

Higgs masses and couplings in the general 2HDM with unitarity bounds

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We investigate the general two Higgs doublet model imposing both the unitarity conditions and the bounded-from-below conditions. Both types of conditions restrict the ranges of the parameters of the scalar potential. We study the model in the Higgs basis, *i.e.* in the basis for the scalar doublets where only one doublet has vacuum expectation value. We use the experimental bounds on the oblique parameter T , to produce scalar particles with masses and cubic and quartic couplings of the Higgs in agreement with the phenomenology. The numerical calculations show that the cubic coupling may be up to 1.6 times larger than in the Standard Model, but it may also be zero or even negative. The quartic coupling is always positive and may be up to four times larger than in the Standard Model.

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1. The model

The Standard Model (SM) predicts a boson h_1 which is a scalar and it predicts its cubic and quartic couplings g_3 and g_4 , which we define through

$$\mathcal{L} = \dots - g_3 (h_1)^3 - g_4 (h_1)^4, \quad (1.1)$$

to be $g_3^{\text{SM}} \approx 32 \text{ GeV}$ and $g_4^{\text{SM}} \approx 0.032$, respectively. But Nature could be more complicated than the SM and then g_3 and g_4 might have different values. In this paper we analyse the cubic g_3 and quartic g_4 couplings of the model with two scalar gauge- $SU(2)$ doublets ϕ_1 and ϕ_2 having the same weak hypercharge. This is usually known as 2HDM. The most general scalar potential is

$$\begin{aligned} V = & \mu_1 \phi_1^\dagger \phi_1 + \mu_2 \phi_2^\dagger \phi_2 + \left(\mu_3 \phi_1^\dagger \phi_2 + \text{H.c.} \right) \\ & + \frac{\lambda_1}{2} \left(\phi_1^\dagger \phi_1 \right)^2 + \frac{\lambda_2}{2} \left(\phi_2^\dagger \phi_2 \right)^2 + \lambda_3 \phi_1^\dagger \phi_1 \phi_2^\dagger \phi_2 + \lambda_4 \phi_1^\dagger \phi_2 \phi_2^\dagger \phi_1 \\ & + \left[\frac{\lambda_5}{2} \left(\phi_1^\dagger \phi_2 \right)^2 + \lambda_6 \phi_1^\dagger \phi_1 \phi_1^\dagger \phi_2 + \lambda_7 \phi_2^\dagger \phi_2 \phi_1^\dagger \phi_2 + \text{H.c.} \right], \end{aligned} \quad (1.2)$$

where $\mu_{1,2}$ and $\lambda_{1,2,3,4}$ are real.

We use the Higgs basis for the scalar doublets where only ϕ_1^0 has VEV

$$\phi_1 = \begin{pmatrix} G^+ \\ v + (H + iG^0) / \sqrt{2} \end{pmatrix}, \quad \phi_2 = \begin{pmatrix} C^+ \\ (\sigma_1 + i\sigma_2) / \sqrt{2} \end{pmatrix}. \quad (1.3)$$

Here v is the VEV, which is real and positive, and G^+ and G^0 are (unphysical) Goldstone bosons. In the second equation, σ_1 and σ_2 are real fields and C^+ is the physical charged scalar of the 2HDM.

The unitarity constraints lead to upper bounds on the parameters of the potential. The idea of these constraints is that the scalar-scalar scattering amplitudes at tree-level must respect unitarity. For the 2HDM there are five two-particle scattering channels having different values of the electric charge Q and of the third component of weak isospin T_3 . In order to derive the unitarity conditions one must write the scattering matrices for pairs of one incoming state and one outgoing state with the same Q and T_3 . In order to satisfy those conditions, the eigenvalues of all the scattering matrices should be smaller, in modulus, than 4π . These conditions were first derived in ref. [1] and the expressions of matrices for our model are presented in ref. [2].

The potential (1.2) must be positive in all the field-space directions for large values of the fields; this means that the scalar potential has to be bounded from below (BFB), which provides stability for the potential. Necessary and sufficient conditions for the scalar potential of the 2HDM to be BFB were derived in ref. [3, 4], which we have implemented in our numerical calculations [2]. Also we apply conditions [5], which guarantee that the vacuum state has a lower value of the potential than all the other possible stability points of the potential.

