

Study of Radiative and Electroweak Penguin Decays at Belle

Nanae Taniguchi*

KEK, High Energy Accelerator Research Organization, Tsukuba, Japan

E-mail: nanae@post.kek.jp

The recent studies on radiative and electroweak penguin decays by the Belle experiment are reported. These new measurements are based on the Belle full data set of $772 \times 10^6 B\bar{B}$. All the measurements of $B \rightarrow K^* \gamma$ are the most precise to date and provide the first evidence of isospin violation in $b \rightarrow s \gamma$ decay with a significance of 3.1σ . Lepton flavor dependent angular analysis of $B \rightarrow K^* \ell^+ \ell^-$ decays, where ℓ is either e or μ , and results of search for lepton flavor violating process of $B^0 \rightarrow K^{*0} \mu e$ and the rare decays of $B \rightarrow h^{(*)} \nu \bar{\nu}$, are presented.

XIV International Conference on Heavy Quarks and Leptons (HQL2018)

May 27- June 1, 2018

Yamagata Terrsa, Yamagata, Japan

*Speaker.

1. Introduction

Loop (penguin) processes are good probe for physics beyond the Standard Model (BSM), since new heavy particles could contribute to the loops. A BSM effect may be observed as a deviation from the SM values. The radiative penguin and electroweak penguin are defined as a process where a charged particle emits an external real photon and an emitted virtual photon or Z^0 that produces a pair of leptons, respectively.

2. Evidence of Isospin violation in $B \rightarrow K^* \gamma$

$B \rightarrow K^* \gamma$ is experimentally the cleanest exclusive decay of $b \rightarrow s \gamma$ process. The branching fraction is $\sim 4 \times 10^{-5}$ and it is corresponding to 12% of inclusive $B \rightarrow X_s \gamma$ decays. The prediction of branching fraction suffers from large uncertainties in form factors, while the isospin (Δ_{0+}) and direct CP asymmetries (A_{CP}) are theoretically clean observables thanks to cancellation of these uncertainties [1]. The Δ_{0+} and A_{CP} are defined as

$$\Delta_{0+} = \frac{\Gamma(B^0 \rightarrow K^{*0} \gamma) - \Gamma(B^+ \rightarrow K^{*+} \gamma)}{\Gamma(B^0 \rightarrow K^{*0} \gamma) + \Gamma(B^+ \rightarrow K^{*+} \gamma)}, \quad (2.1)$$

$$A_{CP} = \frac{\Gamma(\bar{B} \rightarrow \bar{K}^* \gamma) - \Gamma(B \rightarrow K^* \gamma)}{\Gamma(\bar{B} \rightarrow \bar{K}^* \gamma) + \Gamma(B \rightarrow K^* \gamma)} \quad (2.2)$$

where Γ denotes the partial width.

Direct CP asymmetry arises due to the interference of decay amplitudes with different weak and strong CP phases. The $b \rightarrow s \gamma$ process is also possible via annihilation diagram. However this amplitude is suppressed in the SM. Since only one major diagram contributes, A_{CP} of $B \rightarrow K^* \gamma$ is zero in the SM. If there are BSM contributions from annihilation diagram or due to the CP phase, A_{CP} could be detectable. Isospin asymmetry between the neutral and charged decays arises due to the contribution from weak annihilation diagrams. Predictions of the isospin asymmetry range from 2% to 8% with a typical uncertainty of 2% in the SM [1–4]. In the analysis of Ref. [5], $B^0 \rightarrow K^{*0} \gamma$ and $B^{*+} \rightarrow K^* \gamma$ are reconstructed, where K^* is formed from $K^+ \pi^-$, $K_S^0 \pi^0$, $K^+ \pi^0$ or $K_S^0 \pi^+$ combination. The signal yield is extracted from an unbinned maximum likelihood fitting to the beam energy constrained mass distribution (Figure 1), $M_{bc} \equiv \sqrt{(E_{\text{beam}}^*/c^2)^2 - (p_B^*/c)^2}$, where E_{beam}^* is the beam energy and p_B^* is the momentum of the B meson candidate in the c.m. frame is used. The results are

