Theoretical Studies on Pion Photoproduction on Deuterons

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The study of nuclear reactions between elementary particles and atomic nuclei plays an important role in understanding the interdisciplinary area of Nuclear Physics and Particle Physics. The study of photoproduction of mesons has a long history going back to 1950’s. It was in the next decade, studies on photoproduction of π meson on deuteron started. Since then coherent and incoherent photoproduction of π meson on deuteron have been studied theoretically and experimentally. The study of photoproduction of pions describes the coupling among photon, meson and nucleon fields and also gives information about strong interactions that indirectly hold the nucleus together. A thorough investigation of the photoproduction process is firmly believed to give first hand information on two important aspects, one being the threshold of π photoproduction amplitude and the other being propagation of low-energy pions in nuclear medium. The purpose of the present contribution is to theoretically study pion photoproduction on deuterons using model independent irreducible tensor formalism developed earlier to study the photodisintegration of deuterons[1].

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1. Introduction

The study of photoproduction of pions describe the coupling among photon, meson and nucleon fields and also gives information about strong interactions that indirectly hold the nucleus together. In the recent years, the study on pion photoproduction reaction on deuterons has attracted special attention from both theorists [2,3] and experimentalists [4,5]. It allows us to study various properties which are under the influence of nuclear environment like, elementary production amplitude, pion photoproduction on/off-shell nucleon and $N – \Delta$ interaction in nuclear medium. The coherent pion photoproduction on deuteron reaction, $\gamma + d \rightarrow d + \pi^0$ is a source of information on elastic $\pi^0$ scattering but unfortunately the life time of neutral pion is very short making it impossible to study reactions like these. A thorough investigation of the photoproduction process is firmly believed to give first hand information on two important aspects, one being the threshold of $\pi^0$ photoproduction amplitude and the other being propagation of low-energy pions in nuclear medium, since neutral pion produced coherently is found to be sensitive to the pion wave function in the nuclear interior and to entire nuclear matter distribution.

On the theoretical side, the reaction $\gamma + d \rightarrow d + \pi^0$ has been studied by several authors using Impulse Approximation as early as in the 1950’s [6]. Employing the well known CGLN amplitudes [7] for pion photoproduction on nucleons, good agreement was obtained [8] with the then existing experimental data. The differential cross section leading to the different final spin states with $m = 0, \pm1$ has also been calculated and it was found that the forward cross section for $m = 0$ state predominates. Several model calculations like the Maniz Unitary Isobar Model (MAID) [9] and Scattering Analysis Interactive Dial-in (SAID) [10] partial wave analysis, Effective Lagrangian Approach model [11] have also been carried out in the intervening years.

On the Experimental side, a review on the measurement of production of mesons using photon beams is presented[12]. In this review a brief description of various experimental facilities like Jefferson lab, CLAS, ESRF, MAMI and many other are also presented. Rachek et. al., [4] have measured the tensor analyzing power $T_{20}$ in the reaction $\gamma + d \rightarrow d + \pi^0$ at the VEPP–3 storage ring at Budker Institute, Novosibirsk in the energy range $200 < E_\gamma < 500 \text{MeV}$. They have observed that below 350 MeV there is a good agreement between theory and experiment and noted that ‘The quality of agreement between theory and experiment decreases at higher photon energies, hence an improvement in theoretical model seems to be needed’. Recently, theoretical calculations on the tensor target spin asymmetries in coherent $\pi^0$– photoproduction on the deuteron were carried out [3], when their measurements were compared with the experimental data [13], there was a discrepancy for the $T_{21}$ and $T_{22}$ asymmetries.

In view of these experimental and theoretical developments, we are theoretically studying coherent pion photoproduction on deuterons, $\gamma + d \rightarrow d + \pi^0$ at near threshold energies using model independent irreducible tensor formalism developed earlier[1,14,15]. We discuss the Tensor Deuteron target spin asymmetry $A_q^2$ for the reaction $\gamma + d \rightarrow d + \pi^0$.

