



# **Electroweak Physics at the FCC-ee**

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The perspectives of precise measurements of electroweak observables at the Future Circular Collider with  $e^+e^-$  beams (FCC-ee) are briefly discussed.

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## 1. Introduction

The physics program realized at LEP accelerator [1] comprised a detailed test of the Standard Model of electroweak interactions which was passed with flying colours. In view of the Higgs boson discovery [2] in 2012, even more precise determination of electroweak observables has become one of the primordial goals of particle physics. The latter, together with recent advances in accelerator technologies has reinvigorated the interest in circular electron-positron colliders. Two projects of such accelerators have been proposed in the recent years: the Circular Electron Positron Collider (CEPC) [3] in China and Future Cicular Collider (FCC) [4] at CERN. The FCC design study assumes the construction and exploitation of a ~ 100 km tunnel infrastructure in the Geneva area. to host a giant, 100 TeV proton-proton collider as the flagship part of the project. However, the commissioning of the  $e^+e^-$  collider, so-called FCC-ee, is envisioned as the first step. This high-luminosity accelerator would span the range of centre-of-mass (CM) energies from below of the Z pole, through the WW and HZ threshold, till the  $t\bar{t}$  threshold and above. The FCC-ee collider will profit from recent innovations implemented at *b*-factories (nanobeam scheme, continuous injection, crab-crossing collisions etc.) and its most relevant parameters [7] are presented in Table 1 and compared with LEP phase-2 (LEP2) collider [1].

Parameter	LEP2	FCC-ee					
$E_b$ [GeV]	104	45.6	80	120	182.5		
<i>I</i> [mA]	4	1390	147	29	5.4		
# of bunches/beam	4	16640	2000	328	48		
$eta_x^*$ [m]	1.5	0.15	0.2	0.3	1.0		
$\boldsymbol{\beta}_y^*$ [mm]	50	0.8	1	1	1.6		
$\varepsilon_x$ [nm]	30-50	0.27	0.84	0.63	1.46		
$\boldsymbol{\varepsilon}_{y}$ [pm]	250	1	1.7	1.3	2.9		
$P_{\rm SR}$ [MW]	22	100	100	100	100		
$L [10^{34} \text{ cm}^{-2} \text{s}^{-1}/\text{IP}]$	0.012	230	28	8.5	1.55		
Statistics		$\sim 5 \times 10^{12} Z$	$\sim 3 \times 10^7 WW$	$\sim 10^6 ZH$	$\sim 10^6 t \bar{t}$		
(2 expts)		(4 years)	(2 years)	(3 years)	(5 years)		

**Table 1:** Selected parameters of the FCC-ee accelerator. The characteristics of LEP2 collider have been supplemented for comparison. The following parameters are presented: beam energy ( $E_b$ ), beam current (I), number of bunches, horizontal (vertical) beta function at the interaction point (IP)  $\beta_x^*(\beta_y^*)$ , horizontal (vertical) emittance ( $\varepsilon_x(\varepsilon_y)$ ), synchrotron radiation power ( $P_{SR}$ ), luminosity (L) and expected statistics of relevant events (from [6].

In the next two chapters the precision electroweak measurements around the Z pole and WW production threshold are briefly discussed.

#### 2. Electroweak precision measurements at the Z pole

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The FCC-ee offers a very attractive opportunity to go back to the electroweak physics studies at the Z pole. This time, however, the expected statistics of  $5 \times 10^{12}$  Z events (cf. Table 1),

would correspond to the increase by five orders of magnitude to compare with LEP1. As a result, the overall precision of electroweak observables will be generally limited by systematic uncertainties. Special attention should be thus paid to the reduction of them. The accuracy of CM energy calibration, realized via resonant depolarisation will be of particular importance. It would lead to the systematic error of  $\leq 100$  keV on both the mass and width of the Z boson to compare with  $\Delta M_Z = 2.1$  MeV and  $\Delta \Gamma_Z = 2.3$  MeV at LEP1. Correspondingly, the gain in precision by a factor of 20 is envisioned for the evaluation of the normalized partial widths  $R_l = \Gamma_{had} / \Gamma_{l\bar{l}}, \ l = e, \mu, \tau$ and  $R_q = \Gamma_{q\bar{q}}/\Gamma_{had}, q = b, c$  ( $\Gamma_{had}$  is the total hadronic width). The FCC-ee electroweak physics program at the Z pole envisions also new measurements of the relevant set of forward-backward and polarisation asymmetries. Among them, the  $\tau$  lepton pairs offer a unique option of measuring also the polarisation of the final state fermion through the angular distributions and momenta of the decay products. The asymmetries, together with Z partial widths, allow for an extraction of the vector  $(g_V^f)$  and axial  $(g_A^f)$  couplings of the neutral current to fermions with the expected relative precisions collected in Table 2 [8]. Finally, the effective Weinberg electroweak mixing angle, defined as:  $\sin^2 \theta_{W,\text{eff}} = 1/4(1 - g_V^f/g_A^f)$ , can be determined with the gain in the precision of 75 (100) for the tau polarisation (muon forward-backward asymmetry) measurements, respectively.