$$Br(B^0 \rightarrow K^{*0} \gamma) = (3.96 \pm 0.07 \pm 0.14) \times 10^{-5}, \quad (2.3)$$

$$Br(B^+ \rightarrow K^{*+} \gamma) = (3.76 \pm 0.10 \pm 0.12) \times 10^{-5}, \quad (2.4)$$

$$A_{CP}(B^0 \rightarrow K^{*0} \gamma) = (-1.3 \pm 1.7 \pm 0.4)\%, \quad (2.5)$$

$$A_{CP}(B^+ \rightarrow K^{*+} \gamma) = (+1.1 \pm 2.3 \pm 0.3)\%, \quad (2.6)$$

$$A_{CP}(B \rightarrow K^* \gamma) = (-0.4 \pm 1.4 \pm 0.3)\%, \quad (2.7)$$

$$\Delta_{0+} = (+6.2 \pm 1.5 \pm 0.6 \pm 1.2)\%, \quad (2.8)$$

$$\Delta A_{CP} = (+2.4 \pm 2.8 \pm 0.5)\% \quad (2.9)$$

where the first uncertainty is statistical, the second is systematic, and the third for Δ_{0+} is the uncertainty from f_{+-}/f_{00} , the ratio of the branching fraction $\Upsilon(4S) \rightarrow B^+ B^-$ to that of $\Upsilon(4S) \rightarrow B^0 \bar{B}^0$.

The first evidence of isospin violation in $b \rightarrow s\gamma$ decay is found with a significance of 3.1σ [6]. ΔA_{CP} is difference of A_{CP} between charged and neutral B mesons defined as $\Delta A_{CP} = A_{CP}(B^+ \rightarrow K^{*+}\gamma) - A_{CP}(B^0 \rightarrow K^{*0}\gamma)$. ΔA_{CP} is consistent with zero in the first measurement for $B \rightarrow K^*\gamma$. Figures 2 and 3 show a comparison with previous results and SM predictions for isospin and direct CP asymmetries, respectively. They are consistent with the previous results and prediction of the SM [1–4].

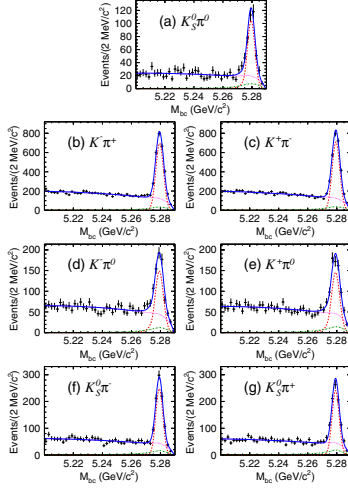


Figure 1: M_{bc} distributions for $K_S^0\pi^0$ (a), $K^-\pi^+$ (b), $K^+\pi^-$ (c), $K^-\pi^0$ (d), $K^+\pi^0$ (e), $K_S^0\pi^-$ (f) and $K_S^0\pi^+$ (g).

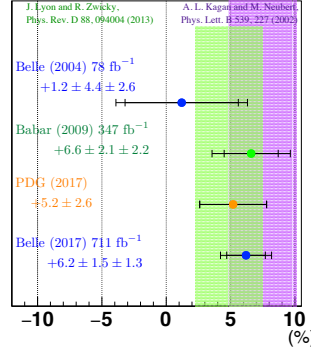


Figure 2: Comparison with previous results and SM predictions for isospin asymmetry.

