



# Jet measurements with ALICE: substructure, dead cone, charm jets

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A selection of recent jet measurements are presented from the ALICE experiment at the CERN LHC in proton-proton collisions at  $\sqrt{s} = 13$  TeV, focusing on substructure results for inclusive and charmed jets. The groomed jet momentum fractions  $(z_g)$  of inclusive full jets are shown for various jet resolution parameters, and  $z_g$ , the groomed splitting radius  $(\theta_g)$  as well as the number of soft drops  $(n_{SD})$  of inclusive and charmed charged-particle jets are compared. The first direct measurement of the dead cone in heavy-flavor jets is also presented. Furthermore, the parallel momentum fractions of charmed D<sup>0</sup> mesons and  $\Lambda_c^+$  baryons are shown. Besides serving as a reference for jet structure modification measurements in heavy-ion collisions, these results provide new insight to QCD parton shower properties and flavor-dependent fragmentation processes.

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# 1. Introduction

Measurement of jets in small collision systems serve as fundamental tests of pQCD and hadronization models [1]. Furthermore, these measurements provide a baseline for modification of jet production rates and structures in heavy-ion collisions by their interactions with the medium that is present in such collisions [2]. Identification of jets with heavy flavor allows for the investigation of flavor-dependent production stemming from mass and color-charge effects, and the understanding of mass-dependent fragmentation. In this contribution, recent groomed substructure measurements of inclusive and heavy-flavor jets are presented, as well as the first direct measurement of the dead cone, and the parallel momentum fraction of charmed D<sup>0</sup> mesons and  $\Lambda_c^+$  baryons. Approximately 59 pb<sup>-1</sup> integrated luminosity is used from pp collisions at  $\sqrt{s} = 13$  TeV. Charged-particle jets are reconstructed in the ALICE [3] central barrel in the central pseudorapidity region  $|\eta| < 0.9$  from tracks identified in the Time Projection Chamber and the Inner Tracking System (ITS). Full jets are reconstructed in a more limited acceptance within the pseudorapidity and azimuthal angle range  $|\eta| < 0.7$  and  $1.4 < \varphi < \pi$ , also using information from the Electromagnetic Calorimeter. Heavyflavor hadrons are fully reconstructed from their decays (in the D<sup>0</sup>/D<sup>0</sup>  $\rightarrow$ K<sup>±</sup> $\pi^{\mp}$  and  $\Lambda_c^+ \rightarrow$ pK<sup>0</sup><sub>S</sub> channels), aided by ITS based on statistical selection of tracks originating from a secondary vertex.

#### 2. Groomed substructure of inclusive jets

Measurements of groomed jet substructures allow for access to the hard parton structure of a jet, while mitigating the effects of the underlying event and hadronization [4]. Ideally, it provides a direct interface with QCD calculations. Soft-drop grooming is a novel technique that is able to remove wide-angle soft radiation (such as initial-state radiation) as well as that of the underlying event [5]. In this method, the jets that had previously been reconstructed with the anti- $k_T$  algorithm [6] are first declustered and then reclustered using the Cambridge-Aachen algorithm [7] to form a clustering tree that follows angular ordering. Then the soft branches are iteratively removed if not fulfilling the so-called soft-drop condition,

$$z > z_{\text{cut}}\theta^{\beta}$$
, where  $z = \frac{p_{\text{T},2}}{p_{\text{T},1} + p_{\text{T},2}}$  and  $\theta = \frac{\Delta R_{1,2}}{R}$  (1)

are the momentum fraction taken by the second prong ( $p_{T,1}$  and  $p_{T,2}$  being the momenta of the two prongs), and the splitting radius (defined as the ratio of the  $\Delta R_{1,2}$  splitting angle between the two prongs and the resolution parameter of the anti- $k_T$  clustering). The soft threshold is set to  $z_{cut} = 0.1$ . The angular exponent  $\beta$  is responsible for rejecting soft radiation. The jet groomed momentum fraction  $z_g$  and the groomed radius  $\theta_g$ , defined as the values of z and  $\theta$  corresponding to the first hard splitting fulfilling the soft-drop condition. Fig. 1 shows  $z_g$  (left panel) and  $\theta_g$  (right panel) for charged-particle based jets for different choices of  $\beta$ . For smaller  $\beta$  values, more soft splittings are groomed away, leading to more collimated jets. This is clearly visible in the case of  $\theta_g$  where the weight of the distribution shifts toward smaller angles with decreasing  $\beta$ .

