

TMD parton distributions from parton showers

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We present the determination of Transverse Momentum Dependent (TMD) parton distributions from Monte Carlo parton showers. We investigate the effective TMD distributions obtained from the PYTHIA8 and HERWIG6 parton showers and compare them to the TMD distributions determined within the Parton Branching method.

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1. Introduction

The Parton Branching (PB) method is introduced in Refs [1, 2, 3] and gives an iterative solution for the evolution of both collinear and transverse momentum dependent parton distributions [4]. The advantage of this method is that the kinematics at every splitting process can be treated exactly since the solution is fully exclusive. It allows one to determine transverse momentum dependent PDFs (TMD).

In this report the determination of TMD parton densities from general parton shower event generators as PYTHIA8 [5] and HERWIG6 [6, 7] is presented.

2. Parton Branching method

The PB method gives the evolution equation for the momentum-weighted TMD parton density $\tilde{\mathcal{A}}_a = x\mathcal{A}_a$,

$$\begin{aligned} \tilde{\mathcal{A}}_a(x, \mathbf{k}, \mu^2) = & \Delta_a(\mu^2) \tilde{\mathcal{A}}_a(x, \mathbf{k}, \mu_0^2) + \sum_b \int \frac{d^2\mathbf{q}'}{\pi\mathbf{q}'^2} \frac{\Delta_a(\mu^2)}{\Delta_a(\mathbf{q}'^2)} \Theta(\mu^2 - \mathbf{q}'^2) \Theta(\mathbf{q}'^2 - \mu_0^2) \\ & \times \int_x^{z_M} dz P_{ab}^{(R)}(\alpha_s, z) \tilde{\mathcal{A}}_b\left(\frac{x}{z}, \mathbf{k} + (1-z)\mathbf{q}', \mathbf{q}'^2\right) \end{aligned} \quad (2.1)$$

where a and b denote the flavour indices, $P_{ab}^{(R)}$ are the real emission splitting kernels, z_M is the resolution scale that separates the region for resolvable emissions from the region for non-resolvable emissions, and the Sudakov form factor $\Delta_a(\mu^2)$ gives the probability for parton a not to have resolvable branchings from scale μ_0 to scale μ . By integrating the PB evolution equation over transverse momenta, in the limit $z_M \rightarrow 1$ one recovers the DGLAP [8, 9, 10, 11] evolution equations for collinear PDFs.

The full solution of the PB evolution equation by iteration calculates the whole chain in the evolution containing the information about all partons and their momenta.

3. Obtaining TMDs from final state

In the following, we want to study how a parton shower can be used to produce effectively a transverse momentum dependent (TMD) parton density. We use Monte Carlo event generators including parton shower, determine the momentum fraction x , the transverse momentum k_\perp and the scale μ , and determine a parton distribution. It allows to visualize the TMD distributions and to compare them directly to TMD PDF sets. To perform the parton shower study a simple toy process is defined that allows to obtain TMD distributions as easily as possible. It contains two incoming partons k_1 and k_2 and a produced particle q with the momentum $k_1 + k_2 = q$. One of the incoming partons k_1 has a fixed momentum fraction $x = 0.98$ and no transverse momentum in order to make the calculation of the kinematics easier. The other parton is subject to the parton shower (or TMD distributions). The transverse momentum of k_2 can then be calculated from

$$\vec{k}_{\perp 2} = \vec{q}_\perp - \vec{k}_{\perp 1} \quad (3.1)$$

This process is not a physical one but useful to investigate the parton distribution from the parton that has no fixed momentum fraction. We label it in the following as PS2TMD.

The partonic process is generated using PYTHIA8¹ [5]. The events produced by PYTHIA are stored in the LHE format [12] which can be read by parton shower event generators. These events are analyzed with Rivet [13], the kinematics are calculated and an effective TMD is determined. Thus a TMD is obtained from standard Monte Carlo parton showers by calculating the cumulative effect of the parton shower.

To validate the method, TMDs are used instead of a parton shower (using the CASCADE package [14, 15] (version 3.0.x)) and the results are compared to the input TMD distributions.

For generating the events two different collinear and two different TMD PDF sets are used, PB-NLO-HERAI+II-2018-set1 and PB-NLO-HERAI+II-2018-set2 (obtained in Ref. [3]). They differ in the choice for the renormalization scale for the argument in α_s . For Set1 the scale is the evolution scale μ^2 and for Set2 it is the transverse momentum q_t .

