

Combined Variables in Isolated Photon Production in Deep Inelastic Scattering at HERA

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Isolated photons with high transverse energy have been studied in ep deep inelastic scattering with the ZEUS detector at HERA, using an integrated luminosity of 326 pb^{-1} . The kinematic region includes photon virtualities $10 < Q^2 < 350 \text{ GeV}^2$. Photons with transverse energy $4 < E_\gamma^T < 15 \text{ GeV}$ and pseudorapidity $0.7 < \eta_\gamma < 0.9$ were measured with accompanying jets having transverse energy and pseudorapidity $2.5 < E_{jet}^T < 35 \text{ GeV}$ and $1.5 < \eta_{jet} < 1.8$. Differential cross sections are presented for several different correlated variables: the fraction of the incoming photon energy and momentum that is transferred to the photon and the jet, the fraction of proton energy taken by the parton that interacts with the photon, the azimuthal angle and the pseudorapidity difference between the prompt photon and the jet and between the prompt photon and the scattered electron. Comparisons with different theoretical calculations are made.

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1. Introduction

The subject of this talk is a continuation to the previous study of prompt photon production in ep deep inelastic scattering (DIS). Outgoing photons with high transverse energy, $E_T > 4\text{GeV}$, may be radiated either by a scattered lepton (LL photons) or by a parton (QQ photons), see Fig.1. A photon isolation requirement reduces the backgrounds from decay products of a high energy hadron and photons from an outgoing high- E_T quark which is problematic to calculate [15].

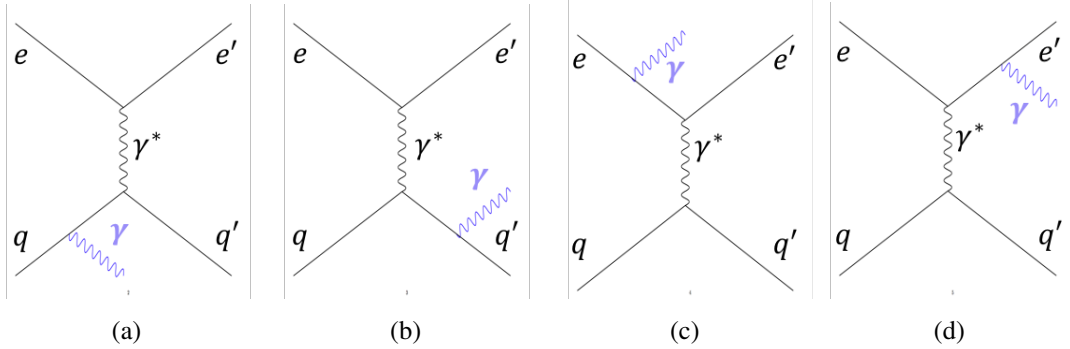


Figure 1: Lowest-order diagrams for photon production in ep scattering. (a) - (b): quark radiative diagrams; (c) - (d): lepton radiative diagrams.

Prompt photons are produced directly during the partonic interaction and in this way they provide a good tool for both QCD and QED tests giving opportunities for additional check of the proton's parton distribution functions. Compared to the studies of final hadronic states, photons do not change their properties before detection since they are not affected by the hadronisation, unlike partons forming a jet. Study of the dynamics of prompt photon emission can be used to probe different theoretical models. One also can observe how the dynamics changes with virtuality scale by comparing new results at different kinematical ranges with the corresponding results in photoproduction [1]. Additionally photons are considered as a possible background to new physics processes and therefore their production should be well understood.

2. Apparatus

For the measurement, a combined data sample, recorded using the ZEUS detector at HERA, with both e^+p and e^-p events and integrated luminosity of 326 pb^{-1} was used. The lepton and proton beam energy were respectively $E_e = 27.5\text{ GeV}$ and $E_p = 920\text{ GeV}$. The detailed description of ZEUS detector can be found elsewhere [2].

