

Charm production at HERA, proton structure, the charm mass, and Higgs Yukawa couplings

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Charm production at HERA (H1 and ZEUS) and its embedding into the Standard Model in the context of results from other experiments are discussed. The latest results on charm production at HERA (D^* and inclusive cross sections, fragmentation fractions, ψ' vs. J/ψ) and their impact on measurements of the proton structure are summarised. Related measurements of the charm quark mass and a determination of its QCD running are presented. Finally, measurements of the charm, beauty and top quark masses from HERA, LEP and LHC are reinterpreted in terms of the corresponding Higgs Yukawa couplings. A graphical representation of the running of the Yukawa couplings derived from measurements is obtained for the first time.

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1. Introduction

The charm quark is one of the fundamental ingredients of the Standard Model, and directly appears in the Standard Model Lagrangian. Measurements of charm quark production and its properties therefore have a strong impact on our understanding of the Standard Model. This contribution focuses on recent measurements, in particular of charm quark production at the HERA collider by the H1 and ZEUS experiments, and on the relation of these measurements to other recent measurements and their interpretation. Measurements of the charm quark mass and its running in the context of QCD are presented and discussed. These measurements, together with corresponding measurements of the beauty quark mass at LEP and HERA and top quark masses at LHC, are then reinterpreted in the context of the full Standard Model, in which masses are represented by the corresponding more fundamental Higgs Yukawa couplings. The status of the experimental evidence for the running of these couplings is worked out and graphically represented for the first time, including a comparison to the Yukawa couplings directly measured at LHC. Figures are included in this writeup only for results which are not yet finally published.

2. Recent measurements of charm production at HERA

A review of all measurements related to open charm production at HERA up to 2015 can be found in [1], with some extensions in [2]. A review of charmonium production can be found in [3], and of diffractive charm production in [4]. Here we will focus on the most recently published HERA charm measurements.

The final cross sections for differential D^* production in Deep Inelastic Scattering (DIS) from H1 and ZEUS were combined [5] and compared to next-to-leading order (NLO) QCD predictions. Differential cross sections in quantities like transverse momentum (p_T) , pseudorapidity (η) or fractional energy flow (z) were obtained with ultimate HERA precision. Their advantage with respect to cross sections depending only on inclusive quantities [6] like photon virtuality (Q^2) and inelasticity (y) is the sensitivity to kinematic details of charm production not accessible otherwise. Their disadvantage is the larger uncertainty of the corresponding NLO theory predictions, typically 3-5 times larger than the one of the data. It is demonstrated that within these large uncertainties it is possible to find a fixed set of theory parameters with which all data distributions can be simultaneously described using experimental uncertainties only. The conclusion [5] is that the data are described very reasonably, but that more work is needed on the theory side (e.g. NNLO, improved fragmentation) to reduce the uncertainties to a level competitive with the data.

Measurements of the fragmentation fractions of charm quarks to different charm hadron final states at HERA have been compared to and combined [7] with measurements from LEP/SLC, B factories and LHC. The HERA measurements, especially those in photoproduction [8], are very competitive with the other measurements, and substantially contribute to the world average. The good agreement between the measurements from different types of collisions in a large range of energies confirms the universality of fragmentation.

In a ZEUS measurement [9] of the ratio of the exclusive production of ψ' and J/ψ mesons in DIS, $ep \to ep\psi$, many experimental and systematic uncertainties cancel. The measurement of the dependence of this ratio on Q^2 is in good agreement with earlier measurements by H1 [10], with

significantly improved precision. It is also consistent with an H1 measurement in photoproduction [11]. The observed rise of this ratio with Q^2 is qualitatively consistent with theoretical models based on different assumptions about the ψ wave functions, and starts to quantitatively discriminate between them [9].

3. Charm and proton structure

Various fits of the proton parton density functions (PDF) at both NLO and NNLO were obtained by H1 and ZEUS from the recent final combination of the H1 and ZEUS inclusive DIS data [12]. All these HERA-data-only fits use values of the charm mass constrained by the combined inclusive HERA charm data [6]. One of the fit variants also explicitly includes these charm data (at NLO), as well as HERA jet data. Good consistency [12] is observed between all these fits, and with various fits obtained by global PDF fitting groups. The inclusive HERA charm data [6] contribute to the reduction of the PDF uncertainties, mainly by constraining the charm quark mass. They are a standard ingredient of all recent global PDF fits [13, 14, 15, 16, 17].

