

New way to access the quark fragmentation functions in electron-positron annihilation

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The description of the polarized quark hadronization process is one of the most challenging problems in the physics of strong interactions. The various single hadron and dihadron fragmentation functions, that quantify this process, are determined by analyzing the inclusive production of hadrons in electron-positron semi-inclusive annihilation process (SIA). These, in turn, are used to extract the transverse momentum dependent parton distribution functions from the experiments on semi-inclusive deep inelastic scattering process (SIDIS), elucidating the spin-orbit correlations in nucleon.

Here we present a framework for new measurements of polarized quark fragmentation functions in electron-positron annihilation and deep inelastic semi-inclusive processes. These measurements offer a number of exciting opportunities to improve our understanding of the polarized quark hadronization and test the universality of the fragmentation functions. Such measurements can be performed at the upcoming BELLE II, JLab 12 and the future EIC experiments.

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1. Introduction

The understanding of the internal structure of the nucleon remains one of the most intriguing issues of the modern physics. At high energies, the semi-inclusive reactions with large momentum transfer are described within the framework of perturbative quantum chromodynamics (pQCD) by convolutions of the transverse momentum dependent (TMD) parton distribution functions (PDFs) and fragmentation functions (FFs). Thus, to have an access to TMD PDFs and study their quark flavor and kinematic dependences, the knowledge of FFs and their flavor dependence is needed. Recently, we have examined the modeling and the phenomenology of one and two hadron TMD FFs in the series of articles [1–7].

The measurements of two-hadron azimuthal asymmetries in semi-inclusive process have been recently used to access the quark transversity PDF of the nucleon [8–10], using a combined analysis of two-hadron production in SIDIS and back-to-back two hadron pair creation in e^+e^- annihilation. A key role here is played by the so-called interference DiFF (IFF), that describes a correlation between the transverse spin polarization of a fragmenting quark with the relative transverse momentum of the produced hadron pair. Similarly, the helicity-dependent DiFF, G_1^\perp , describes a correlation between the longitudinal polarization of a fragmenting quark and the relative transverse momenta of the pair of hadrons. The importance of G_1^\perp lies, to a considerable extent, in its relationship to the phenomenon of longitudinal jet handedness, which was predicted 25 years ago [11], but has not yet been observed experimentally. It is also of interest because it has no analogue in single unpolarized hadron production.

In general, the expressions for the cross-sections of SIDIS and SIA contain many terms with spin (in)dependent azimuthal modulations. In the leading order of pQCD formalism these terms contain complicated transverse momentum convolutions of TMD PDFs and FFs. The idea of using transverse-momentum weighting to break up the transverse momentum convolutions and to single out the PDFs and FFs in the single hadron azimuthal asymmetries was first proposed a number of years ago in [12, 13] and later in [14]. The first experimental results motivated by that work have recently been published by the COMPASS collaboration [15, 16].

In this manuscript we present our recent results for the calculations of various weighted asymmetries of dihadron production in SIDIS and SIA. These weighted asymmetries allow us to disentangle the transverse momentum convolutions between the various PDFs and FFs entering the corresponding expressions, and allow us to study the various DiFFs.

2. The weighted asymmetries for 2h+2h production in SIA

The first proposed azimuthal asymmetry for measuring helicity-dependent DiFFs G_1^\perp , in back-to-back two hadron pair creation in $e^+ + e^- \rightarrow (h_1 h_2) + (\bar{h}_1 \bar{h}_2) + X$, was made over ten years ago in [17]. A subsequent experimental search for this asymmetry by BELLE collaboration did not yield a signal [18, 19], while their previous measurements for IFF signal were sizable [20]. The model calculations of G_1^\perp entering this asymmetry was recently calculated in [3], producing a result that is smaller than that for IFF calculated within the same model, but still non-negligible.

