

## A Bethe-Heitler 5D polarized photon-to-e+e-pair conversion event generator

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**Denis Bernard\***

*LLR, Ecole polytechnique, 91128 Palaiseau, France*

*E-mail: [denis.bernard@in2p3.fr](mailto:denis.bernard@in2p3.fr)*

I present the exact, polarized, event generator of  $\gamma$ -ray conversion to  $e^+e^-$  pairs developed in the context of the gas-TPC HARPO project.

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\*Speaker.

## 1. Introduction

This work originated from a theoretical and experimental development of low density, homogeneous detectors such as a gas time-projection chamber as a high-performance  $\gamma$ -ray telescope and polarimeter in the  $\gamma \rightarrow e^+e^-$  regime, that is, above 1 MeV [1, 2]. In contrast with past and current pair telescopes, the single-track angular resolution of such detectors can be so good that

- at low energy, the single-photon angular resolution is now dominated by the contribution of the momentum of the recoiling ion, that escapes measurement [3];
- polarimetry has been predicted to be possible despite the dilution of the polarization asymmetry induced by multiple scattering [4] and has actually been demonstrated by the characterization of a TPC prototype on beam [6, 7, 2].

Therefore the simulation of modern, high-performance telescopes requires the use of an exact event generator,

- that describes correctly the recoil momentum distribution and
- that is able to simulate the conversion of linearly polarized  $\gamma$  rays.

In addition, polarimetry is extremely demanding in terms of sample statistics, and the ability of gas detectors to detect low energy photons (see Fig. 1 of Ref. [5]) close to threshold ( $2m_e c^2$  for “nuclear” conversion, i.e. in the field of a nucleus) is absolutely critical given the strongly decreasing dependence of the flux of cosmic sources with energy: for a  $\Gamma = 2$  spectral index source, the  $\gamma \rightarrow e^+e^-$  signal peaks at a couple of MeV (Fig. 2 of Ref. [4]). A generator that is free from high-energy approximation is needed. Last, upon conversion to a pair the final state is defined by a set of five variables (see Fig. 3 of Ref. [4]) and given the strong correlation between them in the differential cross section [4], the use of a generator that samples the 5D probability density function (pdf) is highly desirable.

Two mainframes are available for the simulation of high-energy physics (HEP) detectors and in particular of  $\gamma$ -ray telescopes: Geant4 [8] and EGS5 [9]. Their  $\gamma$ -ray conversion generators have been developed so as to reproduce the properties of electromagnetic showers as precisely as possible, not for the simulation of high-performance  $\gamma$ -ray telescopes. None of these mainframes include an event generator that samples the 5D pdf : instead products of 1D pdfs are used. Most often, the electron and the positron are generated back-to-back and their polar angles are generated independently: there happens to be some recoil induced from the momentum mismatch then, that is bound to lie in the conversion plane, but that has nothing to do with the true, QED, recoil momentum distribution [11]. Energy-momentum is obviously not conserved by these generators. The polarized physics model available in Geant4 does not simulate the conversion of linearly polarized  $\gamma$ -rays to pairs at low energy where it is most needed; at higher energy it provides an erratic value of the polarization asymmetry (Fig. 11 of [11]).

## 2. The generator

I have written an event generator that samples exactly (no high-energy nor small-angle approximation) the 5D differential cross section, including for photons with a non-zero linear polarization

fraction [4]. This event generator is using the BASES/SPRING instantiation [14] of the VEGAS method [25]. The differential cross section can be computed by either of two methods :

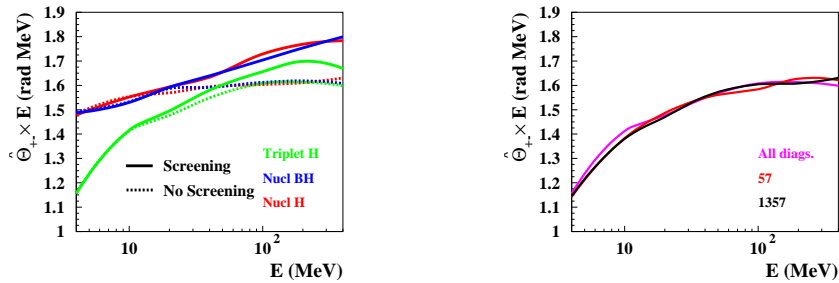
- by the computation of all Feynman diagrams (See Fig. 3 of Ref. [4]) that contribute to first order of the Born approximation, using the HELAS software [15];
- or using the analytical expression obtained by Bethe and Heitler in the unpolarized case [16, 17] from the computation of the two diagrams that dominate nuclear conversion, and by [12, 13, 18] in the fully polarized case (complete expressions, including for partially polarized photons, can be found in our [10]).