We emphasize that both the unitarity conditions and the bounded-from-below conditions for the 2HDM are invariant under a change of the basis used for the two doublets. Therefore, one may implement those conditions directly in the Higgs basis.

The mass terms of the scalars are

$$V = \dots + \frac{1}{2} \begin{pmatrix} H & \sigma_1 & \sigma_2 \end{pmatrix} M \begin{pmatrix} H & \sigma_1 & \sigma_2 \end{pmatrix}^T, \quad (1.4)$$

with mass matrix

$$M = \begin{pmatrix} 2\lambda_1 v^2 & 2v^2 \Re\lambda_6 & -2v^2 \Im\lambda_6 \\ 2v^2 \Re\lambda_6 & M_C + (\lambda_4 + \Re\lambda_5)v^2 & -v^2 \Im\lambda_5 \\ -2v^2 \Im\lambda_6 & -v^2 \Im\lambda_5 & M_C + (\lambda_4 - \Re\lambda_5)v^2 \end{pmatrix} \quad (1.5a)$$

$$= R^T \times \text{diag}(M_1, M_2, M_3) \times R, \quad (1.5b)$$

where R is a 3×3 orthogonal matrix that may be parameterized as

$$R = \bar{R}_{23}(\vartheta_2) \times \bar{R}_{12}(-\vartheta_1) \times \bar{R}_{23}(-\vartheta_3), \quad (1.6)$$

by using rotation matrices \bar{R}_{12} and \bar{R}_{23} . The squared mass $M_1 = (125 \text{ GeV})^2$ and $M_2 < M_3$. We do not impose any lower limit on $M_{2,3}$; we allow them to be lower than M_1 .

According to equation (1.2) the three-Higgs vertex is given by

$$g_3 = \frac{v}{\sqrt{2}} \left[\lambda_1 x_1^3 + (\lambda_3 + \lambda_4)x_1(1 - x_1^2) + x_1(x_2^2 - x_3^2) \Re\lambda_5 - 2x_1 x_2 x_3 \Im\lambda_5 \right. \\ \left. + 3x_1^2(x_2 \Re\lambda_6 - x_3 \Im\lambda_6) + (1 - x_1^2)(x_2 \Re\lambda_7 - x_3 \Im\lambda_7) \right], \quad (1.7)$$

and the four-Higgs vertex is given by

$$g_4 = \frac{1}{8} \left[\lambda_1 x_1^4 + \lambda_2(1 - x_1^2)^2 + 2(\lambda_3 + \lambda_4)x_1^2(1 - x_1^2) + 2x_1^2(x_2^2 - x_3^2) \Re\lambda_5 \right. \\ \left. - 4x_1^2 x_2 x_3 \Im\lambda_5 + 4x_1^3(x_2 \Re\lambda_6 - x_3 \Im\lambda_6) + 4x_1(1 - x_1^2)(x_2 \Re\lambda_7 - x_3 \Im\lambda_7) \right], \quad (1.8)$$

where $x_k = R_{1k}$ for $k = 1, 2, 3$ and R is the matrix in equation (1.6).

2. Numerical analysis

The potential (1.2) of the 2HDM, in the Higgs basis, has 11 real parameters: v , λ_1 , λ_2 , λ_3 , λ_4 , $|\lambda_5|$, $|\lambda_6|$, $|\lambda_7|$, $\arg(\lambda_6^* \lambda_7)$, $\arg(\lambda_5^* \lambda_6 \lambda_7)$, and μ_2 . The other three real parameters of the potential depend on these: $\mu_1 = \lambda_1 v^2$ and $\mu_3 = \lambda_6 v^2$ (μ_3 is complex, therefore it represents two real parameters). By putting these 11 parameters into the mass matrix (1.5a) we can compute its eigenvalues $M_{2,3}$, diagonalizing matrix R , and the Higgs couplings g_3 and g_4 . A detailed description of this calculation method is presented in ref. [2]. It turns out that this method prefers to produce largish values of the masses and misses lower masses.