The same inclusive HERA charm data have also been used [18] in conjunction with LHCb charm data [19] and corresponding beauty data [20, 21] to constrain the NLO gluon PDF at very low x. In particular the LHCb charm data explicitly constrain the gluon distribution down to fractional momenta of order few 10^{-6} , extending the low x reach by more than an order of magnitude with good precision. Other analyses [22, 23] have subsequently obtained similar results. These results are particularly relevant in the context of predictions [22, 24, 25] of the prompt contribution to the cosmic ray neutrino spectrum, as measured e.g. by Ice Cube [26].

4. Measurements of the charm quark mass and its running

As mentioned earlier, the inclusive (or differential) HERA charm data in DIS allow one to constrain the charm quark mass and hence to measure its value within different QCD factorization schemes. The cleanest (in the sense of no extra parameters in the QCD matrix elements beyond those appearing in the QCD Lagrangian) such measurements are obtained [27, 6, 28, 29] in the fixed flavour number scheme, and using the so-called \overline{MS} mass $m_c(m_c)$ (Table 1). For the latter, the cross section has better perturbative convergence behaviour [30] than using the pole mass. Measurements have been obtained both at NLO and approximate NNLO. Similar measurements using the S-ACOT [31] and FONLL C [29] variable flavour number schemes, both of which have one extra parameter in the merging of the massless and massive parts of the calculation, have also been obtained (Table 1). The latter treats the light flavour contributions at NNLO, and the massive part of the heavy flavour contributions at NLO. Although all use the same HERA charm data [6], the measurements partially differ in the choice of input data sets used for the PDF part of the fit, since charm mass and PDFs are fitted simultaneously. The charm mass can also be fitted simultaneously with the strong coupling constant α_s [32]. Preliminary even more precise results are being obtained by adding more data [2, 33] and/or improving the NNLO approximation [34].

All these measurements are consistent with each other, and consistent with the world average [35] 1.27 ± 0.03 GeV, obtained mainly from lattice calculations and e^+e^- data with time-like kine-

scheme	$m_c(m_c)$ [GeV]
FONLL [29]	$1.335 \pm 0.043 (\exp)^{+0.019}_{-0.000} (param)^{+0.011}_{-0.008} (mod)^{+0.033}_{-0.008} (th)$
FFN [29]	$1.318 \pm 0.054 (\exp)_{-0.010}^{+0.011} (param)_{-0.019}^{+0.015} (mod)_{-0.004}^{+0.045} (th)$
FFN (HERA) [6]	$1.26 \pm 0.05 (\exp) \pm 0.03 (\text{mod}) \pm 0.02 (\text{param}) \pm 0.02 (\alpha_s)$
FFN (Alekhin et al.) [28]	$1.24 \pm 0.03 (\exp)^{+0.03}_{-0.02} (\text{scale})^{+0.00}_{-0.07} (\text{th}) (\text{approx. NNLO})$
	$1.15 \pm 0.04 (\exp)^{+0.04}_{-0.00} (\text{scale}) \text{ (NLO)}$
FFN (Alekhin et al.) [34]	$1.252 \pm 0.018 (exp) \pm 0.025 (scale) \pm 0.010 (th)$ (approx. NNLO)
S-ACOT-χ (CT10) [31]	$1.12^{+0.05}_{-0.11}$ (strategy 1)
	$1.18^{+0.05}_{-0.11}$ (strategy 2)
	$1.19_{-0.15}^{+0.06}$ (strategy 3)
	$1.24_{-0.15}^{+0.06}$ (strategy 4)
World average [35]	1.27 ± 0.03

Table 1: List of recent determinations of $m_c(m_c)$ from fits to DIS data adjusted from Ref. [29]. The updated 2016 PDG world average value is also reported for reference.

matics. This demonstrates the consistency of the different QCD approaches, and the usefulness of the HERA charm data to perform such measurements.

One of these measurements [6] has been refined further in order to extract information about the charm mass running. In QCD, the \overline{MS} running mass varies (runs) with scale μ , conceptually similar to the running of α_s . All the measurements above were expressed in terms of the charm quark mass $m_c(\mu)$ at the reference scale $\mu = m_c$ itself, as is customary, just like α_s is often expressed at the scale of the Z boson mass $\alpha_s(M_Z)$, by convention. The running of α_s is usually studied by separately evaluating it for subprocesses measured at different scales μ (e.g. e^+e^- center of mass or jet p_T). Here, the charm data set [6] is subdivided into several subsets, centered on different scales $\mu = \sqrt{Q^2 + m_c^2}$. Each of these subsets is then fitted separately, using PDFs refitted from inclusive data only with the appropriate charm mass setting in a way that minimises the very small residual dependence of the charm mass fit on charm effects in the PDF fit [36, 37]. This results in the graphical representation [36, 1] of the charm mass running shown in Fig. 1. The running is consistent with QCD expectations.