Recently, motivated by the findings in [4], we re-derived the cross section for dihadron production in e^+e^- annihilation [6], and found a number of disagreements with the previous calculations. The two most important conclusions were the resolution of the apparent inconsistencies between the definitions of IFF entering the SIDIS and SIA processes and the realization that the originally proposed azimuthal asymmetry for determining G_1^\perp in e^+e^- annihilation should vanish.

The coordinate system used in the analysis is defined in the center of mass frame of colliding e^+ and e^- by taking the \hat{z} opposite to the total 3-momentum \vec{P}_h of the hadron pair $(h_1 h_2)$, and the components of the vectors perpendicular to \hat{z} are denoted with a subscript \perp , see Fig. 1. In this frame, the perpendicular component of the virtual photon's three-momentum $\mathbf{q}_\perp = 0$. It is also useful to define a reference frame where the total momenta of both hadron pairs are collinear [17], and the components of 3-vectors perpendicular to them is denoted by T . In this frame the virtual photon has a transverse momentum component $\mathbf{q}_T = -\mathbf{P}_{h\perp}/z$, where z is the light-cone momentum fraction of P_h with respect to q . The difference between the T and \perp components of the observed hadron momenta is of the order $1/Q$ and will be neglected here, where $Q^2 = q^2$ is much larger than the typical hadronic scale. Note, that the azimuthal angle of \mathbf{q}_T and \mathbf{P}_T are related as $\varphi_q = \varphi_h + \pi$.

The cross section expression for this process at leading twist was originally derived in [17]. The new derivation, in [6], corrected several errors in the previous result. The part of the cross section giving access to $G_1^{\perp a}$ is given by

$$\frac{d\sigma_{U,L}^{e^+e^- \rightarrow (h_1 h_2)(\bar{h}_1 \bar{h}_2)X}}{d^2\mathbf{q}_T d\varphi_R d\varphi_{\bar{R}} d^7V} = C \sum_{a,\bar{a}} e_a^2 z^2 \bar{z}^2 A(y) \left\{ \mathcal{F} \left[D_1^a \bar{D}_1^{\bar{a}} \right] - \mathcal{F} \left[\frac{(\mathbf{R}_T \times \mathbf{k}_T)_3}{M_h^2} \frac{(\bar{\mathbf{R}}_T \times \bar{\mathbf{k}}_T)_3}{\bar{M}_h^2} G_1^{\perp a} \bar{G}_1^{\perp \bar{a}} \right] \right\}, \quad (2.1)$$

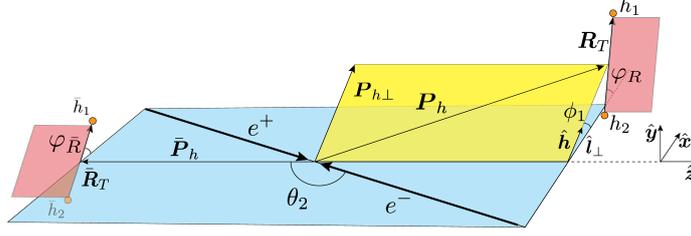


Figure 1: The kinematics of two back-to-back dihadron pair creation in e^+e^- annihilation.

where D_1^a denotes the unpolarized DiFF, $C \equiv \frac{3\alpha^2}{\pi Q^2}$, while $\varphi_R, \varphi_{\bar{R}}$ are the azimuthal angles of \mathbf{R}_T and $\bar{\mathbf{R}}_T$, and $d^7V \equiv dzd\xi dR_T d\bar{z}d\bar{\xi}d\bar{R}_T dy$ denotes the remaining phase space element. The subscript 3 denotes the z component of the vector, and M_h and \bar{M}_h are the invariant masses of the two hadron pairs. The electromagnetic coupling constant is denoted as α , while the charge of a quark of flavor a is e_a , and $A(y) = \frac{1}{2} - y + y^2$. The convolution \mathcal{F} is defined as

$$\mathcal{F}[wD^a\bar{D}^{\bar{a}}] = \int d^2\mathbf{k}_T d^2\bar{\mathbf{k}}_T \delta^2(\mathbf{k}_T + \bar{\mathbf{k}}_T - \mathbf{q}_T) w(\mathbf{k}_T, \bar{\mathbf{k}}_T, \mathbf{R}_T, \bar{\mathbf{R}}_T) D^a D^{\bar{a}}. \quad (2.2)$$