In order to have control on the values of the masses, in our numerical work we can use as input different quantities. Besides $v = 174 \text{ GeV}$ and $M_1 = (125 \text{ GeV})^2$, we input M_2 , M_3 , M_C , ϑ_1 , ϑ_2 , λ_2 , λ_3 , $\Re(\lambda_6 \lambda_7^*)$, and $\Re(\lambda_5^* \lambda_6 \lambda_7)$. The angle ϑ_1 is in either the first quadrant or the fourth quadrant, with a requirement that $c_1 \equiv \cos \vartheta_1 > 0.9$ so that the $h_1 W^+ W^-$ coupling is within 10% of its SM value. The angle ϑ_2 is in the first quadrant, corresponding to a choice of the signs of the fields of the new scalars h_2 and h_3 .

We compute the oblique parameter T [2, 6] and check that it is in its experimentally allowed domain [7] $-0.04 < T < 0.20$. From eq. (1.5) we then compute parameters λ_1 , λ_4 , λ_5 , λ_6 , and λ_7 . Now we have all the parameters of the model and we require them to satisfy both the unitarity conditions and the BFB conditions. Numerically diagonalizing the mass matrix (1.5a), we find the

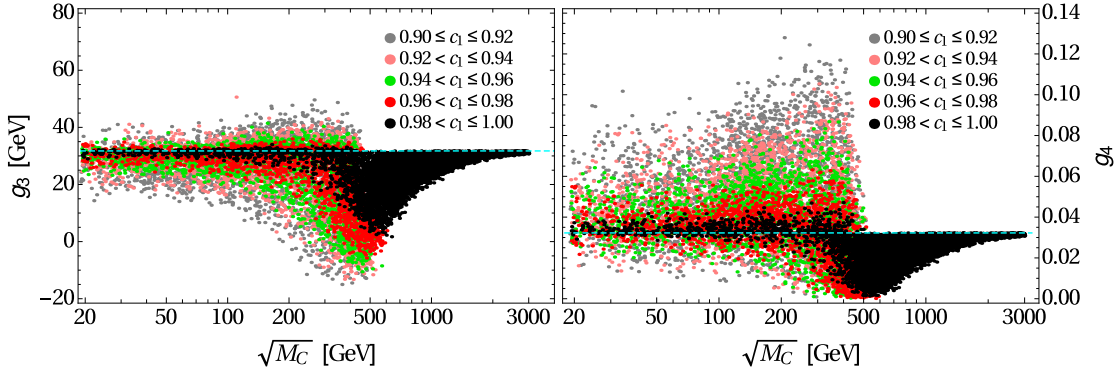


Figure 1: The three-Higgs coupling g_3 (left panel) and of the four-Higgs coupling g_4 (right panel) versus $\sqrt{M_C}$ in the 2HDM for various values of c_1 . The dashed lines mark the values of the couplings in the SM.

diagonalization matrix R and we choose the overall sign of R such that $R_{11} \equiv c_1 > 0$. Having matrix elements R_{11} , R_{12} , and R_{13} we compute the cubic (1.7) and quartic (1.8) Higgs couplings.

In our numerical work we have found that for $c_1 \lesssim 0.99$, the masses of the new scalar particles ($\sqrt{M_C}$, $\sqrt{M_2}$, and $\sqrt{M_3}$) of the 2HDM can be no larger than ~ 700 GeV. For $c_1 \lesssim 0.95$, they can be no larger than ~ 550 GeV. When c_1 becomes close to 1, the masses of all three new scalars may grow strongly to the order of TeV and in that case they become almost identical.

In figure 1 we have plotted the three- and four-Higgs couplings g_3 and g_4 against the mass of the charged scalar $\sqrt{M_C}$. One sees that g_3 may be up to 1.6 times larger than in the SM, but it may also be zero or even negative. On the other hand, g_4 is always positive because of the boundedness from below of the potential and may be up to four times larger than in the SM. For the masses of the new scalars up to 125 GeV, g_3 has values in the range $0.3 \lesssim g_3/g_3^{\text{SM}} \lesssim 1.6$ while g_4 has values in the range $0 \lesssim g_4/g_4^{\text{SM}} \lesssim 3$. For the masses up to 500 GeV the couplings reach their maximal values but when the masses of the new scalars grow beyond 1 TeV the couplings g_3 and g_4 approach their SM values.

For a more detailed description of this model and for a comparison of the couplings g_3 and g_4 in the 2HDM with couplings calculated in other extensions of the SM (Standard Model with the addition of one real singlet, SM with two real singlets, and 2HDM with the addition of one real singlet), we suggest a look at ref. [2].

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