# Event-plane dependent away-side jet-like correlation shape in Au + Au collisions at $\sqrt{s_{N N}}=200 \mathrm{GeV}$ from STAR 

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We employ a data-driven method to subtract the flow background of all harmonics by calculating the difference of the two-particle correlations between the close-region and far-region, determined depending on the pseudo-rapidity $(\eta)$ distance from the region where an enhanced recoil transverse momentum $\left(P_{x}\right)$ from a high- $p_{T}$ trigger particle is selected. We analyze the correlation shape as a function of the trigger particle azimuthal angle relative to the event-plane (EP) reconstructed from the beam-beam counters (BBCs) which are displaced by several units in $\eta$ from the mid-rapidity region. The large $\eta$ gap can effectively eliminate the auto-correlation between trigger particles and EP. We correct for the relatively large resolution effect from the BBC EP determination via an unfolding procedure. The width of unfolded away-side jet-like correlation increases with longer path-length, which is an indication of jet-medium interactions.

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## 1. Introduction

A strongly coupled quark gluon plasma (QGP) is believed to be created in relativistic heavyion collisions [1]. Jet-like correlations are a good probe of the energy loss mechanism of hard partons traversing the QGP medium [2, 3, 4]. They are often analyzed by calculating the azimuthal angle difference $(\Delta \phi)$ between high transverse momentum $\left(p_{T}\right)$ trigger particles and associated particles. While the near-side $(|\Delta \phi|<\pi / 2)$ correlations (in the trigger particle hemisphere) are not much modified, indicating surface bias of these correlations [2], the away-side ( $|\Delta \phi-\pi|<\pi / 2$ ) correlations recoiling from the trigger particles are significantly modified: suppressed at high $p_{T}$ and broadened at low $p_{T}[3,4,5]$. For non-central $\mathrm{Au}+\mathrm{Au}$ collisions, the in-medium path length that the recoil (away-side) parton traverses is expected to depend on its emission angle with respect to the reaction plane (RP) $[6,7]$, spanned by the impact parameter and beam directions and which is approximated by the final state event plane (EP). In these proceedings, we investigate the EP dependence of the away-side jet-like correlation shape.

## 2. Analysis Method

Measurements of jet-like correlations in heavy-ion collisions are complicated by the large underlying background [3]. A novel method to subtract all harmonic flow backgrounds without assumptions on their amplitude and shape [8] is used in this analysis. We first select events with a large recoil transverse momentum $\left(P_{x}\right)$ to a high $-p_{T}$ trigger particle to enhance the away-side jet population for a specific forward or backward pseudo-rapidity $(\eta)$ region ( $-1<\eta<-0.5$ or $0.5<\eta<1$ ). $P_{x}$ is given by

$$
\left.P_{x}\right|_{\eta_{1}} ^{\eta_{2}}=\sum_{\eta_{1}<\eta<\eta_{2},\left|\phi-\phi_{t r i g}\right|>\pi / 2} p_{T} \cos \left(\phi-\phi_{t r i g}\right) \frac{1}{\varepsilon}
$$

where all charged particles $\left(0.15<p_{T}<10 \mathrm{GeV} / c\right)$ in the opposite hemisphere of the trigger particle within a given $\eta$ range are included. We use the inverse of single-particle tracking efficiency $(\varepsilon)$ to correct for particle detection efficiency. Then two $\eta$ regions ( $-0.5<\eta<0$ and $0<\eta<0.5$ ) are defined as the close-region and far-region, respectively, depending on the distance to the $\eta$ region where the $P_{x}$ is calculated. We analyze the two-particle correlations between the trigger and associated particles in the close-region and far-region separately. The anisotropic flow contributions to these two regions are nearly equal because these two regions are symmetric about mid-rapidity. Therefore, the flow contributions to the close-region and far-region are cancelled out in the correlation difference. The away-side jet contribution to the close-region should be significantly larger than that to the far-region because of the different $\eta$ distances. The difference between the closeand far-region two-particle correlations, therefore, contains predominantly the contribution from away-side jet-like correlations, hence is a good measure of the correlation shape.

The $2^{\text {nd }}$ order harmonic EP [9] is reconstructed with the beam-beam counters (BBCs). The $\eta$ ranges of the BBCs are $3.3<|\eta|<5.2$. The trigger and associated particles are detected by the Time Projection Chamber (TPC) at mid-rapidity $(|\eta|<1)$. The large $\eta$ gap between the TPC and BBCs can effectively eliminate the auto-correlation between trigger particles and EP. The resolution of the reconstructed EP from the BBCs is calculated with the two sub-event method [9], and
is found to be $0.135 \pm 0.002$ (stat.) in $20-60 \% \mathrm{Au}+\mathrm{Au}$ collisions at $\sqrt{s_{N N}}=200 \mathrm{GeV}$. This is a measure of its accuracy in representing the true EP, and is relatively poor. Future measurement by STAR's recently installed Event Plane Detector will improve the EP resolution.

## 3. Results



Figure 1: Two-particle azimuthal correlations in the close-region (red solid circles) and far-region (blue open crosses) for different $\phi_{s}$ bins for $3<p_{T}^{\text {trig }}<10 \mathrm{GeV} / c$ and $1<p_{T}^{\text {assoc }}<2 \mathrm{GeV} / c$ in $20-60 \% \mathrm{Au}+\mathrm{Au}$ collisions at $\sqrt{s_{N N}}=200 \mathrm{GeV}$.