As a first step, TMDs are obtained from final state events using PB-TMDs to perform a consistency check. In Fig. 1 the TMD parton densities as a function of k_\perp are shown for gluons at $x = 0.01$ and

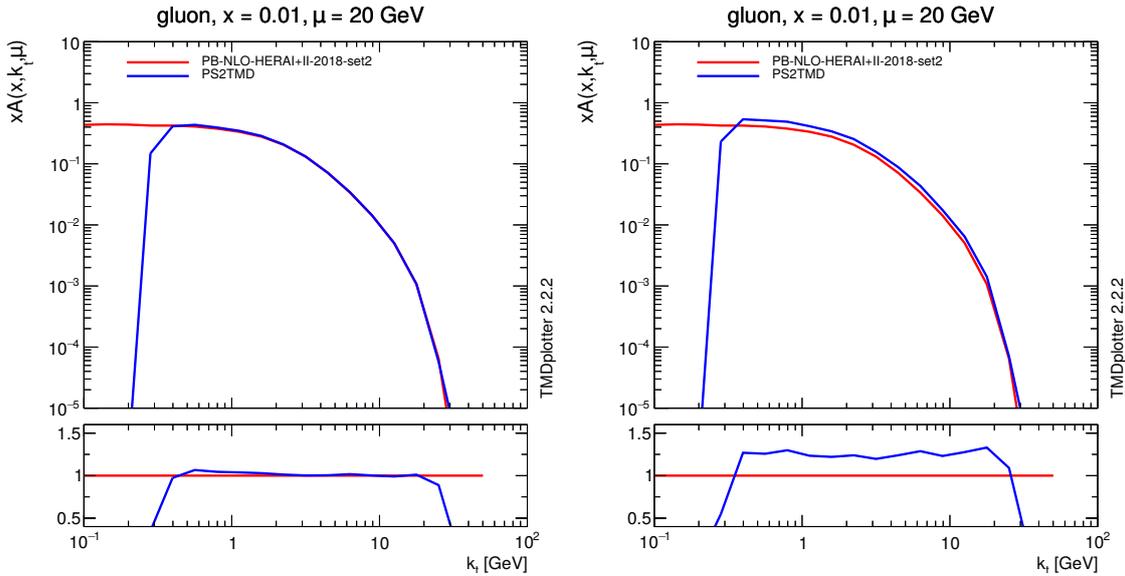


Figure 1: TMD parton density as a function of k_\perp for gluons at $\mu = 20$ GeV and $x = 0.01$ obtained from final state events (blue curve) and obtained from the input TMD PB-Set2 (red curve). Left: Collinear PDF: PB-Set2, TMD PDF: PB-Set2. Right: Collinear PDF: PB-Set1, TMD PDF: PB-Set2

$\mu = 20$ GeV. The distribution obtained from the final state of the PS2TMD process is presented (blue curve) and compared to the PB-Set2 (red curve). In Fig. 1 (left) the collinear PB-Set2 and the TMD PB-Set2 are used for the event generation. The distribution obtained from PS2TMD agrees well with the input distribution, showing that indeed one can obtain a TMD distribution by analysing the final state of a MC event.

¹Thanks to T. Sjöstrand for his help with the setup.

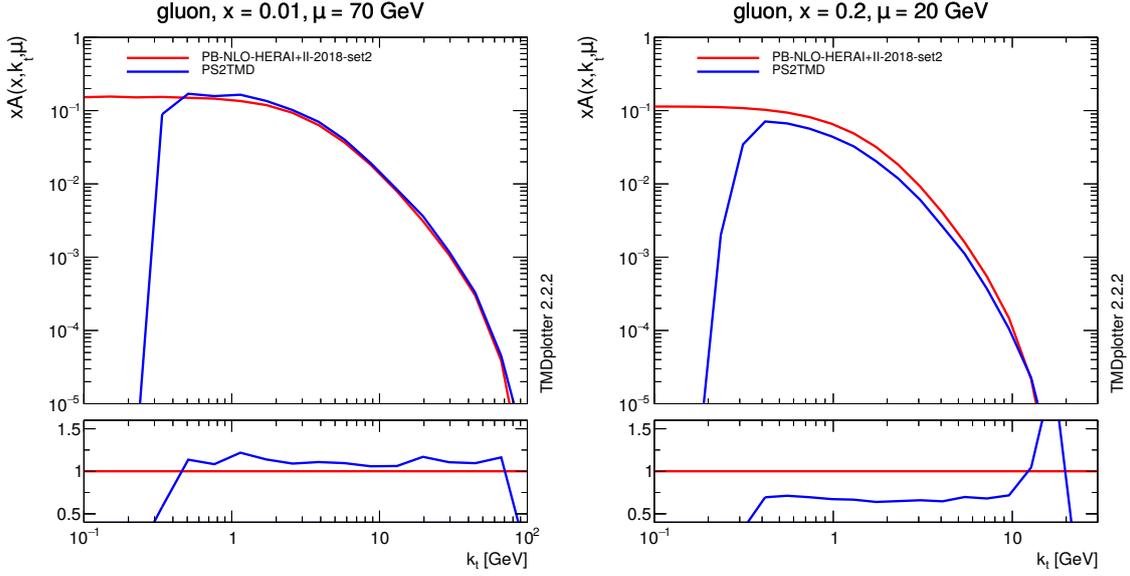


Figure 2: TMD parton density as a function of k_{\perp} for gluons at $\mu = 70$ GeV at $x = 0.01$ (left) and at $\mu = 20$ GeV at $x = 0.2$ (right) obtained from final state events (blue curve) and obtained from the input TMD PB-Set2 (red curve). Collinear PDF: PB-Set1, TMD PDF: PB-Set2