The fully unintegrated DiFFs entering these expressions depend on the relative azimuthal angles between $\mathbf{k}_T, \mathbf{R}_T$ and $\bar{\mathbf{k}}_T, \bar{\mathbf{R}}_T$, respectively: $D^a(z, \xi, \mathbf{k}_T^2, \mathbf{R}_T^2, \mathbf{k}_T \cdot \mathbf{R}_T)$, $D^{\bar{a}}(\bar{z}, \bar{\xi}, \bar{\mathbf{k}}_T^2, \bar{\mathbf{R}}_T^2, \bar{\mathbf{k}}_T \cdot \bar{\mathbf{R}}_T)$.

In order to gain information about the DiFFs entering the cross section, it is helpful to use their decomposition into Fourier cosine series with respect to the relative azimuthal angle $\varphi_{KR} \equiv \varphi_k - \varphi_R$

$$D^a(z, \xi, \mathbf{k}_T^2, \mathbf{R}_T^2, \cos(\varphi_{KR})) = \frac{1}{\pi} \sum_{n=0}^{\infty} \frac{\cos(n \cdot \varphi_{KR})}{1 + \delta^{0,n}} D^{a,[n]}(z, \xi, \mathbf{k}_T^2, \mathbf{R}_T^2), \quad (2.3)$$

The cross section, when integrated over $\mathbf{q}_T, \varphi_R, \varphi_{\bar{R}}, \xi, \bar{\xi}$, contains only the unpolarized DiFFs

$$\langle 1 \rangle = \int d\sigma_{U,L}^{e^+e^- \rightarrow (h_1 h_2)(\bar{h}_1 \bar{h}_2)X} \times 1 = CA(y) \sum_{a,\bar{a}} e_a^2 D_1^a(z, M_h^2) \bar{D}_1^{\bar{a}}(\bar{z}, \bar{M}_h^2), \quad (2.4)$$

where the integrated zeroth Fourier moment of the unpolarized DiFF is defined as

$$D_1^a(z, M_h^2) = z^2 \int d^2\mathbf{k}_T \int d\xi D_1^{a,[0]}(z, \xi, \mathbf{k}_T^2, \mathbf{R}_T^2). \quad (2.5)$$

It is easy to derive from Eq. (2.1), that the previously proposed measurement of an asymmetry containing G_1^\perp in [17] simply vanishes

$$\langle \cos(2(\varphi_R - \varphi_{\bar{R}})) \rangle = 0. \quad (2.6)$$

Moreover, it is easy to demonstrate, that $\langle f(\varphi_R, \varphi_{\bar{R}}) \rangle = 0$ for an arbitrary $f(\varphi_R, \varphi_{\bar{R}})$.

In [6] we proposed a new weight to access the helicity-dependent DiFFs

$$\left\langle \frac{q_T^2 (3 \sin(\varphi_{qR}) \sin(\varphi_{q\bar{R}}) + \cos(\varphi_{qR}) \cos(\varphi_{q\bar{R}}))}{M_h \bar{M}_h} \right\rangle = 4C \sum_{a,\bar{a}} e_a^2 (G_1^{\perp a,[0]} - G_1^{\perp a,[2]}) (\bar{G}_1^{\perp \bar{a},[0]} - \bar{G}_1^{\perp \bar{a},[2]}), \quad (2.7)$$

where the dimensionless integrated n -th moments are defined as

$$G_1^{\perp a,[n]}(z, M_h^2) \equiv z^2 \int d^2\mathbf{k}_T \int d\xi \left(\frac{k_T^2}{2M_h^2} \right) \frac{|\mathbf{R}_T|}{M_h} G_1^{\perp a,[n]}(z, \xi, \mathbf{k}_T^2, \mathbf{R}_T^2). \quad (2.8)$$

