

Cosmic ray fluxes and the role of sub-dominant source populations to the positron excess

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Interactions of cosmic-ray nuclei is the most well-understood contribution to the observed positron flux on Earth. Various scenarios have been proposed, including sources in the spiral arms of the galaxies and nearby isolated sources, which can contribute to the measured flux. In our work we focus on sources such as the gamma-ray novae which can contribute to the observed positron flux, although sub-dominant to the supernova remnants or pulsar-wind nebulae overall, but crucial in the 10s of GeV energy range.

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1. Introduction

Our Galaxy is a continuous source of most of the cosmic-ray charged particles observed on Earth. These cosmic rays, which can be classified into primary and secondary particles as well as the matter particles (proton, Helium, electron, etc.) and antimatter particles (antiproton, positron, etc.), convey a wealth of information about their sources in the Galaxy and intervening medium. Recent data from the balloon-borne detector *High Energy Antimatter Telescope* (HEAT) [1], satellite-borne detectors *Payload for Antimatter Matter Exploration and Light-nuclei Astrophysics* (PAMELA) [2], *Fermi-Large Area Telescope* (LAT) [3] as well as the *Alpha Magnetic Spectrometer* (AMS) [4] onboard the International Space Station has led to significant development in understanding the properties of the cosmic ray sources and interstellar medium (ISM). Observations of cosmic-ray positron flux by HEAT [5], PAMELA [6], AMS [4] and Fermi-LAT [7] have shown an excess in the ≈ 10 -500 GeV energy range, different than expected from propagation and interactions of cosmic-ray protons in the ISM. This excess of positron fraction in the data has been used to understand its origin via dark matter particles [8–11] and via astrophysical sources [12–21]. An astrophysical interpretation is widely favored over dark matter annihilation models at present, although it requires additional mechanisms and/or sources of positrons.

2. Model setup

In our model we considered a diffusive reacceleration with no convection scenario for propagation of cosmic rays using the three dimensional model in the DRAGON code. We did-not consider convection process during the cosmic-ray transport in our Galaxy, as this process lowers the B/C ratio in the 1-10 GeV energy range which needs to be recovered by using large values of Alfvén speed. The magnetic field structure has been selected in the DRAGON code of type Pshirkov, in which the disk component has a value $2\mu\text{G}$ and the halo component can be in a range 2-5 μG based on the radio synchrotron emission of the Galaxy [22]. We have used the Galactic magnetic field disk component value as $2\mu\text{G}$, halo component value as $4\mu\text{G}$ and the turbulent component value as $10\mu\text{G}$, respectively. We have taken the exponential mode of diffusion coefficient in the DRAGON code which varies with particle rigidity ρ and the Galaxy vertical height z [23].

In this propagation scenario we calculate the secondary nuclei, namely the Boron to Carbon ratio [24, 25] to fix the diffusion coefficient and the size of the Galactic halo [26]. In a recent AMS result it has been shown that above rigidity $\rho = 65$ GV, δ takes a value 0.333 ± 0.014 (stat) ± 0.005 (syst) [25]. Below 10 GeV, the effect of solar modulation on the charged particles can be adopted to fit the observational data [27].

3. Fluxes of nuclei and cosmic-ray powers

We have plotted the Boron to Carbon flux ratios in the left panel of Fig. 1, while in the right panel we have plotted the proton flux along with observed data. The proton flux detected by AMS-02 and PAMELA data can be explained by different activity of the solar medium¹. The solid (dotted) lines are fluxes after (before) taking into account solar modulation. Our choice of model parameters

¹All data are taken from the database [28] unless otherwise specified.

Parameters	δ	$D_0(\text{cm}^2/\text{s})$	$v_A(\text{km/s})$	η	$\gamma(p)$	$\gamma(e^-)$	$\gamma(e^\pm)$
Value	0.45	2.8×10^{28}	19	-0.005	2.0/2.32	1.91/2.71/2.36	1.85/2.32

Table 1: Galaxy parameters for a halo of size x,y,z ($= 12$ kpc, 12 kpc, 8 kpc), used in the solution for the transport equation of cosmic rays. All the nuclei follow the same spectrum with a rigidity break at 7 GV, the primary electron has two breaks at 6.2 GV and 85 GV and the additional population has a break at 3.8 GV.

fits the B/C ratio data from both the AMS-02 [25] and PAMELA [24] quite well in the whole observed energy range. This requires the index of the diffusion coefficient to be $\delta = 0.45$, close to the Kraichnan turbulence spectrum rather than the Kolmogorov turbulence spectrum of ≈ 0.33 reported in [25].