These two methods were extensively compared to each other, which proved to be extremely useful, in particular during the debugging phase of the project.

### 3. Generator validations

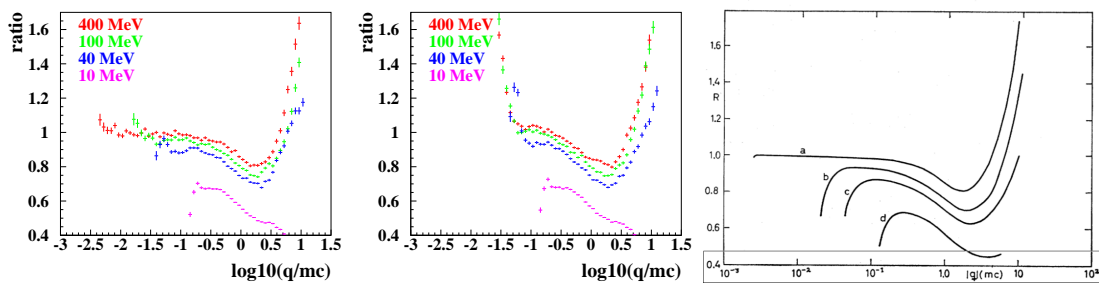
The initial publication of the generator contained a number of validations [4]

- Distributions of individual variables that define the final state, for the HELAS (Feynman amplitudes) and Bethe-Heitler models were compared;
- The value of the most probable value of the opening angle was compared to the analytical calculation by Olsen [19] (Fig. 4 of Ref. [4], here Fig. 1);

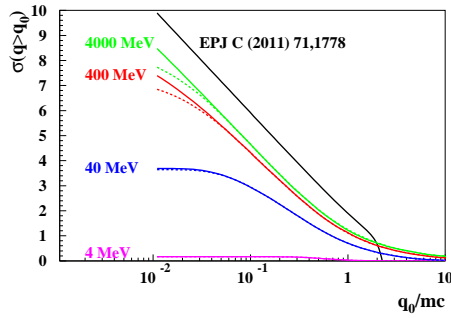


**Figure 1:** Position of the maximum  $\hat{\theta}_{+-} \times E$  of the distribution of  $\theta_{+-} \times E$  (rad MeV), as a function of photon energy. Left: for nuclear and triplet, with/without screening, for the H/BH pdf. Right: triplet with all diagrams, without exchange diagrams (1357) and with BH diagrams only (57).

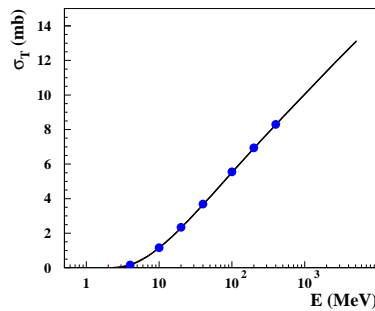
- The ratio of the recoil momentum distribution for triplet to nuclear conversions for various values of the energy of the incident photon was compared to the analytical calculation by Mork [20] (Fig. 5 of Ref. [4], here Fig. 2);
- The variation of the triplet cross section  $\sigma(q > q_0)$  above a recoil momentum threshold,  $q_0$ , as a function of  $q_0$  for several energies was compared with the high-energy asymptotic expression of Ref. [26] (Fig. 6 of Ref. [4], here Fig. 3);
- The total cross section for triplet conversion was compared to the computation by Mork [20] (Fig. 28, supplementary data of Ref. [4], here Fig. 4);



**Figure 2:** Ratio of the recoil momentum distribution for triplet conversion to nuclear conversion. Without (left) and with (center) screening. Right : analytical computation by Mork [20].