For convenience, we define the integrated helicity-dependent DiFF as

$$G_1^{\perp a}(z, M_h^2) \equiv G_1^{\perp a,[0]}(z, M_h^2) - G_1^{\perp a,[2]}(z, M_h^2). \quad (2.9)$$

Note, that this definition differs from that in [17]. Our recent model calculations of Fourier cosine moments of G_1^\perp in [3] suggest a sizable analyzing power for such combination, although in the model calculations we used $k_T R_T$ weighting when defining the Fourier cosine moments of G_1^\perp instead of $k_T^2 R_T$ used in Eq. (2.8).

The corresponding azimuthal asymmetry, which is the ratio of the weighted moment in Eq. (2.7) to the unweighted one in Eq. (2.4), can be expressed as

$$A_{e^+e^-}^{\rightarrow}(z, \bar{z}, M_h^2, \bar{M}_h^2) = 4 \frac{\sum_{a,\bar{a}} G_1^{\perp a}(z, M_h^2) G_1^{\perp \bar{a}}(\bar{z}, \bar{M}_h^2)}{\sum_{a,\bar{a}} D_1^a(z, M_h^2) \bar{D}_1^{\bar{a}}(\bar{z}, \bar{M}_h^2)}. \quad (2.10)$$

3. The weighted asymmetries for 2h production in SIDIS

We next consider the SIDIS process with two observed final state hadrons $l + N \rightarrow l' + h_1 h_2 + X$. For this process, using the the expression for the fully unintegrated cross section of this process, derived in [21], we suggested a new weighted asymmetry for accessing G_1^\perp in [6]

The cross section for this process can be decomposed into various terms according to the polarization of the incident lepton beam and the target nucleon. The two cases of interest here are σ_{UU} and σ_{UL} , describing the unpolarized and target longitudinal polarization dependent parts of the cross section, respectively. Here we only show the explicit dependence of these two terms on the relevant azimuthal angles

$$\frac{d\sigma_{UU}}{d^2\mathbf{P}_{h\perp} d\varphi_R d^6V'} = C' A'(y) \sum_a e_a^2 \mathcal{G} \left[f_1^a D_1^a \right], \quad (3.1)$$

and

$$\frac{d\sigma_{UL}}{d^2\mathbf{P}_{h\perp} d\varphi_R d^6V'} = -S_L \sum_a \frac{\alpha^2 e_a^2}{\pi y Q^2} A'(y) \mathcal{G} \left[\frac{(\mathbf{R}_T \times \mathbf{k}_T)_3}{M_h^2} g_{1L}^a G_1^{\perp a} \right], \quad (3.2)$$

where $C' = \frac{\alpha^2}{\pi y Q^2}$, $d^6V' \equiv dz d\xi dM_h^2 dx dy d\varphi_S$, φ_S is the azimuthal angle of the initial nucleon's transverse polarization \mathbf{S}_T , $A'(y) = 1 - y + y^2/2$, S_L is the longitudinal polarization of the nucleon, and the SIDIS convolution is defined as

$$\mathcal{G}[w f^a D^a] \equiv \int d^2\mathbf{p}_T \int d^2\mathbf{k}_T \delta^2\left(\mathbf{k}_T - \mathbf{p}_T + \frac{\mathbf{P}_{h\perp}}{z}\right) w(\mathbf{p}_T, \mathbf{k}_T, \mathbf{R}_T) f^a(x, p_T) D^a(z, \xi, k_T^2, R_T^2, \mathbf{k}_T \cdot \mathbf{R}_T). \quad (3.3)$$

The unweighted integral of the cross section over $\mathbf{P}_{h\perp}$, φ_R , ξ , φ_S yields a product of the unpolarized PDF and the DiFF

$$\langle 1 \rangle = \int d\sigma_{UU} \times 1 = 2C' A'(y) \sum_a e_a^2 f_1^a(x) z^2 D_1^a(z, M_h^2). \quad (3.4)$$