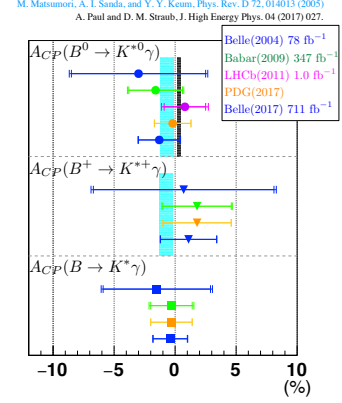


Figure 3: Comparison with previous results and SM predictions for direct CP asymmetry.

3. Lepton flavor dependent angular analysis

LHCb reported 3σ deviation in an angular observable P_5' from a full angular analysis of $B \rightarrow K^*\mu^+\mu^-$. They also reported anomaly in lepton flavor universality [7, 8]. An independent measurement is desired and lepton flavor dependence in angular analysis should also be checked. In the analysis of Ref. [9], angular observables are measured and a test of lepton flavor universality (LFU) is performed in the $B \rightarrow K^*\ell^+\ell^-$ decay, where $\ell = e, \mu$. Figure 4 shows the M_{bc} distribution of $B \rightarrow K^*\ell^+\ell^-$ candidates and P_5' and Q_5 observables for combined electron and muon modes. 2.6σ deviation is seen in P_5' of muon mode. The Q_5 observable is shown for the first time.

4. Search for lepton flavor violating $B^0 \rightarrow K^{*0}\mu e$ decays

Violation of lepton universality is accompanied by lepton flavor violation (LFV) in many models of BSM. The $B^0 \rightarrow K^{*0}\mu^\pm e^\mp$ decay is a promising place to search for LFV. In the analysis of Ref. [13], $B^0 \rightarrow K^{*0}\mu^\pm e^\mp$ is studied using the Belle full data samples, which is more than three times larger than that of the previous results by BABAR [14]. Figure 5 shows the distribution of M_{bc} for the $B^0 \rightarrow K^{*0}\mu^+e^-$, $B^0 \rightarrow K^{*0}\mu^-e^+$ data, as well as for both decays combined. No statistically