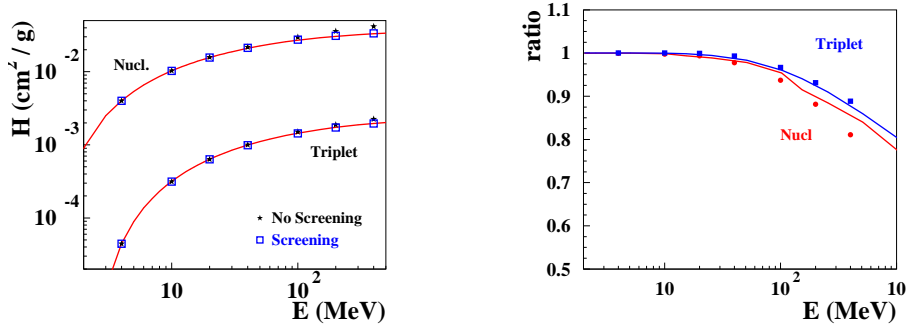


**Figure 3:** Variation of the triplet cross section  $\sigma(q > q_0)$  above a recoil momentum threshold,  $q_0$ , as a function of  $q_0$  in units of  $mc$ , compared with the high-energy asymptotic expression of EPJ C (2011) 71,1778 [26] (thick line). With (dashed line) and without (solid line) form factor applied.



**Figure 4:** Comparison of the total cross section for triplet conversion (bullets, no screening) with the computation by Mork (solid line) [20].

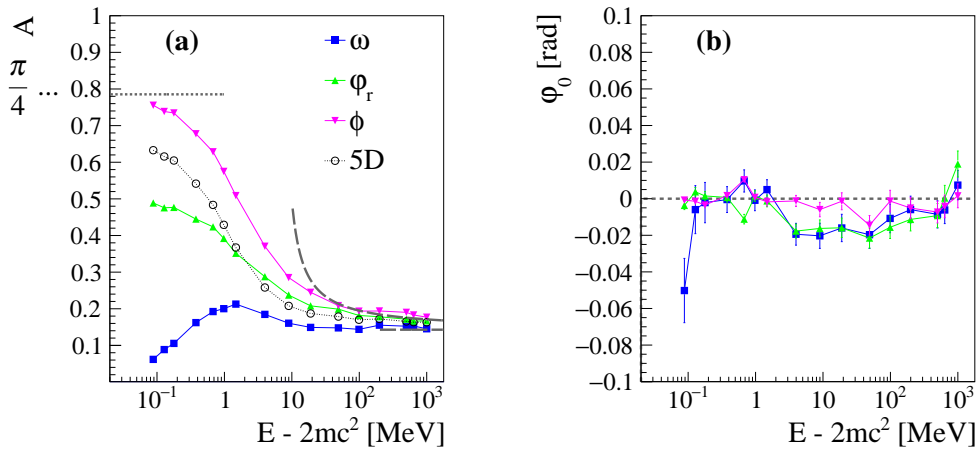
- The decrease of the total mass attenuation coefficients when screening is applied (coherent form factor for nuclear conversion [21], incoherent form factor for triplet conversion [22]) was compared with calculations by Hubbell [23] (Fig. 29, supplementary data of Ref. [4], here Fig. 5);



**Figure 5:** Left: Comparison of the total mass attenuation coefficients (argon; with (squares) and without (bullets) screening) with the data from NIST [24] (curves). Right: ratio of the cross sections with/without screening (symbols) with the results from Ref. [23] (curves).

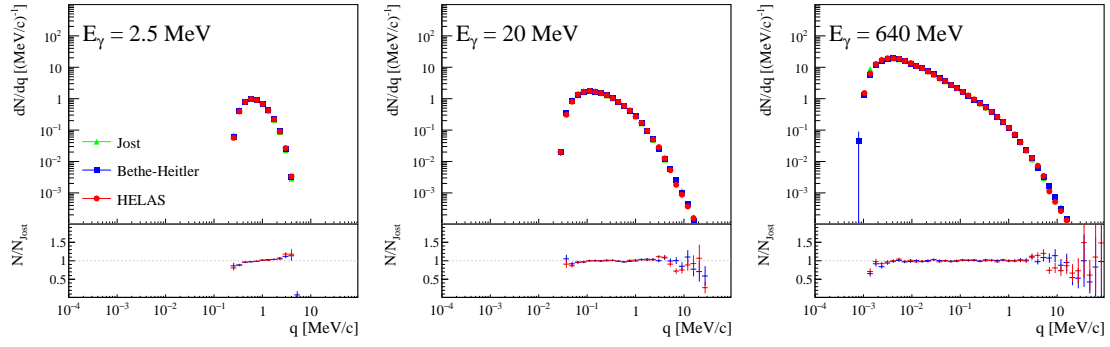
#### 4. Validations : revisited

- Validation of the value of the polarisation asymmetry, compared to the analytical expressions of the high-energy (eq. (14) of Ref.[10]) and low-energy asymptote (eq. (16) of Ref.[10]) of the polarisation asymmetry (Fig. 6).