For the polarization-dependent term in Eq. (3.2), again by weighting this modulation by a factor of $P_{h\perp}$, we can break up the transverse momentum convolution in (3.3) into a product of two independent terms

$$\left\langle \frac{P_{h\perp} \sin(\varphi_h - \varphi_R)}{M_h} \right\rangle = \int d\sigma_{UL} \frac{P_{h\perp} \sin(\varphi_h - \varphi_R)}{M_h} = 2S_L C' A'(y) \sum_a g_{1L}^a(x) z G_1^{\perp a}(z, M_h^2), \quad (3.5)$$

where

$$g_{1L}^a(x) = \int d^2\mathbf{p}_T g_{1L}^a(x, p_T^2), \quad (3.6)$$

is the nucleon collinear helicity PDF, while the combination of the Fourier cosine moments of $G_1^{\perp a}(z, M_h^2)$ is exactly that appearing in the e^+e^- annihilation asymmetry in Eq. (2.9).

Thus, the proposed azimuthal asymmetry can be expressed as

$$A_{SIDIS}^{\vec{\alpha}}(x, z, M_h^2) = S_L \frac{\sum_a g_{1L}^a(x) z G_1^{\perp a}(z, M_h^2)}{\sum_a f_1^a(x) D_1^a(z, M_h^2)}. \quad (3.7)$$

4. The weighted asymmetries for h+2h production in SIA

Recently we proposed a new measurement in e^+e^- annihilation [7], exploiting the production of a single inclusive hadron (generically noted as Λ) back-to-back to a hadron pair, where the relevant cross section should involve convolutions of FFs for the single hadron in one jet and the DiFFs for the hadron pair produced in the opposite jet. The purpose of such a measurement is two-fold, and will leverage our knowledge of the single hadron FFs. First, the absolute cross section measurements will provide a wider basis for extracting the quark flavor dependence of the DiFFs, especially when analyzed together with the inclusive hadron pair measurements in the same jet [22]. Second, by studying various azimuthal asymmetries we can better determine the polarized DiFFs, and also access their sign.

Here we use the e^+e^- center-of-mass coordinate system, similar to that in Sec. 2, where the \hat{z} axis is chosen to point opposite to the Λ 's 3-momentum \mathbf{P}_Λ , see Fig. 2. Some of the results for weighted asymmetries from [7] are:

1. The weighted asymmetry induced by the Collins effect for unpolarized hadron Λ :

$$A^{Coll} = \frac{B(y) \sum_a e_a^2 H_1^{\langle, a \rightarrow h_1 h_2}(z, M_h^2) H_1^{\perp \bar{a}, [1]}(\vec{z})}{A(y) \sum_a e_a^2 D_1^{a \rightarrow h_1 h_2}(z, M_h^2) \bar{D}_1^{\bar{a} \rightarrow \Lambda}(\vec{z})}, \quad (4.1)$$

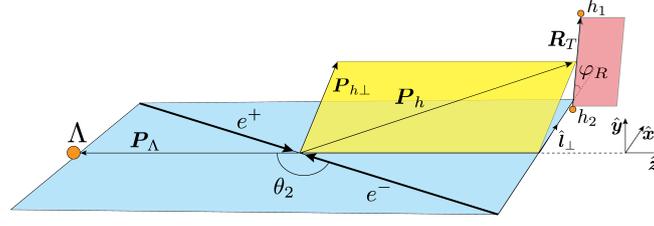


Figure 2: The kinematics of $e^+e^- \rightarrow h_1h_2 + \Lambda + X$ process.

where $B(y) = y(1-y)$.

