

## Current challenges and future prospects for $\gamma$ from $B \rightarrow Dhh'$ decays

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Decays of the type  $B \rightarrow Dhh'$ , where a  $b$  hadron decays to a neutral charm meson that can be an admixture of  $D^0$  and  $\bar{D}^0$  states together with two light particles that are typically a kaon and a pion, have demonstrated potential to enable precise determinations of the angle  $\gamma$  of the CKM Unitarity Triangle. The current status and future prospects of these measurements are reviewed.

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The angle  $\gamma$  of the CKM Unitarity Triangle is a key parameter of quark flavour physics. It is the only  $CP$ -violating parameter that can be measured using only tree-level decays, and as such is a benchmark Standard Model reference point (for a detailed review, see Ref. [1]). Its precise determination is essential in order to be able to disentangle possible contributions from physics beyond the Standard Model to other  $CP$ -violating observables that enter the global CKM fit.

The channels that are most commonly used to determine  $\gamma$  are of the type  $B \rightarrow DK$ , where a  $b$  hadron decays to a neutral charm meson together with a kaon. When the final state is accessible to both  $D^0$  and  $\bar{D}^0$  decays, the neutral  $D$  meson is an admixture of the flavour eigenstates. Since these are produced through  $b \rightarrow u$  and  $b \rightarrow c$  transitions, their interference is sensitive to the relative weak phase  $\gamma$  [2–5]. By measuring the rates and  $CP$  asymmetries of such decays,  $\gamma$  can be determined with negligible theoretical uncertainty [6].

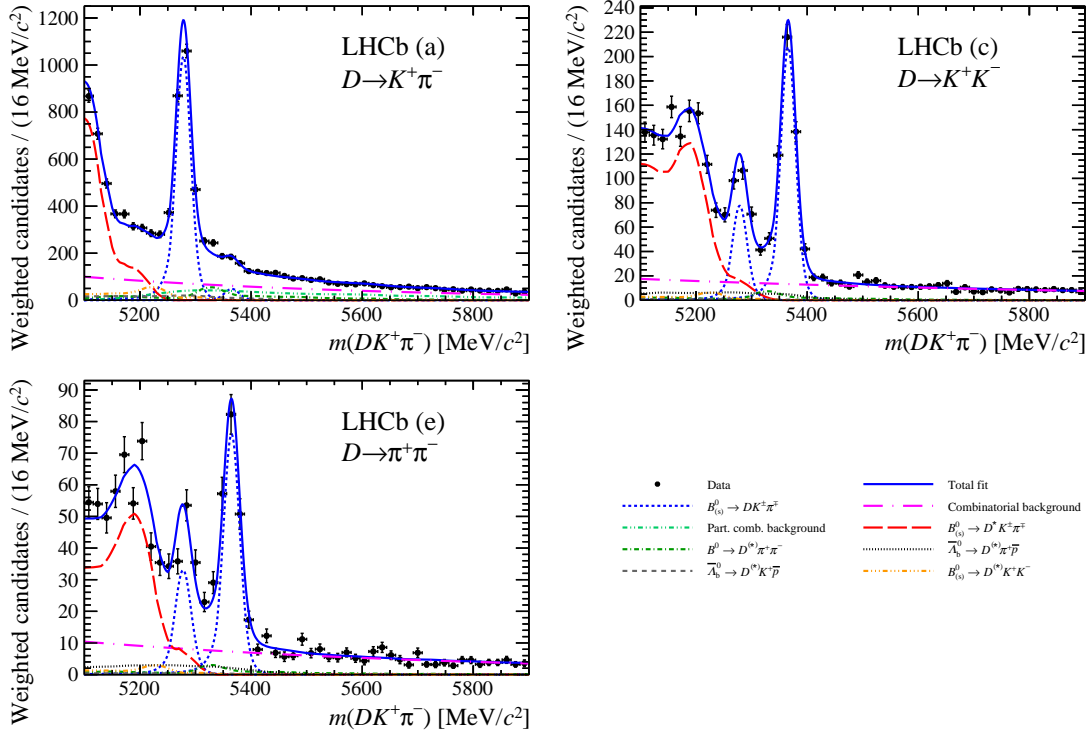
The method can be extended to  $B \rightarrow DK\pi$  decays. In this case, amplitude analysis of the  $B$  decay Dalitz plot [7] provides direct information about the relative phases, and therefore can be used to obtain precise information about  $\gamma$  without ambiguities in the solution. In particular, in the Dalitz plot analysis of  $B^0 \rightarrow DK^+\pi^-$  decays,<sup>1</sup> interference between  $B^0 \rightarrow DK^*(892)^0$  and  $B^0 \rightarrow D_2^*(2460)^-K^+$  amplitudes can be used to obtain more information about  $\gamma$  than is available in a quasi-two-body analysis [8, 9]. A key point is that the  $D_2^*(2460)^-K^+$  amplitude is flavour-tagged and therefore does not depend on the  $D$  decay final state. The method also allows the determination of additional hadronic parameters such as coherence factors that enter the formalism of the quasi-two-body approach [10].

