] The ATLAS Collaboration, Cosmics and Calibration Approved Pixel Plots, PIX-2016-005,
178 <http://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PLOTS/PIX-2016-005>.
- 179 [16] The ATLAS Collaboration, Cosmics and Calibration Approved Pixel Plots, PIX-2015-008,
180 <http://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PLOTS/PIX-2015-008>.
- 181 [17] The ATLAS Collaboration, Cosmics and Calibration Approved Pixel Plots, PIX-2016-010,
182 <http://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PLOTS/PIX-2016-010>.
- 183 [18] The ATLAS Collaboration, Cosmics and Calibration Approved Pixel Plots, PIX-2016-006,
184 <http://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PLOTS/PIX-2016-006>.