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CKM angle γ measurement at LHCb

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Data collected at the LHCb experiment have been used to determine the CKM angle γ with $B^- \to DK^-$, $D \to K_S^0 \pi^+ \pi^-$, $K_S^0 K^+ K^-$ decays, leading to the most precise measurement of γ from a single analysis. Decay-time-dependent *CP* asymmetries in $B^0 \to D^{\mp} \pi^{\pm}$ decays have been measured for the first time at a hadron collider and used to place constraints on $|\sin(2\beta + \gamma)|$ and γ that are consistent with world average values. Combining these new results with other LHCb measurements leads to the most precise determination of γ from a single experiment. Decay modes with future sensitivity to γ , such as $B_s^0 \to \overline{D}^0 K^+ K^-$ and $B_s^0 \to \overline{D}^{*0} \phi$, have been observed for the first time, and the most precise determinations of the branching fractions of $B^0 \to \overline{D}^0 K^+ K^-$ and $B_s^0 \to \overline{D}^0 \phi$ decays have been obtained.

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The angle $\gamma = arg(-(V_{ud}V_{ub}^*)/(V_{cd}V_{cb}^*))$ of the CKM unitarity triangle can be measured directly with tree-level *b*-hadron decays, using methods that exploit the weak-phase difference of γ between $b \rightarrow u$ and $b \rightarrow c$ quark transitions. The current world average value of γ , determined from direct measurements, is $\gamma = (73.5^{+4.2}_{-5.1})^{\circ}$ [1]. Theoretically, within the Standard Model, the value of γ determined in this way is very clean [2] but new physics phenomena, beyond the Standard Model, could affect its value at a level just below current experimental accuracy [3].

The LHCb detector [4] at the Large Hadron Collider is specifically designed for the study of particles containing *b* or *c* quarks. Of particular relevance for γ measurements is its high-precision tracking system, which provides a measurement of the momentum of charged particles with a relative uncertainty that varies from 0.5% at low momentum to 1.0% at 200 GeV/*c*, and a decay time resolution of approximately 50 fs [5]. The impact parameter, the minimum distance between a track and a primary proton-proton interaction vertex, is measured with a resolution of $(15+29/p_T) \mu m$, where p_T is the component of the momentum transverse to the beam, in GeV/*c*. Different types of charged hadrons, such as kaons and pions, are distinguished using information from two ring-imaging Cherenkov detectors.

1. GGSZ analysis of $B^- \rightarrow DK^-$ with $D \rightarrow K_S^0 \pi^+ \pi^-$, $K_S^0 K^+ K^-$ decays

Time-integrated measurements of γ can be made using $B^- \rightarrow DK^-$ decays¹, where *D* represents a superposition of D^0 and \overline{D}^0 mesons decaying to the same final state. Alongside γ , the related parameters r_B (the magnitude of the ratio of amplitudes of the interfering decays) and δ_B (the strong-phase difference between them) are also measured.

The GGSZ method [6] allows γ to be measured from the difference between the distributions of B^- and B^+ candidate decays across the $D \to K_S^0 h^+ h^-$ phase space ($h = \pi$ or K). The resonant structure of the multi-body D decay must therefore be taken into account. In particular, knowledge of the strong-phase difference between D^0 and \overline{D}^0 decays across the phase space is required. One possible "model-independent" approach uses direct measurements of the strong-phase difference in binned regions of phase space, obtained using quantum-correlated charm threshold data [7].

Using proton-proton collision data corresponding to an integrated luminosity of 2.0 fb⁻¹ recorded by LHCb at a centre-of-mass energy of 13 TeV, $B^- \rightarrow DK^-$ with $D \rightarrow K_S^0 \pi^+ \pi^-$, $K_S^0 K^+ K^-$ decays have been used to measure the *CP* observables $x_{\pm} = r_B \cos(\delta_B \pm \gamma)$ and $y_{\pm} = r_B \sin(\delta_B \pm \gamma)$ with the model-independent method [8],

 $\begin{aligned} x_{+} &= (-7.7 \pm 1.9 \text{ (stat.)} \pm 0.7 \text{ (syst.)} \pm 0.4 \text{ (extl.)}) \times 10^{-2}, \\ y_{+} &= (-1.0 \pm 1.9 \text{ (stat.)} \pm 0.4 \text{ (syst.)} \pm 0.9 \text{ (extl.)}) \times 10^{-2}, \\ x_{-} &= (-9.0 \pm 1.7 \text{ (stat.)} \pm 0.7 \text{ (syst.)} \pm 0.4 \text{ (extl.)}) \times 10^{-2}, \\ y_{-} &= (-2.1 \pm 2.2 \text{ (stat.)} \pm 0.5 \text{ (syst.)} \pm 1.1 \text{ (extl.)}) \times 10^{-2}, \end{aligned}$

where the third uncertainties arise from external strong-phase difference input measurements. These results constitute the first observation of *CP* violation in $B^- \rightarrow DK^-$ with $D \rightarrow K_S^0 \pi^+ \pi^-$, $K_S^0 K^+ K^-$ decays.