In Fig. 1 (right) the events are generated using the collinear PB-Set1 but TMD PB-Set2. In this case a significant deviation is observed. It shows that an inconsistent use of the collinear and TMD PDF set does not reproduce the input TMD distribution. The distributions for the same configuration but at larger $\mu = 70$ GeV are presented on the left in Fig. 2. At this scale, the deviation between the distributions is smaller. In Fig. 2 (right) the distributions are shown at larger $x = 0.2$ at $\mu = 20$ GeV. The difference between PB-Set2 and the obtained distribution from PS2TMD is larger. By using TMDs for the new method PS2TMD to obtain effective TMD distributions from parton showers, the concept of the method could be validated. The comparisons show that it is essential to use the collinear and TMD PDF sets consistently.

In the following, effective TMDs are determined from the PYTHIA8 shower with the Monash 2013 tune [16]. PYTHIA uses p_T -ordering for the shower. To generate the events and to produce the shower the collinear PB-Set2 is used. In Fig. 3 the PS2TMD distributions as a function of k_{\perp} obtained from the PYTHIA8 shower are shown. They are presented for gluons at $\mu = 100$ GeV at $x = 0.001$ (left) and $x = 0.01$ (right), respectively, and compared to TMD PB-Set2. A large difference to the TMD PB-Set2 is observed which is more pronounced at the higher value of $x = 0.01$. It has been shown explicitly in Ref. [1] that the p_T -ordering condition leads to flatter transverse momentum distributions. The jumps in the low- k_{\perp} region come from statistical fluctuations.

The HERWIG6 parton shower with default parameter settings is studied next. It follows angular ordering. In Fig. 4 PS2TMDs are obtained from the HERWIG6 parton shower. The parameter settings stay the same as in Fig. 3. On the left $x = 0.001$ and on the right $x = 0.01$. The collinear PB-Set2 is used for the event generation.

The distributions obtained from the HERWIG6 shower are similar to the distributions determined from the PB-Set2 for both values of x . Especially for low k_{\perp} they are close to each other, for larger

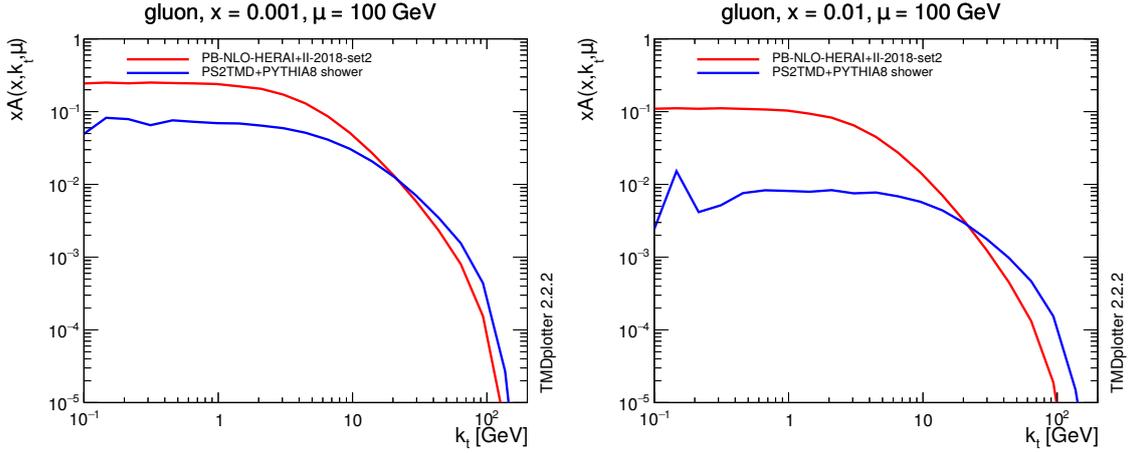


Figure 3: TMD parton density as a function of k_{\perp} for gluons at $\mu = 100$ GeV obtained from final state events of the PYTHIA shower (blue curve) and obtained from the input TMD PB-Set2 (red curve). The collinear PDF PB-Set2 is used for the generation. Left: $x = 0.001$. Right: $x = 0.01$.

