

Electromagnetic properties of neutrinos

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After a brief introduction to neutrino electromagnetic properties we focus on the most important constraints on neutrino magnetic moments, charge radii and millicharges from the terrestrial experiments and astrophysical considerations. The promising new possibilities for constraining neutrino electromagnetic properties in future experiments are also discussed.

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1. Introduction

There is no doubt [1–3] that neutrino electromagnetic properties open a window to *new physics*. The most general form [1] of a neutrino electromagnetic vertex function

$$\Lambda_\mu^{ij}(q) = (\gamma_\mu - q_\mu \not{q}/q^2) \left[f_Q^{ij}(q^2) + f_A^{ij}(q^2) q^2 \gamma_5 \right] - i\sigma_{\mu\nu} q^\nu \left[f_M^{ij}(q^2) + i f_E^{ij}(q^2) \gamma_5 \right],$$

where $\Lambda_\mu(q)$ and form factors $f_{Q,A,M,E}(q^2)$ are 3×3 matrices in the space of massive neutrinos, in the case of coupling with a real photon ($q^2 = 0$) provides four sets of neutrino electromagnetic characteristics: 1) the dipole magnetic moments $\mu_{ij} = f_M^{ij}(0)$, 2) the dipole electric moments $\varepsilon_{ij} = f_E^{ij}(0)$, 3) the millicharges $q_{ij} = f_Q^{ij}(0)$ and 4) the anapole moments $a_{ij} = f_A^{ij}(0)$.

So far, there are no indications in favor of nonzero electromagnetic properties of neutrinos from either data from laboratory experiments with neutrino fluxes from ground-based sources or from astrophysics observations. However, the study of the electromagnetic properties of neutrinos attracts considerable attention.

2. Neutrino magnetic and electric dipole moments

The most well understood and studied are the dipole magnetic and electric moments. In a minimal extension of the Standard Model the diagonal magnetic moment of a Dirac neutrino is given [4] by

$$\mu_{ii}^D = \frac{3eG_F m_i}{8\sqrt{2}\pi^2} \approx 3.2 \times 10^{-19} \left(\frac{m_i}{1 \text{ eV}} \right) \mu_B, \quad (2.1)$$

μ_B is the Bohr magneton. The Majorana neutrinos can have only transition (off-diagonal) magnetic moments $\mu_{i \neq j}^M$. However, in the flavour basis the diagonal magnetic and electric moments can be nonzero even in the case of the Majorana neutrinos (for instance, when the diagonal magnetic moments for the neutrino mass states are not equal, $\mu_{ii} \neq \mu_{jj}$).

The most stringent constraints on the effective neutrino magnetic moment are obtained with the reactor antineutrinos (GEMMA Collaboration [5])

$$\mu_\nu < 2.9 \times 10^{-11} \mu_B, \quad (2.2)$$

and solar neutrinos (Borexino Collaboration [6])

$$\mu_{\nu_e} \leq 2.8 \times 10^{-11} \mu_B. \quad (2.3)$$

It should be noted, that in general in scattering experiments the neutrino is created at some distance from the detector as a flavor neutrino, which is a superposition of massive neutrinos. Therefore, the magnetic and electric moments that is measured in these experiment is not that of a massive neutrino, but it is effective moments which takes into account neutrino mixing and the oscillations during the propagation between source and detector [7, 8]. For the recent and detailed study of the neutrino electromagnetic characteristics dependence on neutrino mixing see [9].

An astrophysical bound (for both Dirac and Majorana neutrinos) is provided [10–12] by observations of the properties of globular cluster stars

$$\left(\sum_{i,j} |\mu_{ij}|^2\right)^{1/2} \leq (2.2-2.6) \times 10^{-12} \mu_B. \quad (2.4)$$

A general and termed model-independent upper bound on the Dirac neutrino magnetic moment, that can be generated by an effective theory beyond a minimal extension of the Standard Model, has been derived in [13]: $\mu_\nu \leq 10^{-14} \mu_B$. The corresponding limit for transition moments of Majorana neutrinos is much weaker [14].

In the theoretical framework with CP violation a neutrino can have nonzero electric moments ε_{ij} . In the laboratory neutrino scattering experiments for searching μ_ν (for instance, in the GEMMA experiment) the electric moment ε_{ij} contributions interfere with those due to μ_{ij} . Thus, these kind of experiments also provide constraints on ε_{ij} . The astrophysical bounds on μ_{ij} are also applicable for constraining ε_{ij} (see [10–12] and [15]).

In what follows below we give a fast flash on less know neutrino electromagnetic characteristics, namely on the neutrino millicharge, charge radius and anapole moment and give some comments on the future prospects of neutrino electromagnetic properties.

3. Neutrino electric millicharge

There are extensions of the Standard Model that allow for nonzero neutrino electric millicharges. This option can be provided by not excluded experimentally possibilities for hypercharge dequantization or another *new physics* related with an additional $U(1)$ symmetry peculiar for extended theoretical frameworks. Neutrino millicharges are strongly constrained on the level $q_\nu \sim 10^{-21} e_0$ (e_0 is the value of an electron charge) from neutrality of the hydrogen atom.

A nonzero neutrino millicharge q_ν would contribute to the neutrino electron scattering in the terrestrial experiments. Therefore, it is possible to get bounds on q_ν in the reactor antineutrino experiments. The most stringent reactor antineutrino constraint

$$q_\nu \leq 1.5 \times 10^{-11} e_0 \quad (3.1)$$

is obtained in [16] (see also [17]) with use of the GEMMA experimental data [5].