2. Theoretical formalism

Using the same notations [1], the reaction matrix for coherent pion photoproduction can be written in the form, $M(\mu) = \sum_{\lambda=0}^{2} (S^4(1,1) \cdot F^4(\mu))$, where, $S^4(1,1)$ of rank $\lambda; \nu = -\lambda, \cdots, \lambda$ are
defined following [14] and $\mu$ denotes photon polarization following [16]. It is important to note that the reaction can be described by only 6 irreducible tensor amplitudes $F^j_{L^j}(\mu)$ with $\lambda = 0, 1, 2$ and $\mu = \pm 1$ at all energies. It is also interesting to note that irreducible tensor amplitudes $F^j_{L^j}(\mu)$; $\nu = -\lambda, \cdots, \lambda$ can be expressed in terms of partial wave $2L$ multipole amplitudes $F^j_{L^j}$ as

$$F^j_{L^j}(\mu) = 4\pi \sqrt{2} \pi \sum_{L=1}^{\infty} \sum_{\lambda=0}^{L+1} \sum_{\mu=\lambda}^{\lambda} (-1)^{j+L-1} W(1LL; j\lambda) [j^2]_L F^j_{L^j} \left(1LL; m_t - \mu \nu \right) Y_{lm}(q).$$

where, $F^j_{L^j} = \frac{1}{2} \left[P_+ M^j_{L^j} + i \mu P_- E^j_{L^j} \right]$ and $P_+ = \frac{1}{2}[1 \pm (-1)^{L-j}]$ and $M^j_{L^j}$ and $E^j_{L^j}$ represent the magnetic and electric multipole amplitudes respectively. The unpolarised differential cross section is given by,

$$\frac{d\sigma_0}{d\Omega} = \frac{1}{6} \sum_{\lambda=0}^{2} (-1)^{j+L} \sum_{\mu} \left(F^j_{L^j}(\mu) \cdot F^{+1\lambda}_{L^j}(\mu)\right),$$

where, $F^j_{L^j}(\mu) = (-1)^{j} F^{j+\lambda}_{L^j}(\mu)$.

### 3. Results and Discussion

The spin of the Deuteron in the initial state is 1 which interacts with a photon to yield a neutral pion and a deuteron in the final state with channel spin $s = 1$. The tensor polarized target is described following [17] in terms of Fano Statistical Tensors $I^k_q$. The spin density matrix of the deuteron can then be written as, $\rho_d = \frac{1}{2} \sum_{k=0}^{2} S_k(1, 1) \cdot t^k$, with initially unpolarized photons and tensor polarized deuterons the differential cross section is given by,

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma_0}{d\Omega} \left[1 + (t^2 \cdot A^2)\right]$$

(1)

where $A^2_q$ denote the Tensor Analyzing Powers which are measurable experimentally.

$$A^2_q = \frac{\sqrt{3}}{2} \sum_{\lambda, \lambda', \mu} \left[[\lambda][\lambda']W(121; \lambda, \lambda') \left(F^j_{L^j}(\mu) \otimes F^{+1\lambda'}_{L^j}(\mu)\right)^2\right]_q$$

(2)

where $q = 0, \pm 1, \pm 2$. $A^2_q$ can explicitly be written as,

$$A^2_q = \sum_{\mu=\pm 1} \frac{1}{\sqrt{2}} \left(F^{0}_{\mu}(\mu) \otimes F^{+1\lambda}_{\mu}(\mu)\right)_q^2 - \frac{1}{\sqrt{2}} \left(F^{1}_{\mu}(\mu) \otimes F^{1\lambda}_{\mu}(\mu)\right)_q^2 - \frac{1}{\sqrt{2}} \left(F^{2}_{\mu}(\mu) \otimes F^{12}_{\mu}(\mu)\right)_q^2$$

$$+ \frac{1}{\sqrt{2}} \left(F^{2}_{\mu}(\mu) \otimes F^{+1\lambda}_{\mu}(\mu)\right)_q^2 - \frac{3}{2\sqrt{2}} \left(F^{1}_{\mu}(\mu) \otimes F^{+2}_{\mu}(\mu)\right)_q^2$$

(3)

These $A^2_q$ are expressed in terms of the irreducible tensor amplitudes. We have presented a model independent analysis of the Tensor Analyzing Powers in the $\gamma + d \rightarrow d + \pi^0$ at near threshold energies. This approach provides an alternative way to analyze the experimental measurements and is valid at both high as well as low energies. The measurements of tensor analyzing powers $A^2_q$ will help us in understanding the coherent pion photoproduction in a better way. In view of the ongoing experimental studies wide scope exists for theoretical studies associated with the reaction $\gamma + d \rightarrow d + \pi^0$. Further work on the polarization observables is in progress.
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