Figure 2: The differences between the close-region and far-region two-particle correlations in Fig. 1. Errors are statistical only. The blue curves are Gaussian fits with the mean value fixed at $\pi$.

Figure 1 shows the close- and far-region two-particle correlations in eight different $\phi_{s}$ bins with the trigger and associated particle $p_{T}$ ranges of $3<p_{T}^{\text {trig }}<10 \mathrm{GeV} / c$ and $1<p_{T}^{\text {assoc }}<2 \mathrm{GeV} / c$ in $20-60 \% \mathrm{Au}+\mathrm{Au}$ collisions at $\sqrt{s_{N N}}=200 \mathrm{GeV}$. Here $\phi_{s}$ is the trigger particle azimuthal angle relative to the reconstructed EP. The near-side correlations are well consistent in all $\phi_{s}$ bins between the close- and far-region. The ratios of the far- to close-region on the near side are approximately unity, with deviations less than $0.5 \%$ (within $2 \sigma$ statistical uncertainty). This remaining deviation is normalized out before taking the correlation difference between the close-region and far-region, shown in Fig. 2. The away-side correlations are different presumably due to away-side jet-like contributions. We use a Gaussian function (with centroid fixed at $\pi$ ) to fit the differences in Fig. 2 to extract the correlation widths. The fits are superimposed as the blue curves. The Gaussian width $(\sigma)$ increases modestly with $\phi_{s}$.

The away-side correlations in different $\phi_{s}$ bins are smeared significantly because of the poor EP resolution. We correct for this smearing effect by an unfolding procedure as follows. We take the measured trigger particle distribution in $\phi_{s}$ and the EP resolution as inputs. The true $\phi_{s}$ distribution is obtained by amplifying the Fourier modulation of the measured $\phi_{s}$ distribution by the inverse of the EP resolution factor [9]. Similarly, the distribution of azimuthal angle difference between the measured EP and true EP is evaluated by the EP resolution [9]. The probability matrix (A) is determined using Monte Carlo simulations, where the element $A_{i j}$ is the probability for the measured $\phi_{s}$ in the $j^{t h}$ bin to come from the true $\phi_{s}$ in the $i^{t h}$ bin. For each $\Delta \phi$ bin, we take the eight amplitudes of the two-particle correlations in eight $\phi_{s}$ bins (as shown in Fig. 1) as the input in the unfolding procedure. We use a least-squares method with Tikhonov regularization [10] as implemented in the TUnfold package [11]. The best value of the regularization strength $\left(\tau^{2}\right)$ is obtained via implementing the L-curve scan in TUnfoldDensity. We set the number of unfolded bins to be half of the input in our analysis. We repeat the unfolding procedure for all $\Delta \phi$ bins and obtain the unfolded correlation results. Figure 3 shows the unfolded two-particle correlations in four $\phi_{s}$ bins. The $\Delta \phi$ bins are rebinned by two to reduce the point-to-point fluctuations. It is found that the unfolded correlation shape in the out-of-plane $\left(3 \pi / 8<\phi_{s}<\pi / 2\right)$ direction is significantly different from the measured correlation shape. This is a result of the poor EP resolution.

Figure 4 shows the differences between the unfolded close- and far-region two-particle correlations. The most in-plane and out-of-plane results have greater uncertainties after unfolding. We also use a Gaussian function to fit the data points to obtain the correlation width. The fits are superimposed as the pink curves.


Figure 3: The unfolded two-particle correlations in the close-region (red solid circles) and far-region (blue open crosses) from those in Fig. 1.


Figure 4: The differences between the unfolded close-region and far-region two-particle correlations in Fig. 3. Errors are statistical only. The pink curves are Gaussian fits with the mean value fixed at $\pi$.

Figure 5 shows the comparison between the raw and unfolded away-side correlation widths
as a functions of $\phi_{s}$. The black and red lines are linear fits to the widths. The slopes of the raw and unfolded results are $0.08 \pm 0.04$ (stat.) and $0.66 \pm 0.27$ (stat.) respectively. Because the errors on the widths of the unfolded correlations are correlated among the $\phi_{s}$ bins, we estimate the statistical error on the unfolded slope as follows: (1) we randomly vary the data points in Fig. 1 using Gaussian sampling according to their statistical errors; (2) we use the same procedure to unfold the varied data points and extract a new Gaussian width after unfolding; (3) we obtain the linear slope of the new Gaussian width as a function of $\phi_{s}$; and (4) we repeat step (1) - (3) many times to obtain a distribution of the slope and take the Gaussian width of the distribution as the statistical uncertainty on the slope. As seen from Fig 5, the unfolded away-side jet-like correlation width increases with $\phi_{s}$, providing a hint of jet-medium interactions.


Figure 5: The raw (black squares) and unfolded (red crosses) away-side correlation widths ( $\sigma$ ) as a function of $\phi_{s}$. The black and red lines are corresponding linear fits.

## 4. Summary

We have applied a data-driven method to subtract flow backgrounds of all harmonics in jetlike correlations relative to high- $p_{T}$ trigger particles ( $3<p_{T}^{\text {trig }}<10 \mathrm{GeV} / c$ ) in $\mathrm{Au}+\mathrm{Au}$ collisions at $\sqrt{s_{N N}}=200 \mathrm{GeV}$. The event-plane dependence of the away-side jet-like correlation shape is reported. The $2^{\text {nd }}$ order EP is reconstructed with BBCs and the EP resolution is corrected via an unfolding procedure. The Gaussian width of the away-side jet-like correlation is found to increase with $\phi_{s}$, providing a hint of jet-medium interactions.

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