The principal component of the ZEUS detector used in this analysis was a high-resolution uranium-scintillator calorimeter (CAL). The calorimeter was divided into three regions, forward (FCAL), barrel (BCAL) and rear(RCAL), and each region consisted of a finely segmented electromagnetic section outside which was a hadronic section with larger cells. The "forward" (Z) direction is the proton beam direction. The fine granularity along the beam axis (5 cm per cell) of the electromagnetic calorimeter of the BCAL (BEMC) allowed a good separation of signal photons from the neutral hadrons decay background. Photons from reactions such as $\pi_0 \rightarrow \gamma\gamma$ have

in general wider spread along the axis and often hit few neighbouring cells while signal photon hits a single cell. The cut on the width of the signal was chosen in the way background could be evaluated.

3. Event reconstruction

The basic event selection and reconstruction were performed as previously [3]. The most important physics selections are as follows:

- **Prompt photon selection**

- $4 < E_T^\gamma < 15$ GeV - transverse energy of the prompt photons candidate
- $0.7 < \eta^\gamma < 0.9$ - pseudorapidity of the prompt photon candidate (detected in BCAL)
- $\frac{E_{EMC}}{E_{EMC} + E_{HAD}} > 0.9$ - at least 90% of the candidate's energy is deposited in BEMC
- $\Delta R(\eta, \phi) < 0.2$ - no other tracks within a narrow cone in the $\eta - \phi$ plane (isolation cut). Here $\Delta R(\eta, \phi) = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2}$
- $\frac{E^\gamma}{E_{jet\ with\ \gamma}} > 0.9$ - at least 90% of the reconstructed jet energy is carried by prompt photon candidate

- **Accompanying jet selection**

- $E_T^{jet} > 2.5$ GeV - transverse energy of the jet
- $1.5 < \eta^\gamma < 1.8$ - pseudorapidity of the jet
- E_{max}^{jet} - the jet that has a highest reconstructed energy is considered as accompanying

- **Kinematics selection**

- $10 < Q^2 < 350$ GeV² - energy of the prompt photons candidate
- $E^e > 10$ GeV - energy of the scattered lepton
- $140^\circ < \theta^e < 180^\circ$ - electron scattering angle
- $35 < E - p_Z < 65$ GeV - calculated from the recorded energy in the detector and corresponds to roughly twice the energy of the incoming electron, showing that the event is well recorded.

Jets were reconstructed using the k_T clustering algorithm with the radius parameter set to 1.0, making use of energy-flow objects (EFO's), constructed from clusters of calorimeter cells and tracks, and excluding the scattered electron.

To extract the photon signal from the background, the quantity $\langle \delta Z \rangle$, defined as the E_T - weighted mean of the distance of the Z position of the electromagnetic cells in the cluster relative to the mean Z of the cluster: $\langle \delta Z \rangle = \sum_i E_i |Z_i - Z_{cluster}| / (w_{cell} \sum_i E_i)$, was used as described in previous ZEUS analyses [3, 4, 5, 6, 7, 8]. Here, Z_i is the Z position of the centre of the i -th cell, $Z_{cluster}$ is