lepton	$\Delta g_V$	$\Delta g_A$	quark	$\Delta g_V$	$\Delta g_A$
е	$2.5  imes 10^{-4}$	$1.5  imes 10^{-4}$	b	$1.0  imes 10^{-2}$	$1.5 \times 10^{-3}$
$\mu$	$2.0  imes 10^{-4}$	$2.5  imes 10^{-5}$	С	$1.0  imes 10^{-2}$	$2.0  imes 10^{-3}$
τ	$3.5  imes 10^{-4}$	$0.5  imes 10^{-4}$			

Table 2: The relative uncertainties on fermion vector and axial neutral couplings to Z.

The formidable experimental precision on the  $A_{FB}^{\mu\mu}$  allows also for a precise, direct determination of the electromagnetic coupling constant at the Z mass scale  $(\alpha_{\text{QED}}(m_Z^2))$ . As demonstrated in [9], the relative uncertainty of the  $\alpha_{\text{QED}}(m_Z^2)$  is proportional to the one of  $A_{FB}^{\mu\mu}$ . By taking the data in two one-year periods of time in the optimal working points at  $\sqrt{s_-} = 87.9$  GeV and  $\sqrt{s_+} = 94.3$  GeV, the electromagnetic coupling constant at the Z mass can be determined, without the necessity of extrapolation from low mass region, with the relative precision of  $3 \times 10^{-5}$ . This value is adequate for future precision electroweak physics fits.

The FCC-ee can also address the question of the number of light neutrino species by measuring the *Z* invisible width ( $\Gamma_{inv}$ ). The first approach will be based on the determination of the hadronic cross section at the *Z* peak ( $\sigma_{had}^{peak,0}$ ):

$$N_{\nu} = \left(\frac{\Gamma_{l\bar{l}}}{\Gamma_{\nu}\bar{\nu}}\right)_{\rm SM} \cdot \left(\sqrt{\frac{12\pi R_l}{M_Z^2 \sigma_{\rm had}^{\rm peak,0}}} - R_l - 3\right)$$
(2.1)

and its accuracy is expected to be of 0.0008(stat.)  $\pm 0.001$ (syst.) (cf. the LEP1 result  $N_v = 2.991 \pm 0.007$  [10]). The second approach would exploit the radiative return process  $e^+e^- \rightarrow Z\gamma$ ,  $Z \rightarrow v\bar{v}$  measured at CM energies above the Z peak. The number of light neutrino species will de determined from the formula

$$N_{\nu} = \left(\frac{e^+e^- \to \gamma Z_{\rm inv}}{e^+e^- \to \gamma Z_{l\bar{l}}}\right)^{\rm meas} / \left(\frac{\Gamma_{\nu\bar{\nu}}}{\Gamma_{l\bar{l}}}\right)^{\rm SM}.$$
(2.2)

Due to the normalization to the number of events observed in the final state  $l\bar{l}\gamma$ , the systematic uncertainties are relatively small and the expected statistical accuracy amounts to  $\times 10^{-3}$  (LEP: 2.92 ± 0.05 [10]).

#### 3. Electroweak physics at the WW production threshold

The FCC-ee aims also at the collection of an enormous sample of WW pairs  $(3 \times 10^7 \text{ events}, \text{ at and above the production threshold})$  [8], which would allow for precise measurements of W boson properties, together with trilinear gauge couplices (TGCs) related to the WWZ and WWY vertices.

The mass and width of the W boson will be extracted from the WW cross-section scan in the close range of the production threshold. The expected precisions are  $\Delta M_W = 1$  MeV and  $\Delta \Gamma_W = 1.5$  MeV. The determination of W branching ratios to hadrons and  $l\bar{v}$ ,  $l = e, \mu, \tau$  will yield important tests of the lepton and lepton-quark universalities, with the accuracy below the permill level to compare with current precision of few percent. The limits TGCs are expected to be improved by a factor of 50.

#### 4. Summary and Acknowledgements

The FCC-ee project offers a vast physics programme, in particular in the sector of electroweak observables. Few selected topics from this field have been briefly discussed.

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