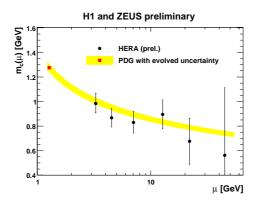


Figure 1: Measured charm mass $m_c(\mu)$ in the \overline{MS} running mass scheme as a function of the scale μ as defined in the text (black points). The red point at scale m_c is the PDG world average [38] and the band is its expected one-loop running [39].

5. Running of Higgs Yukawa couplings

As is well known, in pure QCD, a quark mass term can be introduced as a fundamental contribution to the Lagrangian. In the Standard Model, due to left-right symmetry breaking, such a term violates electroweak gauge invariance. Therefore it has to be 'mimicked' by replacing it with a term coupling to the Higgs field, with a nonzero vacuum expectation value v = 246.22 GeV. The 'constant' part of this coupling is then equivalent to a mass term with

$$m_q := \frac{1}{\sqrt{2}} y_q v \tag{1}$$

where y_q is the Higgs Yukawa coupling for the quarks q = c, b, t which we will consider here. The goal is the extraction of a graphical representation of the running of these Yukawa couplings as a function of scale, based on measurements.

The details of how the masses and couplings evolve with scale are scheme dependent. Here, we use the \overline{MS} scheme for QCD at NLO as in the previous section, and a Higgs/electroweak scheme in which Eq. (1) is taken to be universally valid. In the following, all relevant Higgs and electroweak effects are treated to leading (0-loop) order. Thus, in this scheme, the running of the Yukawa couplings is equivalent to the running of the \overline{MS} mass in QCD.

The direct measurements of the Yukawa couplings for top and beauty from Higgs production and decay at the LHC [40] and at the Tevatron [41] are assumed to be obtained at scale m_H . The QCD \overline{MS} masses which we will use here are all evaluated from real quark pair production cross sections and are assumed to be obtained at scale $\sqrt{Q^2 + 4m_c^2}$ for charm at HERA [6, 36, 37], at scale $\sqrt{Q^2 + 4m_b^2} \simeq 2m_b$ (cross section dominated by low Q^2) for beauty at HERA [20, 37], at scale m_Z for beauty at LEP/SLC [42], and at scale $m_{t\bar{t}}$ for top at LHC [43].

For charm and beauty, Eq. (1), reformulated as $y_q = \sqrt{2} \frac{m_q}{v}$, can be directly applied to the measurements of $m_q(\mu)$ cited above. For q = c, they are exactly the ones explained in the previous section. The measurements for q = b are treated similarly.

A theoretically well-defined experimental scale-dependent evaluation of the running of the \overline{MS} top quark mass can not yet be performed from existing published data. Here, we will apply a preliminary approximate procedure to nevertheless extract such an evaluation, which is on much less solid grounds than those for charm and beauty. It should be replaced by a better-defined procedure as soon as corresponding data become available.

We start from a scale depedendent evaluation [43] of the so-called 'MC mass' (here technically the 'hybrid mass') as a function of the $t\bar{t}$ invariant mass $m_{t\bar{t}}$ performed by CMS. Note that the MC mass is mainly extracted from kinematical quantities in the data. The measurements at different $m_{t\bar{t}}$ are statistically (but not necessarily systematically) independent. This was originally used to demonstrate that, within uncertainties and as expected, the MC masses measured at different scales with different kinematics are consistent with a single value [43].

The top pole mass has been extracted by both ATLAS [44] and CMS [45, 46] as a single value from top cross section measurements, for which the data sets partially overlap with those from which the MC mass is extracted. We note that the top pole and MC masses evaluated from LHC data are numerically almost identical, as to be expected [33]. The theoretical uncertainties of this conversion are poorly known, but can be estimated to be of order 1 GeV or less [38, 33].

The conversion factor from MC mass to pole mass is therefore very close to 1, with an estimated uncertainty of less than 1%.