A neutrino millicharge might have specific phenomenological consequences in astrophysics because of new electromagnetic processes are opened due to a nonzero charge (see [1, 18]). Following this line, the most stringent astrophysical constraint on neutrino millicharges

$$q_\nu \leq 1.3 \times 10^{-19} e_0 \quad (3.2)$$

was obtained in [19]. This bound follows from the impact of the *neutrino star turning* mechanism ($ST\nu$) [19] that can be charged as a *new physics* phenomenon end up with a pulsar rotation frequency shift engendered by the motion of escaping from the star neutrinos along curved trajectories due to millicharge interaction with a constant magnetic field of the star.

4. Neutrino charge radius and anapole moment

Even if a neutrino millicharge is vanishing, the electric form factor $f_Q^{ij}(q^2)$ can still contain nontrivial information about neutrino electromagnetic properties. The corresponding electromagnetic characteristics is determined by the derivative of $f_Q^{ij}(q^2)$ over q^2 at $q^2 = 0$ and is termed neutrino charge radius, $\langle r_{ij}^2 \rangle = -6 \frac{df_Q^{ij}(q^2)}{dq^2} \Big|_{q^2=0}$ (see [1] for the detailed discussions). Note that for a massless neutrino the neutrino charge radius is the only electromagnetic characteristic that can have nonzero value. In the Standard Model the neutrino charge radius and the anapole moment are not defined separately, and there is a relation between these two values: $a = -\frac{\langle r^2 \rangle}{6}$.

A neutrino charge radius contributes to the neutrino scattering cross section on electrons and thus can be constrained by the corresponding laboratory experiments [20]. In all but one previous studies it was claimed that the effect of the neutrino charge radius can be included just as a shift of the vector coupling constant g_V in the weak contribution to the cross section. However, as it has been recently demonstrated in [9] within the direct calculations of the elastic neutrino-electron scattering cross section accounting for all possible neutrino electromagnetic characteristics and neutrino mixing, this is not the fact. The neutrino charge radius dependence of the cross section is more complicated and there are, in particular, the dependence on the interference terms of the type $g_V \langle r_{ij}^2 \rangle$ and also on the neutrino mixing.

5. Conclusions and future prospects

The foreseen progress in constraining neutrino electromagnetic characteristics is related, first of all, with the expected new results from the GEMMA experiment measurements of the reactor antineutrino cross section on electrons at Kalinin Power Plant. The new set of data is expected to arrive next year. The electron energy threshold will be as low as 350 eV (or even lower, up to ~ 200 eV). This will provide possibility to test the neutrino magnetic moment on the level of $\mu_V \sim 0.9 \times 10^{-12} \mu_B$ and also to test the millicharge on the level of $q_V \sim 1.8 \times 10^{-13} e_0$ [16].

The current constraints on the flavour neutrino charge radius $\langle r_{e,\mu,\tau}^2 \rangle \leq 10^{-32} - 10^{-31} \text{ cm}^2$ from the scattering experiments differ only by 1 to 2 orders of magnitude from the values $\langle r_{e,\mu,\tau}^2 \rangle \leq 10^{-33} \text{ cm}^2$ calculated within the minimally extended Standard Model with right-handed neutrinos [20]. This indicates that the minimally extended Standard Model neutrino charge radii could be experimentally tested in the near future.

Note that there is a need to re-estimate experimental constraints on $\langle r_{e,\mu,\tau}^2 \rangle$ from the scattering experiments following new derivation of the cross section [9] that properly accounts for the interference of the weak and charge radius electromagnetic interactions and also for the neutrino mixing.

Recently constraints on charged radii have been obtained [21] from the analysis of the data on coherent elastic neutrino-nucleus scattering obtained in the COHERENT experiment [22, 23]. In addition to the customary diagonal charge radii $\langle r_{e,\mu,\tau}^2 \rangle$ also the neutrino transition (off-diagonal) charge radii have been constrained in [21] for the first time:

$$\left(|\langle r_{\nu e \mu}^2 \rangle|, |\langle r_{\nu e \tau}^2 \rangle|, |\langle r_{\nu \mu \tau}^2 \rangle| \right) < (22, 38, 27) \times 10^{-32} \text{ cm}^2.$$

Quite recently the potential of current and next generation of coherent elastic neutrino-nucleus scattering experiments in probing neutrino electromagnetic interactions has been explored [24]. In particular, the present stage and the next phase of the COHERENT experiment, as well as several other reactor experiments sensitivities to the Majorana neutrino transition magnetic moments is estimated. The conclusion is that future experiments with low-threshold capabilities can improve current limits on transition magnetic moments obtained from Borexino data.

For the future progress in studying (or constraining) neutrino electromagnetic properties a rather promising claim was made in [25, 26]. It was shown that even tiny values of the Majorana neutrino transition moments would probably be tested in future high-precision experiments with the astrophysical neutrinos. In particular, observations of supernova fluxes in the JUNO experiment (see [27–29]) may reveal the effect of collective spin-flavour oscillations due to the Majorana neutrino transition moment $\mu_V^M \sim 10^{-21} \mu_B$. There are indeed other new possibilities for neutrino magnetic moment visualization in extreme astrophysical environments considered recently [30, 31].

In the most recent paper [32] we have proposed an experimental setup to observe coherent elastic neutrino-atom scattering using electron antineutrinos from tritium decay and a liquid helium target. In this scattering process with the whole atom, that has not been observed so far, the electrons tend to screen the weak charge of the nucleus as seen by the electron antineutrino probe. Finally, we study the sensitivity of this apparatus to a possible electron neutrino magnetic moment and we find that it is possible to set an upper limit of about

$$\mu_\nu < 7 \times 10^{-13} \mu_B, \quad (5.1)$$

at 90 % C.L., that is more than one order of magnitude smaller than the current experimental limits (2.2) and (2.3).

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