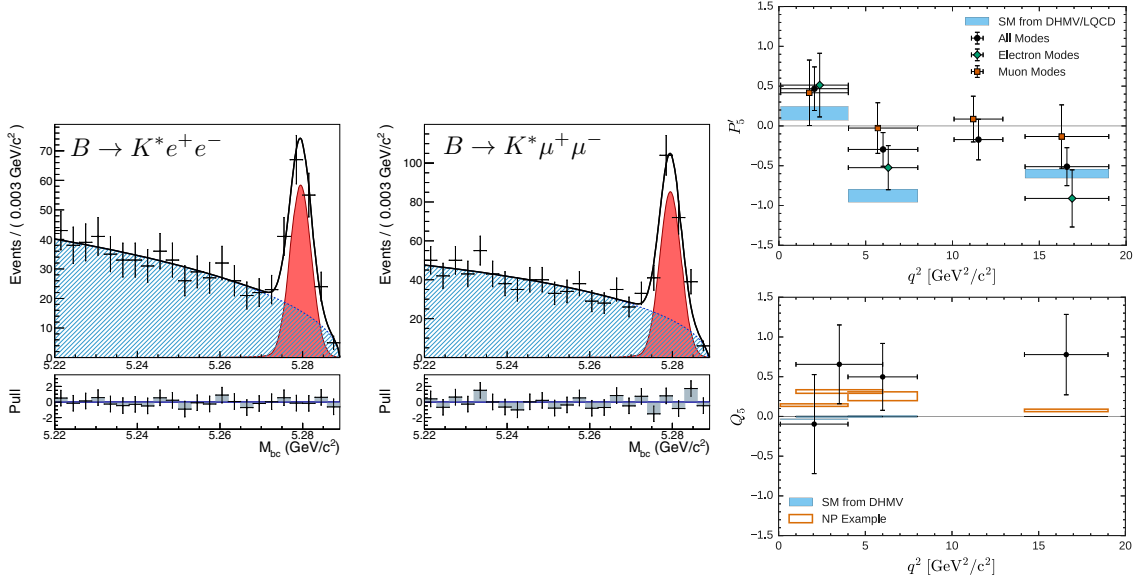


Figure 4: Distribution of the beam energy constrained mass for selected $B \rightarrow K^* e^+ e^-$ and $B \rightarrow K^* \mu^+ \mu^-$ candidates. P_5^{\pm} and Q_5^{\pm} observables for combined electron and muon modes. The SM predictions are provided by DHMC [10, 11] and lattice QCD [12] and displayed as boxes for the muon modes only. Favored NP ‘‘Scenario 1’’ from Ref. [10] for Q_5^{\pm} .

significant signals are seen and 90% confidence level upper limits on the branching fractions are

$$Br(B^0 \rightarrow K^{*0} \mu^+ e^-) < 1.2 \times 10^{-7}, \quad (4.1)$$

$$Br(B^0 \rightarrow K^{*0} \mu^- e^+) < 1.6 \times 10^{-7}, \quad (4.2)$$

$$Br(B^0 \rightarrow K^{*0} \mu^{\pm} e^{\mp}) < 1.8 \times 10^{-7}. \quad (4.3)$$

These results are the most stringent constraints on these LFV decays to date.

5. Search for $B \rightarrow h^{(*)} \nu \bar{\nu}$ decays

The decay $B \rightarrow h \nu \bar{\nu}$ can proceed only via a penguin or a box diagram at leading order in the SM and is highly suppressed. These channels are theoretically clean since they are mediated only by the Z and W bosons, in contrast to $B \rightarrow K^{(*)} \ell \ell$ decays where photon contributes, and there is no charm loop as in $b \rightarrow s \ell \ell$. This decay is also experimentally challenging since there are multiple neutrinos in the final state. Detection of B decays to $h^{(*)}$ and ‘nothing’ is required. $B \rightarrow h \nu \bar{\nu}$ decays have been studied previously by Belle with a hadronic tagging algorithm [15], and by BABAR with both hadronic [16] and semileptonic tagging [17].

In the analysis of Ref. [18], the accompanying B meson (B_{tag}) in the semileptonic decay channels $B \rightarrow D^{(*)} \ell \nu_{\ell}$ ($\ell = e, \mu$) is reconstructed. The neutral (charged) D candidates are reconstructed in 10 (7) different decay channels. Signal B daughter candidates are reconstructed through the decays $K^{*0} \rightarrow K^+ \pi^-$, $K^{*+} \rightarrow K^+ \pi^0$ and $K_S^0 \pi^+$, $\rho^+ \rightarrow \pi^+ \pi^0$, $\rho^0 \rightarrow \pi^+ \pi^-$, $K_S^0 \rightarrow \pi^+ \pi^-$, and $\pi^0 \rightarrow \gamma \gamma$. Figure 6 shows the extra energy (E_{ECL}) distributions for all eight $B \rightarrow h \nu \bar{\nu}$ channels. E_{ECL} is the