Figure 2 shows the full-jet groomed momentum fraction  $z_g$  in the  $30 < p_T^{\text{jet}} < 40 \text{ GeV}/c$  transverse for different *R* values. The difference suggests that jets with small radii tend to split more symmetrically, while in case of larger radii there is a higher sensitivity to non-perturbative effects. Trends observed both in the *R* and the  $\beta$ -dependent groomed jet substructure results are reproduced rather well by Monte-Carlo event generators [8, 9].



**Figure 1:** Charged-particle jet  $z_g$  (left) and  $\theta_g$  (right) in pp collisions at  $\sqrt{s} = 13$  TeV for different  $\beta$  values, compared to PYTHIA 8 simulations.



**Figure 2:** Full jet  $z_g$  for different jet resolution parameter values compared to PYTHIA simulations, in pp collisions at  $\sqrt{s} = 13$  TeV.



**Figure 3:** Ratio of the angular distribution of splittings with different  $k_{\rm T}$  cuts for  $D^0$ -tagged jets over inclusive jets, shown for  $5 < E_{\rm rad} < 15$  GeV.

## 3. Structure and fragmentation of heavy-flavor jets

In gauge theories, charged particles with a mass m > 0 and energy E emit radiation that is suppressed below angles  $\theta \approx m/E$  with respect to the axis of the radiator. This so-called dead-cone effect is expected to be present in jets containing heavy flavor [10, 11]. The ALICE collaboration presented the first direct measurement of the dead cone in heavy-flavor jets, following the iterative declustering method proposed in Ref. [12]. A cut on the relative transverse momentum fraction of the splitting,  $k_{\rm T}$ , is applied to remove non-perturbative effects. Fig. 3 shows the ratio of the angular distribution of splittings for D<sup>0</sup>-tagged jets over inclusive jets for radiator energy of  $5 < E_{\rm rad} < 15$  GeV. The D<sup>0</sup>-tagged jets show a significant suppression toward smaller splitting angles. This suppression becomes stronger if the  $k_{\rm T}$  cut is set to higher values, corresponding to a



**Figure 4:** Parallel momentum fraction  $z_{\parallel}$  of charged-particle jets tagged with D<sup>0</sup> mesons (left) and  $\Lambda_c^+$  baryons (right).



**Figure 5:** Substructure variables  $z_g$  (left),  $\theta_g$  (center) and  $n_{SD}$  (right) of D<sup>0</sup>-tagged charged-particle jets compared to inclusive charged-particle jets, in pp collisions at  $\sqrt{s} = 13$  TeV.

cleaner dead-cone signature with less contamination by non-perturbative effects.

The reconstruction of heavy-flavor hadrons within a jet allows for direct access to the fragmentation properties, and also allows for a comparison of meson and baryon fragmentation. The parallel momentum fraction  $z_{\parallel}$  of D<sup>0</sup> mesons and  $\Lambda_c^+$  baryons, shown in Fig. 4 left and right panels respectively, exhibit similar trends in the chosen momentum range. It is to be noted however, that a quantitative description of the observations still poses a challenge to some of the most popular model calculations [9, 13, 14].

The groomed jet substructure of D<sup>0</sup>-tagged jets has been measured for the first time, and compared to that of inclusive jets. Trends in  $z_g$  (Fig. 5 left) and  $\theta_g$  (Fig. 5 center) are slightly different for charmed and inclusive jets, giving a hint about flavor-dependent jet substructure. A more obvious difference is present in the distribution of the number of splittings fulfilling the softdrop condition,  $n_{SD}$  (Fig. 5 right). The fact that charm jets typically have less hard splittings than inclusive jets is consistent with harder heavy-flavor fragmentation caused by mass and color charge effects.

## 4. Summary

In this contribution, recent jet-related results were presented from the ALICE experiment in pp collisions at  $\sqrt{s} = 13$  TeV. Soft-drop groomed substructure measurements of full and charged jets provide an excellent opportunity to test perturbative QCD and hadronization models, besides serving as a baseline for heavy-ion collisions. We also presented the first direct measurement of the dead cone in heavy-flavor jets. Parallel momentum fractions of charmed D<sup>0</sup> mesons and  $\Lambda_c^+$  baryons provide great discrimination power among models on heavy-flavor fragmentation. Charmed jets have been found to typically have less hard splittings than inclusive jets, suggesting a harder fragmentation of heavy than light flavor. The upcoming Run–3 phase of LHC with higher luminosity will allow for high-precision measurements of jets, charmed baryons as well as beauty-jets, further facilitating model developments and moving toward a deeper understanding of the strong interaction [15].

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