We calculate the total powers in cosmic-ray components, using the spectra in Table 1, by injecting and propagating these particles using the DRAGON code for $T = 64$ Million years in a volume $V = 1.15 \times 10^3 \text{ kpc}^3$ of the Galaxy, as

$$P = \frac{V}{T} \frac{4\pi}{c} \int_{E_1}^{E_2} n(E) E dE. \quad (3.1)$$

Here E_1 and E_2 are the two limiting energies in the spectrum. In our calculations, the approximate total power required in the protons is 10^{56} GeV/Myr, in primary electrons is 2.8×10^{54} GeV/Myr and in the additional e^\pm populations is 1.5×10^{53} GeV/Myr.

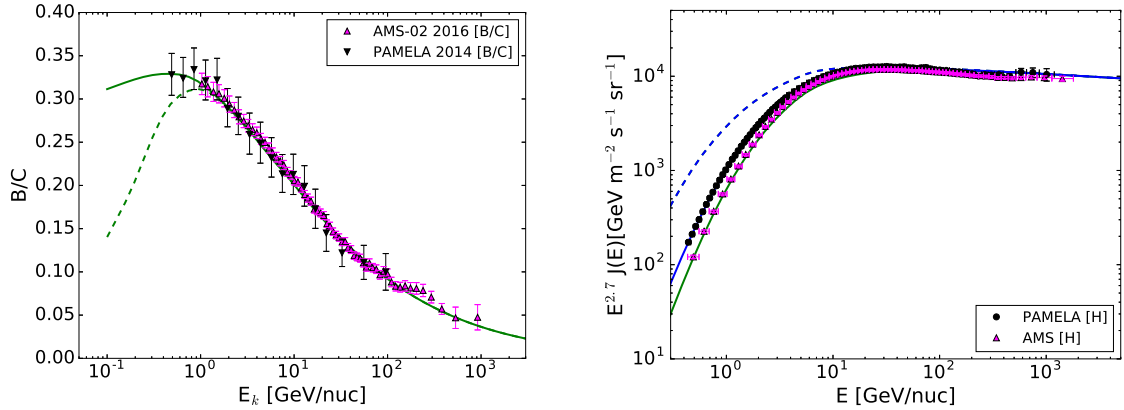


Figure 1: *Left panel:* B/C ratio calculation using the DRAGON code and plotted against the AMS [25] and PAMELA [24] data (From Joshi et al., 2017 [29]). *Right panel:* The proton flux plotted against the PAMELA [30] and AMS data [31]. The PAMELA proton data can be explained with a lower solar modulation $\phi = 0.42$ GV (blue solid line), while the AMS-02 data requires a stronger solar modulation $\phi = 0.65$ GV (green solid line). The solid (dashed) lines are fluxes with (without) solar modulation taken into account.

4. Contributions from nearby sub-kpc pulsars

We have selected nearby pulsars from the Australia Telescope National Facility (ATNF) catalogue² with parameters taken from there as well as from [32, 33] and we have listed those in Table 2.

²<http://www.atnf.csiro.au/people/pulsar/psrcat/>

Name	Distance (kpc)	Age (yr)	\dot{E} (erg/sec)	$E(t)$ (erg)
J0633+1746 (Geminga)	$0.25^{+0.23}_{-0.08}$	3.42×10^5	3.2×10^{34}	4.4×10^{48}
B1055-52	$0.35^{+0.15}_{-0.15}$	5.35×10^5	$3. \times 10^{34}$	1.0×10^{49}

Table 2: Pulsar parameters based on the ATNF catalogue. Geminga and B1055-52 dominantly contribute to the electron-positron fluxes in the 100 GeV-1 TeV range in our model.

If η_{e^\pm} is the efficiency for conversion of spin-down energy of a pulsar into electron-positron pairs, then we can estimate emitted energy in e^\pm fluxes as [13],

$$E(t) \simeq \eta_{e^\pm} \dot{E} \frac{t}{\tau_0}. \quad (4.1)$$

Here $\tau_0 \sim 10^4$ is the luminosity decay time for mature pulsars and \dot{E} is the spin-down luminosity for a given pulsar. We have taken $\eta_{e^\pm} = 0.4$ in Table 2.