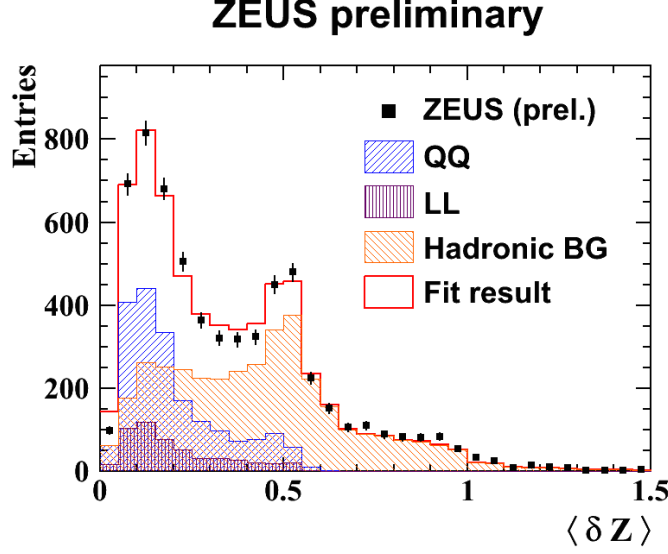


Figure 2: Distribution of $\langle \delta Z \rangle$ for the entire selected data sample. The error bars represent the statistical uncertainties. The red line shows a fit to the data of three components with fixed shapes as described in the text. The blue shaded histogram represents the QQ component of the fit, and the purple histogram - the LL component.

the centroid of the EFO cluster, w_{cell} is the width of the cell in the Z direction, and E_i is the energy recorded in the cell. The sum runs over all BEMC cells in the EFO.

The total $\langle \delta Z \rangle$ distribution is shown in Fig.2 plotted in units of electromagnetic calorimeter cell widths. The MC samples corresponding to different sources of high- E_T photons were scaled to the level of data luminosity. The LL component (generated using ARIADNE 4.12 [9]) of the signal was extracted from the data before the fit, in which only one parameter was varied - the fraction of QQ (generated using PYTHIA 6.416 [10]) component in the data after the extraction of LL component from it. A peak at low $\langle \delta Z \rangle$ indicates that most of the energy is found in one cell - actual prompt photons. The broader background peaks around a value of approximately 0.5, where the cluster energy is shared between two contiguous cells - hadronic background. Fits are performed to the data in each bin of each quantity in the range $0.05 < \langle \delta Z \rangle < 0.8$ whose cross section is to be measured, to extract the photon signal.

4. Combined variables

For this paper a set of six combined variables was studied. They are defined as follows:

- the variable x_γ is a measure of the fraction of the incoming photon energy and momentum that is given to the photon and the jet:

$$x_\gamma = \frac{E^\gamma - p_z^\gamma + E^{\text{jet}} - p_z^{\text{jet}}}{2E_e y_{JB}},$$

where E^γ and E^{jet} denote the energies of the photon and the jet, respectively, p_z denotes the corresponding longitudinal momenta, $E_e = 27.5$ GeV, $y_{JB} = \sum(E - p_z)/2E_e$, summing over all energy-flow objects in the event, each treated as equivalent to a massless particle;

- the kinematic quantity x_p equals to the fraction of proton energy taken by the parton that interacts with the photon:

$$x_p = \frac{E^\gamma + p_z^\gamma + E^{jet} + p_z^{jet}}{2E_p},$$

where $E_p = 920$ GeV. This distribution is sensitive to the partonic structure of the proton;

- $\Delta\phi$ is the azimuthal angle between the prompt photon and the jet: $\Delta\phi = \phi^{jet} - \phi^\gamma$, where ϕ^γ and ϕ^{jet} denote azimuthal angles of the photon and the jet, respectively. This variable is sensitive to the presence of higher-order gluon radiation relative to the outgoing quark, which can generate non-collinearity between the photon and the leading jet;
- $\Delta\eta$ is the pseudorapidity difference between the prompt photon and the jet: $\Delta\eta = \eta^{jet} - \eta^\gamma$, where η^γ and η^{jet} denote pseudorapidity of the photon and the jet, respectively. The pseudorapidity is sensitive to the spin of the exchanged quantum [12];
- $\Delta\phi^{e,\gamma}$ is the azimuthal angle between the prompt photon and the scattered electron: $\Delta\phi^{e,\gamma} = \phi^e - \phi^\gamma$, where ϕ^γ and ϕ^e denote azimuthal angles of the photon and the electron, respectively;
- $\Delta\eta^{e,\gamma}$ is the pseudorapidity difference between the prompt photon and the scattered electron: $\Delta\eta^{e,\gamma} = \eta^e - \eta^\gamma$, where η^γ and η^e denote pseudorapidity of the photon and the scattered electron, respectively;

A similar analysis was previously done for photoproduction [11] for all variables except the last two, since the scattered electron is unobservable in photoproduction.