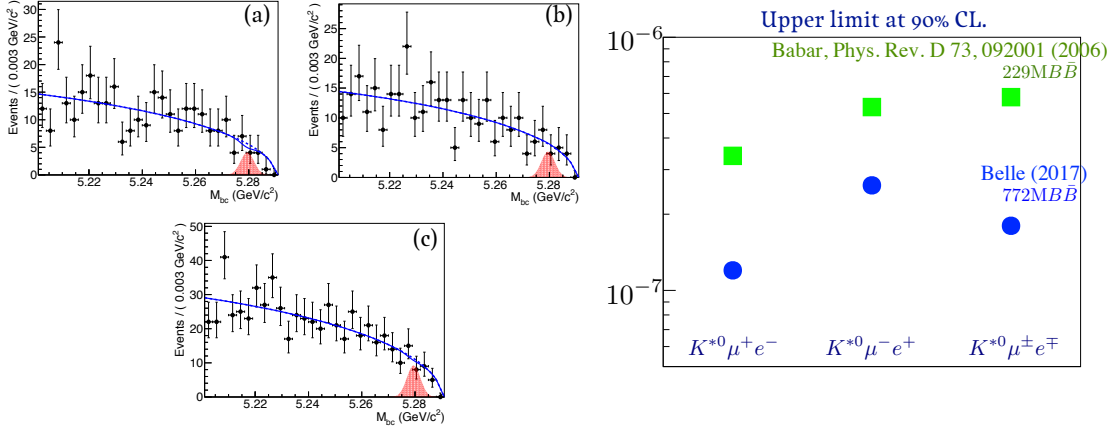


Figure 5: Distribution of the beam energy constrained mass for selected $B^0 \rightarrow K^{*0}\mu^+e^-$ (a), $B^0 \rightarrow K^{*0}\mu^-e^+$ (b), and also both decays combined (c). Dots with error bars are the data, and the blue solid curve is the result of the fit for the signal-plus-background hypothesis, where the blue dashed curve is the background component. The red shaded histogram represents the probability density function for signal with arbitrary normalization. Comparison with the previous results [14] for the upper limit at 90% confidence level.

sum of the energies of all clusters in the electromagnetic calorimeter comprised of CsI(Tl) crystals (ECL) which are not used in reconstruction of the $\Upsilon(4S)$. No statistically significant signal is seen and 90% confidence level upper limits on the branching fractions are

$$Br(B \rightarrow K\nu\bar{\nu}) < 1.6 \times 10^{-5}, \quad (5.1)$$

$$Br(B \rightarrow K^*\nu\bar{\nu}) < 2.7 \times 10^{-5}, \quad (5.2)$$

$$Br(B \rightarrow \pi\nu\bar{\nu}) < 0.8 \times 10^{-5}, \quad (5.3)$$

$$Br(B \rightarrow \rho\nu\bar{\nu}) < 2.8 \times 10^{-5}. \quad (5.4)$$

The limits on the branching fraction for the $B^0 \rightarrow K_S^0\nu\bar{\nu}$, $B^0 \rightarrow K^{*0}\nu\bar{\nu}$, $B^+ \rightarrow \pi^+\nu\bar{\nu}$, $B^0 \rightarrow \pi^0\nu\bar{\nu}$, $B^+ \rightarrow \rho^+\nu\bar{\nu}$ and $B^0 \rightarrow \rho^0\nu\bar{\nu}$ channels are the most stringent to date.

6. Summary

New measurements for radiative and electroweak penguin are performed with the Belle full data set of $772 \times 10^6 B\bar{B}$. All the measurements of $B \rightarrow K^*\gamma$ are the most precise to date and the first evidence of isospin violation in $b \rightarrow s\gamma$ decay is detected with a significance of 3.1σ . The first measurement of ΔA_{CP} for $B \rightarrow K^*\gamma$ is performed. A 2.6σ deviation is seen in P'_5 of muon mode in the lepton flavor dependent angular analysis of $B \rightarrow K^*\ell\ell$. The most stringent constraints on LFV in B decays are obtained by a search for LFV $B \rightarrow K^*\mu e$. Upper limits on $B \rightarrow h^*\nu\bar{\nu}$ channels are also obtained.

References

- [1] M. Matsumori, A. I. Sanda, and Y. Y. Keum, Phys. Rev. D **72**, 014013 (2005).