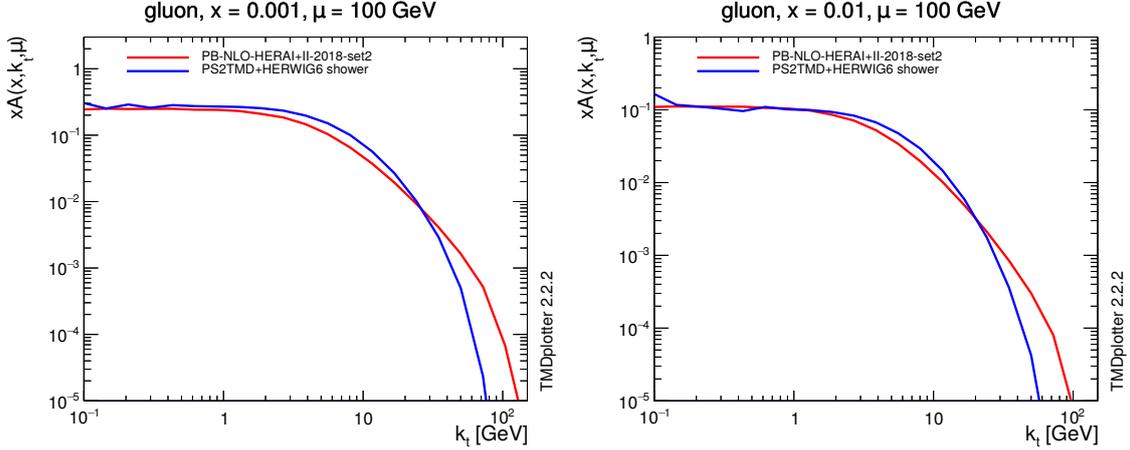


Figure 4: TMD parton density as a function of k_{\perp} for gluons at $\mu = 100$ GeV obtained from final state events of the HERWIG shower (blue curve) and obtained from the input TMD PB-Set2 (red curve). The collinear PDF PB-Set2 is used for the generation. Left: $x = 0.001$. Right: $x = 0.01$.

k_{\perp} they differ.

4. Conclusion

A method was described to determine TMDs from final state events. The concept is proven by using TMDs, the distributions are exactly reproduced by using the collinear and TMD PDF sets consistently. The PS2TMD method is applied to the PYTHIA8 parton shower. Effective TMD distributions were obtained and differences to the PB-TMD distributions were observed coming from different ordering conditions in the PYTHIA shower and the PB method. Effective TMD distributions were also obtained from the HERWIG6 parton shower, and the distributions are closer to the PB-TMD distributions.

References

- [1] F. Hautmann, H. Jung, A. Lelek, V. Radescu, and R. Zlebcik, *Phys. Lett.*, vol. B772, pp. 446-451 (2017). [1704.01757](#)
- [2] F. Hautmann, H. Jung, A. Lelek, V. Radescu, and R. Zlebcik, *JHEP*, vol. 01, p. 070 (2018). [1708.03279](#)
- [3] A. Bermudez Martinez, P. Connor, F. Hautmann, H. Jung, A. Lelek, V. Radescu, and R. Zlebcik, DESY 18-042, (2018). [1804.11152](#)
- [4] R. Angeles-Martinez *et al.*, *Acta Phys. Polon.*, vol. B46, no.12, pp. 2501-2534 (2015). [1507.05267](#)
- [5] T. Sjostrand, S. Mrenna, P. Z. Skands, *Comput. Phys. Commun.*, vol. 178, pp. 852-867 (2008). [0710.3820](#)
- [6] G. Corcella *et al.*, [1512.01178](#)
- [7] G. Marchesini *et al.*, *Comput. Phys. Commun* 67 (1992) 465-508. [1512.01178](#)
- [8] V. N. Gribov and L. N. Lipatov, *Sov. J. Nucl. Phys.* 15, 438. [*Yad. Fiz.*15,781(1972)].
- [9] L. N. Lipatov, *Sov. J. Nucl. Phys.* 20, 94. [*Yad. Fiz.*20,181(1974)].
- [10] G. Altarelli and G. Parisi, *Nucl. Phys.* B126, 298 (1977).
- [11] Y. L. Dokshitzer, *Sov. Phys. JETP* 46, 641 [*Zh. Eksp. Teor. Fiz.*73,1216(1977)]
- [12] J. Alwall *et al.*, *Comput. Phys. Commun.* 176 (2007) 300. [hep/ph0609017](#)
- [13] A. Buckley *et al.*, *Comput. Phys. Commun.* 184 (2013) 2803-2819. [1003.0694](#)
- [14] H. Jung *et al.*, *Eur. Phys. J. C*70 (2010) 1237. [1008.0152](#)
- [15] H. Jung, *Comput. Phys. Commun.* 143 (2002) 100. [hep/ph0109012](#)
- [16] P. Skands, S. Carrazza, J. Rojo, *J. Eur. Phys. J. C.* (2014) 74:3024. [1404.5630](#)