5. Summary

The cross sections were calculated using acceptance corrections estimated on a bin-by-bin basis. Systematic uncertainties are dominated by the photon energy scale and the modelling of the background.

Preliminary resulting DIS cross sections as functions of $0 < x_\gamma < 1$, $0 < x_p < 0.07$, $0^\circ < \Delta\phi < 180^\circ$, $-2.2 < \Delta\eta < 2$, $0^\circ < \Delta\phi_{e,\gamma} < 180^\circ$ and $-3.6 < \Delta\eta_{e,\gamma} < -0.6$ are shown in Fig.3-5. For these figures, the MC QQ cross section was scaled by a factor of 1.6, as in previous publication. Results are provided in the kinematical ranges $10 < Q^2 < 350$ GeV², $10 < Q^2 < 30$ GeV² and $30 < Q^2 < 350$ GeV² respectively. The 30 GeV² point was chosen to split the available data sample in two with roughly the same statistics. After the rescaling Pythia MC shows a good agreement in shape with the data.

The cross sections are additionally compared to two theory models on the whole Q^2 range: the model of Baranov *et al.*, [15], in which unintegrated parton densities are used in a LO QCD framework, see Fig.6 and a NLO QCD calculation of Aurenche, Fontannaz and Guillet (AFG)[13, 14]. The uncertainties for the AFG predictions are set to 10% arising from the QCD scale at this

order. Both predictions are provided by the authors at the parton level and later bin-by-bin corrected for the hadronisation effects. Hadronisation corrections are estimated from the $10 < Q^2 < 350$ GeV range using MC samples produced with Pythia.

There is a reasonable agreement between the experimental data and the shapes of the model distributions (except for x_γ).