A comparison of the measured top cross section as a function of $m_{t\bar{t}}$ [47] with theoretical predictions at leading order plus parton shower, which use the MC mass, and at NLO, which use the pole mass, shows that all three are reasonably consistent. They are also consistent [48] between ATLAS and CMS and with NNLO predictions. This raises the hope (but does not prove) that, if evaluated as a function of $m_{t\bar{t}}$, the conversion factor from MC mass to pole mass would experimentally turn out to be constant within uncertainties. This is not completely obvious, since the MC mass is mainly evaluated from kinematical quantities within each event, while the pole mass is evaluated from the derived cross sections. Here, we will simply make the assumption that the ratio is constant, and neglect the related (currently unknown) additional uncertainties, which is the weakest point of the procedure.

In a first step, we use this constant conversion factor to translate the $m_{t\bar{t}}$ -scale-dependent MC-mass evaluation [43] into a corresponding scale-dependent evaluation of the top pole mass. In the much better defined second step, we translate [39] the top pole mass (semi)independently obtained at each scale $m_{t\bar{t}}$ into the corresponding \overline{MS} running mass $m_t(\mu)$ with $\mu = m_{t\bar{t}}$. Somewhat arbitrarily, we use the two-loop formula. In the final step we again use Eq. (1) to translate each mass into a Yukawa coupling.

The result of this preliminary procedure is shown in Fig. 2, together with various evaluations of α_s at different scales. The blue bands showing the running expected from Standard Model theory are obtained by using the PDG reference values [38] for each of the parameters with their uncertainties, and evolving them using the three-loop QCD formula for α_s , the one-loop QCD formula for the charm and Higgs Yukawa couplings, and the two-loop formula for top, consistent with the way the points have been extracted from the data. No Higgs or electroweak loop corrections are applied. The Yukawa couplings from direct measurements [40, 41] are also shown in the figure, appropriately evaluated at the scale m_H .

Although not yet quite perfect (see some of the caveats explained above) this is, to the best knowledge of the author, the first graphical representation of the running of Higgs Yukawa couplings extracted from data. It is shown together with the well-known running of α_s extracted from τ decays [38], Upsilon decays [49], PETRA [50], LEP [51], HERA [52], S $p\bar{p}$ S [53], Tevatron [54] and LHC [45, 55] data to illustrate the analogy. It will be noted that data and Standard Model expectations are consistent, and that the couplings extracted indirectly from the mass measurements are consistent with those extracted from direct measurements of Higgs production. This nicely illustrates the interplay between QCD and the Higgs sector in the Standard Model.

6. Conclusions

The most recent measurements of charm production at HERA have been presented and put into the context of other measurements. With these examples, the author hopes to have illustrated that, if the interpretation of such measurements is pursued all the way down to the basics of the Standard Model, heavy quark physics (including charm physics) is almost always also both QCD and Higgs physics, since the quark masses almost always play a role. So far, the Higgs Yukawa

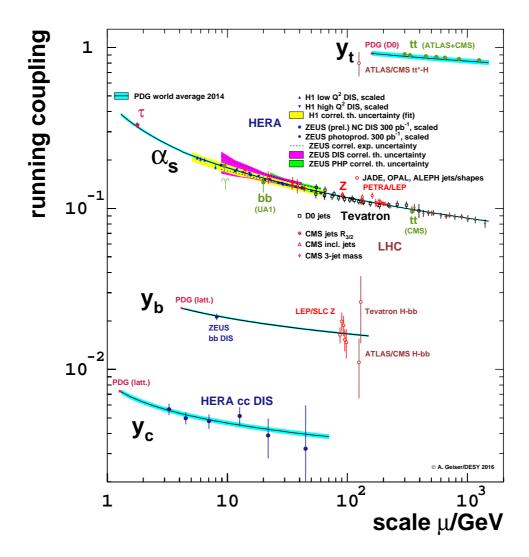


Figure 2: The running of the Higgs Yukawa couplings for charm (y_c) , beauty (y_b) and top (y_t) as a function of energy/momentum scale μ , compared to the running of QCD coupling constant α_s . For references, details of their evaluation and the scheme definitions see text. The brown points at scale m_H are directly extracted Higgs couplings, while all other Yukawa coupling points are obtained indrectly from mass measurements. The red stars are the PDG world averages, and the blue bands represent their expected evolution.

couplings and their running as obtained from measurements of the quark masses are consistent with the Yukawa couplings obtained from direct measurements of processes involving Higgs bosons.

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