2. The acquired longitudinal polarization s_L of Λ :

$$\langle s_L \rangle^{\sin(\varphi_q - \varphi_R)}(z, M_h^2, \bar{z}, y) = \frac{\sum_a e_a^2 G_1^{\perp, a \rightarrow h_1 h_2}(z, M_h^2) G_{1L}^{\bar{a} \rightarrow \Lambda}(\bar{z})}{\sum_a e_a^2 D_1^{a \rightarrow h_1 h_2}(z, M_h^2) \bar{D}_1^{\bar{a} \rightarrow \Lambda}(\bar{z})}, \quad (4.2)$$

3. An asymmetry induced by the correlation of the transverse polarizations of the quark and the anti quark, which couples the Collin DiFFs with the H_{1L}^\perp type FF:

$$\langle \beta_L \rangle_{H_1^\perp H_{1L}^\perp} = \left\langle \frac{q_T}{M_\Lambda} \sin(\varphi_q + \varphi_R) \right\rangle = \frac{3\alpha_{em}^2}{(2\pi)^2 Q^2} B(y) \sum_a e_a^2 H_1^{\leftarrow, a \rightarrow h_1 h_2}(z, M_h^2) H_{1L}^{\perp, \bar{a}, [1]}(\bar{z}), \quad (4.3)$$

4. The asymmetries for a transversely polarized Λ :

$$\frac{\langle s_y \rangle^{\cos(\varphi_q)}(z, M_h^2, \bar{z}, y) - \langle s_x \rangle^{\sin(\varphi_q)}(z, M_h^2, \bar{z}, y)}{M_\Lambda} = 2 \frac{\sum_a e_a^2 D_1^{a \rightarrow h_1 h_2}(z, M_h^2) D_{1T}^{\perp, \bar{a}, [1]}(\bar{z})}{\sum_a e_a^2 D_1^{a \rightarrow h_1 h_2}(z, M_h^2) \bar{D}_1^{\bar{a} \rightarrow \Lambda}(\bar{z})}, \quad (4.4)$$

$$\frac{\langle s_y \rangle^{\cos(\varphi_q)}(z, M_h^2, \bar{z}, y) + \langle s_x \rangle^{\sin(\varphi_q)}(z, M_h^2, \bar{z}, y)}{M_h} = \frac{2B(y)}{A(y)} \frac{\sum_a e_a^2 H_1^{\perp, a \rightarrow h_1 h_2}(z, M_h^2) H_1^{\bar{a} \rightarrow \Lambda}(\bar{z})}{\sum_a e_a^2 D_1^{a \rightarrow h_1 h_2}(z, M_h^2) \bar{D}_1^{\bar{a} \rightarrow \Lambda}(\bar{z})}. \quad (4.5)$$

$$\langle s_T \rangle_x^{\sin(\varphi_R)}(z, M_h^2, \bar{z}, y) = \langle s_T \rangle_y^{\cos(\varphi_R)}(z, M_h^2, \bar{z}, y) = \frac{1}{2} \frac{B(y)}{A(y)} \frac{\sum_a e_a^2 H_1^{\leftarrow, a \rightarrow h_1 h_2}(z, M_h^2) H_1^{\bar{a}}(\bar{z})}{\sum_a e_a^2 D_1^{a \rightarrow h_1 h_2}(z, M_h^2) \bar{D}_1^{\bar{a} \rightarrow \Lambda}(\bar{z})}. \quad (4.6)$$

5. Conclusions

In our recent work, we proposed a new inclusive measurement, where an unpolarized hadron pair is detected back-to-back with an another hadron pair or a single hadron, that may or may not be polarized. A number of new exciting measurements were then discussed. For example:

a) New measurements which will permit us to access the helicity-dependent DiFF, both in back-to-back two hadron pair production in e^+e^- annihilation and in forward two hadron production in SIDIS.

b) A new measurement of the unpolarized cross section will enable us to determine the flavor dependence of the DiFF using our knowledge of the ordinary unpolarized FFs, which is not possible with single hadron e^+e^- measurements. The measurement of the hadron polarization dependent asymmetries gives access to a wide variety of combinations of polarized spin 1/2 baryon FFs and polarized DiFFs.

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