5. Electron-positron flux levels and positron excess

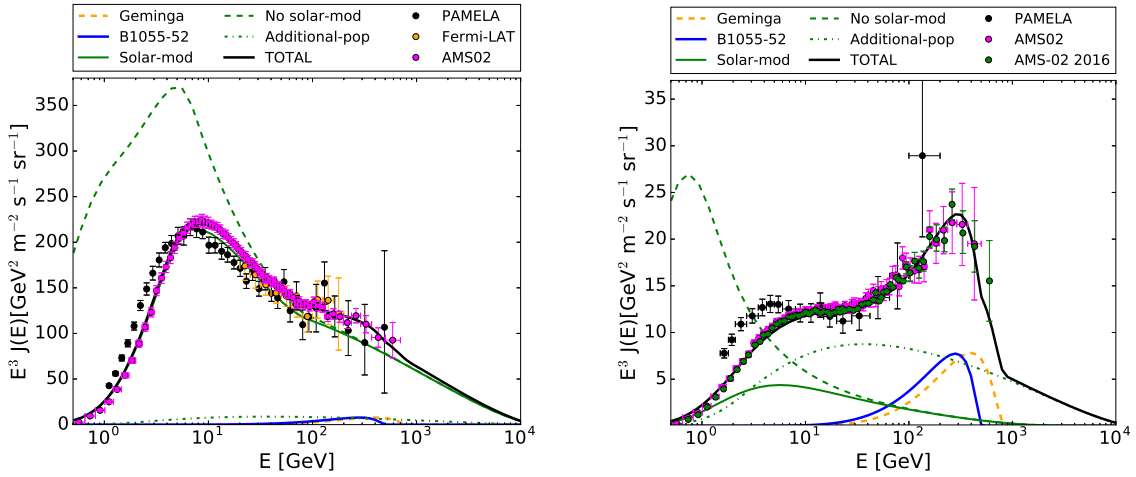


Figure 2: *Left panel:* Electron flux data from PAMELA [34], Fermi-LAT [7] and AMS-02 [35] (From Joshi et al.,2017 [29]). *Right panel:* Positron Flux data from PAMELA [6] and AMS-02 [35]). Also shown are our model fluxes with black solid line for the total flux in both panels. See main text for more details.

In Figure 2 the electron and positron flux coming from Galaxy due to cosmic ray interactions with the ISM, primary electron source contribution, primary e^\pm pair emitters (additional population) and nearby source contribution have shown. The total addition of all processes explains the AMS observations. To fit the positron flux at high energies we require a diffusion coefficient 28 times smaller for the diffusion of electrons and positrons from nearby sources than used for the Galactic diffusion of cosmic rays in the DRAGON code. We have assumed that Geminga and B1055-52 are contributing 40% ($\eta_{e^\pm} = 0.4$) of their spin-down power into electron-positron pair emission, as mentioned earlier. We have also taken an injection spectral index of 1.89 for both of these sources with an exponential folding energy of 1.3 TeV for Geminga and 1.5 TeV for B1055-52, which fit the observational data. Please see the earlier work [29] for a full description.

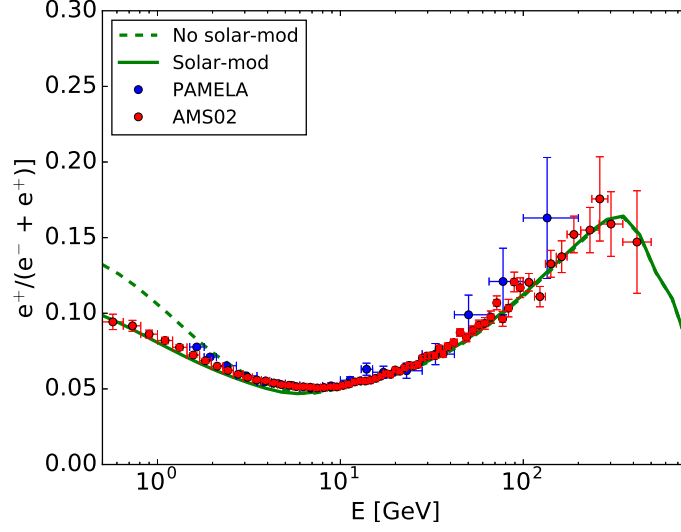


Figure 3: Positron fraction based on our model (solid line) plotted against data from PAMELA [6] and AMS-02 [36].

In Fig. 3 we have plotted the fractional positron flux of the total electron and positron flux. The solid (dashed) line corresponds to the flux model with (without) solar modulation taken into account. Our model fits both the AMS-02 data [35] and PAMELA data [6] rather well.

6. Summary and Discussion

The Proton flux we have calculated using the diffusion-reacceleration model in the DRAGON code can fit the AMS-02 data and PAMELA data with higher ($\phi = 0.65$ GV) and lower ($\phi = 0.42$ GV) solar modulation activity respectively. We found that the observed spectrum of electrons required two breaks for its successful interpretation. The low energy break is due to electrons cooling in the Galactic magnetic field, while the second break, after which the flux is harder, might have an origin in the source population of cosmic rays or in the energy dependent escape of electrons from their sources. We needed an additional population of sources, distributed similar to the conventional source population but producing primary electron-positron fluxes. A requirement for an additional source population to fit data is not new. Sources such as white dwarf pulsars [37] with magnetic field $\gtrsim 10^9$ G (see, e.g. [38]) can accelerate e^\pm to 10 TeV [39]. Other sources such as gamma-ray novae [40] can also possibly accelerate particles to very-high energies [41, 42].

We have also used contributions to the e^- and e^+ fluxes from the nearby pulsars, as is usually considered, e.g., [43, 44, 21] to explain observations in the $\gtrsim 100$ GeV-1 TeV range.

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