¹The inclusion of charge conjugated processes is implied throughout.

Combining them with previous measurements made using LHCb data corresponding to an integrated luminosity of 3.0 fb⁻¹ recorded at centre-of-mass energies of 7 and 8 TeV [9] results in the constraints $\gamma = (80^{+10}_{-9})^{\circ}$, $r_B = 0.080 \pm 0.011$ and $\delta_B = (110 \pm 10)^{\circ}$. Two-dimensional projections of the corresponding confidence regions are shown in Fig. 1, alongside the constraints obtained using the two separate data sets, demonstrating good agreement between results.



Figure 1: Two-dimensional projections of the 68.3% and 95.5% confidence regions onto the (γ, r_B) and (γ, δ_B) planes [8]. The constraints from measurements using an integrated luminosity of 3.0 fb⁻¹ recorded at centre-of-mass energies of 7 and 8 TeV ("Run 1") and 2.0 fb⁻¹ at a centre-of-mass energy of 13 TeV ("2015 & 2016 data") are shown, along with their combination.

2. Time-dependent analysis of $B^0 \rightarrow D^{\mp} \pi^{\pm}$ decays

Measurements of *CP* violation in $B^0 \to D^{\mp} \pi^{\pm}$ decays provide information on the angles $\beta = arg(-(V_{cd}V_{cb}^*)/(V_{td}V_{tb}^*))$ and γ of the unitarity triangle. A total phase difference of $2\beta + \gamma$ arises between the amplitude of a direct $B^0 \to D^{\mp} \pi^{\pm}$ decay and a decay after oscillation to \overline{B}^0 . It is possible to measure decay-time-dependent *CP* asymmetries in $B^0 \to D^{\mp} \pi^{\pm}$ decays using the decay rates of B^0 and \overline{B}^0 mesons of known initial flavour, as a function of their decay time *t*. For an initial B^0 meson, the decay rate to final state $f = D^-\pi^+$ ($\overline{f} = D^+\pi^-$) is $\Gamma_{B^0 \to f(\overline{f})}(t) \propto e^{-\Gamma t} [1 + C_{f(\overline{f})} \cos(\Delta m t) - S_{f(\overline{f})} \sin(\Delta m t)]$, where Γ is the average B^0 decay width and Δm is the $B^0 - \overline{B}^0$ oscillation frequency. It is assumed that there is no *CP* violation in the decay, that |q/p| = 1, where q and p are the coefficients defining the heavy and light mass eigenstates of the B^0 meson system, and that $\Delta \Gamma = 0$, where $\Delta \Gamma$ is the difference in decay width between the two mass eigenstates. Under these assumptions, the *CP* asymmetries $C_{f(\overline{f})}$ and $S_{f(\overline{f})}$ are defined as

$$C_{f} = -C_{\overline{f}} = \frac{1 - r_{D\pi}^{2}}{1 + r_{D\pi}^{2}}, \quad S_{f} = -\frac{2r_{D\pi}\sin[\delta - (2\beta + \gamma)]}{1 + r_{D\pi}^{2}}, \quad S_{\overline{f}} = \frac{2r_{D\pi}\sin[\delta + (2\beta + \gamma)]}{1 + r_{D\pi}^{2}}$$

where $r_{D\pi} = |A(B^0 \to D^+\pi^-)/A(B^0 \to D^-\pi^+)|$ is the ratio of decay amplitudes and δ is the *CP*-conserving phase between them.

A measurement of the *CP* asymmetries S_f and $S_{\overline{f}}$ has been performed using proton-proton collision data collected at LHCb at centre-of-mass energies of 7 and 8 TeV and corresponding to an integrated luminosity of 3.0 fb⁻¹ [10]. In this analysis, terms $\mathscr{O}(r_{D\pi}^2)$ are neglected, due to the small value of $r_{D\pi}$; this fixes $C_f = -C_{\overline{f}} = 1$. Using external measurement constraints for Δm and Γ ,

this first measurement of the asymmetries at a hadron collider results in $S_f = 0.058 \pm 0.020$ (stat.) \pm 0.011 (syst.) and $S_{\overline{f}} = 0.038 \pm 0.020$ (stat.) ± 0.007 (syst.), which agree with, and are more precise than, previous measurements [11]. They are used to place constraints on $|\sin(2\beta + \gamma)|$ and γ , shown in Fig. 2, that are consistent with world average values.



Figure 2: 1-CL (confidence level) as a function of $|\sin(2\beta + \gamma)|$ (left) and γ (right), determined from *CP* asymmetries in $B^0 \rightarrow D^{\mp} \pi^{\pm}$ decays [10].