Determination of  $\gamma$  with this method has recently been achieved, for the first time, by LHCb. In the first step, the Dalitz plot distribution of  $B^0 \rightarrow \bar{D}^0K^+\pi^-$  decays is obtained by fitting a sample reconstructed in the  $\bar{D}^0 \rightarrow K^+\pi^-$  channel (which is flavour-specific, to a good approximation). With the full LHC run I data sample of  $3\text{ fb}^{-1}$  of  $pp$  collision data at centre-of-mass energies of  $\sqrt{s} = 7$  and  $8\text{ TeV}$ ,  $2344 \pm 66$  signal decays are found inside the  $B^0$  signal window [11]. The Dalitz plot analysis provides a model for the  $b \rightarrow c$  transition, and reveals that the largest resonant contributions are from the  $K^*(892)^0$  and  $D_2^*(2460)^-$  states, with additional significant components from  $K\pi$  and  $D\pi$  S-waves. Results on the masses and widths of the  $D_0^*(2400)^-$  and  $D_2^*(2460)^-$  states are also obtained in the analysis.

With the  $b \rightarrow c$  model thus established, the analysis is extended to include decays of the  $D$  meson to the  $CP$ -even  $K^+K^-$  and  $\pi^+\pi^-$  final states, where yields of  $339 \pm 22$  and  $168 \pm 19$  signal events are available inside the  $B^0$  signal window [12], as shown in Fig. 1. A simultaneous Dalitz plot fit, implemented in Laura++ [13] with the *jFit* method [14], is carried out to the samples with  $D \rightarrow K^+\pi^-$ ,  $K^+K^-$  and  $\pi^+\pi^-$  – this is the first such simultaneous Dalitz plot analysis ever performed. The  $B^0 \rightarrow DK^+\pi^-$ ,  $D \rightarrow K^+\pi^-$  sample is fitted with the  $b \rightarrow c$  model, while the model is modified for the  $D \rightarrow K^+K^-$  and  $\pi^+\pi^-$  samples to account for effect of the  $b \rightarrow u$  contributions. Specifically, the complex coefficient  $c_j$  which describes the relative contribution of the resonance  $j$  to the overall amplitude is unchanged for  $D\pi^-$  resonances, since the charge of the pion tags the flavour of the resonance, while amplitudes for  $K^+\pi^-$  resonances receive additional contributions,

$$c_j \longrightarrow \begin{cases} c_j & \text{for a } D\pi^- \text{ resonance,} \\ c_j [1 + x_{\pm,j} + iy_{\pm,j}] & \text{for a } K^+\pi^- \text{ resonance,} \end{cases} \quad (1)$$

<sup>1</sup>The inclusion of charge conjugate processes is implied throughout unless explicitly stated otherwise.

