

1 Measurement of low-mass dileptons in ALICE

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The ALICE experiment is optimized for the study of the quark-gluon plasma (QGP), a state of matter where, due to high temperature and density, chiral symmetry is restored, and quarks and gluons are deconfined. In order to obtain information on its properties, it is particularly valuable to study lepton pair production, as leptons are produced in the plasma at all stages of its evolution and carry information about the medium properties at the time of their emission with negligible final state effects. In particular, the production of dileptons is a promising tool for the understanding of the chiral symmetry restoration and the thermodynamical properties of the QGP.

6 To differentiate possible medium contributions to the dilepton yield in nucleus-nucleus collisions from those originating in hadron decays, studies in pp and p–A collision systems are necessary to obtain a medium-free reference. These measurements can also be used to study charm and beauty production, which are an excellent test of perturbative quantum chromodynamics (pQCD) predictions. We will report on the latest results for low mass dilepton production measured with the ALICE detector in various collision systems at mid- and forward rapidity. All the measurements are compared to the expected dilepton yields from known hadronic sources and model calculations. Finally, we will discuss the prospects for future measurements.

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7 1. Introduction

8 The quark-gluon plasma (QGP) is a state of matter predicted by QCD where quarks and gluons
 9 are deconfined. It is possible to recreate the QGP with ultra-relativistic heavy-ion collisions. The
 10 production of dileptons is a promising tool for the understanding of the chiral symmetry restoration
 11 in the hot medium and the thermodynamical properties of the QGP [1]. Studies of dileptons in
 12 pp and p–Pb collisions provide reference measurements, and a way to study cold nuclear matter
 13 effects in p–Pb collisions. In the following, a selection of some of the latest dilepton measurements
 14 obtained by ALICE in pp, p–Pb, and Pb–Pb collisions are reported. A full description of the ALICE
 15 detector and performance can be found in [2].

16 There are several sources of dileptons. In the low mass range ($0 < m_{ll} < 1.1 \text{ GeV}/c^2$) the
 17 main source are Dalitz decays of pseudo-scalar and vector mesons ($\pi^0, \eta, \omega, \eta', \phi$), and 2-body
 18 decays of light-flavor vector mesons (ρ, ω, ϕ). In particular, the $\phi \rightarrow l^+l^-$ decay allows to study the
 19 strangeness production and ρ is sensitive to the chiral symmetry restoration. In the intermediate
 20 mass region ($1.1 < m_{ll} < 2.7 \text{ GeV}/c^2$), dileptons are coming mainly from the decays of correlated
 21 heavy-flavor hadrons ($c\bar{c} \rightarrow D\bar{D} \rightarrow XYl^+l^-$ and $b\bar{b} \rightarrow B\bar{B} \rightarrow XYl^+l^-$). Their yield is sensitive to
 22 the production cross section of charm and beauty quarks, $\sigma_{c\bar{c}}$ and $\sigma_{b\bar{b}}$. In addition, thermal radiation
 23 contributes to the e+e- spectrum over a broad mass range and provides insight into the temperature
 24 of the medium. However, the measurement of the thermal radiation is difficult in the intermediate
 25 mass range due to the dominant contribution from charm and beauty hadrons.

26 2. Dimuon Spectra

27 In pp collisions, new measurements of the ϕ meson cross section at forward rapidity were
 28 performed at $\sqrt{s} = 5.02, 8$ and 13 TeV [3]. The results are compared with measurements at
 29 $\sqrt{s} = 2.76$ and 7 TeV , and are shown in Figure 1 left. The differential cross section as a function of
 30 p_T shows a hardening of the p_T spectra with increasing collision energy.