3. Combination of LHCb measurements

Using LHCb data, measurements of γ using time-integrated GLW [12], ADS [13], GGSZ [6] and Dalitz [14] approaches have been performed using b-hadron decays such as $B^- \rightarrow D^{(*)}K^{(*)-}$, $B^0 \to DK^{(*)0}, B^0 \to DK^+\pi^-$ and $B^- \to DK^-\pi^+\pi^+$ [15]. Time-dependent measurements [16] have also been performed with $B_s^0 \to D_s^{\pm} K^{\pm}$ and $B^0 \to D^{\pm} \pi^{\pm}$ decays. These results, including the new measurements described in Sections 1 and 2, are combined using a frequentist approach to determine γ with the highest precision from a single experiment, $\gamma = (74.0^{+5.0}_{-5.8})^{\circ}$ [17].

4. Future γ decay modes: $B^0_{(s)} \to \overline{D}^{(*)0} \phi$ and $B^0_{(s)} \to \overline{D}^0 K^+ K^-$

Proton-proton collision data collected at LHCb at centre-of-mass energies of 7 and 8 TeV and corresponding to an integrated luminosity of 3.0 fb⁻¹ have been used to measure and set limits on

the branching fractions of $B^0_{(s)} \to \overline{D}^0 K^+ K^-$, $B^0_{(s)} \to \overline{D}^0 \phi$ and $B^0_s \to \overline{D}^{*0} \phi$ decays [18]. The decays $B^0_s \to \overline{D}^0 K^+ K^-$ and $B^0_s \to \overline{D}^{*0} \phi$ are observed for the first time with measured branching fractions $\mathscr{B}(B^0_s \to \overline{D}^0 K^+ K^-) = (5.7 \pm 0.5 \text{ (stat.)} \pm 0.4 \text{ (syst.)} \pm 0.5 \text{ (norm.)}) \times 10^{-5}$ and $\mathscr{B}(B^0_s \to \overline{D}^{*0} \phi) = (3.7 \pm 0.5 \text{ (stat.)} \pm 0.3 \text{ (syst.)} \pm 0.2 \text{ (norm.)}) \times 10^{-5}$; the third uncertainties arise from the branching fractions of decay modes used for normalisation. The most precise determinations of $\mathscr{B}(B^0 \to \overline{D}^0 K^+ K^-) = (6.1 \pm 0.4 \text{ (stat.)} \pm 0.3 \text{ (syst.)} \pm 0.3 \text{ (norm.)}) \times 10^{-5}$ and $\mathscr{B}(B^0_s \to \overline{D}^0 \phi) = (3.0 \pm 0.3 \text{ (stat.)} \pm 0.2 \text{ (syst.)} \pm 0.2 \text{ (norm.)}) \times 10^{-5}$ are obtained, and an upper limit on the branching fraction of $B^0 \to \overline{D}^0 \phi$ is set, $\mathscr{B}(B^0 \to \overline{D}^0 \phi) < 2.0 \ (2.2) \times 10^{-6}$ at 90% (95%) confidence level. Figure 3 shows the $B_{(s)}^0$ candidate invariant mass distributions fitted to determine these results.

In future, with larger data samples, $B_s^0 \rightarrow D^{(*)}\phi$ will be a promising decay mode to measure γ , and a time-dependent amplitude analysis of $B_s^0 \rightarrow DK^+K^-$ decays will allow constraints to be placed on γ and the *CP*-violating phase ϕ_s of the B_s^0 meson system.



Figure 3: Fitted invariant mass distributions for $B^0_{(s)} \to \overline{D}^0 K^+ K^-$ candidates [18]. The left distribution shows all selected $B^0_{(s)} \to \overline{D}^0 K^+ K^-$ candidates. The right distribution shows the subset of candidates that have been statistically separated as $B^0_{(s)} \to \overline{D}^0 \phi$ using a fit to the $K^+ K^-$ invariant mass distribution.

5. Conclusions and prospects

Using LHCb data, the most precise measurement of γ from a single analysis has been determined with $B^- \to DK^-$ with $D \to K_S^0 \pi^+ \pi^-$, $K_S^0 K^+ K^-$ decays. The first measurement of decaytime-dependent *CP* asymmetries in $B^0 \to D^{\mp} \pi^{\pm}$ decays at a hadron collider has been used to place constraints on $|\sin(2\beta + \gamma)|$ and γ that are consistent with world average values. A combination of LHCb γ measurements, including these new results, leads to the most precise determination of γ from a single experiment. Decay modes with future sensitivity to γ , such as $B_s^0 \to \overline{D}^0 K^+ K^$ and $B_s^0 \to \overline{D}^{*0} \phi$, have been observed for the first time, and the most precise determinations of the branching fractions of $B^0 \to \overline{D}^0 K^+ K^-$ and $B_s^0 \to \overline{D}^0 \phi$ decays have been obtained. Studies of other new decay modes and analysis updates to include additional LHCb data will provide further measurements and constraints on γ in the near future.

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