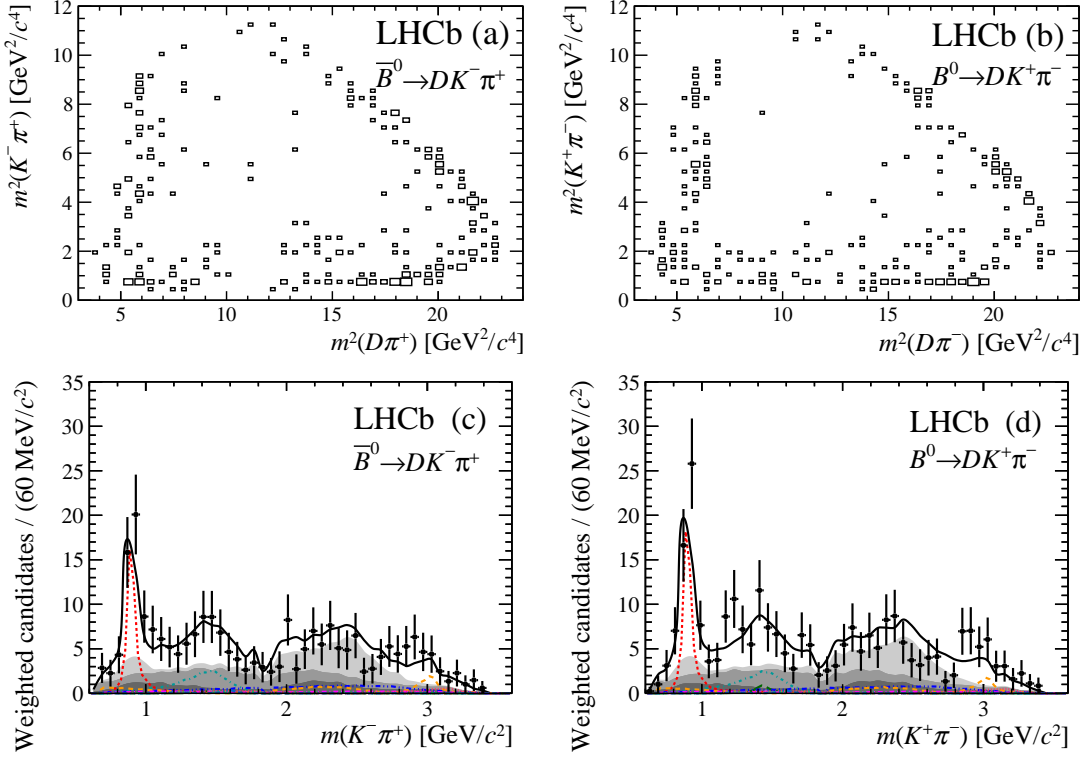


**Figure 1:** Candidates for  $B^0 \rightarrow DK^+\pi^-$  decays in the (top left)  $D \rightarrow K^+\pi^-$ , (top right)  $K^+K^-$  and (bottom left)  $\pi^+\pi^-$  channels [12]. The largest peak in the  $CP$ -eigenstate modes is due to  $B_s^0 \rightarrow \bar{D}^0 K^- \pi^+$  decays, with an associated satellite peak (long-dashed red line) from  $B_s^0 \rightarrow \bar{D}^{*0} K^- \pi^+$  decays. The candidates have been weighted by the signal-to-background fractions in the different samples that are fitted.

with  $x_{\pm,j} = r_{B,j} \cos(\delta_{B,j} \pm \gamma)$  and  $y_{\pm,j} = r_{B,j} \sin(\delta_{B,j} \pm \gamma)$ , where the  $+$  and  $-$  signs correspond to  $B^0$  and  $\bar{B}^0$  decay amplitudes, respectively. Here  $r_{B,j}$  and  $\delta_{B,j}$  are the relative magnitude and strong phase of the  $b \rightarrow u$  and  $b \rightarrow c$  amplitudes for each  $K^+\pi^-$  resonance  $j$ . A component corresponding to the  $B^0 \rightarrow D_{s1}^{*+}(2700)\pi^-$  decay, which is mediated by the  $b \rightarrow u$  amplitude alone, is also included.