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ZEUS preliminary

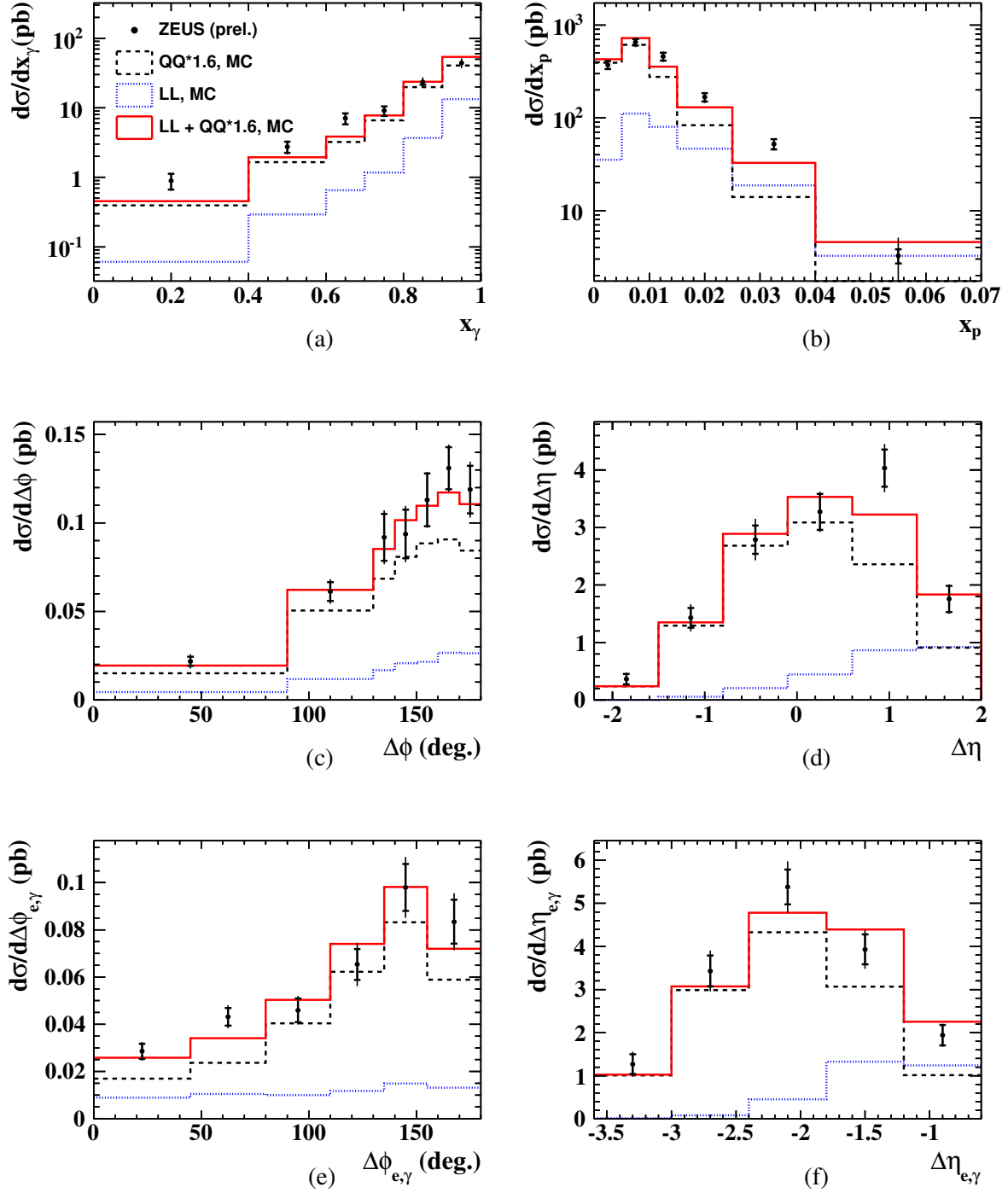


Figure 3: Differential cross sections in (a) x_γ , (b) x_p , (c) $\Delta\phi$ (deg.), (d) $\Delta\eta$, (e) $\Delta\phi_{e,\gamma}$ (deg.), and (f) $\Delta\eta_{e,\gamma}$. The inner and outer error bars show, respectively, the statistical uncertainty and the statistical and systematic uncertainties added in quadrature. The solid histograms are the reweighted Monte Carlo predictions from the sum of QQ photons from PYTHIA normalized by a factor 1.6 plus DJANGO LL photons. The dashed (dotted) lines show the QQ (LL) contributions.