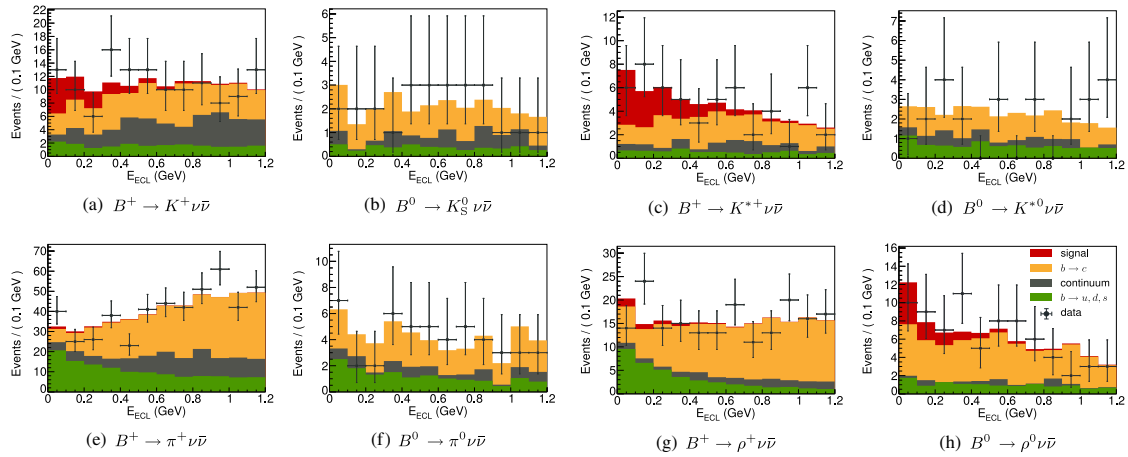


Figure 6: E_{ECL} distributions for all eight $B \rightarrow hv\bar{\nu}$ channels.

- [2] J. Lyon and R. Zwicky, Phys. Rev. D **88**, 094004 (2013).
- [3] M. Beneke, T. Feldmann, and D. Seidel, Eur. Phys. J. C **41**, 173 (2005); P. Ball, G. W. Jones, and R. Zwicky, Phys. Rev. D **75**, 054004 (2007).
- [4] A. L. Kagan and M. Neubert, Phys. Lett. B **539**, 227 (2003).
- [5] T. Horiguchi, A. Ishikawa, H. Yamamoto *et al.* (Belle Collaboration), Phys. Rev. Lett. **119**, 191802 (2017).
- [6] S. Watanuki *et al.* (Belle Collaboration), arXiv:1807.04236 [hep-ex].
- [7] R. Aaij *et al.* (LHCb Collaboration), J. High Energy Phys. 02 (2016) 104.
- [8] R. Aaij *et al.* (LHCb Collaboration), Phys. Rev. Lett. **111**, 191801 (2013).
- [9] S. Wehle, C. Niebuhr, S. Yashchenko *et al.* (Belle Collaboration), Phys. Rev. Lett. **118**, 11801 (2017).
- [10] B. Capdevila, S. Descotes-Genon, J. Matias, and J. Virto, J. High Energy Phys. 10 (2016) 075.
- [11] S. Descotes-Genon, L. Hofer, J. Matias, and J. Virto, J. High Energy Phys. 12 (2014) 125.
- [12] R. R. Horgan, Z. Liu, S. Meinel, and M. Wingate, Proc. Sci., LATTICE2014 (2015) 372.
- [13] S. Sandilya, K. Trabelsi, A. J. Schwartz *et al.* (Belle Collaboration), arXiv:1807.03267 [hep-ex].
- [14] B. Aubert *et al.* (BaBar Collaboration), Phys. Rev. D **73**, 092001 (2006).
- [15] O. Lutz, S. Neubauer, M. Heck, T. Kuhr, A. Zupanc *et al.* (Belle Collaboration), Phys. Rev. D **87**, 111103 (2013).
- [16] J. P. Lees *et al.* (BaBar Collaboration), Phys. Rev. D **87**, 112005 (2013).
- [17] J. P. Lees *et al.* (BaBar Collaboration), Phys. Rev. D **82**, 112002 (2013).
- [18] J. Grygier, P. Goldenzweig, M. Heck *et al.* (Belle Collaboration), Phys. Rev. D **96**, 091101(R) (2017).