The Dalitz plots for candidates in the  $B^0 \rightarrow DK^+\pi^-$ ,  $D \rightarrow K^+K^-$  and  $\pi^+\pi^-$  samples combined are shown in Fig. 2, together with projections of the data and the fit result onto  $m(K^\mp\pi^\pm)$ . Within the available statistics, there is no evidence for  $CP$  violation. The results for the parameters of the  $B^0 \rightarrow DK^*(892)^0$  decay are consistent with those of a quasi-two-body analysis based on the same data sample [15]. The determination of the  $x_{\pm}, y_{\pm}$  parameters of Eq. (1) allows also a comparison with results obtained from the  $B^0 \rightarrow DK^*(892)^0$ ,  $D \rightarrow K_s^0\pi^+\pi^-$  and  $K_s^0K^+K^-$  mode [16, 17]: the  $x_{\pm}$  results from the  $B^0 \rightarrow DK^+\pi^-$  Dalitz plot analysis are slightly more precise, while the  $y_{\pm}$  results are slightly less precise; all results are consistent.

Since the central values of the  $x_{\pm}, y_{\pm}$  parameters are not significantly different from zero, limited precision on  $\gamma$  is obtained using the results of the  $B^0 \rightarrow DK^+\pi^-$  Dalitz plot analysis alone. However, the analysis also yields information about the hadronic parameters needed to interpret results obtained from quasi-two-body analyses. In particular, the coherence factor  $\kappa$ , which would be unity in the case that the  $K^*(892)^0$  selection window contains only contributions from the  $K^*(892)^0$



**Figure 2:** (Top) Dalitz plots for (left)  $\bar{B}^0$  and (right)  $B^0$  candidates, together with (bottom) their projections onto  $m(K^\mp \pi^\pm)$  with results of the fit superimposed [12]. In the projections the shaded areas indicate backgrounds, while the red dotted line is the contribution from the  $\bar{K}^*(892)^0$  resonance.

resonance, is determined to be  $\kappa = 0.958_{-0.010}^{+0.005} {}_{-0.045}^{+0.002}$ , where the uncertainties are statistical and systematic. The results therefore have an important impact on the combined determination of  $\gamma$  using results from all  $B \rightarrow DK$  type processes [18, 19]. The LHCb combination [18] gives for the ratio of magnitudes of  $b \rightarrow u$  and  $b \rightarrow c$  amplitudes,  $r_B(DK^*(892)^0) = 0.218_{-0.047}^{+0.045}$ , smaller than but consistent with the expected value of  $\sim 0.3$ . Analyses with larger data samples will therefore be important to see if this value increases, in which case  $B^0 \rightarrow DK^*(892)^0$  decays will have an even larger impact on the overall combination than now. In addition to increasing the size of the sample, it will be important to improve understanding on  $K\pi$  and  $D\pi$  S-wave amplitudes (for which good progress has been reported recently [20, 21]) and to control background contributions from  $B_s^0 \rightarrow \bar{D}^{*0} K^- \pi^+$  decays (the  $B_s^0 \rightarrow \bar{D}^0 K^- \pi^+$  Dalitz plot has already been studied [22, 23]). More  $D$  meson decay modes can also be added, including the possibility of a model-independent  $B^0 \rightarrow DK^+ \pi^-$ ,  $D \rightarrow K_s^0 \pi^+ \pi^-$  double Dalitz plot analysis [24].

Given the success of the  $B \rightarrow DK\pi$  Dalitz plot analysis, it is reasonable to ask whether similar approaches can be applied for other  $B \rightarrow Dhh'$  modes. The isospin partner  $B^+ \rightarrow DK^+ \pi^0$  would be more challenging experimentally, due to the presence of a neutral pion in the final state. A further challenge in this channel is that  $D\pi^0$  resonances are not flavour tagged by the charge of the pion, so the associated amplitudes can differ depending on the  $D$  meson final state. While this complicates the formalism, it also means that in principle there may be more interference between

$b \rightarrow u$  and  $b \rightarrow c$  amplitudes, and therefore better sensitivity to  $\gamma$ . Although this mode has been investigated in the past [25], a more detailed investigation taking into account the latest knowledge is warranted [26].