ZEUS preliminary

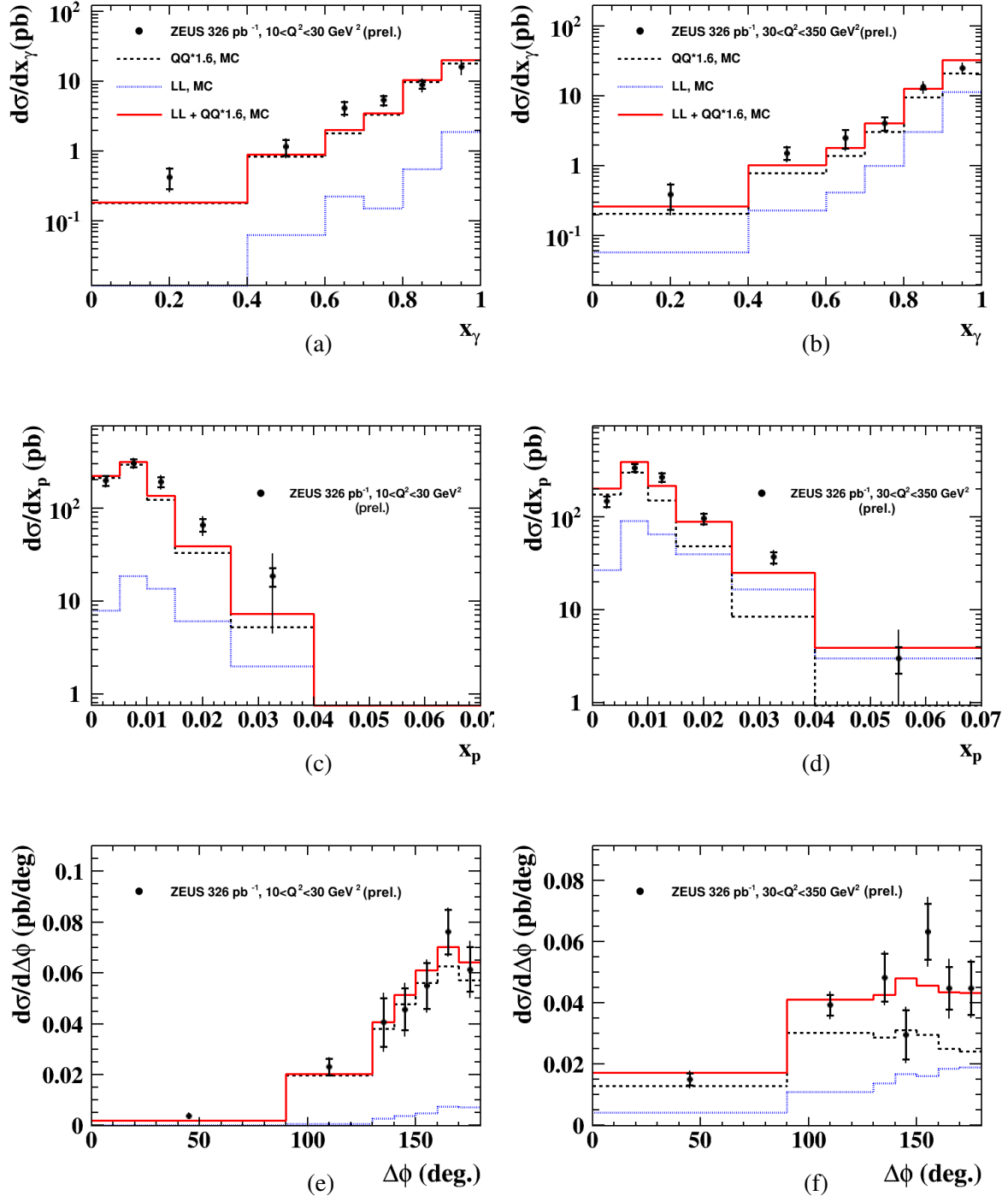


Figure 4: Differential cross sections for the separate regions $10 < Q^2 < 30$ and $30 < Q^2 < 350 \text{ GeV}^2$ of x_γ (a,b), x_p (c,d) and $\Delta\phi$ (deg.) (e,f). The inner and outer error bars show, respectively, the statistical uncertainty and the statistical and systematic uncertainties added in quadrature. The solid histograms are the reweighted Monte Carlo predictions from the sum of QQ photons from PYTHIA normalized by a factor 1.6 plus DJANGO LL photons. The dashed (dotted) lines show the QQ (LL) contributions.