31 In Pb–Pb collisions, the nuclear modification factor is used to quantify the effect of the medium
 32 on the ϕ meson production, and is defined as $R_{AA} = (1/N_{\text{coll}}) \cdot Y_{\text{PbPb}}/Y_{\text{pp}}$, with Y_{PbPb} and Y_{pp} the
 33 production yield in Pb–Pb and pp collisions and N_{coll} the number of binary nucleon-nucleon
 34 collisions. The R_{AA} of the ϕ meson measured at $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ at forward rapidity [4] is
 35 shown in Figure 1 right. A suppression is observed in central collisions, whereas the R_{AA} is
 36 compatible with unity in peripheral collisions. Results at forward and midrapidity are compatible
 37 within uncertainties, which hints towards similar mechanisms driving the interaction of the ϕ meson
 38 with the medium in the two rapidity ranges.

39 3. Dielectron spectra

40 The heavy-flavor cross section in pp collisions at midrapidity was measured recently at $\sqrt{s} =$
 41 5.02 TeV [5]. The cross section is compared with results at $\sqrt{s} = 7$ and 13 TeV [6, 7], derived using
 42 the same event generators, i.e. PYTHIA [8] or POWHEG [9]. The extracted cross sections are all in
 43 agreement with FONLL calculations [10] within uncertainties. The results are shown in Figure 2.
 44 However, a new measurement of the charm quark fragmentation fractions was performed [11],

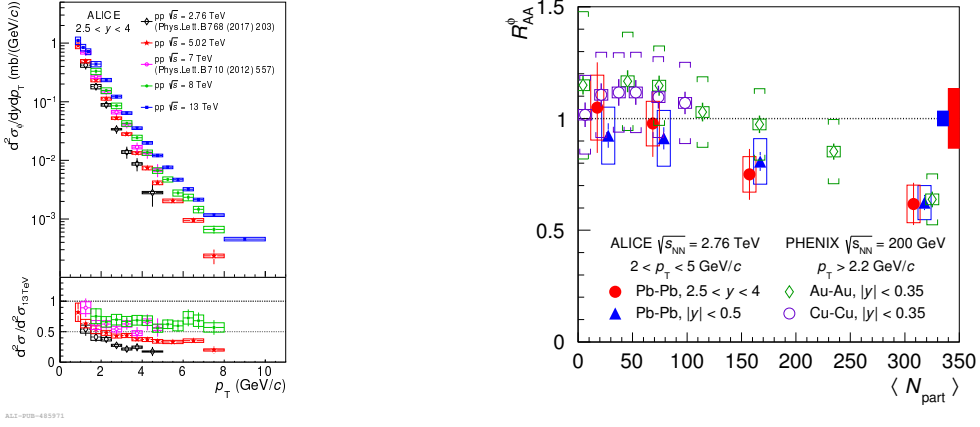


Figure 1: Left: Differential ϕ meson cross section as a function of p_T [3]. Right: R_{AA} of the ϕ meson as a function of $\langle N_{part} \rangle$ [4].

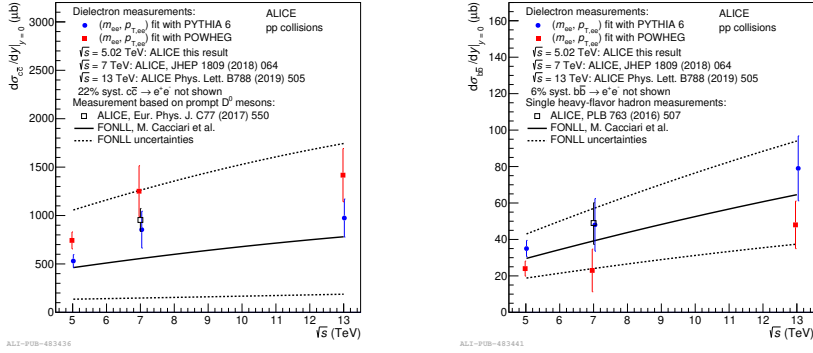


Figure 2: Cross sections at midrapidity for $c\bar{c}$ (left) and $b\bar{b}$ (right) as a function of \sqrt{s} in pp collisions [5]

45 which leads to a modification of the value of $\sigma_{c\bar{c}}$ measured via open-charm hadrons and the
 46 effective branching ratio for $c \rightarrow e$. Consequently the charm production measured via e^+e^- will be
 47 updated accordingly.