A further potential advantage of the  $B^+ \rightarrow DK^+\pi^0$  mode is that for many  $D\pi$  resonances, the relative magnitude of the contributing amplitudes  $r_B$  can be known independently from studies of the  $B^+ \rightarrow D^+K^+\pi^-$  and  $B^+ \rightarrow D^-K^+\pi^+$  decays [27]. (This is not the case for the  $D^*(2007)^0$  resonance, which is below threshold for decay to  $D^+\pi^-$ , but has other advantages [28]. The quasi-two-body approach is preferable for analysis of decays involving this narrow resonance.) Both these modes have recently been observed by LHCb [29,30]. A large  $D_2^*(2460)^0$  component is seen in the Dalitz plot analysis of the favoured mode. In the suppressed mode, the available statistics are not sufficient for amplitude analysis, so instead a novel method involving weighting data by angular moments is used to set a limit on the  $\bar{D}_2^*(2460)^0$  contribution. These results give an upper limit  $r_B(D_2^*(2460)K^+) < 0.30$  (0.36) at 90 (95) % confidence level.

Extending to four-body decays, similar methods could potentially be used to determine  $\gamma$  from the interference of  $b \rightarrow u$  and  $b \rightarrow c$  amplitudes in  $B^+ \rightarrow D_1(2420)K^+$  decays. A possible sign of the  $b \rightarrow c$  decay was seen in early LHCb data [31], in the  $\bar{D}_1(2420)^0 \rightarrow \bar{D}^0\pi^+\pi^-$  channel. In the case that the ground state  $D$  meson is reconstructed as a  $CP$  eigenstate, it is possible that decays of  $D_1(2420)$  to the  $D^{*+}\pi^-$  and  $D(\pi^+\pi^-)$  would allow interference between flavour-tagged and untagged  $D$  mesons in the same final state. However, a full four-body amplitude analysis may be necessary, as there is also a significant contribution from  $B^+ \rightarrow DK_1^+ \rightarrow DK^+\pi^+\pi^-$ , which has been used to determine  $\gamma$  with a quasi-two-body approach [32]. Further studies will be necessary to establish by how much such an analysis would benefit the sensitivity to  $\gamma$ .

All  $b$  hadron species are produced in  $pp$  collisions, and LHCb has recorded large samples of  $B_s^0$  and  $\Lambda_b^0$  decays. Hence, possibilities to determine  $\gamma$  in  $B_s^0 \rightarrow DK^+K^-$  and  $\Lambda_b^0 \rightarrow DpK^-$  decays can also be considered. Both of these modes have been observed in LHCb data [33–35], but with modest yields. Moreover, a full analysis of  $B_s^0 \rightarrow DK^+K^-$  decays requires tagging of the initial  $B$  meson flavour, which leads to a reduction of sensitivity. In the case of  $\Lambda_b^0 \rightarrow DpK^-$  decays, the kinematic boundary of the phase space (due to the proton mass) limits overlap between  $Dp$  and  $pK^-$  amplitudes, and thus it is unclear how much gain in sensitivity may be possible compared to the quasi-two-body analysis. A detailed study of the amplitude structure of  $\Lambda_b^0 \rightarrow DpK^-$  decays, similar to that recently performed for the related  $\Lambda_b^0 \rightarrow Dp\pi^-$  channel [36] will be needed to address this issue.

In summary,  $B \rightarrow Dhh'$  decays provide many interesting ways to determine  $\gamma$ , with Dalitz plot analysis methods being particularly sensitive in certain cases. The results from these methods on the  $B^0 \rightarrow DK^*(892)^0$  mode give competitive sensitivity to those from  $B^+ \rightarrow DK^+$ , with the precision expected to improve further as results with additional  $D$  decay modes become available. Other  $B \rightarrow Dhh'$  decays, which have not yet been used to determine  $\gamma$ , are well worth pursuing, since in addition to helping to improve the overall knowledge of  $CP$  violation, these channels can also provide interesting results in charm meson spectroscopy.

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