ZEUS preliminary

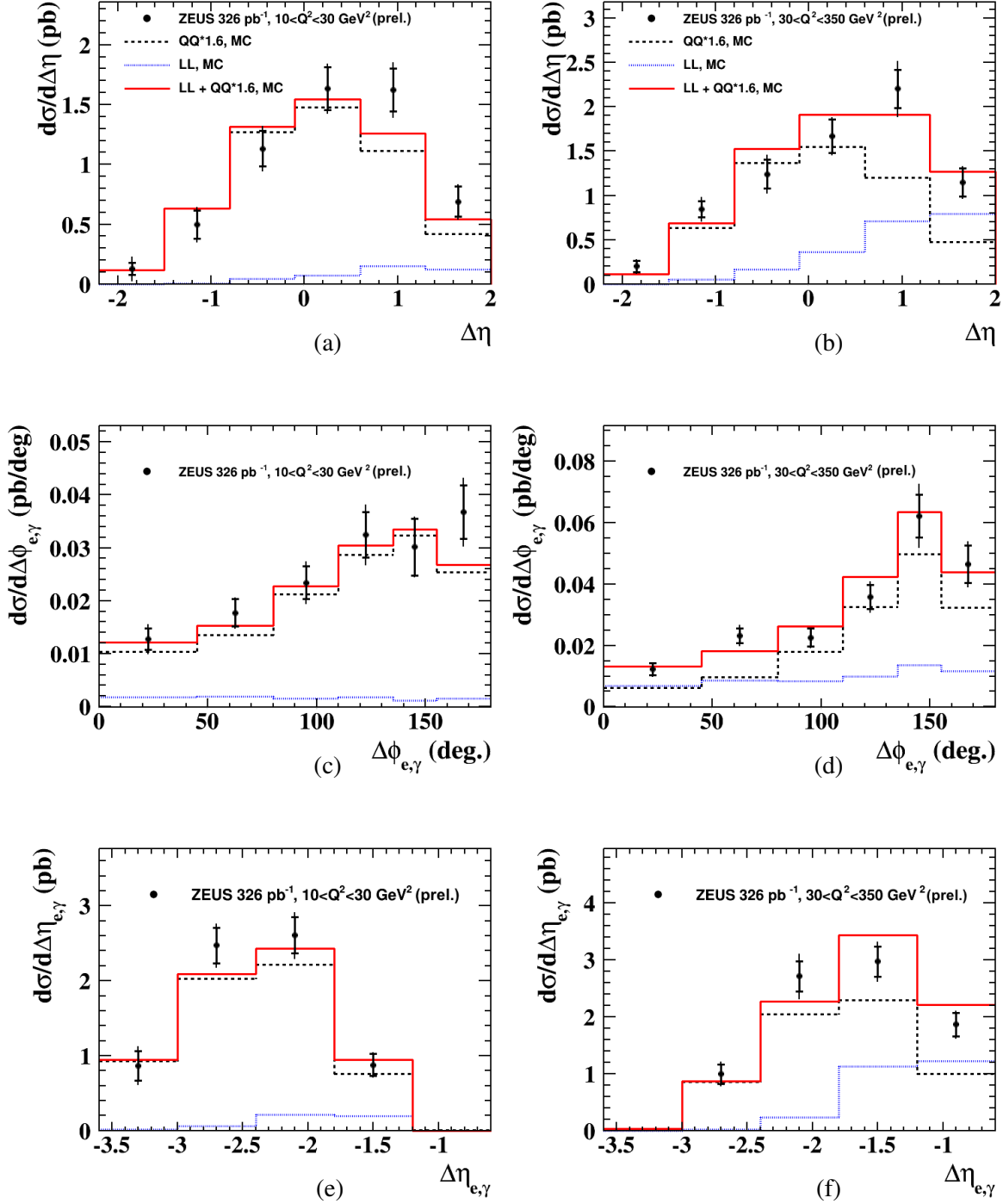


Figure 5: Differential cross sections for the separate regions $10 < Q^2 < 30$ and $30 < Q^2 < 350$ GeV^2 of $\Delta\eta$ (a,b), $\Delta\phi_{e,\gamma}$ (deg.) (c,d), and $\Delta\eta_{e,\gamma}$ (e,f). The inner and outer error bars show, respectively, the statistical uncertainty and the statistical and systematic uncertainties added in quadrature. The solid histograms are the reweighted Monte Carlo predictions from the sum of QQ photons from PYTHIA normalized by a factor 1.6 plus DJANGO LL photons. The dashed (dotted) lines show the QQ (LL) contributions.