48 In p–Pb collisions, a new measurement of the $R_{p\text{-Pb}}$ was performed at $\sqrt{s_{NN}} = 5.02$ TeV [5].
 49 The nuclear modification factor is defined as $R_{p\text{-Pb}} = (1/A) \cdot (d\sigma_{ee}^{p\text{-Pb}}/dm_{ee}) / (d\sigma_{ee}^{pp}/dm_{ee})$, with
 50 $\sigma_{ee}^{p\text{-Pb}}$ and σ_{ee}^{pp} the cross section for dielectron production in p–Pb and pp collisions respectively,
 51 and A denoting the mass number of the Pb nucleus. The result is shown in Figure 3 left. For
 52 $m_{ee} < 1.1$ GeV/ c^2 , $R_{p\text{-Pb}} < 1$ whereas in the intermediate mass region, $R_{p\text{-Pb}}$ is compatible with
 53 unity within uncertainties. This suggests a different scaling behavior of the light-flavor production
 54 from binary NN collision scaling.

55 Finally, the dielectron production was measured recently at midrapidity in pp collisions at
 56 $\sqrt{s} = 13$ TeV with a reduced magnetic field [12], allowing to investigate a low m_{ee} and $p_{T,ee}$ region.
 57 The result is shown in Figure 3 right. Within the uncertainties, the dielectron cross section and
 58 the hadronic cocktail are in good agreement at $m_{ee} < m_{\pi\pi}$ while an excess over the hadronic cocktail
 59 is observed at larger masses, with a significance of 1.6σ . This excess cannot be explained with
 60 contributions from known hadronic decays.

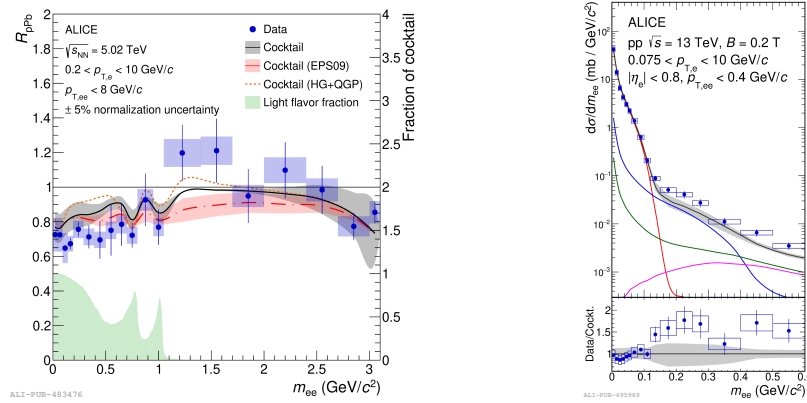


Figure 3: Left : Dielectron $R_{p\text{-Pb}}$ as a function of invariant mass at $\sqrt{s_{NN}} = 5.02$ TeV [11]. Right : Soft Dielectron cross section as a function of invariant mass [12]

61 4. Outlook for Run 3 and 4

62 In the future, dilepton studies will profit from the now installed upgrades of the ALICE
 63 detector [13]. The upgrade of the ITS and TPC will allow to improve vertex resolution and
 64 increase the readout rate in Pb – Pb collisions by a factor 100. In particular, a better separation of
 65 prompt sources from heavy-flavor electrons will be achieved via DCA studies. At the same time,
 66 the installation of the Muon Forward Tracker will add vertexing capability to the ALICE muon
 67 spectrometer and improve the dimuon mass resolution and the signal over background. Together,
 68 these upgrades will significantly increase the precision of dielectron and dimuon, as well as electron-
 69 muon measurements. Finally, runs with low magnetic fields will give access to very soft dielectron
 70 production at the LHC energies.

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