



Latest Results from the Daya Bay Reactor Neutrino Experiment

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Starting in 2011, the Daya Bay Reactor Neutrino Experiment has been observing antineutrinos from six nuclear reactors with eight identically designed underground antineutrino detectors in three experimental halls. It has accumulated the world's largest dataset of electron antineutrino candidates. The measurement of the neutrino mixing angle θ_{13} and the neutrino mass squared difference $|\Delta m_{ee}^2|$ have reached a precision of better than 4%. The large dataset allows study of a variety of topics in neutrino physics, such as absolute reactor flux and spectrum. In this poster, we present the latest results from Daya Bay on several topics.

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1. Introduction

The Daya Bay Reactor Neutrino Experiment studies antineutrinos from six reactors near Shenzhen, China, with eight identically designed antineutrino detectors. Four of the detectors are split between two near halls and the other four are placed in the far hall. The near detectors constrain the reactor antineutrino flux; the far detectors measure antineutrino disappearance. The six reactor cores in Daya Bay produce 17.4 GW_{th} power, producing an estimated 35×10^{20} neutrinos per second. Each antineutrino detector (AD) has three cylindrical zones [1], starting for the center: a 20-ton Gd-doped liquid scintillator, a 22-ton liquid scintillator and a 40-ton mineral oil buffer. The energy resolution [1] in the detectors is $8.5\%/\sqrt{E}$ (MeV). Outside the detectors, water Cherenkov detector is used to shield the ADs from natural radioactivity and to detect cosmogenic muons; together with resistive plate chambers placed on top of the detector, it produces an efficient veto of muons and muon-related backgrounds [2].

2. Oscillation Analysis

The electron antineutrinos have a distinctive inverse beta decay (IBD) signature in the Daya Bay ADs:

$$\bar{\nu}_e + p \to e^+ + n. \tag{2.1}$$

The coincidence of the prompt scintillation from the positron and the delayed neutron capture on Gd (or H) allows for powerful background rejection. For gadolinium-captured neutron events, we allow a subsequent single neutron signal with energy from 6 to 12 MeV in a time window between 1 to 200 μ s after the initial positron signal with energy from 0.7 to 12 MeV. Events with preceding cosmic muons or flashers [3] from the photomultiplier tubes are also rejected. We report our full set of cuts for the IBD selection in Ref. [3].

To further reduce uncertainties in the absolute energy calibration, an additional Flash-ADC readout system was installed in one AD, and a special calibration campaign which used different radioactive source enclosures was performed. Therefore, in the most recent oscillation analysis, this uncertainty was controlled to less than 0.5% for visible energies larger than 2 MeV. The uncertainty for the cosmogenic ⁹Li and ⁸He background was reduced from 45% to 30% for the near detectors. We also studied the spent nuclear fuel history, which improved its uncertainty from 100% to 30%. With 1958 days of data collection, more than 3.9 million antineutrino candidates were observed by all ADs, of which about 0.5 million were observed in the far site detectors. The new dataset provides a more than 60% increase of statistics from our previous publication [3].

With reduced systematic and statistical uncertainties, a χ^2 expression was defined to compare the observation in the far-hall detectors to the prediction based on near-hall detectors measurement. The systematic variation was evaluated using simulation, and statistical variation was analytically estimated; both incorporated into the covariance matrix method (described as "Method A" in Ref. [3]).

Fig. 1 shows the background-subtracted reconstructed prompt energy spectrum at the far site, with the expectations derived from near-site measurements. Fig. 2 shows the fitted 2D contours regarding the 68.3%, 95.5%, and 99.7% C.L. allowed regions. Our latest analysis yields $\sin^2 2\theta_{13} = 0.0856 \pm 0.0029$ and $\Delta m_{32}^2 = (2.471^{+0.068}_{-0.070}) \times 10^{-3} \text{eV}^2$ assuming the normal hierarchy,



Figure 1: The reconstructed spectrum at the far site after background subtraction (black points) and the expectation derived from near-site measurements excluding (red line) and including (blue line) the best-fit oscillation. The ratio of measured data over predictions with no oscillation is shown at the bottom. The top right inset shows the background components with a logarithmic scale. Figure from Ref. [4].



Figure 2: The 68.3%, 95.5%, and 99.7% C.L. allowed regions in the $\Delta m_{ee}^2 - \sin^2 2\theta_{13}$ plane. The onedimensional plots represent the $\Delta \chi^2$ for $\sin^2 2\theta_{13}$ and Δm_{ee}^2 respectively. Figure from Ref .[4]

and $\Delta m_{32}^2 = -(2.575^{+0.068}_{-0.070}) \times 10^{-3} \text{eV}^2$ assuming the inverted hierarchy. This is by far the most precise measurement of $\sin^2 2\theta_{13}$. The measurement of Δm_{32}^2 has comparable precision to that of the accelerator-based experiments [5, 6, 7].

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