ZEUS preliminary

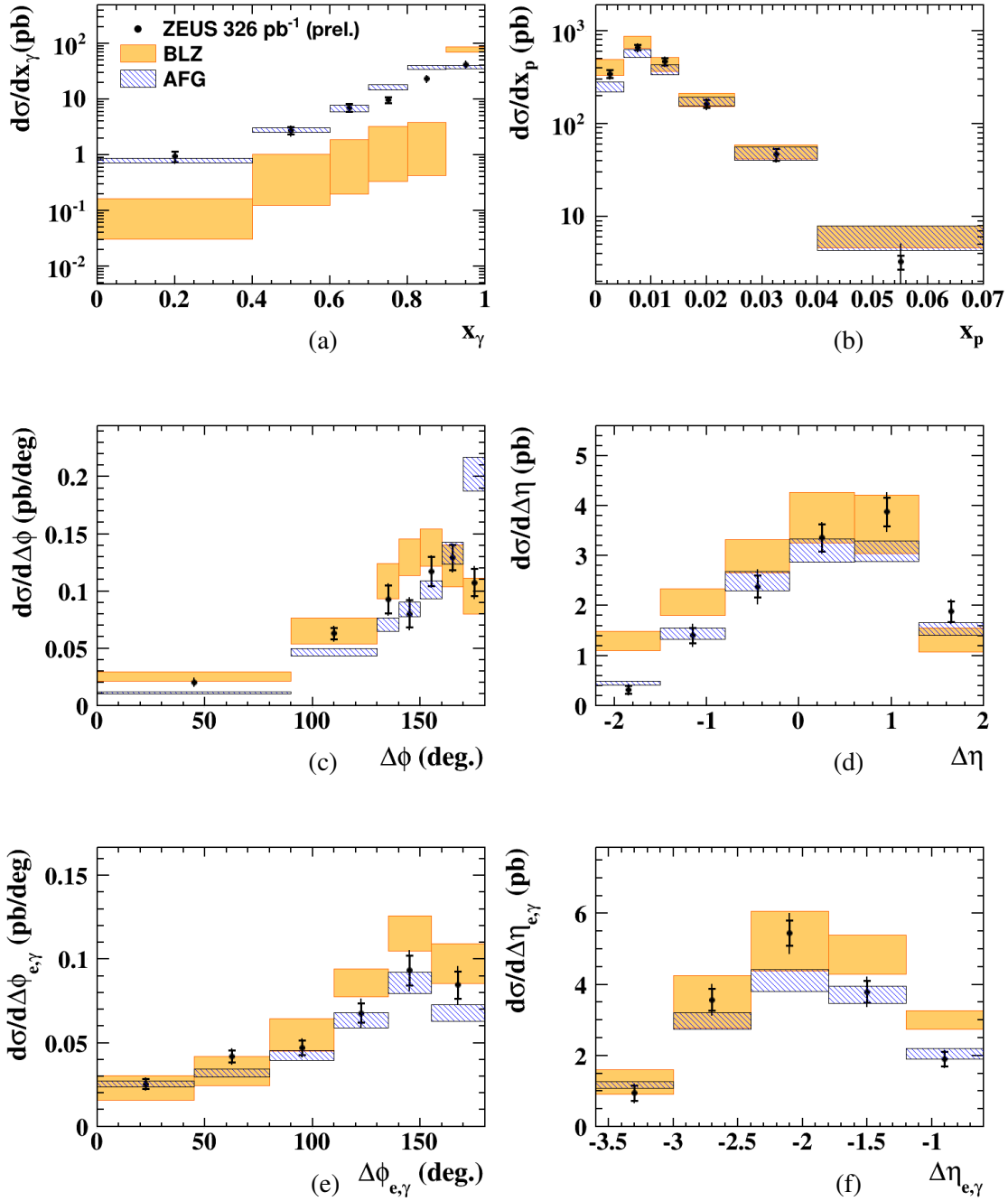


Figure 6: Data points as shown in Fig. 3. Theoretical predictions from BLZ [15] and AFG [13] are shown, with associated uncertainties indicated by the shaded bands.