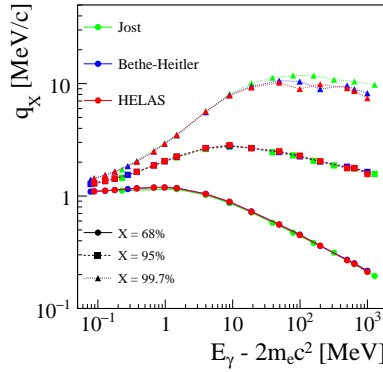


**Figure 6:** Amplitude  $A$  (left) and phase  $\phi_0$  (right) of the modulation of the distribution for various definitions of the azimuthal angle [10] (squares:  $\omega$ , upward triangles:  $\phi_r$ , downward triangles:  $\phi$ ). In addition, the performance of the 5D estimator is shown (circles). Measured value (top row), and uncertainty (middle row). The dotted lines show: (a) the asymptotic values of  $A = \pi/4$  at low energy; The dashed curve shows the high energy asymptotic expression for  $A$  (eq. (16) of Ref.[10]) (a). These results were obtained using simulated samples with  $P = 1$  and  $N = 10^5$  events each. Please note that the same quantity (example  $A$ ) is measured for various definitions of the azimuthal angle using the same event sample at a given photon energy, so that their statistical fluctuations are correlated. The error bars in plot (a) amount to  $\approx \sqrt{2/N} \approx 0.0045$  and are therefore not visible.

- Comparison of  $q_X$ , the value of the recoil momentum  $q$  at  $X\%$  containment, for three values of  $X$  (Fig. 8).
- Comparison of the distributions of  $x_+ = E_+/E_\gamma$  and  $\theta_{+-}$  for the HELAS and BH models (Fig. 9).



**Figure 7:** Examples of distribution of the momentum transfer  $q$  for different energies (2.5 MeV, 20 MeV, 640 MeV), without form factor [11].



**Figure 8:** The containment value  $q_X$  as a function of the photon energy, for different value of  $X$  [11].

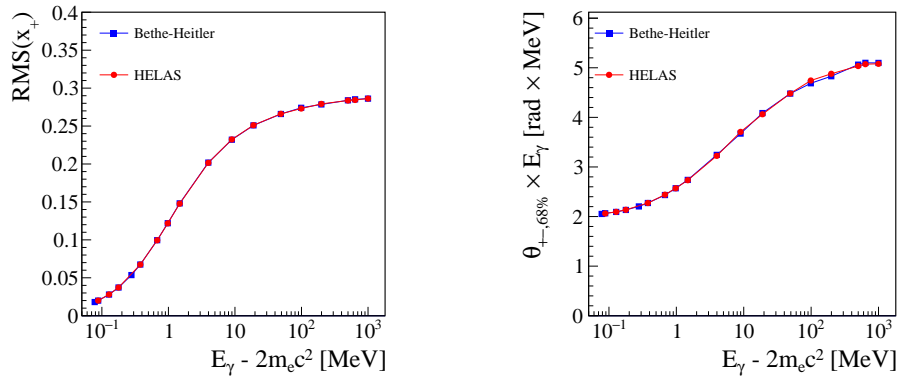
## 5. Perspectives

This event generator is using the BASES/SPRING instantiation [14] of the VEGAS method: at a given energy, for a given target nucleus (nuclear conversion) or atom (triplet conversion), after a 5D grid has been optimized for differential cross section integration precision that is for event generator efficiency, the differential cross section is tabulated and stored, something that needs several seconds of computation. Then zillions of  $\gamma$  conversions can be generated quickly at that energy and for that target.

I am now developing a VEGAS-free version of that code, that will allow the fast generation of the conversion of a photon of a given energy on a given target, with the same other properties as for the VEGAS-based generator, i.e., 5D, exact, polarized and strictly energy-momentum conserving.

## 6. Acknowledgement

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**Figure 9:** Comparison of the variables  $x$  and  $\theta_{+-}$  between BH and HELAS [11].

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